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LOGARITHMIC COEFFICIENTS AND GENERALIZED MULTIFRACTALITY OF WHOLE-PLANE SLE

BERTRAND DUPLANTIER⁽¹⁾, XUAN HIEU HO⁽²⁾, THANH BINH LE⁽²⁾,
AND MICHEL ZINSMEISTER⁽²⁾

ABSTRACT. It has been shown that for f an instance of the whole-plane SLE unbounded conformal map from the unit disk \mathbb{D} to the slit plane, the derivative moments $\mathbb{E}(|f'(z)|^p)$ can be written in a closed form for certain values of p depending continuously on the SLE parameter $\kappa \in (0, \infty)$. We generalize this property to the mixed moments, $\mathbb{E}(\frac{|f'(z)|^p}{|f(z)|^q})$, along integrability curves in the moment plane $(p, q) \in \mathbb{R}^2$ depending continuously on κ , by extending the so-called Beliaev–Smirnov equation to this case. The generalization of this integrability property to the m -fold transform of f is also given. We define a novel generalized integral means spectrum, $\beta(p, q; \kappa)$, corresponding to the singular behavior of the above mixed moments. By inversion, it allows a unified description of the unbounded interior and bounded exterior versions of whole-plane SLE, and of their m -fold generalizations. The average generalized spectrum of whole-plane SLE is found to take four possible forms, separated by five phase transition lines in the moment plane \mathbb{R}^2 . The average generalized spectrum of the m -fold whole-plane SLE is directly obtained from the $m = 1$ case by a linear map acting in the moment plane. We also conjecture the precise form of the universal generalized integral means spectrum.

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Key words and phrases. Whole-plane SLE, logarithmic coefficients, Beliaev–Smirnov equation, generalized integral means spectrum, universal spectrum.

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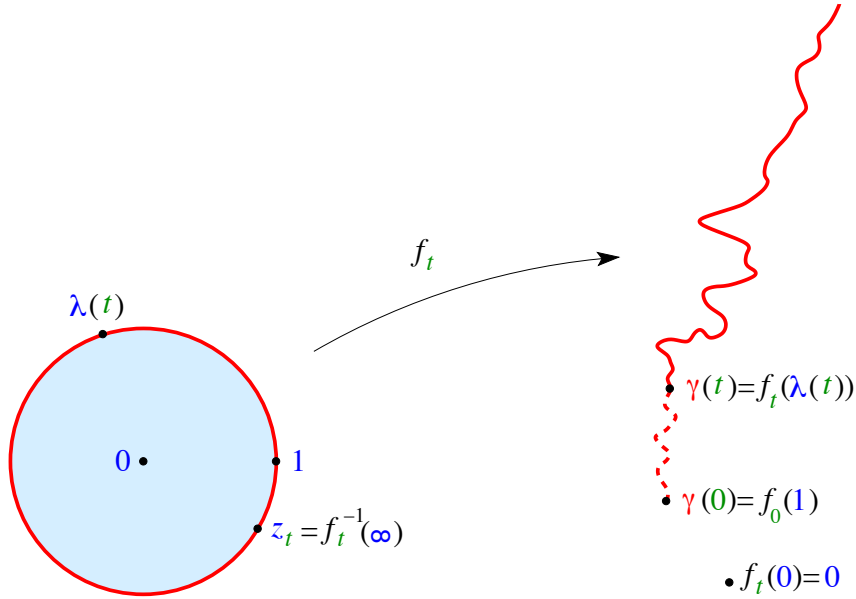


FIGURE 1. Loewner map $z \mapsto f_t(z)$ from \mathbb{D} to the slit domain $\Omega_t = \mathbb{C} \setminus \gamma([t, \infty))$ (here slit by a single curve $\gamma([t, \infty))$) for $\text{SLE}_{\kappa \leq 4}$. One has $f_t(0) = 0, \forall t \geq 0$. At $t = 0$, the driving function $\lambda(0) = 1$, so that the image of $z = 1$ is at the tip $\gamma(0) = f_0(1)$ of the curve. [From Ref. [14]].

1. INTRODUCTION

1.1. **Logarithmic coefficients.** Consider f , a holomorphic function in the unit disk \mathbb{D} ,

$$(1) \quad f(z) = \sum_{n \geq 0} a_n z^n.$$

Bieberbach observed in 1916 [6] that if f is further assumed to be injective, then

$$|a_2| \leq 2|a_1|,$$

and he conjectured that $|a_n| \leq n|a_1|$ for all $n > 2$. This famous conjecture has been proved in 1984 by de Branges [8]. A crucial ingredient of his proof is the theory of growth processes that was developed by Loewner in 1923 [31], precisely in order to solve the $n = 3$ case of the Bieberbach conjecture.

Let $\gamma : [0, \infty) \rightarrow \mathbb{C}$ be a simple curve such that $|\gamma(t)| \rightarrow +\infty$ as $t \rightarrow +\infty$ and such that $\gamma(t) \neq 0, t \geq 0$. Let $\Omega_t = \mathbb{C} \setminus \gamma([t, \infty))$ and $f_t : \mathbb{D} = \{|z| \leq 1\} \rightarrow \Omega_t$ be the Riemann map characterized by $f_t(0) = 0, f_t'(0) > 0$ (See Fig. 1). It is easy to see that $t \mapsto f_t'(0)$ is an increasing continuous function, diverging to $+\infty$ as $t \rightarrow +\infty$. Assuming that $f_0'(0) = 1$, and changing parameterization if necessary, we may set $f_t'(0) = e^t, t \geq 0$. Loewner has shown that f_t satisfies the following

PDE:

$$(2) \quad \frac{\partial}{\partial t} f_t(z) = z \frac{\partial}{\partial z} f_t(z) \frac{\lambda(t) + z}{\lambda(t) - z},$$

where $\lambda : [0, \infty) \rightarrow \partial\mathbb{D}$ is a continuous function on the unit circle. With the sole information that $|\lambda(t)| = 1, \forall t$, he could prove that $|a_3| \leq 3|a_1|$.

Besides Loewner's theory of growth processes, de Branges' proof also heavily relied on the consideration, developed by Grunsky [17] and later Lebedev and Milin [26], of *logarithmic* coefficients. More precisely, if $f : \mathbb{D} \rightarrow \mathbb{C}$ is holomorphic and injective with $f(0) = 0$, we may consider the power series,

$$(3) \quad \log \frac{f(z)}{z} = 2 \sum_{n \geq 1} \gamma_n z^n.$$

The purpose of introducing this logarithm was to prove Robertson's conjecture [35], which was known to imply Bieberbach's. Let f be in the \mathcal{S} class of schlicht functions, i.e., holomorphic and injective in the unit disk, and normalized such that $f(0) = 0, f'(0) = 1$. There is a branch $f^{[2]}$ of $z \mapsto \sqrt{f(z^2)}$ which is an odd function in \mathcal{S} and plays a crucial role in the theory of univalent functions. Let us then write

$$(4) \quad f^{[2]}(z) := z \sqrt{f(z^2)/z^2} = \sum_{n=0}^{\infty} b_{2n+1} z^{2n+1},$$

with $b_1 = 1$. Robertson's conjecture states that:

$$(5) \quad \forall n \geq 0, \sum_{k=0}^n |b_{2k+1}|^2 \leq n + 1.$$

The Lebedev and Milin approach to this conjecture consisted in observing that

$$\log \frac{f^{[2]}(\sqrt{z})}{\sqrt{z}} = \frac{1}{2} \log \frac{f(z)}{z},$$

and consequently that

$$\sum_{n=0}^{\infty} b_{2n+1} z^n = \exp \left(\sum_{n=1}^{\infty} \gamma_n z^n \right).$$

They proved what is now called the second Lebedev-Milin inequality, a combinatorial inequality connecting the coefficients of any power series to those of its exponential, namely

$$(6) \quad \forall n \geq 0, \sum_{k=0}^n |b_{2k+1}|^2 \leq (n + 1) \exp \left(\frac{1}{n + 1} \sum_{m=1}^n \sum_{k=1}^m \left(k |\gamma_k|^2 - \frac{1}{k} \right) \right).$$

This naturally led Milin [33] to conjecture that

$$(7) \quad \forall f \in \mathcal{S}, \forall n \geq 1, \sum_{m=1}^n \sum_{k=1}^m \left(k |\gamma_k|^2 - \frac{1}{k} \right) \leq 0;$$

this conjecture, proved by de Branges in 1984, implies Robertson's, hence Bieberbach's conjecture.

Returning to Loewner's theory, his derivation of Eq. (2) above is only half of the story. There is indeed a converse: given any continuous function $\lambda : [0, +\infty[\rightarrow \mathbb{C}$ with $|\lambda(t)| = 1$ for $t \geq 0$, the the Loewner equation (2), supplemented by the boundary ("initial") condition, $\lim_{t \rightarrow +\infty} f_t(e^{-t}z) = z$, has a solution $(t, z) \mapsto f_t(z)$, such that $(f_t(z))_{t \geq 0}$ is a chain of Riemann maps onto simply connected domains (Ω_t) that are increasing with t .

In 1999, Schramm [38] introduced into the Loewner equation the *random* driving function,

$$(8) \quad \lambda(t) := \sqrt{\kappa} B_t,$$

where B_t is standard one dimensional Brownian motion and κ a non-negative parameter, thereby making Eq. (2) a stochastic PDE, and creating the celebrated *Schramm-Loewner Evolution* SLE_κ .

The associated conformal maps f_t from \mathbb{D} to $\mathbb{C} \setminus \gamma([t, \infty)$, obeying (2) for (8), define the *interior whole-plane* Schramm-Loewner evolution. Their coefficients $a_n(t)$, which are random variables, are defined by a normalized series expansion, as described in the following proposition [14].

Proposition 1.1. *Let $(f_t(z))_{t \geq 0}, z \in \mathbb{D}$, be the interior Schramm-Loewner whole-plane process driven by $\lambda(t) = e^{i\sqrt{\kappa}B_t}$ in Eq. (2). We write*

$$(9) \quad f_t(z) = e^t \left(z + \sum_{n \geq 2} a_n(t) z^n \right).$$

and for its logarithm,

$$(10) \quad \log \frac{e^{-t} f_t(z)}{z} = 2 \sum_{n \geq 1} \gamma_n(t) z^n.$$

Then the conjugate whole-plane Schramm-Loewner evolution $e^{-i\sqrt{\kappa}B_t} f_t(e^{i\sqrt{\kappa}B_t} z)$ has the same law as $f_0(z)$, hence $e^{i(n-1)\sqrt{\kappa}B_t} a_n(t) \stackrel{(\text{law})}{=} a_n(0)$. From this and Eqs. (9), (10), follows the identity $e^{in\sqrt{\kappa}B_t} \gamma_n(t) \stackrel{(\text{law})}{=} \gamma_n(0)$. In the sequel, we set: $a_n := a_n(0)$ and $\gamma_n := \gamma_n(0)$.

The starting point of the present article is the observation, made in Ref. [13], that the SLE_κ process, in its *interior* whole-plane version, has a rich algebraic structure, giving rise to a host of (integrability-like) closed form results. The first hint was the fact that, beyond the coefficient expectations $\mathbb{E}(a_n)$ for Eq. (9), the coefficient squared moments, $\mathbb{E}(|a_n|^2)$, have very simple expressions for

specific values of κ . This has been developed in detail in Refs. [14] and [27] (see also Refs. [28, 29, 30]), by using a PDE obeyed by the derivative moments $\mathbb{E}(|f'(z)|^p)$. Following the work of Rohde–Schramm [36], it was originally derived by Beliaev–Smirnov in Ref. [3] to study the average integral means spectrum of the *exterior* version of the whole-plane SLE_κ map. Note also that similar ideas already appeared in Ref. [21], where A. Kemppainen studied in detail the coefficients associated with the Schramm–Loewner evolution, using a stationarity property of SLE [22]. However, the focus there was on expectations of the moments of those coefficients, rather than on the moments of their moduli.

Here, we study in particular the logarithmic coefficients (10) of whole-plane SLE and the generalizations thereof. The main idea of this work is the introduction for the (unbounded) whole-plane SLE map f of *mixed moments*, $\mathbb{E}(|f'(z)|^p/|f(z)|^q)$, for any $(p, q) \in \mathbb{R}^2$, which are found to obey a two-parameter family of Beliaev–Smirnov-type equations. We also define and study in detail the *generalized integral means spectrum* associated with these mixed moments, which is a priori a function of p and q .

We argue that the general approach proposed here is the natural one for whole-plane SLE. It unifies the earlier SLE integral means studies of Refs. [2, 3, 14, 27, 28, 29] into a much broader framework, which in particular allows one to fully exploit the *inversion symmetry* between the unbounded inner and bounded outer versions, and also to cover the $p = q$ logarithmic case, as well as the integral means of the map itself or of its m -fold transforms.

As we shall see, the (p, q) -plane for whole-plane SLE is structured by *integrable probability lines*, as well as by *phase-transition lines*. The integrability lines yield a continuum of closed forms for mixed moments, that generalize results in Refs. [14, 27, 28, 29]. The phase transition lines mark the breakdown of the standard analysis of B–S-type equations initiated in Ref. [3], as already seen in the case of the standard integral means spectrum in Refs. [2, 14, 27, 28, 29]. In particular, a certain phase transition *point* here manifests the appearance, for the bounded exterior whole-plane SLE, of a subjacent spectrum, thus requiring a novel proof as given in Ref. [2]. To study the phase structure of the whole-plane generalized spectrum, we unify and extend here the non-standard methods initiated in Refs. [2, 14]. Moreover, this field leads to further interesting open questions, which apparently exceed the power of the various methods proposed so far.

1.2. Main results. A first motivation of this article is the proof, originally obtained for small n by the third author [25], of the following.

Theorem 1.2. *Let $f(z) := f_0(z)$ be the time 0 unbounded whole-plane SLE_κ map, in the same setting as in Proposition 1.1, such that*

$$\log \frac{f(z)}{z} = 2 \sum_{n \geq 1} \gamma_n z^n;$$

then, for $\kappa = 2$, which corresponds to the scaling limit of the Loop-Erased Random Walk [24, 38],

$$\begin{aligned}\mathbb{E}(\gamma_1) &= -1/2, \quad \mathbb{E}(\gamma_n) = 0, \quad n \geq 2, \\ \mathbb{E}(|\gamma_n|^2) &= \frac{1}{2n^2}, \quad n \geq 1, \\ \mathbb{E}(\gamma_n \bar{\gamma}_{n+1}) &= -\frac{1}{4n(n+1)}, \quad \mathbb{E}(\gamma_n \bar{\gamma}_{n+k}) = 0, \quad n \geq 1, k \geq 2.\end{aligned}$$

Let us then briefly return to the Lebedev-Milin theory. By Theorem 1.2, we have for SLE_2 ,

$$\mathbb{E} \left(\sum_{m=1}^n \sum_{k=1}^m \left(k |\gamma_k|^2 - \frac{1}{k} \right) \right) = -\frac{1}{2} \sum_{m=1}^n \sum_{k=1}^m \frac{1}{k} = -\frac{n+1}{2} \sum_{k=2}^{n+1} \frac{1}{k},$$

which gives an example of the validity ‘‘in expectation’’ of the Milin conjecture. Recalling Definition (4), we also get, in expectation, a check of Robertson’s conjecture (5):

$$\mathbb{E} \left(\log \frac{1}{n+1} \sum_{k=0}^n |b_{2k+1}|^2 \right) \leq -\frac{1}{2} \sum_{k=2}^{n+1} \frac{1}{k}.$$

The idea behind the proof of Theorem 1.2 is to differentiate (3),

$$z \frac{d}{dz} \log \frac{f(z)}{z} = z \frac{f'(z)}{f(z)} - 1 = 2 \sum_{n \geq 1} n \gamma_n z^n,$$

such that

$$(11) \quad \left| z \frac{f'(z)}{f(z)} \right|^2 = 1 + 2 \sum_{n \geq 1} n \gamma_n (z^n + \bar{z}^n) + 4 \sum_{n \geq 1} \sum_{m \geq 1} n m \gamma_n \bar{\gamma}_m z^n \bar{z}^m,$$

and to compute $\mathbb{E} \left(\left| z \frac{f'(z)}{f(z)} \right|^2 \right)$. We indeed prove:

Theorem 1.3. *Let f be the interior whole-plane SLE_κ map, in the same setting as in Theorem 1.2; then for $\kappa = 2$,*

$$\mathbb{E} \left(\left| z \frac{f'(z)}{f(z)} \right|^2 \right) = \frac{(1-z)(1-\bar{z})}{1-z\bar{z}}.$$

This gives

$$\mathbb{E} \left(\left| z \frac{f'(z)}{f(z)} \right|^2 \right) = 1 - \sum_{n \geq 0} (z^{n+1} \bar{z}^n + z^n \bar{z}^{n+1}) + 2 \sum_{n \geq 1} z^n \bar{z}^n,$$

and Theorem 1.2 follows by direct identification to Eq. (11).

Theorem 1.3 is actually a consequence of Theorems 3.1 and 4.2 of Sections 3 and 4 below, which give expressions in closed form for the mixed moments,

$$(12) \quad (a) \quad \mathbb{E} \left(\frac{(f'(z))^{p/2}}{(f(z))^{q/2}} \right); \quad (b) \quad \mathbb{E} \left(\frac{|f'(z)|^p}{|f(z)|^q} \right),$$

along an integrability curve \mathcal{R} , which is a *parabola* in the (p, q) plane depending on the SLE parameter κ . In fact, we establish a general integrability result along the \mathcal{R} parabola for the SLE *two-point function*:

$$G(z_1, \bar{z}_2) := \mathbb{E} \left(z_1^{\frac{q}{2}} \frac{(f'(z_1))^{\frac{p}{2}}}{(f(z_1))^{\frac{q}{2}}} \overline{\left[z_2^{\frac{q}{2}} \frac{(f'(z_2))^{\frac{p}{2}}}{(f(z_2))^{\frac{q}{2}}} \right]} \right).$$

The mixed moments (12) can also be seen respectively as the value $G(z, 0)$ of this SLE two-point function at $(z_1 = z, z_2 \rightarrow 0)$ for (a), and the value $G(z, \bar{z})$ at coinciding points, $z = z_1 = z_2$, for (b). These integrability theorems, which provide full generalizations of the results of Refs. [14] and [27], give rise to a host of new algebraic identities concerning the whole-plane SLE_κ random map.

These integrability results are generalized in Section 4.4 to the so-called m -fold symmetric transforms of the whole-plane SLE map [14], $f^{[m]} : z \mapsto \sqrt[m]{f(z^m)}$, $m \in \mathbb{N} \setminus \{0\}$, defined for $f \in \mathcal{S}$ as the holomorphic branch whose derivative is equal to 1 at 0. These are functions in \mathcal{S} whose Taylor series are of the form $f(z) = \sum_{k \geq 0} a_{mk+1} z^{mk+1}$, the $m = 2$ case corresponding to odd functions. To extend the definition to negative m , the m -fold transform of the outer map is conjugate by inversion of the $(-m)$ -fold transform of the inner map, $f^{[m]}(\zeta) = 1/f^{[-m]}(1/\zeta)$, for $m \in \mathbb{Z} \setminus \mathbb{N}$ and $\zeta \in \mathbb{D}_- := \mathbb{C} \setminus \bar{\mathbb{D}}$. In particular, the map, $\zeta \in \mathbb{D}_- \mapsto f^{[-1]}(\zeta) = 1/f(1/\zeta)$, is just the exterior whole-plane map considered by Beliaev and Smirnov in Ref. [3], from \mathbb{D}_- to the slit plane, a domain with bounded boundary. Note similarly that for all $m < 0$, $f^{[m]}(\mathbb{D}_-)$ has bounded boundary. Interestingly enough, a linear map in the (p, q) -moment plane allows one to directly relate the generalized mixed moments of the m -fold map to those of f , thus yielding the integrability Theorem 4.6.

From the consideration in Eq. (12) of the mixed moments (b) of moduli of f' and f for whole-plane SLE, we are led to introduce the following definition.

Definition 1.4. For a function $f \in \mathcal{S}$, the (*average*) *generalized integral means spectrum* $\beta_f(p, q)$, depending on p and q , is defined as

$$\beta_f(p, q) := \limsup_{r \rightarrow 1^-} \frac{\log \int_{r\partial\mathbb{D}} \mathbb{E} \left(\frac{|f'(z)|^p}{|f(z)|^q} \right) |dz|}{\log \left(\frac{1}{1-r} \right)}.$$

In the case where the limit exists, this can be written as,

$$\int_{r\partial\mathbb{D}} \mathbb{E} \left(\frac{|f'(z)|^p}{|f(z)|^q} \right) |dz| \stackrel{(r \rightarrow 1^-)}{\asymp} (1-r)^{-\beta_f(p, q)},$$

in the sense of the equivalence of the logarithms of both terms.

Remark 1.5. *The $q = 0$ case yields the standard spectrum $\beta_f(p) := \beta_f(p, 0)$, which is defined for any univalent function. For a map $f \in \mathcal{S}$, the introduction of the q -parameter is aimed at the analysis of the behavior of f at infinity. Indeed, when f is bounded, the generalized integral means spectrum does not depend on q and coincides with the standard one.*

Remark 1.6. *Exterior-Interior Duality. For $\hat{f} := f^{[-1]}$ and $0 < r < 1$, we have*

$$(13) \quad \int_{r^{-1}\partial\mathbb{D}} \mathbb{E} \left(\frac{|\hat{f}'(\zeta)|^p}{|\hat{f}(\zeta)|^q} \right) |d\zeta| = r^{2p-2} \int_{r\partial\mathbb{D}} \mathbb{E} \left(\frac{|f'(z)|^p}{|f(z)|^{q'}} \right) |dz|,$$

for $q + q' = 2p$. In particular, the $(p, 0)$ standard integral means for the exterior [2, 3] or interior [14, 28, 29] whole-plane maps correspond to the $(p, 2p)$ generalized integral means of their respective inverted maps.

The introduction of this generalized spectrum in the whole (p, q) -plane also allows for a unified description of the *standard* spectra $\beta_{f^{[m]}}(p)$ of the whole collection of m -fold transforms of a given map f , in terms of its generalized spectrum $\beta_f(p, q)$. We indeed have the identity (Section 5.5):

$$(14) \quad \beta_{f^{[m]}}(p) = \beta_f(p, (1 - 1/m)p), \quad m \in \mathbb{Z} \setminus \{0\}.$$

In particular, for $m = -1$, the ray $q = 2p$ describes the standard integral means spectrum in the exterior (bounded) case, as studied in Refs. [2, 3].

In Section 5, we study this generalized spectrum, $\beta(p, q; \kappa)$, for the interior whole-plane SLE $_{\kappa}$ and $(p, q) \in \mathbb{R}^2$. We show that it takes four possible forms, $\beta_0(p)$, $\beta_{\text{tip}}(p)$, $\beta_{\text{lin}}(p)$ and $\beta_1(p, q)$. The first three spectra are independent of q , and are respectively given by the bulk, the tip and the linear SLE average spectra appearing in the work by Beliaev and Smirnov [3], recently revisited and partially corrected in Ref. [2] by Beliaev and two of the present authors. The bulk case corresponds to the earlier harmonic measure multifractal spectrum derived by the first author in Refs. [9, 10] (see also [7, 11, 12] and [4, 5, 37]), and recently established in an almost sure sense by Gwynne, Miller and Sun [18]; the tip harmonic measure spectrum was predicted by Hastings [19] and proved in an almost sure sense by Johansson Viklund and Lawler [20]. The fourth spectrum, $\beta_1(p, q)$, is the extension to non-vanishing q of a novel integral means spectrum, which was discovered and studied in Refs. [14] and [27, 29], and which is due to the unboundedness of whole-plane SLE. As shown in Ref. [14], this spectrum is also closely related to the (non-standard) SLE tip exponents obtained by quantum gravity techniques in Ref. [10], and to the so-called radial SLE derivative exponents of Ref. [23]. As shown in Ref. [2] for the exterior bounded case, the corresponding spectrum, $\beta_1(p, 2p)$, due to the SLE seed or ‘second tip’, is subjacent in between the bulk and the tip spectra.

Five different phase transition lines partition the (p, q) -plane into four domains, where the generalized integral means spectrum takes four different forms, as given by the following theorem.

Theorem 1.7. *Define the functions:*

$$\begin{aligned}\beta_{\text{tip}}(p; \kappa) &:= -p - 1 + \frac{1}{4}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa p}), \\ \beta_0(p; \kappa) &:= -p + \frac{4 + \kappa}{4\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa p}), \\ \beta_{\text{lin}}(p; \kappa) &:= p - \frac{(4 + \kappa)^2}{16\kappa}, \\ \beta_1(p, q; \kappa) &:= 3p - 2q - \frac{1}{2} - \frac{1}{2}\sqrt{1 + 2\kappa(p - q)}.\end{aligned}$$

Consider the ‘green’ parabola \mathcal{G} defined in the (p, q) -plane by the parametric equations,

$$\begin{aligned}p = p_{\mathcal{G}}(\gamma) &:= \frac{(4 + \kappa)^2}{8\kappa} - \frac{\kappa}{2}\gamma^2, \\ q = q_{\mathcal{G}}(\gamma) &:= \frac{(4 + \kappa)^2}{8\kappa} + \gamma - \kappa\gamma^2, \quad \gamma \in \mathbb{R},\end{aligned}$$

and the ‘blue’ quartic \mathcal{Q} ,

$$\begin{aligned}p = p_{\mathcal{Q}}(\gamma) &:= \frac{\kappa}{16} + \left(1 + \frac{\kappa}{4}\right)\gamma - \frac{\kappa}{2}\gamma^2 - \frac{1}{8}\Delta^{\frac{1}{2}}(\gamma), \\ q = q_{\mathcal{Q}}(\gamma) &:= p_{\mathcal{Q}}(\gamma) + \gamma - \frac{\kappa}{2}\gamma^2, \\ \Delta(\gamma) &:= 4\kappa^2\gamma^2 - 2\kappa(4 + \kappa)\gamma + \frac{1}{4}(8 + \kappa)^2 + 4\kappa, \quad \gamma \in \mathbb{R}.\end{aligned}$$

Phase transition lines for the generalized integral means spectrum $\beta(p, q)$ of whole-plane SLE_{κ} are in the (p, q) -plane (Fig. 2):

- (i) the vertical half-line D_0 above $P_0 = (p_0, q_0)$, with $p_0 = 3(4 + \kappa)^2/32\kappa$, $q_0 = (4 + \kappa)(8 + \kappa)/16\kappa$, where $\beta_0(p; \kappa) = \beta_{\text{lin}}(p; \kappa)$;
- (ii) the unit slope half-line D_1 originating at P_0 , whose equation is $q - p = (16 - \kappa^2)/32\kappa$ with $p \geq p_0$, and where $\beta_{\text{lin}}(p; \kappa) = \beta_1(p, q; \kappa)$;
- (iii) the section of green parabola \mathcal{G} , with parametric coordinates $(p_{\mathcal{G}}(\gamma), q_{\mathcal{G}}(\gamma))$ for $\gamma \in [1/4 + 1/\kappa, 1 + 2/\kappa]$, between P_0 and $Q_0 = (p'_0, q'_0)$, with $p'_0 = -1 - 3\kappa/8$, $q'_0 = -2 - 7\kappa/8$, and where $\beta_0(p; \kappa) = \beta_1(p, q; \kappa)$;
- (iv) the vertical half-line D'_0 above point Q_0 , where $\beta_{\text{tip}}(p; \kappa) = \beta_0(p; \kappa)$;
- (v) the branch of the blue quartic \mathcal{Q} from Q_0 to ∞ , with parametric coordinates $(p_{\mathcal{Q}}(\gamma), q_{\mathcal{Q}}(\gamma))$ for $\gamma \in [1 + 2/\kappa, +\infty)$, where $\beta_{\text{tip}}(p; \kappa) = \beta_1(p, q; \kappa)$.

The analytic form of $\beta(p, q; \kappa)$ changes across lines (i), (ii), (iv) and (v) according to Figure 2; in case (ii), $\beta(p, q; \kappa) = \beta_{\text{lin}}(p; \kappa)$ above D_1 , whereas $\beta(p, q; \kappa) \leq \beta_1(p, q; \kappa)$ below D_1 .

Remark 1.8. *The proof of Theorem 1.7, as given in Section 5.4, establishes the analytic form of $\beta(p, q; \kappa)$ in an upper semi-infinite domain, whose interior*

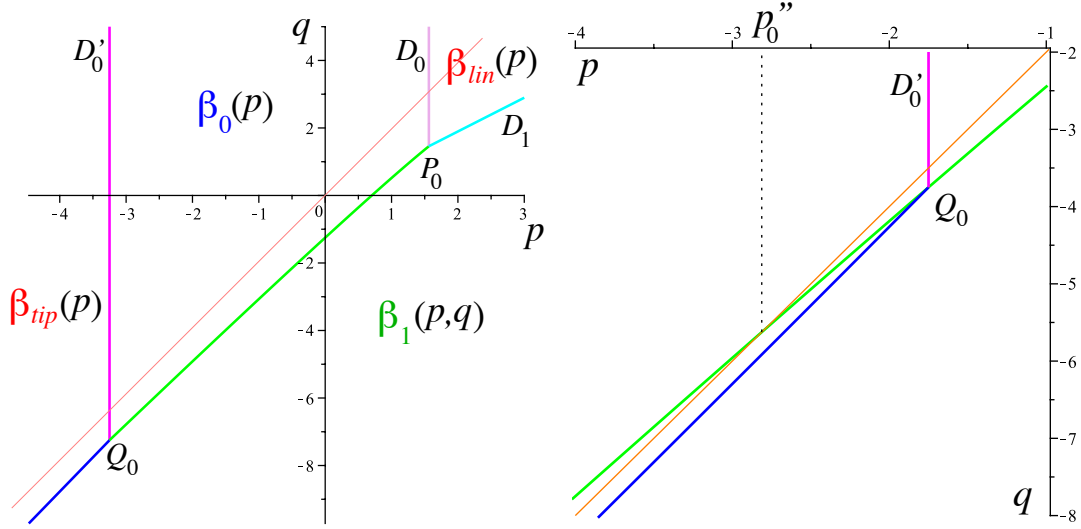


FIGURE 2. *Left: Phase diagram for the SLE generalized integral means spectrum $\beta(p, q)$ illustrating Theorem 1.7. The $q = 2p$ (coral) straight line yields the standard spectrum of the exterior version of whole-plane SLE, as studied in Refs. [2, 3]. Right: [Zoom below Q_0] This line, being asymptotically parallel to the blue quartic \mathcal{Q} , does not intersect the lower domain where the form β_1 holds, but intersects the green parabola \mathcal{G} below point Q_0 at abscissa p_0'' .*

strictly includes D_0 , D_0' and the branches of quartic and parabola of the theorem's statement, and whose frontier includes D_1 .

Remark 1.9. *The fact that the line $q = 2p$ intersects the extension of the green parabola \mathcal{G} below Q_0 at $p_0'' = -\frac{1}{128}(4+\kappa)^2(8+\kappa)$ (Fig. 2, right) signals the presence of an underlying phase transition from the bulk spectrum to the β_1 spectrum. As recently shown in Ref. [2], the original proof of Ref. [3] breaks down there and requires for $p \leq p_0''$ a novel argument.*

The generalization of this four-domain structure to the generalized integral means spectrum of m -fold transforms, $f^{[m]}$, $m \in \mathbb{N} \setminus \{0\}$, is obtained in a straightforward way from the above mentioned linear map in (p, q) co-ordinates (Theorem 5.20). The structure obtained appears so robust that the *universal* generalized spectrum $B(p, q)$, i.e., the maximum of $\beta_f(p, q)$ over all univalent functions $f \in \mathcal{S}$, shows a similar partition of the mixed moment plane. In Section 6, we give a precise result for its four forms, which incorporates known results and conjectures on the standard universal spectra for univalent, not necessarily bounded, functions [16, 32, 34].

1.3. Synopsis. This article is organized as follows. Section 2 deals with logarithmic derivatives. It sets up the martingale techniques needed for dealing with

mixed moments. Section 3 uses them for the study of the complex one-point function (a) in (12), which is shown to obey a simple differential equation in complex variable z . This leads to Theorem 3.1, which establishes a closed form for this function along the integrability parabola \mathcal{R} in the (p, q) -plane. Section 4 is concerned with the moduli one-point function (b) in (12), and more generally, with the SLE two-point function $G(z_1, \bar{z}_2)$. A PDE in (z_1, \bar{z}_2) is derived for $G(z_1, \bar{z}_2)$, which yields a proof of Theorem 4.2 establishing closed form expressions for G for all $(p, q) \in \mathcal{R}$. Section 4.4 deals with the generalization of the previous integrability results to the m -fold symmetric transforms $f^{[m]}$, $m \in \mathbb{Z} \setminus \{0\}$, of the whole-plane SLE map f . Section 5 is devoted to the study of the averaged generalized spectrum $\beta(p, q; \kappa)$ of the whole-plane SLE_κ random map f , as well as to the averaged generalized integral means spectrum $\beta^{[m]}(p, q; \kappa)$ of the m -fold transform $f^{[m]}$ for $m \in \mathbb{N} \setminus \{0\}$. Of particular interest are the five phase transition lines separating the four different analytic expressions of β (or $\beta^{[m]}$) in the moment plane, given in Theorems 1.7 and 5.20. In the final Section 6, we give a full description of the expected form for the universal generalized integral means spectrum, $B(p, q)$, in terms of known or conjectured results on the standard universal spectrum for univalent functions.

Acknowledgments. It is a pleasure to thank Kari Astala for extended discussions about the universal generalized integral means spectrum.

2. EXPECTATIONS OF LOGARITHMIC DERIVATIVES

In this section, we prove the following simple result:

Proposition 2.1. *Let $f(z) = f_0(z)$ be the interior whole-plane SLE_2 map at time 0, in the same setting as in Proposition 1.1; we then have*

$$\mathbb{E} \left(z \frac{f'(z)}{f(z)} \right) = 1 - z.$$

The method explained here will allow us to swiftly move to more complicated cases in the next sections. Let us then introduce

$$(15) \quad G(z) := \mathbb{E} \left(z \frac{f'(z)}{f(z)} \right),$$

and, following Refs. [36] and [3], aim at finding a partial differential equation satisfied by G . For the benefit of the reader not familiar with Refs. [3, 36], let us detail the strategy of these papers that we will apply in various contexts here.

The starting point is to consider the radial SLE_κ , solution to the ODE,

$$\partial_t g_t(z) = g_t(z) \frac{\lambda(t) + g_t(z)}{\lambda(t) - g_t(z)}, \quad z \in \mathbb{D},$$

with the initial condition $g_0(z) = z$, and where $\lambda(t) = e^{i\sqrt{\kappa}B_t}$. The map g_t conformally maps a subdomain of the unit disk onto the latter. As we shall see

shortly, the whole-plane map f is rather related to the map g_t^{-1} , but this last function satisfies, by Loewner's theory, a PDE not well-suited to Itô calculus. To overcome this difficulty, one runs backward the ODE of radial SLE, i.e., one compares g_t^{-1} to g_{-t} . This is the purpose of Lemma 1 in [3] (an analog of Lemma 3.1 in Ref. [36]), which states that, for $t \in \mathbb{R}$, $g_{-t}(z)$ has the same law as the process $\tilde{f}_t(z)$, defined as follows.

Definition 2.2. The (conjugate, inverse) radial SLE process \tilde{f}_t is defined, for $t \in \mathbb{R}$, as

$$(16) \quad \tilde{f}_t(z) := g_t^{-1}(z\lambda(t))/\lambda(t).$$

The lemma then results from the simple observation that

$$\tilde{f}_s(z) = \hat{g}_{-s}(z),$$

where, for fixed $s \in \mathbb{R}$, the new process $\hat{g}_t(z) := g_{s+t} \circ g_s^{-1}(z\lambda(s))/\lambda(s)$ can be shown to be a radial SLE. This lemma implies in particular that \tilde{f}_t is solution to the ODE:

$$(17) \quad \partial_t \tilde{f}_t(z) = \tilde{f}_t(z) \frac{\tilde{f}_t'(z) + \lambda(t)}{\tilde{f}_t(z) - \lambda(t)}, \quad \tilde{f}_0(z) = z.$$

To apply Itô's stochastic calculus, one then uses Lemma 2 in Ref. [3], which is a version of the SLE's Markov property,

$$\tilde{f}_t(z) = \lambda(s) \tilde{f}_{t-s}(\tilde{f}_s(z)/\lambda(s)).$$

To finish, one has to relate the whole-plane SLE to the (modified) radial one. This is done through Lemma 3 in [3], which is in our present setting (with a change of an e^{-t} convergence factor there to an e^t factor here, when passing from the exterior to the interior of the unit disk \mathbb{D}):

Lemma 2.3. *The limit in law, $\lim_{t \rightarrow +\infty} e^t \tilde{f}_t(z)$, exists, and has the same law as the (time zero) interior whole-plane random map $f_0(z)$:*

$$\lim_{t \rightarrow +\infty} e^t \tilde{f}_t(z) \stackrel{(\text{law})}{=} f_0(z).$$

Let us now turn to the proof of Proposition 2.1.

Proof. Let us introduce the auxiliary, time-dependent, radial variant of the SLE one-point function $G(z)$ (15) above,

$$(18) \quad \tilde{G}(z, t) := \mathbb{E} \left(z \frac{\tilde{f}_t'(z)}{\tilde{f}_t(z)} \right),$$

where \tilde{f}_t is a modified radial SLE map at time t as in Definition 2.2. Owing to Lemma (2.3), we have

$$(19) \quad \lim_{t \rightarrow +\infty} \tilde{G}(z, t) = G(z).$$

We then use a martingale technique to obtain an equation satisfied by $\tilde{G}(z, t)$. For $s \leq t$, define $\mathcal{M}_s := \mathbb{E} \left(\frac{\tilde{f}'_t(z)}{\tilde{f}_t(z)} | \mathcal{F}_s \right)$, where \mathcal{F}_s is the σ -algebra generated by $\{B_u, u \leq s\}$. $(\mathcal{M}_s)_{s \geq 0}$ is by construction a martingale. Because of the Markov property of SLE, we have [3]

$$\begin{aligned} \mathcal{M}_s &= \mathbb{E} \left(\frac{\tilde{f}'_t(z)}{\tilde{f}_t(z)} | \mathcal{F}_s \right) = \mathbb{E} \left(\frac{\tilde{f}'_s(z) \tilde{f}'_{t-s}(\tilde{f}_s(z)/\lambda(s))}{\lambda(s) \tilde{f}_{t-s}(\tilde{f}_s(z)/\lambda(s))} | \mathcal{F}_s \right) \\ &= \frac{\tilde{f}'_s(z)}{\lambda(s)} \mathbb{E} \left(\frac{\tilde{f}'_{t-s}(\tilde{f}_s(z)/\lambda(s))}{\tilde{f}_{t-s}(\tilde{f}_s(z)/\lambda(s))} | \mathcal{F}_s \right) \\ &= \frac{\tilde{f}'_s(z)}{\tilde{f}_s(z)} \tilde{G}(z_s, \tau), \end{aligned}$$

where $z_s := \tilde{f}_s(z)/\lambda(s)$, and $\tau := t - s$.

We have from Eq. (17)

$$\begin{aligned} (20) \quad \partial_s \log \tilde{f}'_s &= \frac{\partial_z \left[\tilde{f}'_s \frac{\tilde{f}_s + \lambda(s)}{\tilde{f}_s - \lambda(s)} \right]}{\tilde{f}'_s} = \frac{\tilde{f}_s + \lambda(s)}{\tilde{f}_s - \lambda(s)} - \frac{2\lambda(s)\tilde{f}_s}{(\tilde{f}_s - \lambda(s))^2} \\ &= 1 - \frac{2}{(1 - z_s)^2}, \end{aligned}$$

$$(21) \quad \partial_s \log \tilde{f}_s = \frac{\partial_s \tilde{f}_s}{\tilde{f}_s} = \frac{z_s + 1}{z_s - 1},$$

$$(22) \quad dz_s = z_s \left[\frac{z_s + 1}{z_s - 1} - \frac{\kappa}{2} \right] ds - iz_s \sqrt{\kappa} dB_s.$$

The coefficient of the ds -drift term of the Itô derivative of \mathcal{M}_s is obtained from the above as,

$$(23) \quad \frac{\tilde{f}'_s(z)}{\tilde{f}_s(z)} \left[-\frac{2z_s}{(1 - z_s)^2} + z_s \left(\frac{z_s + 1}{z_s - 1} - \frac{\kappa}{2} \right) \partial_z - \partial_\tau - \frac{\kappa}{2} z_s^2 \partial_z^2 \right] \tilde{G}(z_s, \tau),$$

and vanishes by the (local) martingale property. Because \tilde{f}_s is univalent, \tilde{f}'_s does not vanish in \mathbb{D} , therefore the bracket above vanishes.

Owing to the existence of the limit (19), we now take the limit as $\tau \rightarrow +\infty$ in the above, and obtain the ODE,

$$\begin{aligned} (24) \quad \mathcal{P}(\partial)[G(z)] &:= -\frac{2z}{(1 - z)^2} G(z) + z \left(\frac{z + 1}{z - 1} - \frac{\kappa}{2} \right) G'(z) - \frac{\kappa}{2} z^2 G''(z) \\ &= \left[-\frac{2z}{(1 - z)^2} + z \left(\frac{z + 1}{z - 1} \right) \partial_z - \frac{\kappa}{2} (z \partial_z)^2 \right] G(z) = 0. \end{aligned}$$

In the above, the exchange of the $\tau \rightarrow +\infty$ limit and of partial derivation of $\tilde{G}(z, \tau)$ with respect to z is justified by the fact that the τ -family $e^\tau \tilde{f}_\tau(z)$ in Lemma (2.3) and all its z -derivatives are normal, i.e., uniformly bounded in any

compact of \mathbb{D} , so that the spatial derivatives of $\tilde{G}(z, \tau)$ form an equicontinuous family. A further requirement is that $\lim_{\tau \rightarrow +\infty} \frac{\partial}{\partial \tau} \tilde{G} = 0$. Use of the Schramm-Loewner equation (17) for \tilde{f}_τ shows that

$$\frac{\partial}{\partial \tau} \left(\frac{\tilde{f}'_\tau(z)}{\tilde{f}_\tau(z)} \right) = -\frac{2\tilde{f}'_\tau(z)}{(\tilde{f}_\tau(z) - \lambda(\tau))^2}.$$

Since $z \mapsto e^\tau \tilde{f}_\tau(z)$ belongs to the \mathcal{S} class, classical Koebe distortion theorems then show that the right-hand side is bounded by $C(z)e^{-\tau}$, with C defined in \mathbb{D} ; this insures both the validity of the exchange of expectation and τ -derivation, and the vanishing of the limit above.

Following Ref. [14], we now look for solutions to Eq. (24) of the form $\varphi_\alpha(z) := (1 - z)^\alpha$. We have

$$\mathcal{P}(\partial)[\varphi_\alpha] = A(2, 2, \alpha)\varphi_\alpha + B(2, \alpha)\varphi_{\alpha-1} + C(2, \alpha)\varphi_{\alpha-2},$$

where, in anticipation of the notation that will be introduced in Section 3 below,

$$\begin{aligned} A(2, 2, \alpha) &:= \alpha - \frac{\kappa}{2}\alpha^2, \\ B(2, \alpha) &:= 2 - \left(3 + \frac{\kappa}{2}\right)\alpha + \kappa\alpha^2, \\ C(2, \alpha) &:= -2 + \left(2 + \frac{\kappa}{2}\right)\alpha - \frac{\kappa}{2}\alpha^2, \end{aligned}$$

with, identically, $A + B + C = 0$. The linear independence of $\varphi_\alpha, \varphi_{\alpha-1}, \varphi_{\alpha-2}$ thus shows that $\mathcal{P}(\partial)[\varphi_\alpha] = 0$ is equivalent to $A = B = C = 0$, which yields $\kappa = 2, \alpha = 1$, and $1 - z$ is a solution. That it must be G follows from two facts. First, G is holomorphic, since it is clearly continuous by Lebesgue theorem and holomorphicity then follows from Morera's theorem. Second, the space of formal power series that are solutions to Eq. (24) is one-dimensional. This is analogous to Lemma 3.1 in Ref. [14], and follows here from the fact that $\mathcal{P}(\partial)(z^n) = -(\frac{\kappa}{2}n^2 + n)z^n + O(z^{n+1})$. \square

3. SLE ONE-POINT FUNCTION

Let us now turn to the natural generalization of Proposition 2.1.

Theorem 3.1. *Let $f(z) = f_0(z)$ be the interior whole-plane SLE_κ map at time zero, in the same setting as in Proposition 1.1. Consider the curve \mathcal{R} , defined parametrically by*

$$(25) \quad p = -\frac{\kappa}{2}\gamma^2 + \left(2 + \frac{\kappa}{2}\right)\gamma, \quad 2p - q = \left(1 + \frac{\kappa}{2}\right)\gamma, \quad \gamma \in \mathbb{R}.$$

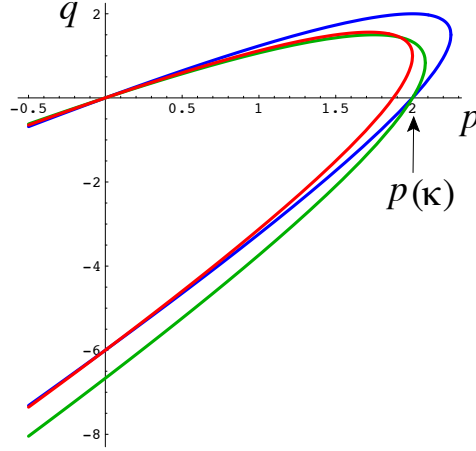


FIGURE 3. Integral curves \mathcal{R} of Theorem 3.1, for $\kappa = 2$ (blue), $\kappa = 4$ (red), and $\kappa = 6$ (green). In addition to the origin, the $q = 0$ intersection point with the p -axis is at $p(\kappa) := (6 + \kappa)(2 + \kappa)/8\kappa$, with $p(2) = p(6) = 2$ [14, 27].

On \mathcal{R} , the whole-plane SLE_κ one-point function has the integrable form,

$$\mathbb{E}\left(\frac{(f'(z))^{\frac{p}{2}}}{(f(z)/z)^{\frac{q}{2}}}\right) = (1-z)^\gamma.$$

Remark 3.2. Eq. (25) describes a parabola in the (p, q) plane (see (Fig. 3)), which is given in Cartesian coordinates by

$$(26) \quad 2\kappa \left(\frac{2p-q}{2+\kappa}\right)^2 - (4+\kappa) \frac{2p-q}{2+\kappa} + p = 0,$$

with two branches,

$$(27) \quad \begin{aligned} \gamma &= \gamma_0^\pm(p) := \frac{1}{2\kappa} \left(4 + \kappa \pm \sqrt{(4 + \kappa)^2 - 8\kappa p}\right), \quad p \leq \frac{(4 + \kappa)^2}{8\kappa}, \\ q &= 2p - \left(1 + \frac{\kappa}{2}\right) \gamma_0^\pm(p). \end{aligned}$$

or, equivalently,

$$(28) \quad 2p = q + \frac{2 + \kappa}{8\kappa} \left(6 + \kappa \pm \sqrt{(6 + \kappa)^2 - 16\kappa q}\right), \quad q \leq \frac{(6 + \kappa)^2}{16\kappa}.$$

Proof. Our aim is to derive an ODE satisfied by the whole-plane SLE one-point function,

$$(29) \quad G(z) := \mathbb{E}\left(z^{\frac{q}{2}} \frac{(f'(z))^{\frac{p}{2}}}{(f(z))^{\frac{q}{2}}}\right),$$

which, by construction, stays finite at the origin and such that $G(0) = 1$.

Let us introduce the shorthand notation,

$$(30) \quad X_t(z) := \frac{(\tilde{f}'_t(z))^{\frac{p}{2}}}{(\tilde{f}_t(z))^{\frac{q}{2}}},$$

where \tilde{f}_t is the conjugate, reversed radial SLE process in \mathbb{D} , as introduced in Definition 2.2, and such that by Lemma 2.3, the limit, $\lim_{t \rightarrow +\infty} e^t \tilde{f}_t(z) \stackrel{(\text{law})}{=} f_0(z)$, is the same in law as the whole-plane map at time zero. Applying the same method as in the previous section, we consider the time-dependent function

$$(31) \quad \tilde{G}(z, t) := \mathbb{E} \left(z^{\frac{q}{2}} X_t(z) \right),$$

such that

$$(32) \quad \lim_{t \rightarrow +\infty} \exp \left(\frac{p-q}{2} t \right) \tilde{G}(z, t) = G(z).$$

Consider now the martingale $(\mathcal{M}_s)_{t \geq s \geq 0}$, defined by

$$\mathcal{M}_s = \mathbb{E}(X_t(z) | \mathcal{F}_s).$$

By the SLE Markov property we get, setting $z_s := \tilde{f}_s(z)/\lambda(s)$,

$$(33) \quad \mathcal{M}_s = X_s(z) \tilde{G}(z_s, \tau), \quad \tau := t - s.$$

As before, the partial differential equation satisfied by $\tilde{G}(z_s, \tau)$ is obtained by expressing the fact that the ds -drift term of the Itô differential of Eq. (33),

$$d\mathcal{M}_s = \tilde{G} dX_s + X_s d\tilde{G},$$

vanishes. The differential of X_s is simply computed from Eqs. (20) and (21) above as:

$$(34) \quad \begin{aligned} dX_s(z) &= X_s(z) F(z_s) ds, \\ F(z) &:= \frac{p}{2} \left[1 - \frac{2}{(1-z)^2} \right] - \frac{q}{2} \left[1 - \frac{2}{1-z} \right]. \end{aligned}$$

The Itô differential $d\tilde{G}$ brings in the ds terms proportional to $\partial_{z_s} \tilde{G}$, $\partial_{z_s}^2 \tilde{G}$, and $\partial_\tau \tilde{G}$; therefore, in the PDE satisfied by \tilde{G} , the latter terms are exactly the same as in the PDE (23). We therefore directly arrive at the vanishing condition of the overall drift term coefficient in $d\mathcal{M}_s$,

$$(35) \quad X_s(z) \left[F(z_s) + z_s \left(\frac{z_s + 1}{z_s - 1} - \frac{\kappa}{2} \right) \partial_z - \partial_\tau - \frac{\kappa}{2} z_s^2 \partial_z^2 \right] \tilde{G}(z_s, \tau) = 0.$$

Since $X_s(z)$ does not vanish in \mathbb{D} , the bracket in (35) must identically vanish:

$$(36) \quad \left[F(z_s) + z_s \frac{z_s + 1}{z_s - 1} \partial_z - \partial_\tau - \frac{\kappa}{2} (z_s \partial_z)^2 \right] \tilde{G}(z_s, \tau) = 0,$$

where we used $z \partial_z + z^2 \partial_z^2 = (z \partial_z)^2$.

To derive the ODE satisfied by $G(z)$ (29), we first recall its expression as the limit (32), which further implies

$$\lim_{\tau \rightarrow +\infty} \exp\left(\frac{p-q}{2}\tau\right) \partial_\tau \tilde{G}(z, \tau) = -\frac{p-q}{2}G(z),$$

with an exchange of derivative and limit similar to that in the proof of Proposition 2.1. Multiplying the PDE (35) satisfied by \tilde{G} by $\exp(\frac{p-q}{2}\tau)$ and letting $\tau \rightarrow +\infty$, we get

$$\begin{aligned} \mathcal{P}(\partial)[G(z)] &:= \left[-\frac{\kappa}{2}(z\partial_z)^2 - \frac{1+z}{1-z}z\partial_z + F(z) + \frac{p-q}{2} \right] G(z) \\ (37) \quad &= \left[-\frac{\kappa}{2}(z\partial_z)^2 - \frac{1+z}{1-z}z\partial_z - \frac{p}{(1-z)^2} + \frac{q}{1-z} + p - q \right] G(z) = 0. \end{aligned}$$

Again, the space of holomorphic solutions to (37) is one-dimensional, because $F(z) + (p-q)/2$ vanishes at $z = 0$. For the boundary condition $G(0) = 1$, we now look for special solutions of the form $\varphi_\alpha(z) = (1-z)^\alpha$. This function satisfies the simple differential operator algebra [14]

$$(38) \quad \mathcal{P}(\partial)[\varphi_\alpha] = A(p, q, \alpha)\varphi_\alpha + B(q, \alpha)\varphi_{\alpha-1} + C(p, \alpha)\varphi_{\alpha-2},$$

$$(39) \quad A(p, q, \alpha) := p - q + \alpha - \frac{\kappa}{2}\alpha^2,$$

$$(40) \quad B(q, \alpha) := q - \left(3 + \frac{\kappa}{2}\right)\alpha + \kappa\alpha^2,$$

$$(41) \quad C(p, \alpha) := -p + \left(2 + \frac{\kappa}{2}\right)\alpha - \frac{\kappa}{2}\alpha^2,$$

such that, identically, $A + B + C = 0$. Because $\varphi_\alpha, \varphi_{\alpha-1}, \varphi_{\alpha-2}$ are linearly independent, the condition $\mathcal{P}(\partial)[\varphi_\alpha] = 0$ is equivalent to the system $A = C = 0$, hence $C(p, \alpha) = 0$ and $A(p, q, \alpha) - C(p, \alpha) = 2p - q - (1 + \kappa/2)\alpha = 0$. It yields precisely the parabola parametrization (25) given in Theorem 3.1, and has for solution (27). \square

4. SLE TWO-POINT FUNCTION

4.1. Beliaev–Smirnov type equations. In this section, we will determine the mixed moments of moduli, $\mathbb{E}\left(\frac{|f'(z)|^p}{|f(z)|^q}\right)$, for (p, q) belonging to the same parabola \mathcal{R} as in Theorem 3.1, and where $f = f_0$ is the (time zero) interior whole-plane SLE_κ map.

In contradistinction to the method used in Refs. [3, 14] for writing a PDE obeyed by $\mathbb{E}(|f'(z)|^p)$, we shall use here a slightly different approach, building on the results obtained in Section 2.1. We shall study the SLE two-point function for $z_1, z_2 \in \mathbb{D}$,

$$(42) \quad G(z_1, \bar{z}_2) := \mathbb{E}\left(\frac{z_1^{\frac{q}{2}} \frac{(f'(z_1))^{\frac{p}{2}}}{(f(z_1))^{\frac{q}{2}}}}{\left[z_2^{\frac{q}{2}} \frac{(f'(z_2))^{\frac{p}{2}}}{(f(z_2))^{\frac{q}{2}}} \right]}\right).$$

As before, we define a time-dependent, auxiliary two-point function,

$$(43) \quad \begin{aligned} \tilde{G}(z_1, \bar{z}_2, t) &:= \mathbb{E} \left(z_1^{\frac{q}{2}} \frac{(\tilde{f}'_t(z_1))^{\frac{p}{2}}}{(\tilde{f}_t(z_1))^{\frac{q}{2}}} \overline{\left[z_2^{\frac{q}{2}} \frac{(\tilde{f}'_t(z_2))^{\frac{p}{2}}}{(\tilde{f}_t(z_2))^{\frac{q}{2}}} \right]} \right) \\ &= \mathbb{E} \left(z_1^{\frac{q}{2}} X_t(z_1) \overline{z_2^{\frac{q}{2}} X_t(z_2)} \right), \end{aligned}$$

where as above \tilde{f}_t is the reverse radial SLE $_{\kappa}$ process 2.2, and where we used the shorthand notation (30). This time, the two-point function (42) is the limit

$$(44) \quad \lim_{t \rightarrow +\infty} e^{(p-q)t} \tilde{G}(z_1, \bar{z}_2, t) = G(z_1, \bar{z}_2).$$

Let us define the two-point martingale $(\mathcal{M}_s)_{t \geq s \geq 0}$, with

$$\mathcal{M}_s := \mathbb{E}(X_t(z_1) \overline{X_t(z_2)} | \mathcal{F}_s).$$

By the Markov property of SLE,

$$(45) \quad \mathbb{E}(X_t(z_1) \overline{X_t(z_2)} | \mathcal{F}_s) = X_s(z_1) \overline{X_s(z_2)} \tilde{G}(z_{1s}, \bar{z}_{2s}, \tau), \quad \tau := t - s,$$

where

$$(46) \quad z_{1s} := \tilde{f}_s(z_1)/\lambda(s); \quad \bar{z}_{2s} := \overline{\tilde{f}_s(z_2)/\lambda(s)} = \overline{\tilde{f}_s(z_2)}\lambda(s).$$

Their Itô differentials, dz_{1s} and $d\bar{z}_{2s}$, are as in (22),

$$(47) \quad \begin{aligned} dz_{1s} &= z_{1s} \left[\frac{z_{1s} + 1}{z_{1s} - 1} - \frac{\kappa}{2} \right] ds - i\sqrt{\kappa} z_{1s} dB_s, \\ d\bar{z}_{2s} &= \bar{z}_{2s} \left[\frac{\bar{z}_{2s} + 1}{\bar{z}_{2s} - 1} - \frac{\kappa}{2} \right] ds + i\sqrt{\kappa} \bar{z}_{2s} dB_s. \end{aligned}$$

As before, the partial differential equation satisfied by $\tilde{G}(z_{1s}, z_{2s}, \tau)$ is obtained by expressing the fact that the ds -drift term of the Itô differential of Eq. (45),

$$(48) \quad d\mathcal{M}_s = [dX_s(z_1) \overline{X_s(z_2)} + X_s(z_1) d\overline{X_s(z_2)}] \tilde{G} + X_s(z_1) \overline{X_s(z_2)} d\tilde{G},$$

vanishes.

The differentials of $X_s, \overline{X_s}$ are as in Eq. (34) above:

$$(49) \quad \begin{aligned} dX_s(z_1) &= X_s(z_1) F(z_{1s}) ds, \quad d\overline{X_s(z_2)} = \overline{X_s(z_2)} F(\bar{z}_{2s}) ds, \\ F(z) &:= \frac{p}{2} - \frac{q}{2} - \frac{p}{(1-z)^2} + \frac{q}{1-z}. \end{aligned}$$

We thus obtain the simple expression

$$(50) \quad d\mathcal{M}_s = X_s(z_1) \overline{X_s(z_2)} \left[[F(z_{1s}) + F(\bar{z}_{2s})] \tilde{G} ds + d\tilde{G} \right],$$

and the vanishing of the ds -drift term in $d\mathcal{M}_s$ requires that of the drift term in the right-hand side bracket in (50), since $X_s(z)$ does not vanish in \mathbb{D} .

The Itô differential of $\tilde{G}(z_{1s}, \bar{z}_{2s}, \tau)$ can be obtained from Eqs. (47) and Itô calculus as

$$(51) \quad d\tilde{G}(z_{1s}, \bar{z}_{2s}, \tau) = \partial_1 \tilde{G} dz_{1s} + \bar{\partial}_2 \tilde{G} d\bar{z}_{2s} - \partial_\tau \tilde{G} ds \\ - \frac{\kappa}{2} z_{1s}^2 \partial_1^2 \tilde{G} ds - \frac{\kappa}{2} \bar{z}_{2s}^2 \bar{\partial}_2^2 \tilde{G} ds + \kappa z_{1s} \bar{z}_{2s} \partial_1 \bar{\partial}_2 \tilde{G} ds,$$

where use was made of the shorthand notations, $\partial_1 := \partial_{z_1}$ and $\bar{\partial}_2 := \partial_{\bar{z}_2}$. We observe that the only coupling between the z_{1s}, \bar{z}_{2s} variables arises in the last term of (51), the other terms simply resulting from the independent contributions of the z_{1s} and \bar{z}_{2s} parts.

Using again the Itô differentials (47), we can rewrite (51) as

$$(52) \quad d\tilde{G} = -i\sqrt{\kappa} (z_{1s}\partial_1 - \bar{z}_{2s}\bar{\partial}_2) \tilde{G} dB_s \\ + \frac{z_{1s} + 1}{z_{1s} - 1} z_{1s} \partial_1 \tilde{G} ds + \frac{\bar{z}_{2s} + 1}{\bar{z}_{2s} - 1} \bar{z}_{2s} \bar{\partial}_2 \tilde{G} ds - \partial_\tau \tilde{G} ds \\ - \frac{\kappa}{2} (z_{1s}\partial_1 - \bar{z}_{2s}\bar{\partial}_2)^2 \tilde{G} ds,$$

where we used the obvious formal identity

$$(53) \quad (z_1 \partial_1)^2 + (\bar{z}_2 \bar{\partial}_2)^2 - 2z_1 \partial_1 \bar{z}_2 \bar{\partial}_2 = (z_1 \partial_1 - \bar{z}_2 \bar{\partial}_2)^2.$$

At this stage, comparing the computations (50) and (52) above with those in the one-point martingale study in Section 2.1, it is clear that the PDE obeyed by $\tilde{G} = \tilde{G}(z_{1s}, \bar{z}_{2s}, \tau)$ is obtained as two duplicates of Eq. (36), completed as in (53) by the derivative coupling between variables z_{1s}, \bar{z}_{2s} :

$$(54) \quad \left[F(z_{1s}) + z_{1s} \frac{z_{1s} + 1}{z_{1s} - 1} \partial_1 + F(\bar{z}_{2s}) + \bar{z}_{2s} \frac{\bar{z}_{2s} + 1}{\bar{z}_{2s} - 1} \bar{\partial}_2 - \partial_\tau - \frac{\kappa}{2} (z_{1s}\partial_1 - \bar{z}_{2s}\bar{\partial}_2)^2 \right] \tilde{G} = 0.$$

The existence of the limit (44) further implies that of

$$\lim_{\tau \rightarrow \infty} e^{(p-q)\tau} \partial_\tau \tilde{G}(z_1, \bar{z}_2, \tau) = -(p-q)G(z_1, \bar{z}_2).$$

Multiplying the PDE (54) satisfied by \tilde{G} by $\exp((p-q)\tau)$ and letting $\tau \rightarrow +\infty$, then gives the expected PDE for $G(z_1, \bar{z}_2)$. It can be most compactly written in terms of the ODE (37) as

$$(55) \quad [\mathcal{P}(\partial_1) + \mathcal{P}(\bar{\partial}_2) + \kappa z_1 \partial_1 \bar{z}_2 \bar{\partial}_2] G(z_1, \bar{z}_2) = 0,$$

and its fully explicit expression is

$$(56) \quad \mathcal{P}(D)[G(z_1, \bar{z}_2)] = -\frac{\kappa}{2} (z_1 \partial_1 - \bar{z}_2 \bar{\partial}_2)^2 G - \frac{1+z_1}{1-z_1} z_1 \partial_1 G - \frac{1+\bar{z}_2}{1-\bar{z}_2} \bar{z}_2 \bar{\partial}_2 G \\ + \left[-\frac{p}{(1-z_1)^2} - \frac{p}{(1-\bar{z}_2)^2} + \frac{q}{1-z_1} + \frac{q}{1-\bar{z}_2} + 2p - 2q \right] G = 0.$$

4.2. Moduli one-point function. Note that one can take the $z_1 = z_2 = z$ case in Definition (42) above, thereby obtaining the moduli one-point function,

$$(57) \quad G(z, \bar{z}) = \mathbb{E} \left(|z|^q \frac{|f'(z)|^p}{|f(z)|^q} \right).$$

Because of Eq. (56), it obeys the corresponding ODE,

$$(58) \quad \begin{aligned} \mathcal{P}(D)[G(z, \bar{z})] &= -\frac{\kappa}{2}(z\partial - \bar{z}\bar{\partial})^2 G - \frac{1+z}{1-z} z\partial G - \frac{1+\bar{z}}{1-\bar{z}} \bar{z}\bar{\partial} G \\ &+ \left[-\frac{p}{(1-z)^2} - \frac{p}{(1-\bar{z})^2} + \frac{q}{1-z} + \frac{q}{1-\bar{z}} + 2p - 2q \right] G = 0, \end{aligned}$$

which is the generalization to $q \neq 0$ of the Beliaev–Smirnov equation studied in Refs. [14] and [27].

4.3. Integrable case.

Lemma 4.1. *The space of formal series $F(z_1, \bar{z}_2) = \sum_{k, \ell \in \mathbb{N}} a_{k, \ell} z_1^k \bar{z}_2^\ell$, with complex coefficients and that are solutions of the PDE (56), is one-dimensional.*

Proof. We assume that F is a solution to (56) with $F(0, 0) = 0$; it suffices to prove that, necessarily, $F = 0$. We argue by contradiction: If not, consider the minimal (necessarily non constant) term $a_{k, \ell} z_1^k \bar{z}_2^\ell$ in the series of F , with $a_{k, \ell} \neq 0$ and $k + \ell$ minimal (and non vanishing). Then $\mathcal{P}(D)[F]$ (56) will have a minimal term, equal to $-a_{k, \ell} \left[\frac{\kappa}{2}(k - \ell)^2 + k + \ell \right] z_1^k \bar{z}_2^\ell$, which is non-zero, contradicting the fact that $\mathcal{P}(D)[F]$ vanishes. \square

As a second step, following Ref. [14], let us consider the action of the operator $\mathcal{P}(D)$ of (56) on a function of the factorized form $\varphi(z_1)\varphi(\bar{z}_2)P(z_1, \bar{z}_2)$, which we write, in a shorthand notation, as $\varphi\bar{\varphi}P$. By Leibniz's rule, it is given by

$$\begin{aligned} \mathcal{P}(D)[\varphi\bar{\varphi}P] &= -\frac{\kappa}{2}\varphi\bar{\varphi}(z_1\partial_1 - \bar{z}_2\bar{\partial}_2)^2 P - \kappa(z_1\partial_1 - \bar{z}_2\bar{\partial}_2)(\varphi\bar{\varphi})(z_1\partial_1 - \bar{z}_2\bar{\partial}_2)P \\ &+ \kappa(z_1\partial_1\varphi)(\bar{z}_2\bar{\partial}_2\bar{\varphi})P - \varphi\bar{\varphi}\frac{1+z_1}{1-z_1}z_1\partial_1 P - \varphi\bar{\varphi}\frac{1+\bar{z}_2}{1-\bar{z}_2}\bar{z}_2\bar{\partial}_2 P \\ &- \left[\frac{\kappa}{2}\bar{\varphi}(z_1\partial_1)^2\varphi + \frac{\kappa}{2}\varphi(\bar{z}_2\bar{\partial}_2)^2\bar{\varphi} + \bar{\varphi}\frac{1+z_1}{1-z_1}z_1\partial_1\varphi + \varphi\frac{1+\bar{z}_2}{1-\bar{z}_2}\bar{z}_2\bar{\partial}_2\bar{\varphi} \right] P \\ &+ \left[-\frac{p}{(1-z_1)^2} - \frac{p}{(1-\bar{z}_2)^2} + \frac{q}{1-z_1} + \frac{q}{1-\bar{z}_2} + 2p - 2q \right] \varphi\bar{\varphi}P. \end{aligned}$$

Note that the operator $z_1\partial_1 - \bar{z}_2\bar{\partial}_2$ is antisymmetric with respect to z_1, \bar{z}_2 ; therefore, if we choose a symmetric function, $P(z_1, \bar{z}_2) = P(z_1\bar{z}_2)$, the first line of $\mathcal{P}(D)[\varphi\bar{\varphi}P]$ above identically vanishes.

One then looks for solutions to (56) of the particular form,

$$G(z_1, \bar{z}_2) = \varphi_\alpha(z_1)\varphi_\alpha(\bar{z}_2)P(z_1\bar{z}_2),$$

where, as before, $\varphi_\alpha(z) = (1 - z)^\alpha$. The action of the differential operator then takes the simple form,

$$\begin{aligned} \mathcal{P}(D)[\varphi_\alpha \bar{\varphi}_\alpha P] &= z_1 \bar{z}_2 \varphi_{\alpha-1} \bar{\varphi}_{\alpha-1} (\kappa \alpha^2 P - 2(1 - z_1 \bar{z}_2) P') \\ &\quad + \mathcal{P}(\partial_1)[\varphi_\alpha] \bar{\varphi}_\alpha P + \mathcal{P}(\partial_2)[\bar{\varphi}_\alpha] \varphi_\alpha P, \end{aligned}$$

where P' is the derivative of P with respect to $z_1 \bar{z}_2$, and $\mathcal{P}(\partial)$ is the so-called boundary operator (37) [14].

The ODE, $\kappa \alpha^2 P(x) - 2(1 - x)P'(x) = 0$ with $x = z_1 \bar{z}_2$ and $P(0) = 1$, has for solution $P(z_1 \bar{z}_2) = (1 - z_1 \bar{z}_2)^{-\kappa \alpha^2 / 2}$. It is then sufficient to pick for α the value $\gamma = \gamma_0^\pm(p)$ (27) such that $\mathcal{P}(\partial)[\varphi_\gamma] = 0$, as obtained in the proof of Theorem 3.1, to get a solution of the PDE, $\mathcal{P}(D)[\varphi_\gamma \bar{\varphi}_\gamma P] = 0$ (56). By uniqueness of the solution with $G(0, 0) = 1$, it gives the explicit form of the SLE two-point function,

$$G(z_1, \bar{z}_2) = \varphi_\gamma(z_1) \varphi_\gamma(\bar{z}_2) (1 - z_1 \bar{z}_2)^{-\kappa \gamma^2 / 2}.$$

We thus get:

Theorem 4.2. *Let $f(z) = f_0(z)$ be the interior whole-plane SLE_κ map in the setting of Proposition (1.1); then, for (p, q) belonging to the parabola \mathcal{R} defined in Theorem 3.1 by Eqs. (25) or (26) or (27), and for any pair $(z_1, z_2) \in \mathbb{D} \times \mathbb{D}$,*

$$\mathbb{E} \left(z_1^{\frac{q}{2}} \frac{(f'(z_1))^{\frac{p}{2}}}{(f(z_1))^{\frac{q}{2}}} \overline{z_2^{\frac{q}{2}} \frac{(f'(z_2))^{\frac{p}{2}}}{(f(z_2))^{\frac{q}{2}}}} \right) = \frac{(1 - z_1)^\gamma (1 - \bar{z}_2)^\gamma}{(1 - z_1 \bar{z}_2)^\beta}, \quad \beta = \frac{\kappa}{2} \gamma^2.$$

Corollary 4.3. *In the same setting as in Theorem 4.2, we have for $z \in \mathbb{D}$,*

$$\mathbb{E} \left(|z|^q \frac{|f'(z)|^p}{|f(z)|^q} \right) = \frac{(1 - z)^\gamma (1 - \bar{z})^\gamma}{(1 - z \bar{z})^\beta}, \quad \beta = \frac{\kappa}{2} \gamma^2,$$

for

$$\begin{aligned} \gamma &= \gamma_0^\pm(p) := \frac{1}{2\kappa} \left(4 + \kappa \pm \sqrt{(4 + \kappa)^2 - 8\kappa p} \right), \quad p \leq \frac{(4 + \kappa)^2}{8\kappa}, \\ q &= 2p - \left(1 + \frac{\kappa}{2} \right) \gamma_0^\pm(p). \end{aligned}$$

Let us stress some particular cases of interest. First, the $p = 0$ case gives some integral means of f .

Corollary 4.4. *The interior whole-plane SLE_κ map has the integrable moments,*

$$\begin{aligned} \mathbb{E} \left(\left[\frac{f(z_1)}{z_1} \right]^{\frac{(2+\kappa)(4+\kappa)}{4\kappa}} \overline{\left[\frac{f(z_2)}{\bar{z}_2} \right]^{\frac{(2+\kappa)(4+\kappa)}{4\kappa}}} \right) &= \frac{(1 - z_1)^{\frac{4+\kappa}{\kappa}} (1 - \bar{z}_2)^{\frac{4+\kappa}{\kappa}}}{(1 - z_1 \bar{z}_2)^{\frac{(4+\kappa)^2}{2\kappa}}}, \\ \mathbb{E} \left(\left| \frac{f(z)}{z} \right|^{\frac{(2+\kappa)(4+\kappa)}{2\kappa}} \right) &= \frac{(1 - z)^{\frac{4+\kappa}{\kappa}} (1 - \bar{z})^{\frac{4+\kappa}{\kappa}}}{(1 - z \bar{z})^{\frac{(4+\kappa)^2}{2\kappa}}}. \end{aligned}$$

Second, taking $p = q$ yields the logarithmic integral means we started with:

Corollary 4.5. *The interior whole-plane SLE_κ map has the integrable logarithmic derivative two-point function,*

$$\mathbb{E}\left(\left[\begin{array}{c} z_1 \frac{f'(z_1)}{f(z_1)} \\ z_1 \frac{f'(z_1)}{f(z_1)} \end{array}\right]^{\frac{2+\kappa}{2\kappa}} \left[\begin{array}{c} \bar{z}_2 \frac{f'(z_2)}{f(z_2)} \\ \bar{z}_2 \frac{f'(z_2)}{f(z_2)} \end{array}\right]^{\frac{2+\kappa}{2\kappa}}\right) = \frac{(1-z_1)^{\frac{2}{\kappa}}(1-\bar{z}_2)^{\frac{2}{\kappa}}}{(1-z_1\bar{z}_2)^{\frac{2}{\kappa}}},$$

$$\mathbb{E}\left(\left|z \frac{f'(z)}{f(z)}\right|^{\frac{2+\kappa}{\kappa}}\right) = \frac{(1-z)^{\frac{2}{\kappa}}(1-\bar{z})^{\frac{2}{\kappa}}}{(1-z\bar{z})^{\frac{2}{\kappa}}}.$$

Theorem 1.3 describes the $\kappa = 2$ case of the latter result.

4.4. Generalization to processes with m -fold symmetry. The moments, $\mathbb{E}(|(f^{[m]})'(z)|^p)$ (for $m \in \mathbb{N} \setminus \{0\}$), as well as their associated integral means spectra were studied in Ref. [14]. Using Itô calculus, a PDE satisfied by these moments was derived for each value of m . The introduction of mixed (p, q) moments allows us to circumvent these calculations in a unified approach for $m \in \mathbb{Z} \setminus \{0\}$. To see this, notice that

$$(f^{[m]})'(z) = z^{m-1} f'(z^m) f(z^m)^{\frac{1}{m}-1}.$$

As a consequence,

$$\frac{|z|^q |(f^{[m]})'(z)|^p}{|f^{[m]}(z)|^q} = |z|^{q+p(m-1)} \frac{|f'(z^m)|^p}{|f(z^m)|^{p+\frac{q-p}{m}}},$$

so that we identically have

$$(59) \quad \mathbb{E}\left(|z|^q \frac{|(f^{[m]})'(z)|^p}{|f^{[m]}(z)|^q}\right) = G(z^m; p, q_m),$$

$$(60) \quad q_m = q_m(p, q) := p + \frac{q-p}{m},$$

with the notation,

$$(61) \quad G(z; p, q) := G(z, \bar{z}) = \mathbb{E}\left(|z|^q \frac{|f'(z)|^p}{|f(z)|^q}\right),$$

where we have made explicit the dependence on the (p, q) parameters of the SLE moduli one-point function (57) introduced in Section 3. From Theorem 4.2, we immediately get the following.

Theorem 4.6. *Let $f^{[m]}$ be the m -fold whole-plane SLE_κ map, $m \in \mathbb{Z} \setminus \{0\}$, with $z \in \mathbb{D}$ for $m > 0$ and $z \in \mathbb{C} \setminus \bar{\mathbb{D}}$ for $m < 0$. Then,*

$$\mathbb{E}\left(|z|^q \frac{|(f^{[m]})'(z)|^p}{|f^{[m]}(z)|^q}\right) = \frac{(1-z^m)^\alpha (1-\bar{z}^m)^\alpha}{(1-(z\bar{z})^m)^{\frac{\kappa}{2}\alpha^2}},$$

for (p, q) belonging to the m -dependent parabola $\mathcal{R}^{[m]}$, given in parametric form by

$$(62) \quad p = \left(2 + \frac{\kappa}{2}\right) \alpha - \frac{\kappa}{2} \alpha^2, \quad q = \left(m + 2 + \frac{\kappa}{2}\right) \alpha - \frac{\kappa}{2} (m + 1) \alpha^2, \quad \alpha \in \mathbb{R}.$$

In Cartesian coordinates, an equivalent statement is

$$\alpha = \frac{(m + 1)p - q}{m \left(1 + \frac{\kappa}{2}\right)},$$

with

$$q = (m + 1)p - m \frac{2 + \kappa}{4\kappa} \left(4 + \kappa \pm \sqrt{(4 + \kappa)^2 - 8\kappa p}\right), \quad p \leq \frac{(4 + \kappa)^2}{8\kappa},$$

or,

$$p = \frac{q}{m + 1} + \frac{m}{(m + 1)^2} \frac{2 + \kappa}{4\kappa} \left(2m + 4 + \kappa \pm \sqrt{(2m + 4 + \kappa)^2 - 8(m + 1)\kappa q}\right),$$

$$q \leq \frac{(2m + 4 + \kappa)^2}{8(m + 1)\kappa}.$$

As for logarithmic coefficients, first observe that trivially,

$$(63) \quad \log \frac{f^{[m]}(z)}{z} = \frac{1}{m} \log \frac{f(z^m)}{z^m}.$$

From this, and Theorem 1.2, we thus get

Corollary 4.7. *Let $f^{[m]}(z)$ be the m -fold whole-plane SLE₂ map and*

$$(64) \quad \log \frac{f^{[m]}(z)}{z} = 2 \sum_{n \geq 1} \gamma_n^{[m]} z^n;$$

then

$$\mathbb{E}(|\gamma_n^{[m]}|^2) = \begin{cases} \frac{1}{2n^2} & n = mk, \quad k \geq 1, \\ 0 & \text{otherwise.} \end{cases}$$

We can also see this result as a corollary of Theorem 4.6, which, for the logarithmic case $p = q$, and for any value of m , yields $p = q = 2$ for $\kappa = 2$ as the only integrable case.

5. GENERALIZED INTEGRAL MEANS SPECTRUM

In this section we aim at generalizing to the setting of the present work the integral means spectrum analysis of Refs. [3], [14] and [2] (see also [27, 28, 29]) concerning the whole-plane SLE. The original work by Beliaev–Smirnov [3] and Ref. [2] deal with the exterior case, whereas Ref. [14] and this work are concerned with the interior case, both being related by duality (13).

5.1. Modified One-Point Function. Let us first consider the *modified* SLE moduli one-point function,

$$(65) \quad F(z, \bar{z}) := \frac{1}{|z|^q} G(z, \bar{z}) = \mathbb{E} \left(\frac{|f'(z)|^p}{|f(z)|^q} \right).$$

Because of Eq. (58), it obeys the modified PDE,

$$(66) \quad \begin{aligned} \mathcal{P}(D)[F(z, \bar{z})] = & -\frac{\kappa}{2}(z\partial - \bar{z}\bar{\partial})^2 F - \frac{1+z}{1-z} z\partial F - \frac{1+\bar{z}}{1-\bar{z}} \bar{z}\bar{\partial} F \\ & + \left[-\frac{p}{(1-z)^2} - \frac{p}{(1-\bar{z})^2} + 2p - q \right] F(z, \bar{z}) = 0, \end{aligned}$$

which, of course, differs from Eq. (58). We can rewrite it as

$$(67) \quad \begin{aligned} \mathcal{P}(D)[F(z, \bar{z})] = & -\frac{\kappa}{2}(z\partial - \bar{z}\bar{\partial})^2 F - \frac{1+z}{1-z} z\partial F - \frac{1+\bar{z}}{1-\bar{z}} \bar{z}\bar{\partial} F \\ & - p \left[\frac{1}{(1-z)^2} + \frac{1}{(1-\bar{z})^2} + \sigma - 1 \right] F = 0, \end{aligned}$$

in term of the important new parameter,

$$(68) \quad \sigma := q/p - 1.$$

This PDE then exactly coincides with Eq. (106) in Ref. [14], where σ was meant to represent ± 1 , whereas here $\sigma \in \mathbb{R}$.

The value $\sigma = +1$ corresponds to the original Beliaev–Smirnov case, where the integral means spectrum successively involves three functions [2, 3, 7, 9, 18, 19, 20]:

$$(69) \quad \beta_{\text{tip}}(p, \kappa) := -p - 1 + \frac{1}{4}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa p}),$$

$$(70) \quad \text{for } p \leq p'_0(\kappa) := -1 - \frac{3\kappa}{8};$$

$$(71) \quad \begin{aligned} \beta_0(p, \kappa) := & -p + \frac{4 + \kappa}{4\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa p}), \\ & \text{for } p'_0(\kappa) \leq p \leq p_0(\kappa); \end{aligned}$$

$$(72) \quad \beta_{\text{lin}}(p, \kappa) := p - \frac{(4 + \kappa)^2}{16\kappa},$$

$$(73) \quad \text{for } p \geq p_0(\kappa) := \frac{3(4 + \kappa)^2}{32\kappa}.$$

As shown in Refs. [14, 27, 28, 29] in the $\sigma = -1$ interior case, because of the unboundedness of the interior whole-plane SLE map, there exists a phase

transition at $p = p^*(\kappa)$, with

$$(74) \quad \begin{aligned} p^*(\kappa) &:= \frac{1}{16\kappa} \left((4 + \kappa)^2 - 4 - 2\sqrt{2(4 + \kappa)^2 + 4} \right) \\ &= \frac{1}{32\kappa} \left(\sqrt{2(4 + \kappa)^2 + 4} - 6 \right) \left(\sqrt{2(4 + \kappa)^2 + 4} + 2 \right). \end{aligned}$$

The integral means spectrum is afterwards given by

$$(75) \quad \beta(p, \kappa) := 3p - \frac{1}{2} - \frac{1}{2}\sqrt{1 + 2\kappa p}, \text{ for } p \geq p^*(\kappa).$$

Since $p^*(\kappa) < p_0(\kappa)$ (73), this transition precedes and supersedes the transition from the bulk spectrum (71) towards the linear behavior (72).

The search in Ref. [14, Section 4.2.1] for exact boundary solutions to Eq. 66 (in the $\sigma = \pm 1$ cases) led to the introduction of the σ -dependent spectrum function [14, Eq. (180)]

$$(76) \quad \beta_+^\sigma(p, \kappa) = (1 - 2\sigma)p - \frac{1}{2}(1 + \sqrt{1 - 2\sigma\kappa p}).$$

For $\sigma = -1$, it recovers for interior whole-plane SLE the integral means spectrum (75) mentioned above. For $\sigma = +1$ it introduces a new spectrum,

$$(77) \quad \beta_+^{(+1)}(p, \kappa) = -p - \frac{1}{2}(1 + \sqrt{1 - 2\kappa p}),$$

the relevance of which for the exterior whole-plane SLE case is analyzed in a joint work by D. Beliaev and two of the present authors [2].

For general real values of σ (68), we can rewrite (76) as a function of (p, q, κ) ,

$$(78) \quad \beta_+^\sigma(p, \kappa) = \beta_1(p, q; \kappa) := 3p - 2q - \frac{1}{2} - \frac{1}{2}\sqrt{1 + 2\kappa(p - q)}.$$

We claim that the spectrum generated by the integral means in Definition (1.4) in the general (p, q) case will involve the standard multifractal spectra (69), (71), (72), that are independent of q , and also the new (p, q) -dependent multifractal spectrum (78). Phase transitions between these spectra will occur along lines drawn in the (p, q) plane. As we shall see in Section 5.4, the analysis of the integral means spectrum performed in Ref. [14], in particular that concerning the range of validity of the form (75) of (76) for $\sigma = -1$, as well as that given in Ref. [2] for the range of validity of the form (77) for $\sigma = +1$, and their corresponding proofs, can be extended to general values of the σ parameter via Propositions 5.18 and 5.9 to establish Theorem 1.7.

We first describe the corresponding partition of the (p, q) plane into the respective domains of validity of the four spectra above. For this, we need to determine boundary curves where pairs (possibly triplets) of these spectra coincide, that are marking the onset of the respective phase transitions.

5.2. Phase transition lines. Following Ref. [14], let us first introduce the analytical form of the various multifractal spectra based on the use of functions A (39), B (40) and C (41). It will be convenient to use the notation [14],

$$(79) \quad A^\sigma(p, \gamma) := -\frac{\kappa}{2}\gamma^2 + \gamma - \sigma p,$$

such that for $\sigma = q/p - 1$ (68),

$$(80) \quad A^\sigma(p, \gamma) = A(p, q; \gamma) := p - q + \gamma - \frac{\kappa}{2}\gamma^2,$$

as well as,

$$(81) \quad B(q, \gamma) = q - \left(3 + \frac{\kappa}{2}\right)\gamma + \kappa\gamma^2,$$

$$(82) \quad C(p, \gamma) = -\frac{\kappa}{2}\gamma^2 + \left(2 + \frac{\kappa}{2}\right)\gamma - p,$$

$$(83) \quad \beta(\gamma) := \beta(p, \gamma) := \frac{\kappa}{2}\gamma^2 - C(p, \gamma) = \kappa\gamma^2 - \left(2 + \frac{\kappa}{2}\right)\gamma + p,$$

where the last function, $\beta(p, \gamma)$, is the so-called ‘‘spectrum function’’ of Ref. [14]. Recall also that this function possesses an important duality property [14],

$$(84) \quad \beta(p, \gamma) = \beta(p, \gamma'), \quad \gamma + \gamma' := \frac{2}{\kappa} + \frac{1}{2}.$$

Remark 5.1. *The B - S parameter γ_0 , and bulk spectrum (71) $\beta_0 := \beta(p, \gamma_0)$, (corresponding to Eqs. (11) and (12) in Ref. [3]) are obtained from the equations (see Ref. [14]),*

$$(85) \quad C(p, \gamma_0) = 0; \quad \beta_0 = \beta(p, \gamma_0) = \kappa\gamma_0^2/2.$$

The two solutions to (85) are $\gamma_0^\pm(p)$ as in Eq. (27), where the lower branch $\gamma_0 := \gamma_0^-$ is the one selected for the bulk spectrum, $\beta_0(p) = \frac{1}{2}\kappa\gamma_0^-(p)^2$.

This spectrum (71) is defined only to the left of a vertical line in the (p, q) plane, as given by (see Fig. 4)

$$(86) \quad \Delta_0 := \left\{ p = \frac{(4 + \kappa)^2}{8\kappa}, q \in \mathbb{R} \right\}.$$

Remark 5.2. *The σ -dependent spectrum (76) is obtained from the equations*

$$(87) \quad A^\sigma(p, \gamma) = 0; \quad \beta(p, \gamma) = \kappa\gamma^2/2 - C(p, \gamma).$$

The solutions to Eq. (87) are

$$(88) \quad \gamma_\pm^\sigma(p) = \frac{1}{\kappa}(1 \pm \sqrt{1 - 2\sigma\kappa p}),$$

$$(89) \quad \beta_\pm^\sigma(p) = (1 - 2\sigma)p - \frac{\kappa}{2}\gamma_\pm^\sigma(p) = (1 - 2\sigma)p - \frac{1}{2}(1 \pm \sqrt{1 - 2\sigma\kappa p}).$$

The multifractal spectrum (76) is then given by the upper branch $\beta_+^\sigma(p)$ [14]. Note also that this spectrum is defined only for $2\sigma\kappa p \leq 1$, hence for points in the (p, q) plane below the oblique line (Fig. 4):

$$(90) \quad \Delta_1 := \left\{ (p, q) \in \mathbb{R}^2, q = p + \frac{1}{2\kappa} \right\}.$$

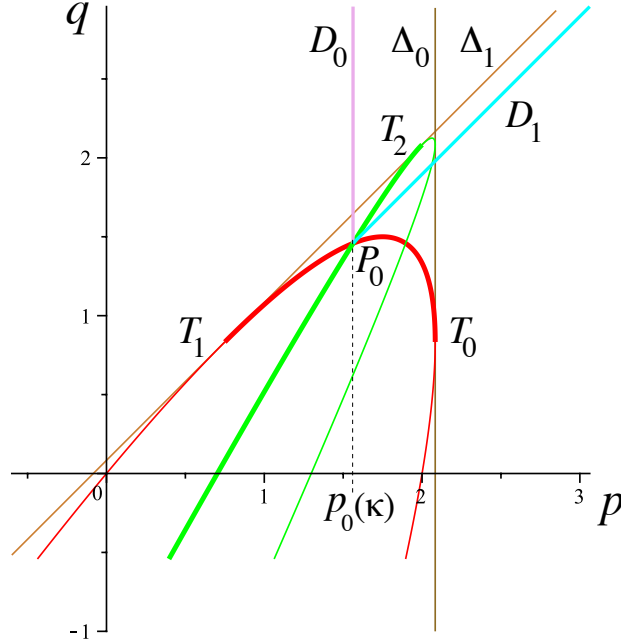


FIGURE 4. Red parabola \mathcal{R} (92) and green parabola \mathcal{G} (97) (for $\kappa = 6$). From the intersection point P_0 (100) originate the two (half)-lines D_0 (102) and D_1 (103). The bulk spectrum $\beta_0(p)$ and the generalized spectrum $\beta_1(p, q)$ coincide along the arc (94) of red parabola between its tangency points T_0 and T_1 with Δ_0 and Δ_1 (thick red line). They also coincide along the infinite left branch (99) of the green parabola, up to its tangency point T_2 to Δ_1 (thick green line). The $\beta_0(p)$ spectrum and the linear one $\beta_{\text{lin}}(p)$ coincide along D_0 , whereas $\beta_1(p, q)$ and $\beta_{\text{lin}}(p)$ coincide along D_1 .

5.2.1. ‘Red’ Parabola. The parabola \mathcal{R} of Theorems 3.1 and 4.2, which we shall hereafter call (and draw in) **red** (see Fig. 4), is given by the simultaneous conditions,

$$(91) \quad A^\sigma(p, \gamma) = A(p, q, \gamma) = 0, \quad C(p, \gamma) = 0,$$

hence also $B(q, \gamma) = 0$, which recovers the parametric form (25)

$$(92) \quad \begin{aligned} p &= p_{\mathcal{R}}(\gamma) := \left(2 + \frac{\kappa}{2}\right) \gamma - \frac{\kappa}{2} \gamma^2, \\ q &= q_{\mathcal{R}}(\gamma) := \left(3 + \frac{\kappa}{2}\right) \gamma - \kappa \gamma^2, \quad \gamma \in \mathbb{R}. \end{aligned}$$

By construction, the associated spectrum $\beta(p, \gamma)$ is therefore both of the B–S type, $\beta_0^\pm(p)$, and of the novel type, $\beta_\pm^\sigma(p)$. We successively have:

$$(93) \quad \gamma = \gamma_-^\sigma(p) = \gamma_0^-(p); \quad \beta_-^\sigma(p) = \beta_0^-(p), \quad \gamma \in (-\infty, 1/\kappa],$$

$$(94) \quad \gamma = \gamma_+^\sigma(p) = \gamma_0^-(p); \quad \beta_+^\sigma(p) = \beta_0^-(p), \quad \gamma \in [1/\kappa, 2/\kappa + 1/2],$$

$$(95) \quad \gamma = \gamma_+^\sigma(p) = \gamma_0^+(p); \quad \beta_+^\sigma(p) = \beta_0^+(p), \quad \gamma \in [2/\kappa + 1/2, +\infty),$$

where the change of analytic branch from the first to the second line corresponds to a tangency at T_1 of the red parabola to the boundary line Δ_1 , whereas the change from second to third corresponds to a tangency at T_0 to the vertical boundary line Δ_0 . The interval where the multifractal spectra coincide, i.e., when $\beta_+^\sigma(p) = \beta_0^-(p)$, is thus given by line (94) in the equations above.

In Cartesian coordinates, the red parabola \mathcal{R} (92) has for equation (26).

5.2.2. *‘Green’ Parabola.* A second parabola in the (p, q) plane, hereafter called **green** (see Fig. 4) and denoted by \mathcal{G} , is such that the multifractal spectra $\beta_0^-(p)$ and $\beta_+^\sigma(p) = \beta(p, q; \kappa)$ coincide on part of it. We use the *duality* property (84) of the spectrum function [14], and set the simultaneous seed conditions,

$$(96) \quad \begin{aligned} A^\sigma(p, \gamma') &= A(p, q, \gamma') = 0, \quad C(p, \gamma'') = 0, \\ \gamma' + \gamma'' &= 2/\kappa + 1/2, \end{aligned}$$

where γ' and γ'' are *dual* of each other and such that $\beta(p, \gamma') = \beta(p, \gamma'')$.

Eqs. (39) and (41) immediately give the parametric form for the green parabola,

$$(97) \quad \begin{aligned} p &= p_{\mathcal{G}}(\gamma') := \frac{(4 + \kappa)^2}{8\kappa} - \frac{\kappa}{2} \gamma'^2, \\ q &= q_{\mathcal{G}}(\gamma') := \frac{(4 + \kappa)^2}{8\kappa} + \gamma' - \kappa \gamma'^2, \quad \gamma' \in \mathbb{R}. \end{aligned}$$

Along this locus, we successively have:

$$(98) \quad \gamma' = \gamma_-^\sigma(p), \gamma'' = \gamma_0^+(p); \quad \beta_-^\sigma(p) = \beta_0^+(p), \quad \gamma' \in (-\infty, 0],$$

$$\gamma' = \gamma_-^\sigma(p), \gamma'' = \gamma_0^-(p); \quad \beta_-^\sigma(p) = \beta_0^-(p), \quad \gamma' \in [0, \kappa^{-1}],$$

$$(99) \quad \gamma' = \gamma_+^\sigma(p), \gamma'' = \gamma_0^-(p); \quad \beta_+^\sigma(p) = \beta_0^-(p), \quad \gamma' \in [\kappa^{-1}, +\infty),$$

where the changes of branches correspond to a tangency of the green parabola to Δ_0 followed by a tangency to Δ_1 . The multifractal spectra coincide when $\beta_+^\sigma(p) = \beta_0^-(p)$, which corresponds to the third line (99) in the equations above, i.e., to the domain where $\gamma' \geq 1/\kappa$.

5.2.3. *Quadruple point.* The intersection of the red and green parabolae (92) and (97) can be found by combining the seed equations (91) and (96). We find either $\gamma = \gamma' = 1/\kappa + 1/4$, or $\gamma = 2/\kappa + 1/4, \gamma' = -1/4$, which lead to the two intersection points,

$$(100) \quad P_0 : p_0 = p_0(\kappa) = \frac{3(4 + \kappa)^2}{32\kappa}, \quad q_0 = \frac{(4 + \kappa)(8 + \kappa)}{16\kappa},$$

$$(101) \quad P_1 : p_1 = \frac{(8 + \kappa)(8 + 3\kappa)}{32\kappa}, \quad q_0 = \frac{(4 + \kappa)(8 + \kappa)}{16\kappa}.$$

Note that these points have same ordinate, while the abscissa of the left-most one, P_0 , is $p_0(\kappa)$ (73), where the integral means spectrum transits from the bulk form (71) to its linear form (72).

Through this intersection point P_0 further pass two important straight lines in the (p, q) plane.

Definition 5.3. D_0 and D_1 are, respectively, the vertical line and the slope one line passing through point P_0 , of equations

$$(102) \quad D_0 := \{(p, q) : p = p_0\},$$

$$(103) \quad D_1 := \left\{ (p, q) : q - p = q_0 - p_0 = \frac{16 - \kappa^2}{32\kappa} \right\}.$$

On D_0 , one has $\beta_0(p, \kappa) = \beta_{\text{lin}}(p, \kappa)$. A key property of D_1 is the following. The difference,

$$(104) \quad \beta_1(p, q; \kappa) - \beta_{\text{lin}}(p, \kappa) = \frac{1}{\kappa} \left(\frac{\kappa}{4} - \sqrt{1 + 2\kappa(p - q)} \right)^2,$$

is always positive, and vanishes only on line D_1 , where

$$(105) \quad \forall (p, q) \in D_1, \quad \beta_1(p, q; \kappa) = \beta_{\text{lin}}(p, \kappa) = p - \frac{(4 + \kappa)^2}{16\kappa}.$$

5.2.4. *'Blue' Quartic.* A third locus, the **blue** quartic \mathcal{Q} , will also play an important role, that is where the tip-spectrum, $\beta_{\text{tip}}(p; \kappa)$ (69), coincides with the novel spectrum, $\beta_+^\sigma(p) = \beta_1(p, q; \kappa)$. The tip spectrum is given by $\beta_{\text{tip}}(p; \kappa) = \beta(p, \gamma_0) - 2\gamma_0 - 1$, where γ_0 is solution to $C(p, \gamma_0) = 0$ and such that the tip contribution is positive, $2\gamma_0 + 1 \leq 0$ [3, 14]; this corresponds to the tip condition (70) [3]. In the (p, q) plane, this describes the domain to the left of the straight line D'_0 (Fig. 5), defined by

$$(106) \quad D'_0 := \{(p, q) : p = p'_0(\kappa) = -1 - 3\kappa/8\}.$$

The generalized spectrum is given by $\beta_+^\sigma(p) = \beta(p, \gamma)$ where γ is solution to $A^\sigma(p, \gamma) = 0$. We therefore look for simultaneous solutions to the seed equations,

$$(107) \quad \begin{aligned} \beta(p, \gamma) &= \beta(p, \gamma_0) - 2\gamma_0 - 1, \quad 2\gamma_0 + 1 \leq 0, \\ A^\sigma(p, \gamma) &= 0, \quad C(p, \gamma_0) = 0. \end{aligned}$$

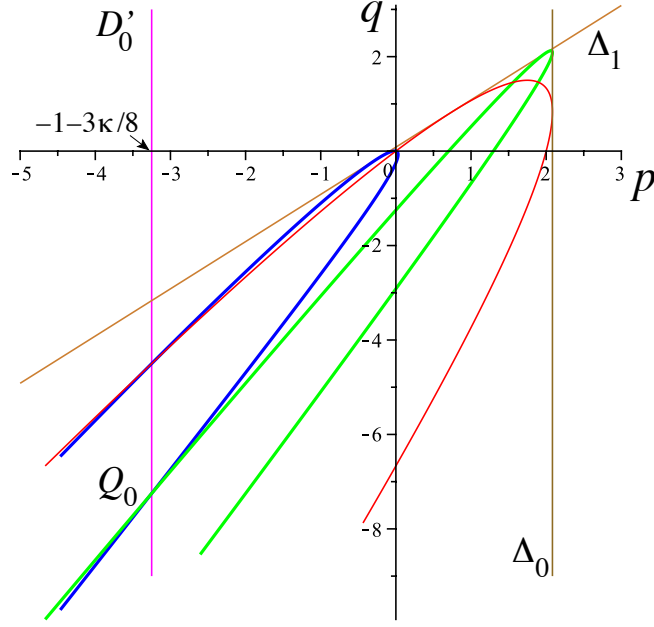


FIGURE 5. The blue quartic (113) for $\kappa = 6$. It intersects the green parabola at point Q_0 (117) and the red parabola at point Q_1 (116) (not marked), both of abscissa $p'_0(\kappa) = -1 - 3\kappa/8$.

Using Eq. (80), we first find, as for the red and green parabolae,

$$(108) \quad q - p = \gamma - \frac{\kappa}{2}\gamma^2,$$

and from (83) and (82), by substitution in the above,

$$(109) \quad 2p - q + \frac{1}{2} = \frac{\kappa}{4}(\gamma + \gamma_0),$$

$$(110) \quad \frac{4 + \kappa}{2}\gamma - \kappa\gamma^2 - 1 = \frac{8 + \kappa}{2}\gamma_0 - \kappa\gamma_0^2.$$

Solving for γ_0 in terms of γ gives

$$(111) \quad \gamma_0 = \gamma_0^\pm := \frac{8 + \kappa}{4\kappa} \pm \frac{1}{2\kappa}\Delta^{\frac{1}{2}}(\gamma),$$

$$(112) \quad \Delta(\gamma) := 4\kappa^2\gamma^2 - 2\kappa(4 + \kappa)\gamma + \frac{1}{4}(8 + \kappa)^2 + 4\kappa,$$

with $\Delta(\gamma) > 0, \forall \gamma \in \mathbb{R}$. The tip relevance inequality in (107), $2\gamma_0 + 1 \leq 0$, implies the choice of the negative branch in (111): $\gamma_0 = \gamma_0^-$. We thus get the

desired explicit parameterization of that branch of the quartic,

$$(113) \quad \begin{aligned} p &= p_{\mathcal{Q}}(\gamma) := \frac{\kappa}{16} + \left(1 + \frac{\kappa}{4}\right) \gamma - \frac{\kappa}{2} \gamma^2 - \frac{1}{8} \Delta^{\frac{1}{2}}(\gamma), \\ q &= q_{\mathcal{Q}}(\gamma) := p_{\mathcal{Q}}(\gamma) + \gamma - \frac{\kappa}{2} \gamma^2, \quad \gamma \in \mathbb{R}. \end{aligned}$$

Remark 5.4. *Note that because of the very choice to parameterize the parabola and the quartic by γ , such that A (80) vanishes, Eq. (108) holds for each of the pairs of parametric equations.*

We successively have along the branch (113) of the blue quartic:

$$(114) \quad \begin{aligned} \gamma &= \gamma_{-}^{\sigma}(p); \quad \beta_{-}^{\sigma}(p) = \beta_{\text{tip}}(p), \quad \gamma \in (-\infty, 1/\kappa], \\ \gamma &= \gamma_{+}^{\sigma}(p); \quad \beta_{+}^{\sigma}(p) = \beta_{\text{tip}}(p) < \beta_{0}^{-}(p), \quad \gamma \in [1/\kappa, 1 + 2/\kappa], \end{aligned}$$

$$(115) \quad \gamma = \gamma_{+}^{\sigma}(p); \quad \beta_{+}^{\sigma}(p) = \beta_{\text{tip}}(p) \geq \beta_{0}^{-}(p), \quad \gamma \in [1 + 2/\kappa, +\infty),$$

The intersection of the **blue** quartic (113) with the **red** parabola \mathcal{R} (92) is located at

$$(116) \quad Q_1 : p'_0 = -1 - \frac{3\kappa}{8}, q = -\frac{1}{2}(3 + \kappa); \quad \gamma = \gamma_0 = -\frac{1}{2},$$

followed by a second intersection at the origin, $p = q = 0$, for $\gamma = \frac{2}{\kappa}$ and $\gamma_0 = 0$.

The intersection of the **blue** quartic (113) with the **green** parabola \mathcal{G} (97) is located at

$$(117) \quad Q_0 : p'_0 = -1 - \frac{3\kappa}{8}, q'_0 := -2 - \frac{7\kappa}{8}; \quad \gamma = \gamma' = 1 + \frac{2}{\kappa}, \gamma_0 = -\frac{1}{2}.$$

Notice that these two intersection points have same abscissae, $p'_0(\kappa)$ (70), where the transition for $\gamma_0 = -\frac{1}{2}$ from the bulk spectrum β_0 to the tip spectrum β_{tip} takes place. They are found by combining Eqs. (91) or Eqs. (96) with (107).

The tip spectrum and the generalized one coincide in both γ -intervals (114) and (115), which together parameterize the branch of the quartic located below its contact with Δ_1 (see Fig. 5). Because of the tip relevance condition (70), only the interval (115) describing the lower infinite branch of the quartic located to the *left* of Q_0 will matter for the integral means spectrum.

5.3. Whole-plane SLE_{κ} generalized spectrum.

5.3.1. *Phase diagram.* Let us briefly summarize the results of Section 5.2. We know from Eq. (94) that the bulk spectrum $\beta_0(p)$ and the mixed spectrum $\beta_1(p, q)$ coincide along the finite sector of parabola \mathcal{R} located between tangency points T_0 and T_1 (Fig. 4). From Eq. (99), we also know that they coincide along the infinite left branch of parabola \mathcal{G} below the tangency point T_2 (Fig. 4).

The linear bulk spectrum $\beta_{\text{lin}}(p)$ coincides with $\beta_0(p)$ along line D_0 and supersedes the latter to the right of D_0 (Fig. 4). We know from (105) that $\beta_{\text{lin}}(p)$ and $\beta_1(p, q)$ coincide along the line D_1 (Fig. 4).

The tip spectrum $\beta_{\text{tip}}(p)$ coincides with $\beta_0(p)$ along line D'_0 , and supersedes it to the left of D'_0 . We finally know from Eq. (115) that this tip spectrum $\beta_{\text{tip}}(p)$ coincides with $\beta_1(p, q)$ along the lower branch of the blue quartic located below point Q_0 (117) (Fig. 5).

The only possible scenario which thus emerges to construct the average generalized integral means spectrum by a continuous matching of the 4 different spectra along the phase transition lines described above, is the partition of the (p, q) plane in 4 different regions as indicated in Fig. 2:

- a part (I) to the left of D'_0 and located above the blue quartic up to point Q_0 , where the average integral means spectrum is $\beta_{\text{tip}}(p)$;
- an upper part (II) bounded by lines D'_0, D_0 , and located above the section of the green parabola between points Q_0 and P_0 , where the spectrum is given by $\beta_0(p)$;
- an infinite wedge (III) of apex P_0 located between the upper half-lines D_0 and D_1 , where the spectrum is given by $\beta_{\text{lin}}(p)$;
- a lower part (IV) whose boundary is the blue quartic up to point Q_0 , followed by the arc of green parabola between points Q_0 and P_0 , followed by the half-line D_1 above P_0 where the spectrum is $\beta_1(p, q)$.

The two wings T_1P_0 and P_0T_0 of the red parabola (Fig. 4), where we know from Theorem 4.2 that the average spectrum is given by $\beta_0(p) = \beta_1(p, q)$, can thus be seen as the respective extensions of region IV into II and of region II into IV. The validity of this geometrical analysis of the phase diagram of Figure 2, associated with the generalized integral means spectrum of whole-plane SLE_κ , is established in Theorem 1.7.

5.3.2. *The B-S line.* As mentioned above, the whole-plane SLE case studied by Beliaev and Smirnov corresponds to the $q = 2p$ line. Because of Eq. (26), it intersects the red parabola \mathcal{R} only at $p = 0$. The green parabola \mathcal{G} (97) has for Cartesian equation,

$$(118) \quad \frac{\kappa}{2}(2p - q)^2 - \frac{1}{8}(4 + \kappa)^2(2p - q) + p + \frac{1}{128}(4 + \kappa)^2(8 + \kappa) = 0,$$

which shows that it intersects the B-S line at [2]

$$(119) \quad p = p_0''(\kappa) := -\frac{1}{128}(4 + \kappa)^2(8 + \kappa),$$

which is to the *left* of the tip transition line at $p'_0(\kappa) = -1 - \frac{3}{8}\kappa$ (70). The quartic \mathcal{Q} (113) obeys

$$(120) \quad \left[\left(2p - q - \frac{\kappa}{16} \right)^2 - \frac{c}{4} \right] \left(2p - q - 1 - \frac{\kappa}{8} \right) (2p - q) = \frac{\kappa}{2}(p - q) \left(2p - q - \frac{1}{4} - \frac{\kappa}{8} \right)^2$$

$$c = c(\kappa) := \frac{1}{64}(8 + \kappa)^2 + \frac{\kappa}{4},$$

which immediately shows that the B–S line $q = 2p$ intersects \mathcal{Q} only at the origin and stays above its lower branch.

The B–S line therefore does not intersect the segment of green parabola \mathcal{G} between P_0 and Q_0 , nor the quartic \mathcal{Q} below Q_0 (Fig. 2). Thus, as shown rigorously in Ref. [2], the novel spectrum β_1 *does not* a priori appear in the version of whole-plane SLE considered in Ref. [3]. The B–S line nevertheless intersects \mathcal{G} at p'_0 (119) to the left of Q_0 , in a domain lying above the quartic and where the integral mean receives a non-vanishing contribution from the SLE tip. But if that integral mean is restricted to avoid a neighborhood of $z = 1$, whose image is the tip, only the bulk spectrum remains, and a phase transition will take place from β_0 to β_1 when the line $q = 2p$ crosses \mathcal{G} , as shown in Ref. [2]. As we shall see in Section 5.5, the β_1 spectrum can also directly appear in the averaged integral means spectra of higher m -fold transforms of the B–S version of whole-plane SLE.

5.3.3. *The Koebe $\kappa \rightarrow 0$ limit.* In this limit, the whole-plane SLE map tends to the Koebe function. Eq. (26) for the red parabola \mathcal{R} becomes $3p - 2q = 0$, Eq. (118) for the green one \mathcal{G} becomes $3p - 2q - 1 = 0$, and Eq. (120) for the quartic factorizes into that of four parallel lines, among which $q = 2p$ gives the relevant lower branch. Point P_0 moves up to infinity, whereas $Q_0 \rightarrow (-1, -2)$. The phase diagram is thus made of only three different regions, I, where $\beta_{\text{tip}}(p) = -p - 1$, II, where $\beta_0 = 0$, and IV, where $\beta_1(p, q) = 3p - 2q - 1$.

5.4. **Proof of Theorem 1.7.** Let us first consider the following geometrical elements (Fig. 4). The red parabola \mathcal{R} (92) partitions the half-plane below Δ_1 into an open interior \mathcal{I} of \mathcal{R} , an open exterior \mathcal{E}_- located to the left of tangency point T_1 , and an open exterior \mathcal{E}_+ to the right of T_1 .

Lemma 5.5. *The average generalized integral means spectrum $\beta(p, q)$ of whole-plane SLE is bounded below as $\beta(p, q) \geq \beta_1(p, q)$ in $\mathcal{E}_- \cup \mathcal{I}$, whereas $\beta(p, q) \leq \beta_1(p, q)$ in \mathcal{E}_+ .*

Proof. The red parabola \mathcal{R} is parameterized by γ in (92) such that $C(p, \gamma) = 0$, where owing to (94), (93) and (95), $\gamma = \gamma_-^\sigma(p)$ before the tangency point T_1 , and $\gamma = \gamma_+^\sigma(p)$ after it (Fig. 4). In the open interior \mathcal{I} of \mathcal{R} , $C(p, \gamma_+^\sigma) > 0$ and $C(p, \gamma_-^\sigma) < 0$; in the open exterior \mathcal{E}_- to the left of the tangency point T_1 , $C(p, \gamma_\pm^\sigma) > 0$; in the open exterior \mathcal{E}_+ to the right of T_1 , $C(p, \gamma_\pm^\sigma) < 0$. According to [14, Section 4.2.5], and the generalization thereof to PDE (67), there exists then in $\mathcal{E}_- \cup \mathcal{I}$ a supersolution to Eq. (67) of critical exponent (83) $\beta(\gamma_+^\sigma) = \beta_+^\sigma(p) = \beta_1(p, q)$, such that the average integral means spectrum $\beta(p, q)$ is bounded below as $\beta(p, q) \geq \beta_1(p, q)$, whereas there exists in \mathcal{E}_+ a subsolution to (67) with the same critical exponent $\beta(\gamma_+^\sigma) = \beta_1(p, q)$, such that now $\beta(p, q) \leq \beta_1(p, q)$. \square

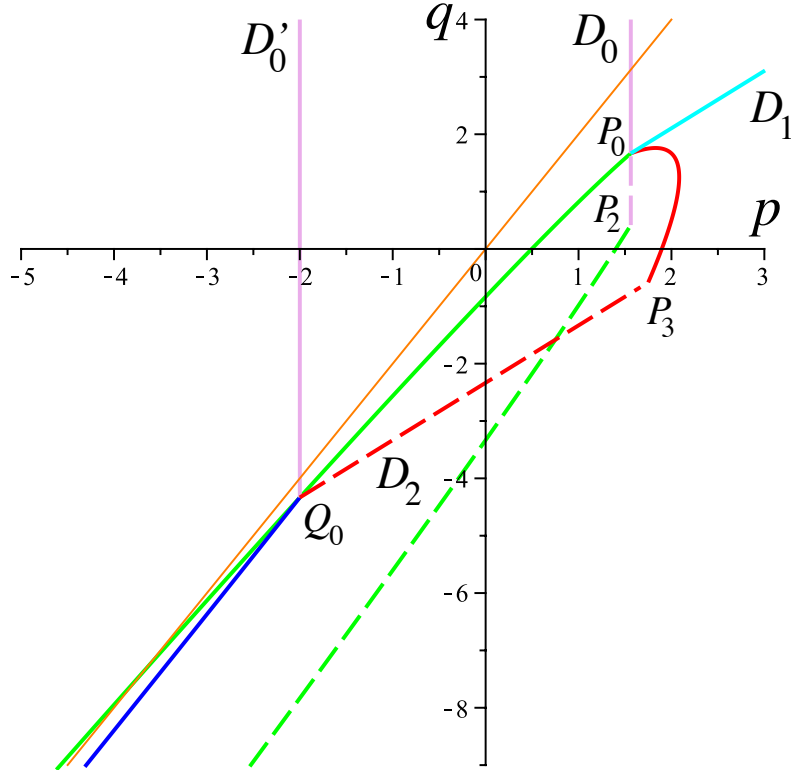


FIGURE 6. (Case $\kappa = 8/3$ shown here.) Various geometrical elements appearing in Section 5.4, when establishing the respective domains of validity of integral means spectra $\beta_{\text{tip}}(p)$, $\beta_0(p)$, $\beta_{\text{lin}}(p)$, and $\beta_1(p, q)$. The green parabola \mathcal{G} intersects D_0 at points P_0 (100) and $P_2 = (\frac{3(4+\kappa)^2}{32\kappa}, \frac{4+\kappa}{16})$. The dashed red line D_3 corresponds to Eq. (140), and intersects D'_0 at Q_0 and the red parabola \mathcal{R} at point $P_3 = (1 + \frac{2}{\kappa}, \frac{4-\kappa^2}{2\kappa})$. The $q = 2p$ continuous straight line in coral, corresponding to the whole-plane SLE version of Ref. [3], does not intersect the blue quartic, but intersects the green parabola at a point of abscissa (119).

i) *Linear spectrum.*

Proposition 5.6. *Consider in the (p, q) plane the upward wedge-like domain delimited by lines D_0 and D_1 intersecting at P_0 (Figs. 2 and 6). In this domain, the average generalized integral means spectrum has the linear form β_{lin} (72).*

Proof. For this purpose, we consider the convex function,

$$(121) \quad \mathfrak{b}(p, q) := \beta(p, q) - \beta_{\text{lin}}(p, q),$$

where β is the true average generalized integral means spectrum. As all integral means spectra, $\beta(p, q)$, hence $\mathbf{b}(p, q)$, are convex functions in \mathbb{R}^2 , as can be seen by using Hölder's inequality (see also Section 6). We know that the function \mathbf{b} (121) is equal to 0 on the boundary part of the wedge sector along D_0 , since $\beta = \beta_0 = \beta_{\text{lin}}$ there (recall Definition 5.3). On the other hand, from Lemma 5.5, we know that on the boundary part of the same sector along D_1 , $\beta \leq \beta_1 = \beta_{\text{lin}}$, so that $\mathbf{b} \leq 0$ there. We now use the Beliaev–Smirnov result asserting that $\mathbf{b} = 0$ on the intersection of the sector with the B–S line $q = 2p$. We will then conclude with the maximum principle satisfied by the convex function \mathbf{b} in the wedge sector. To be precise, let us consider the bounded triangular domain, intersection of the wedge with the half-plane $q \leq M$, with M a large positive number. We already know that $\mathbf{b} \leq 0$ on the vertical and oblique parts of the triangle's boundary. On the horizontal part, we know that at the end-points the convex function satisfies $\mathbf{b} \leq 0$, and is equal to 0 at the interior point of intersection of the horizontal side with the line $q = 2p$. It follows that \mathbf{b} is identically 0 on that side. We can then apply the maximum principle in the interior of the bounded triangle to conclude that $\mathbf{b} = 0$ there; letting $M \rightarrow +\infty$ gives the result in the whole domain. \square

ii) *Duplantier–Hastings–Beliaev–Smirnov spectrum.*

Proposition 5.7. *Consider the infinite domain of the (p, q) plane above the infinite upper branch (99) of parabola \mathcal{G} (96) located below point P_0 , and to the left of the half-line D_0 above P_0 (Fig. 6). The average generalized integral means spectrum $\beta(p, q)$ is given in this domain by the standard bulk spectrum $\beta_0(p)$ for $p \geq -1 - 3\kappa/8$ or by the tip version $\beta_{\text{tip}}(p)$ in the opposite case.*

Proof. We use the analysis performed in Ref. [14], Section 4.2.4, in the proof of Proposition 4.1 there. Two important quantities are defined as

$$(122) \quad a(p) := \gamma_0(p) - \gamma_+^\sigma(p) = \gamma_0(p) - \frac{1}{\kappa} - \frac{1}{\kappa} \sqrt{1 - 2\sigma\kappa p},$$

$$(123) \quad b(p) := \gamma_0(p) - \gamma_-^\sigma(p) = \gamma_0(p) - \frac{1}{\kappa} + \frac{1}{\kappa} \sqrt{1 - 2\sigma\kappa p},$$

where $\gamma_0(p) = \gamma_0^-(p)$ is the B–S parameter of Eq. (27),

$$(124) \quad \gamma_0^\pm(p) = \frac{1}{2\kappa} \left(4 + \kappa \pm \sqrt{(4 + \kappa)^2 - 8\kappa p} \right), \quad p \leq \frac{(4 + \kappa)^2}{8\kappa}.$$

The analysis in Refs. [3] and [14] was originally meant to apply to $\sigma = \pm 1$, but it identically goes through for $\sigma \in \mathbb{R}$. It is shown there that two (necessary and sufficient) conditions for the validity of the B–S method, hence of the corresponding spectra β_0 and β_{tip} , are respectively:

$$(125) \quad \frac{1}{2} - a(p) - b(p) \geq 0, \quad \frac{1}{2} - b(p) \geq 0.$$

The first condition is independent of σ and simply amounts to $p \leq p_0(\kappa) = \frac{3(4+\kappa)^2}{32\kappa}$, i.e., to points (p, q) belonging to the half-plane located to the left of D_0 . The second condition can be rewritten as the inequality,

$$(126) \quad \frac{1}{2} + \frac{2}{\kappa} - \gamma_0(p) - \gamma_+^\sigma(p) \geq 0.$$

The equality case corresponds to the duality condition (84) obeyed by $\gamma_0(p)$ and $\gamma_+^\sigma(p)$, as in Eq. (96), the solution of which is precisely given by the branch (99) of the green parabola below P_0 . Since at the origin, $a_0(0) = b_0(0) = -1/\kappa$, the inequalities (125) hold together in the unbounded domain to the left of the union of the left branch of \mathcal{G} below P_0 and of the half-line D_0 above P_0 . We thus conclude that the average generalized integral means spectrum $\beta(p, q)$ is given in this domain by $\beta_0(p)$ for $p \geq -1 - 3\kappa/8$ or by $\beta_{\text{tip}}(p)$ in the opposite case. \square

iii) Inside the green parabola. In Ref. [2], a new method of proof was given, which establishes the validity of the D–H–BS spectrum along the $\sigma = +1$ line $q = 2p$, below the transition point (119). There $\beta_0(p) \leq \beta_1(p, 2p) \leq \beta_{\text{tip}}(p)$, and the integral means spectrum is still given by β_{tip} ; however avoiding the tip neighborhood in the integral mean reveals the existence of a phase transition to the β_1 spectrum. We now generalize this method to $\sigma \in \mathbb{R}$.

Definition 5.8. The domain \mathcal{D} , defined by the inequalities,

$$(127) \quad 0 < \frac{1}{2} - a(p) - b(p), \quad 0 < b(p) - \frac{1}{2} < \frac{2}{\kappa},$$

is the interior part of the green parabola \mathcal{G} located to the left of D_0 (102).

For general σ in the (p, q) plane, the first inequality in (127) corresponds to the left of D_0 , while the second set corresponds to the interior of the \mathcal{G} parabola. Indeed, Eqs. (123) and (124) show that $b = (\gamma_0^- + \gamma_+^\sigma) - 2/\kappa = 1 + 4/\kappa - (\gamma_0^+ + \gamma_-^\sigma)$. Thus, the r.h.s. of (127) is equivalent to $\gamma_0^- + \gamma_+^\sigma \geq 1/2 + 2/\kappa$ and $\gamma_0^+ + \gamma_-^\sigma > 1/2 + 2/\kappa$. The first inequality is saturated on the top branch (99) of the green parabola (96) (the phase transition line), while the second one is saturated on the lower branch (98). The domain \mathcal{D} is closed to the right by the segment of line D_0 in between the intersection points P_0 (100) and $P_2 := \left(\frac{3(4+\kappa)^2}{32\kappa}, \frac{4+\kappa}{16}\right)$ of D_0 with the green parabola (Fig. 6).

Proposition 5.9. *In domain \mathcal{D} , the average generalized integral means spectrum is given by $\beta(p, q) = \max\{\beta_{\text{tip}}(p), \beta_1(p, q)\}$.*

Corollary 5.10. *Since in \mathcal{D} the blue quartic \mathcal{Q} below Q_0 is the separatrix for $\beta_{\text{tip}}(p) = \beta_1(p, q)$, the average generalized integral means spectrum is given by $\beta_{\text{tip}}(p)$ in the infinite thin wedge \mathcal{W} in \mathcal{D} of apex Q_0 , located to the left of line D'_0 , inbetween the green parabola \mathcal{G} and the blue quartic \mathcal{Q} , and by $\beta_1(p, q)$ in the remaining part $\mathcal{D} \setminus \mathcal{W}$ of \mathcal{D} (Fig. 6).*

Proof. The extension to $\sigma \in \mathbb{R}$ of the proof given in Ref. [2] for $\sigma = +1$ and for $|z| > 1$, involves, now for $z \in \mathbb{D}$ and $u := |1 - z|^2$, the function (see [2, Eq. (4.1)]):

$$(128) \quad \psi := \varsigma\psi_0 + \psi_1 := \varsigma g_0(u)u^{\gamma_0}(1 - z\bar{z})^{-\beta_0} + u^{\gamma_1}(1 - z\bar{z})^{-\beta_1},$$

where we set $\gamma_0 = \gamma_0^-(p)$ (124) with $\beta_0 := \beta_0(p)$ (71), together with $\gamma_1 := \gamma_+^\sigma(p)$ (88) and $\beta_1 = \beta_+^\sigma(p) = \beta_1(p, q)$ (89). The function g_0 is a peculiar combination of two hypergeometric functions, parameterized by a (122) and b (123), with $\varsigma := \text{sgn } g_0(0)$ (see [2, Eqs. (3.5)–(3.7)]),

$$(129) \quad g_0(u) = C_0 {}_2F_1(a, b, c, u/4) - C'_0 (u/4)^{1/2-a-b} {}_2F_1(a', b', c', u/4),$$

where

$$(130) \quad \begin{aligned} a &= \gamma_0 - \gamma_1 = \gamma_0 - \gamma_+^\sigma, & b &= \gamma_0 - \gamma_-^\sigma, & c &= \frac{1}{2} + a + b, \\ a' &= \frac{1}{2} - a, & b' &= \frac{1}{2} - b, & c' &= \frac{1}{2} + a' + b'. \end{aligned}$$

The hypergeometric functions are singular at $u = 4$, but the coefficients C_0 and C'_0 are chosen in such a way that the function g_0 is smooth at $z = -1$ [3]:

$$(131) \quad g_0(0) = C_0 = \frac{\Gamma(3/2 - a - b)}{\Gamma(1/2 - a)\Gamma(1/2 - b)}, \quad C'_0 = \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)},$$

such that near $u = 4$, $g_0(u) = \frac{1}{\sqrt{\pi}} \left(\frac{1}{2} - a - b\right) + O(4 - u)$. This, together with the condition $1/2 - a - b > 0$ in \mathcal{D} , insures that the function g_0 stays bounded on the whole interval $u \in [0, 4]$.

One further defines

$$(132) \quad \ell_\delta := (-\log(1 - z\bar{z}))^\delta, \quad \delta \in \mathbb{R}.$$

As in the original Ref. [3], the proof then consists in showing that $\psi\ell_\delta$ provides a positive sub- or supersolution of the differential operator (58), depending on the sign of δ .

The key steps are Lemmas 4.1 and 4.2 in Ref. [2]. *Mutatis mutandis*, these Lemmas and their proofs are as follows for general σ in the (p, q) plane.

Lemma 5.11. *For $(p, q) \in \mathcal{D}$, there is $r_0 < 1$ such that $\psi = \varsigma\psi_0 + \psi_1 > 0$ for all z such that $r_0 < |z| < 1$.*

Proof. According to Ref. [2, Lemma 4.1], the condition for this to hold is $\beta_1 > \beta_0$; this is satisfied in the intersection of the interior domains of the green parabola \mathcal{G} and of the red parabola \mathcal{R} , and \mathcal{D} is contained in this intersection. \square

Along the top branch (99) of the green parabola (96) (the onset of the phase transition), one has $b = 1/2$, and this precisely corresponds to $g_0(0) = 0$ in Eq. (131). Actually, this vanishing, which is critical to the proof below, also happens for half-integer values $b = n + 1/2, n \in \mathbb{N}$. Proposition 5.13 below shows that

there indeed exist a finite set of integers \mathcal{J}_κ , and a discrete set \mathcal{T}_κ of parabolic trajectories in the (p, q) -plane, where this is realized. We need the following

Definition 5.12. For $n \in \mathbb{N}$, define \mathcal{P}_n as the parabola given in parametric equations by

$$(133) \quad \begin{aligned} p = p_n(\gamma) &= -\frac{\kappa}{2}\gamma^2 + \kappa n\gamma + \frac{\kappa}{2} \left(\frac{(4 + \kappa)^2}{4\kappa^2} - n^2 \right), \quad \gamma \in \mathbb{R}, \\ q = q_n(\gamma) &= -\kappa\gamma^2 + (\kappa n + 1)\gamma + \frac{\kappa}{2} \left(\frac{(4 + \kappa)^2}{4\kappa^2} - n^2 \right); \end{aligned}$$

note that for $n = 0$, $\mathcal{P}_0 = \mathcal{G}$. Define then \mathcal{P}_n^+ as the branch corresponding to $\gamma \geq \max\{n, \kappa^{-1}\}$.

For $n = 0$, \mathcal{P}_0^+ coincides with the upper branch of the green parabola \mathcal{G} stemming from its tangency point to Δ_1 . For $\kappa n < 1$, the branch originates from the point of tangency of \mathcal{P}_n to Δ_1 (90), whereas for $\kappa n > 1$, it originates from the point of tangency to Δ_0 (86). For $\kappa n = 1$, it degenerates into the bisector of Δ_0 and Δ_1 , $q = 2p - p(\kappa)$ with $p(\kappa) = (6 + \kappa)(2 + \kappa)/8\kappa$. (See Fig. 5.4.)

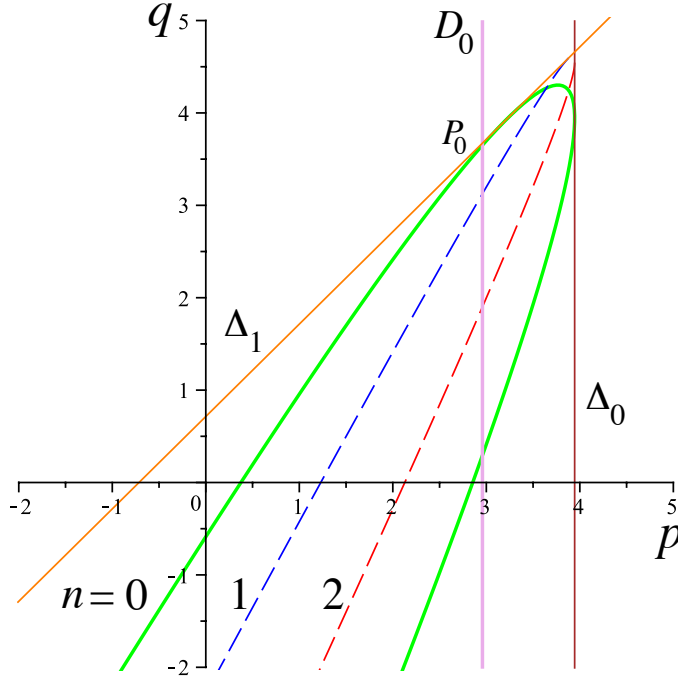


FIGURE 7. Green parabola $\mathcal{G} = \mathcal{P}_0$, and parabolic branches \mathcal{P}_n^+ for $n \in \{0, 1, 2\}$, and $\kappa = 0.7$. They cross into domain \mathcal{D} at the intersection points (134) with D_0 .

Lemma 5.13. *The set of equations, $b = n + \frac{1}{2}$, $n \in \mathbb{N}$, where $g_0(0) = 0$, is realized on the set of parabola branches \mathcal{P}_n^+ as defined in Def. 5.12. Inside domain \mathcal{D} , this yields the set $\mathcal{T}_\kappa := \cup_{n \in \mathcal{J}_\kappa} \mathcal{P}_n^+ \cap \mathcal{D}$, where $\mathcal{J}_\kappa := \{n \in \mathbb{N}, 0 \leq n \leq \lfloor 2\kappa^{-1} \rfloor\}$.*

Proof. From definition (123) or (130) of b , we have

$$b = \gamma_0^- + \gamma_+^\sigma - \frac{2}{\kappa} = \gamma_0 + \gamma_1 - \frac{2}{\kappa},$$

and we get the condition $\gamma_1 + \gamma_0 = n + \frac{1}{2} + \frac{2}{\kappa}$, together with, as in (96), $A^\sigma(p, \gamma_1) = 0, C(p, \gamma_0) = 0$. This yields the parabola equations (133) for \mathcal{P}_n , where $\gamma = \gamma_1$. The root determinations $\gamma_1 = \gamma_+^\sigma$ and $\gamma_0 = \gamma_0^-$ then determine the range $\gamma = \gamma_1 > \max\{n, 1/\kappa\}$ of branch \mathcal{P}_n^+ . The restriction on the values of n in \mathcal{J}_κ directly comes from definition (127) of \mathcal{D} . \square

Remark 5.14. *Note that for $\kappa > 2$, \mathcal{T}_κ is reduced to the $n = 0$ phase transition line on \mathcal{G} , whereas positive values of n exist in \mathcal{J}_κ only for $0 < \kappa \leq 2$.*

Remark 5.15. *The parabolic branch \mathcal{P}_n^+ intersects D_0 for $\gamma = \gamma^{(n)} := n + \frac{4+\kappa}{4\kappa}$, at*

$$(134) \quad p_n = p_0 = \frac{3}{32\kappa}(4 + \kappa)^2, q_n = \frac{(4 + \kappa)(8 + \kappa)}{16\kappa} - \frac{\kappa}{2}n \left(n + \frac{1}{2} \right),$$

its intersection with \mathcal{D} corresponding to the range $\gamma > \gamma^{(n)}$.

Remark 5.16. *The parabolic branch \mathcal{P}_n^+ intersects the B-S line $q = 2p$ only when $0 \leq n \leq \lfloor \kappa^{-1} \rfloor$, and their intersection is given by*

$$(135) \quad p_n''(\kappa) := -\frac{(1 + 2n)(8 + \kappa - 2n\kappa)(4 + \kappa + 2n\kappa)(4 + \kappa - 2n\kappa)}{128(1 - n\kappa)^2},$$

in agreement with Ref. [2, Eq. (3.18)]. The $n = 0$ case corresponds, for any value of κ , to the phase transition point $p_0''(\kappa)$ (119), as studied in Ref. [2].

Lemma 5.17. *For $(p, q) \in \mathcal{D}$ and $(p, q) \notin \mathcal{T}_\kappa$, there is $r_0 < 1$ such that $\mathcal{P}(D)[\psi l_\delta]$ for ψ (128) has a constant sign in the annulus $r_0 < |z| < 1$, which depends only on that of δ .*

Proof. We follow the proof of the similar [2, Lemma 4.2], which distinguishes three cases according to the relative rates at which $r := |z| \rightarrow 1$ and $u = |1 - z|^2 \rightarrow 0$, with $1 - r \leq u^{1/2}$. In case I, u is bounded away from zero, and the only requirement is $\beta_0(p) < \beta_1(p, q)$, which here holds in domain \mathcal{D} . In case II, one assumes that $(1 - r)^{2-\varepsilon} < u < u_0$, with $\varepsilon > 0$, and $u_0 > 0$ chosen such that $\varsigma g_0(u) > 0$ for $0 < u < u_0$. This u_0 exists here thanks to $(p, q) \notin \mathcal{T}_\kappa$ and Lemma 5.13. In case III, one assumes that $1 - r > u^{1/2+\varepsilon}$, for some $\varepsilon > 0$ to be determined such that the following bound holds, $\psi_1 \leq \varsigma \psi_0 u^{1/2}$ [2, Eq. (4.3)]. For that, the proof of [2, Lemma 4.2] makes a crucial use of the inequality $\beta_1 < \beta_{\text{tip}}$. Here, this is precisely valid in the wedge $\mathcal{W} \subset \mathcal{D}$ (recall Corollary 5.10), which thus establishes Lemma 5.17 in \mathcal{W} . In $\mathcal{D} \setminus \mathcal{W}$, we need a different, more general argument.

For functions A and B of r , we shall use the short-hand notations, $A \lesssim B$ for $A \leq cB$ with c some positive constant, and $A \approx B$ when both $A \lesssim B$ and $A \gtrsim B$ hold. We start with

$$\psi_1 \approx (1-r)^{-\beta_1} u^{\gamma_1} = (1-r)^{-\beta_0} u^{\gamma_0} u^{\gamma_1-\gamma_0} (1-r)^{\beta_0-\beta_1} \approx \varsigma \psi_0 u^{\gamma_1-\gamma_0} (1-r)^{\beta_0-\beta_1}.$$

Note that the last estimate requires that $g_0(0) \neq 0$, hence the condition $(p, q) \notin \mathcal{T}_\kappa$ of Lemma 5.17. We expect $u^{\gamma_1-\gamma_0} (1-r)^{\beta_0-\beta_1}$ to be bounded by some positive power of u . Recall that in \mathcal{D} , $\beta_0 - \beta_1 < 0$. Hence, for case III where $1-r > u^{1/2+\varepsilon}$,

$$u^{\gamma_1-\gamma_0} (1-r)^{\beta_0-\beta_1} < u^{\gamma_1-\gamma_0+(1/2+\varepsilon)(\beta_0-\beta_1)}.$$

Recalling notation (83), we have

$$\beta_0 - \beta_1 = \beta(\gamma_0) - \beta(\gamma_1) = (\gamma_0 - \gamma_1)(\kappa(\gamma_0 + \gamma_1) - (2 + \kappa/2)),$$

which yields $u^{\gamma_1-\gamma_0} (1-r)^{\beta_0-\beta_1} < u^\alpha$, with

$$\begin{aligned} \alpha &:= (\gamma_1 - \gamma_0) [1 - (1/2 + \varepsilon)(\kappa(\gamma_0 + \gamma_1) - (2 + \kappa/2))] \\ &= (\gamma_1 - \gamma_0) [1 - \kappa(1/2 + \varepsilon)(b - 1/2)]. \end{aligned}$$

Recall that $\gamma_1 - \gamma_0 > 0$ and $b - 1/2 < 2/\kappa$ by Definition (5.8) of \mathcal{D} , so that for sufficiently small ε , we have $\alpha > 0$. Hence for this ε , we have that

$$(136) \quad \psi_1 \lesssim \varsigma \psi_0 u^\alpha, \quad \alpha > 0.$$

As in the proof of [2, Lemma 4.2], this estimate insures that Lemma 5.17 holds in case III. \square

To conclude the proof of Proposition 5.9, Lemmas 5.11 and 5.17 show that for $(p, q) \notin \mathcal{T}_\kappa$, $\psi \ell_\delta$ provides both positive sub- and supersolutions depending on the sign of δ , and the generalized spectrum $\beta(p, q)$ can then be read off from the explicit form (128) of ψ . The spectrum function $\beta(p, q)$ is extended by convexity, hence continuity on \mathcal{T}_κ , which by Lemma 5.13 has no interior. \square

iv) Between the green and red parabolae. We now concentrate on the infinite domain $\tilde{\mathcal{D}}$ located between the top branch (99) of the green parabola \mathcal{G} stemming from point P_0 and the right branch of the red parabola \mathcal{R} (91) starting from P_0 (Fig. 6). Ref. [14, Theorem 1.4], states the validity of the spectrum $\beta_1(p, 0)$ on the $q = 0$ line, i.e., for $\sigma = -1$, between the transition point $p^*(\kappa)$ (74), which is the (left) intersection point of the green parabola \mathcal{G} with that line, up to the point $\min\{\hat{p}(\kappa), p(\kappa)\}$, where $\hat{p}(\kappa) := 1 + \kappa/2$, and where $p(\kappa) := (2 + \kappa)(6 + \kappa)/8\kappa$ is the rightmost intersection point of the red parabola \mathcal{R} with the $q = 0$ line.

Proposition 5.18. *Define $\tilde{\mathcal{D}}$ as the infinite domain located between the left branch (99) of the green parabola \mathcal{G} stemming from point P_0 and the right branch of the red parabola \mathcal{R} (91) starting from P_0 . Define $\hat{\mathcal{D}} \subset \tilde{\mathcal{D}}$ as the racket-shaped finite domain located above the straight line D_3 of equation $p - q = \hat{p}(\kappa) = 1 + \frac{\kappa}{2}$,*

which intersects \mathcal{G} at point Q_0 and \mathcal{R} at point $P_3 := (1 + \frac{2}{\kappa}, \frac{4-\kappa^2}{2\kappa})$ (Fig. 6). In $\hat{\mathcal{D}}$, the average generalized integral means spectrum is $\beta_1(p, q)$.

Proof. The proof given for $\sigma = -1$ in [14, Section 4.2.7], was based on the key *duality* property (84) of function (83), i.e., $\beta(\gamma) = \beta(\gamma')$ for $\gamma' = 1/2 + 2/\kappa - \gamma$. By using the general σ -formalism set up in Ref. [14, Sections 4.2.1 – 4.2.3], the same proof applies word for word to the $\sigma \in \mathbb{R}$ case. It suffices to replace everywhere the quantities there, $\gamma_{\pm} = \frac{1}{\kappa}(1 \pm \sqrt{1 + 2\kappa p})$, by (88), $\gamma_{\pm}^{\sigma} = \frac{1}{\kappa}(1 \pm \sqrt{1 - 2\kappa\sigma p})$. Instead of the test function (128), (129) as above, use is made of

$$(137) \quad \psi := g(u)(1 - z\bar{z})^{-\beta(\gamma)},$$

where g is the weighted combination of two hypergeometric functions,

$$(138) \quad g(u) = (u/4)^{\gamma} {}_2F_1(a, b, c, u/4) - (C'_0/C_0) (u/4)^{\gamma'} {}_2F_1(a', b', c', u/4),$$

where now γ is a free parameter, taken such that $\gamma > \gamma_1 = \gamma_+^{\sigma}(p)$ (88), and

$$(139) \quad \begin{aligned} a &= \gamma - \gamma_1 = \gamma - \gamma_+^{\sigma}, & b &= \gamma - \gamma_-^{\sigma}, & c &= \frac{1}{2} + a + b, \\ a' &= \frac{1}{2} - a, & b' &= \frac{1}{2} - b, & c' &= \frac{1}{2} + a' + b'. \end{aligned}$$

At the end of the argument, one lets $a \rightarrow 0^+$, $\gamma \rightarrow \gamma_1 = \gamma_+^{\sigma}$, so that $\beta(\gamma) \rightarrow \beta(\gamma_1) = \beta_1(p, q)$ (89).

The proof given in Section 4.2.7 of Ref. [14] holds provided that two simultaneous conditions are fulfilled, $C(p, \gamma_+) > 0$ and $C(p, \gamma'_+) < 0$, where $\gamma'_+ := \frac{1}{2} + \frac{2}{\kappa} - \gamma_+$ is the dual value of γ_+ . These two conditions precisely gave $p^*(\kappa) < p < p(\kappa)$. By definition, on the branch (94), (95) of the red parabola \mathcal{R} (91), hence on its right branch starting from P_0 , we have $C(p, \gamma_+^{\sigma}) = 0$. On the branch (99) of the green parabola \mathcal{G} (96), hence on its left branch starting from P_0 , we have $C(p, \gamma_+^{\sigma'}) = 0$, where $\gamma_+^{\sigma'}$ is the dual value of γ_+^{σ} . By continuity, since $C(p, \gamma_+^{\sigma}) > 0$ to the right of \mathcal{G} on the $q = 0$, $\sigma = -1$ line, and $C(p, \gamma_+^{\sigma'}) < 0$ to the left of \mathcal{R} on the same line, the generalized proof conditions $C(p, \gamma_+^{\sigma}) > 0$ and $C(p, \gamma_+^{\sigma'}) < 0$ simultaneously hold in the domain $\tilde{\mathcal{D}}$ of the (p, q) -plane, as defined in Proposition 5.18.

Lastly, the rationale for the occurrence of $\min\{\hat{p}(\kappa), p(\kappa)\}$ in [14, Theorem 1.4] was the technical requirement that the dual value γ'_+ of γ_+ is such that there is no dual tip contribution, i.e., that it satisfies $2\gamma'_+ + 1 \geq 0$, yielding the auxiliary condition $p \leq \hat{p}(\kappa)$. Since $\sigma = q/p - 1$, we simply have $\gamma_{\pm}^{\sigma}(p) = \gamma_{\pm}(p - q)$, and the condition in the (p, q) -plane is now for a general value of σ ,

$$(140) \quad p - q \leq \hat{p}(\kappa) = 1 + \frac{\kappa}{2}.$$

This restricts the domain of validity of the argument to the part $\hat{\mathcal{D}}$ of $\tilde{\mathcal{D}}$ located above the straight line D_3 of equation $p - q = \hat{p}(\kappa)$ (Fig. 6), and concludes the proof of the validity of the spectrum $\beta_1(p, q)$ (78) in $\hat{\mathcal{D}}$. \square

Proof of Theorem 1.7.

Proof. Propositions 5.6 in item *i*), 5.7 in *ii*), 5.9 and its corollary 5.10 in *iii*), and 5.18 in *iv*), altogether establish the validity of the various spectra given in Theorem 1.7 and Fig. 2, in the following maximal domain of the (p, q) -plane. It is in Figure 6 the infinite domain above the frontier line made by the union of the lower branch of the green parabola up to its intersection point with D_3 , obtained for $\gamma' = -1$ in Eq. (97) as $((4 - \kappa)(4 + 3\kappa)/8\kappa, (16 - 7\kappa^2)/8\kappa)$, of the segment of D_3 from that point up to P_3 , of the branch of the red parabola \mathcal{R} between P_3 and P_0 , and finally of D_1 from P_0 up to infinity. This concludes the proof of Theorem 1.7. \square

5.5. m -fold spectrum. For $m \geq 1$, the generalized integral means spectrum $\beta^{[m]}(p, q; \kappa)$, associated with the m -fold transform $f^{[m]}$ of the SLE whole-plane map, can be directly derived from the analysis given in Section 4.4. Definition 1.4 and identities (59) and (60) immediately imply that

$$(141) \quad \begin{aligned} \beta^{[m]}(p, q; \kappa) &= \beta^{[1]}(p, q_m; \kappa), \\ q_m &= q_m(p, q) = (1 - 1/m)p + q/m, \end{aligned}$$

where $\beta^{[1]}(p, q; \kappa) := \beta(p, q; \kappa)$ is the $m = 1$ average generalized integral means spectrum of whole-plane SLE_κ studied above. As we shall see in Section 5.5.3, Definition 1.4 can be formally extended to $m \leq -1$, and the resulting $\beta^{[m]}$ still obeys (141).

5.5.1. Standard m -fold integral means spectrum. Let us first focus on the m -fold *standard* spectrum, i.e., on the $q = 0$ case, for which

$$(142) \quad q_m(p, 0) = (1 - 1/m)p, \quad m \in \mathbb{Z} \setminus \{0\}, p \in \mathbb{R}.$$

This defines a line $q = q_m(p, 0)$ in the original (p, q) plane. For the β_1 part of the spectrum, this yields

$$(143) \quad \beta_1(p, (1 - 1/m)p; \kappa) = \left(1 + \frac{2}{m}\right)p - \frac{1}{2} - \frac{1}{2}\sqrt{1 + \frac{2\kappa p}{m}},$$

in agreement with the result obtained in [14, Eq. (22)]. Note, however, that for $m = -2$ this function is negative and cannot appear in the spectrum.

Let us consider the point P_0 whose coordinates (p_0, q_0) are given by (100). The line OP_0 has slope $\frac{q_0}{p_0} = 1 - \frac{1}{3} \frac{\kappa - 4}{\kappa + 4}$, a quantity that decreases from $4/3$ to $2/3$ as κ runs from 0 to ∞ and takes the value 1 for $\kappa = 4$.

In the case $\kappa > 4$, it is thus possible to find $m > 0$ such that $\frac{q_0}{p_0} \leq 1 - \frac{1}{m} < 1$. This holds for $m \geq m(\kappa)$, with the definition,

$$(144) \quad \begin{aligned} m(\kappa) &:= 3 \frac{\kappa + 4}{\kappa - 4}; \quad m(16/\kappa) = -m(\kappa), \\ \kappa > 4, m(\kappa) &> 3; \quad 0 < \kappa < 4, m(\kappa) < -3. \end{aligned}$$

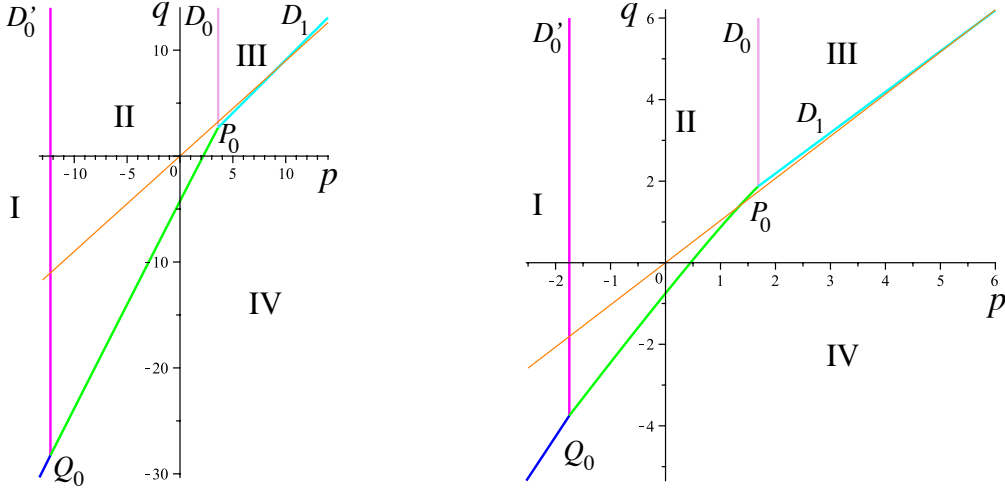


FIGURE 8. Phase diagram for the m -fold whole-plane SLE_κ in the (p, q_m) -plane, along trajectories (142) (coral color). Left: For $m = +10$ and $\kappa = 30$, the line successively crosses domains I, II, III and IV. Right: For $m = -30$ and $\kappa = 2$, it successively crosses domains I, II, IV and III.

The line (142) then first intersects the vertical line $p = p_0$ above point P_0 , and the line D_1 of unit slope afterwards (Fig. 8, left). So in this case the standard m -fold integral means spectrum $\beta^{[m]}(p) := \beta^{[m]}(p, 0)$, for $p \in \mathbb{R}$, has four phases, $\beta_{\text{tip}}, \beta_0, \beta_{\text{lin}}$ and β_1 (143). In the second case, $\kappa < 4$, it is possible to find $m < 0$ such that $1 < 1 - \frac{1}{m} \leq \frac{q_0}{p_0}$, which is equivalent to $m \leq m(\kappa) < -3$. The line (142) then crosses the line $p = p_0$ at a point below point P_0 . It follows that it first crosses the green parabola and afterwards crosses the line D_1 , after which the spectrum becomes linear (Fig. 8, right). So the standard m -fold spectrum again has four phases, but now in the order $\beta_{\text{tip}}, \beta_0, \beta_1$ (143) and β_{lin} for $p \in \mathbb{R}$. For the critical case $\kappa = 4$, $\frac{q_0}{p_0} = 1$, and for $m \geq 1$, we observe the ordered set $\beta_{\text{tip}}, \beta_0, \beta_1$, or for $m \leq -1$, the set $\beta_{\text{tip}}, \beta_0$ and β_{lin} .

We finally obtain the following table for the successive expressions taken by $\beta^{[m]}(p)$ for $p \in \mathbb{R}$.

$\kappa > 4$	$m \leq -1$	$1 \leq m \leq m(\kappa)$	$3 < m(\kappa) \leq m$
	$\beta_{\text{tip}}, \beta_0, \beta_{\text{lin}}$	$\beta_{\text{tip}}, \beta_0, \beta_1$	$\beta_{\text{tip}}, \beta_0, \beta_{\text{lin}}, \beta_1$
$\kappa = 4$	$m \leq -1$	$1 \leq m$	
	$\beta_{\text{tip}}, \beta_0, \beta_{\text{lin}}$	$\beta_{\text{tip}}, \beta_0, \beta_1$	
$\kappa < 4$	$m \leq m(\kappa) < -3$	$m(\kappa) \leq m \leq -1$	$1 \leq m$
	$\beta_{\text{tip}}, \beta_0, \beta_1, \beta_{\text{lin}}$	$\beta_{\text{tip}}, \beta_0, \beta_{\text{lin}}$	$\beta_{\text{tip}}, \beta_0, \beta_1$

Remark 5.19. *In the three cases $m \in \{-1, -2, -3\}$, and for any $\kappa > 0$, only the usual spectra β_{tip} , β_0 , and β_{lin} appear in the m -fold standard integral means spectrum $\beta^{[m]}(p)$.*

The discussion shows that the spectrum β_1 (at (p, q_m)) (143) may appear in $\beta^{[m]}(p)$ even if the boundary of the SLE image domain is *bounded*. This indeed happens for $m \leq -4$ and κ small enough: for negative m , the m -fold transform of the outer whole-plane SLE is then conjugate by $z \mapsto 1/z$ of the $(-m)$ -fold transform of the inner whole-plane SLE, which gives rise to a univalent function map onto a domain with bounded boundary. In this case, the appearance of the β_1 spectrum is due to a higher $|m|$ -fold branching at the origin for $\kappa < 4$ and $|m| \geq |m(\kappa)|$.

5.5.2. *Generalized m -fold integral means spectrum ($m \geq 1$).* Let now $T_m = \begin{pmatrix} 1 & 0 \\ -1/m & 1/m \end{pmatrix}$ be the endomorphism of \mathbb{R}^2 defined by $T_m(p, q) = (p, q_m)$, with inverse $T_m^{-1} = \begin{pmatrix} 1 & 0 \\ 1-m & m \end{pmatrix}$. The separatrix lines for the m -fold case are the images by T_m^{-1} of those for $m = 1$. Theorem 1.7 then yields:

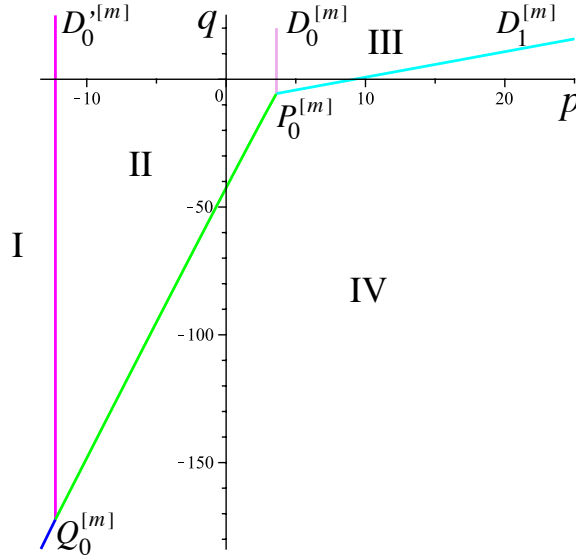


FIGURE 9. Phase diagram for the m -fold whole-plane SLE_κ and domains of validity of spectra β_{tip} (I), β_0 (II), β_{lin} (III), and β_m (IV). For $m = +10$ and $\kappa = 30$, the $q = 0$ line successively crosses domains I, II, III and IV.

Theorem 5.20. (Figure 9) *Separatrix curves for the generalized integral means spectrum $\beta^{[m]}(p, q; \kappa)$ of the m -fold whole-plane SLE_κ are given, for $m \geq 1$, by the same as in Theorem 1.7 for $m = 1$, provided that one replaces there,*

- D_0 by $D_0^{[m]}$, P_0 by $P_0^{[m]} = (p_0, q_0^{[m]})$, q_0 by $q_0^{[m]} := p_0 + m(16 - \kappa^2)/32\kappa$;

- D_1 by $D_1^{[m]}$ of equation $q - p = m(16 - \kappa^2)/32\kappa$;
- $q_{\mathcal{G}}(\gamma)$ by $p_{\mathcal{G}}(\gamma) + m(\gamma - \kappa\gamma^2/2)$, Q_0 by $Q_0^{[m]} = (p'_0, q_0^{[m]})$, with $q_0^{[m]} := p'_0 - m(1 + \kappa/2)$;
- D'_0 by the vertical half-line $D_0^{[m]}$ above $Q_0^{[m]}$;
- $q_{\mathcal{Q}}(\gamma)$ by $p_{\mathcal{Q}}(\gamma) + m(\gamma - \kappa\gamma^2/2)$.

The figures of the m -separatrices are easily deduced from Theorem 5.20; in particular, the transformed quartic $\mathcal{Q}^{[m]}$ is asymptotic for $p \rightarrow -\infty$ to a straight-line,

$$(145) \quad q = (m + 1)p - m(2 + \kappa)/8,$$

whose direction is also that of the axis of the transformed green parabola $\mathcal{G}^{[m]}$. In region IV, the m -fold integral means spectrum is thus given by

$$(146) \quad \beta^{[m]}(p, q; \kappa) = \beta_1(p, q_m; \kappa) = \left(1 + \frac{2}{m}\right)p - \frac{2}{m}q - \frac{1}{2} - \frac{1}{2}\sqrt{1 + \frac{2\kappa}{m}(p - q)}.$$

5.5.3. *The $m \leq -1$ case.* In this range, $f^{[m]}$ is an exterior map from \mathbb{D}_- , whose continuation to $\partial\mathbb{D}$ vanish on $-m$ points, and the image domain is the plane slit by $-m$ curves in a star configuration at the origin. When defined via a formal extension of Definition 1.4, this generalized means spectrum $\beta^{[m]}(p, q; \kappa)$ now describes for $q > 0$ the singular behavior of the map near $(f^{[m]})^{-1}(0)$. It still obeys identity (141), and the same conclusions as in Theorem 5.20 hold formally, except that, because $\det T_m^{-1} = m < 0$, the vertical positions of the respective domains of validity of the spectra are all *reversed*, the vertical separatrix lines $D_0^{[m]}$ and $D_1^{[m]}$ being now half-lines going downwards from $P_0^{[m]}$ and $Q_0^{[m]}$ to $-\infty$, and the domain IV lying *above* the half-line $D_1^{[m]}$, the transformed green parabola and the transformed quartic. The concavity of the separatrix curves is correspondingly inverted (Fig. 10).

Note the slope inversion in (145) when going from $m \geq 1$ to $m \leq -1$. The $m = -1$ case, i.e., the B–S exterior version of whole-plane SLE, is thus peculiar: the parabola’s axis and the quartic’s linear asymptote are both horizontal, and $\mathcal{G}^{[-1]}$ intersects the p -axis at $p''_0(\kappa) \leq p'_0(\kappa)$, in agreement with Section 5.3.2. (See Fig. 10, left.) As hinted above, the $m = -2$ case is special, and one can check that the green parabola $\mathcal{G}^{[-2]}$ stays strictly above the p -axis for any $\kappa > 0$. Hence, this axis never enters region IV, rendering the (degenerate) negative spectrum (143) irrelevant.

In Theorem 5.20, the coordinate $q_0^{[m]}$ of point $P_0^{[m]}$ can become *negative*. This corresponds to the fact that the $q = 0$ axis intersects all four different regions in the m -fold spectrum (Figs. 9 and 10). This happens for $\frac{m}{3} \frac{\kappa - 4}{\kappa + 4} \geq 1$, which gives either $m \geq m(\kappa) > 3$ for $\kappa > 4$ or $m \leq m(\kappa) < -3$ for $\kappa < 4$, in agreement with the discussion of Fig. 8 above.

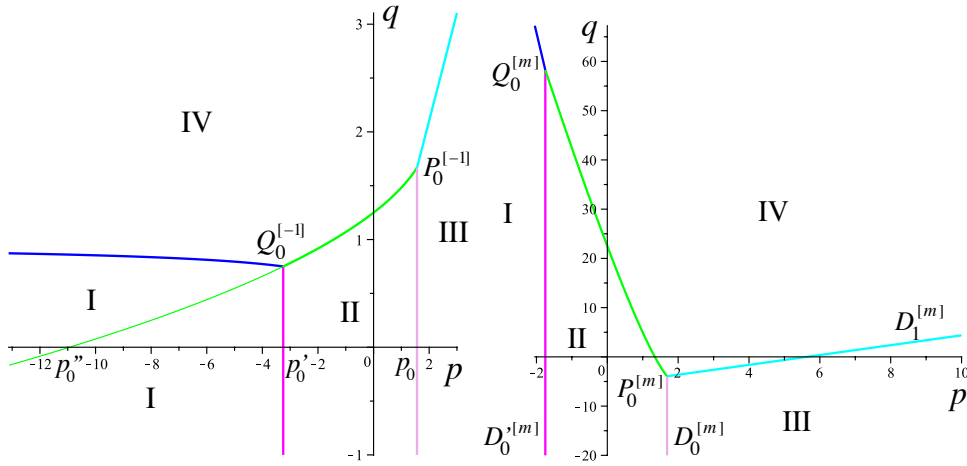


FIGURE 10. *Left: Phase diagram for $m = -1$ and $\kappa = 6$. Right: Phase diagram for $m = -30$ and $\kappa = 2$; the $q = 0$ line successively crosses domains I, II, IV and III.*

6. UNIVERSAL SPECTRUM

It is worthwhile to compare the above results to universal ones. We aim at generalizing the universal spectrum for integral means of derivatives of univalent functions to the case of mixed integrals of the type $\int_0^{2\pi} \frac{|f'(re^{i\theta})|^p}{|f(re^{i\theta})|^q} d\theta$. More precisely, for f injective and holomorphic in the unit disk, we define $\beta_f(p, q)$ as being the smallest number such that

$$\int_0^{2\pi} \frac{|f'(re^{i\theta})|^p}{|f(re^{i\theta})|^q} d\theta \leq O(1-r)^{-\beta_f(p,q)-\varepsilon}, \forall \varepsilon > 0, r \rightarrow 1.$$

The *universal spectrum* $B(p, q)$ is then defined as the supremum of $\beta_f(p, q)$ over all holomorphic and injective f 's on the disk.

It should be first noticed that if one restricts oneself to bounded univalent functions, there will be no change with respect to the usual integral means spectrum, i.e., the denominator in the integrand (and thus q) plays no role in this case. In the general case, we will mimic the Feng-McGregor approach [15].

Theorem 6.1. *Let f be holomorphic and injective in the unit disk. For $p \in \mathbb{R}^+$, $q \in \mathbb{R}$ such that $q < \min\{2, \frac{5}{4}p - \frac{1}{2}\}$, there exists a constant $C > 0$ such that*

$$(147) \quad \int_0^{2\pi} \frac{|f'(re^{i\theta})|^p}{|f(re^{i\theta})|^q} d\theta \leq \frac{C}{(1-r)^{3p-2q-1}}.$$

The universal spectrum is therefore finite and such that $B(p, q) \leq 3p - 2q - 1$, at least in the domain $\mathcal{D}_0 := \{0 \leq p, q < \min\{2, \frac{5}{4}p - \frac{1}{2}\}\}$ of Theorem 6.1. In that domain, the Koebe function, $\mathcal{K}(z) = z(1+z)^{-2}$, saturates the bound, therefore $B(p, q) = 3p - 2q - 1 > 0$ for $(p, q) \in \mathcal{D}_0$.

In order to make the proof lighter, we will neither write the variables in the functions involved, which are of the form $re^{i\theta}$ with r fixed, nor the angular integration interval, which is meant to be $[0, 2\pi]$.

Proof. Let a, b be two reals, to be fixed later, such that $a - b = 1$. Let us first consider the case $p < 2$, for which Hölder's inequality gives

$$(148) \quad \int \frac{|f'|^p}{|f|^q} = \int \frac{|f'|^p}{|f|^{aq}} |f|^{bq} \leq \left(\int \frac{|f'|^2}{|f|^{2aq/p}} \right)^{p/2} \left(\int |f|^{2bq/(2-p)} \right)^{(2-p)/2}.$$

In order to estimate the first integral on the right-hand side, we invoke Hardy's inequality [34]: For any $p' > 0$, there exists a constant $C' > 0$ such that for any function f which is holomorphic and injective in the unit disk,

$$\int |f'|^2 |f|^{p'-2} \leq \frac{C'}{(1-r)^{2p'+1}}.$$

For the rightmost integral in (148), we use the Prawitz inequality [34]: For any $p'' > 1/2$, there exists a constant $C'' > 0$ such that for any function f holomorphic and injective in the unit disk,

$$\int |f|^{p''} \leq \frac{C''}{(1-r)^{2p''-1}}.$$

We then take $p' := 2 - \frac{2aq}{p}$, $p'' := \frac{2bq}{2-p}$, and assume that $p' > 0$ and $p'' > 1/2$; we may then use the two inequalities above and get from (148)

$$(149) \quad \int \frac{|f'|^p}{|f|^q} \leq \frac{C}{(1-r)^{3p-2q-1}},$$

for some $C > 0$ and any f as above. For this, we need to find $a, b \in \mathbb{R}$ such that $a - b = 1$, $p' > 0$, $p'' > 1/2$. The first inequality is equivalent to $p > aq$, and the second one gives $aq > q + \frac{2-p}{4}$. We thus find that the universal bound (149) holds for $q + \frac{1}{2} < \frac{5}{4}p$. Recall then the original condition of validity, $p < 2$, which implies that $q < 2$. The theorem being already proved for $p < 2$, we may now assume that $p \geq 2$. Let then p' be such that $\frac{4}{5}q + \frac{2}{5} < p' < 2 \leq p$. We now invoke Koebe distortion theorem:

$$\forall z \in \mathbb{D}, |f'(z)| \leq 2 \frac{|f'(0)|}{(1-|z|)^3},$$

from which follows, by writing $|f'|^p = |f'|^{p'} |f'|^{p-p'}$ and by using (147) for the couple (p', q) , that for some $C > 0$,

$$(150) \quad \int \frac{|f'|^p}{|f|^q} \leq \frac{C}{(1-r)^{3p'-2q-1+3(p-p')}} = \frac{C}{(1-r)^{3p-2q-1}}.$$

□

Guided by the results obtained above for the generalized integral means spectrum of whole-plane SLE, we will now state a conjecture concerning the universal generalized spectrum. As we shall see, its structure turns out to be very similar, each of the SLE four spectra having its own analogue in the universal case.

Let us first recall that the universal spectrum for **bounded** holomorphic and injective functions, $B(p)$, is known to be equal to $p - 1$ for $p \geq 2$, and equal to $-p - 1$ below a certain threshold $p^\dagger \leq -2$. For $p \in [p^\dagger, 2]$, it is equal to a unknown function, $B_0(p)$. Two famous conjectures are that by Brennan, stating that $B_0(-2) = 1$ and implying that $p^\dagger = -2$, and the broader conjecture by Kraetzer stating that $B_0(p) = p^2/4$ (see Ref. [34] and references therein).

For **unbounded** functions, a classical result by Makarov [32] states that the universal spectrum is simply given by

$$(151) \quad \max\{B(p), 3p - 1\},$$

the second term corresponding to the extremal case of the Koebe function.

Now, in the case of generalized spectra, the universal analogue of the SLE generalized spectrum $\beta_1(p, q; \kappa)$ is naturally the spectrum that we have just obtained in Theorem 6.1, and that corresponds to the Koebe limit of Section 5.3.3,

$$(152) \quad B_1(p, q) := 3p - 2q - 1.$$

The analogue of the SLE bulk spectrum, $\beta_0(p)$, is then naturally given by the function $B_0(p)$ of the bounded universal spectrum above, while the two remaining SLE spectrum functions, $\beta_{\text{tip}}(p)$ and $\beta_{\text{lin}}(p)$, have respectively for universal analogues, $B_{\text{tip}}(p) := -p - 1$ for $p \leq p^\dagger$, and $B_{\text{lin}}(p) := p - 1$ for $p \geq 2$.

We then proceed as for SLE, looking for the sets of points in the (p, q) plane such that $B_1(p, q) = B_{\text{tip}}(p)$, $B_1(p, q) = B_0(p)$, $B_1(p, q) = B_{\text{lin}}(p)$. They turn out to be, in the same order,

- the line $q = 2p$ for $p \leq p^\dagger$,
- the curve $2q = 3p - 1 - B_0(p)$ for $p \in [p^\dagger, 2]$,
- the line $p = q$ for $p \geq 2$.

Note that if Brennan's conjecture holds, $p^\dagger = -2$, and it is equivalent to the fact that the separatrix curve, $2q = 3p - 1 - B_0(p)$, the vertical line, $p = -2$, and the separatrix, $q = 2p$, all meet at point $(-2, -4)$. If Kraetzer's conjecture also holds, the first curve becomes the segment of parabola $2q = 3p - 1 - p^2/4$, with $p \in [-2, 2]$.

In complete analogy with the SLE case (see Fig. 2), we thus obtain a prediction for the universal spectrum $B(p, q)$, with a partition of the plane into four zones corresponding to the four spectra introduced above, as illustrated in Fig. 11.

Observe that the figure contains, for $q = 0$, the universal spectrum for all univalent functions, as well as, along the line $q = 2p$, the spectrum of bounded univalent ones. For $p \leq p^\dagger$ (possibly (-2)), the latter line also appears as a *separatrix* of the (conjectured) universal spectrum. A small departure from it

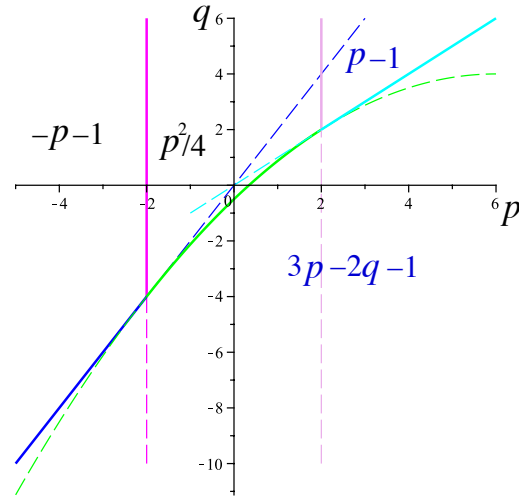


FIGURE 11. The four functions giving the universal generalized spectrum (assuming here the validity of Kraetzer’s conjecture).

triggers a phase transition in the spectrum, which is thus *unstable* along the bounded functions line.

As work done with Kari Astala shows [1], it is actually possible to extend Makarov’s approach [32] to the universal generalized spectrum $B(p, q)$, and to generalize result (151) into

Theorem 6.2. *The universal generalized spectrum is given by*

$$B(p, q) = \max\{B(p), 3p - 2q - 1\},$$

where $B(p)$ is the universal spectrum for bounded univalent functions.

This confirms the conclusions drawn above for the universal generalized spectrum, the unknown remaining the position of p^\dagger and the form of $B_0(p)$ in the standard universal spectrum.

REFERENCES

- [1] K. Astala, B. Duplantier, and M. Zinsmeister, 2015. Unpublished manuscript.
- [2] D. Beliaev, B. Duplantier, and M. Zinsmeister. Integral Means Spectrum of Whole-Plane SLE. 2015. arXiv:1605.03112. *Commun. Math. Phys.* <http://link.springer.com/article/10.1007/s00220-017-2868-z>.
- [3] D. Beliaev and S. Smirnov. Harmonic Measure and SLE. *Commun. Math. Phys.*, 290:577–595, 2009.
- [4] A Belikov, I A Gruzberg, and I. I Rushkin. Statistics of harmonic measure and winding of critical curves from conformal field theory. *J. Phys. A: Math. Theor.*, 41(28):285006, 2008.
- [5] E. Bettelheim, I. Rushkin, I. A. Gruzberg, and P. Wiegmann. Harmonic Measure of Critical Curves. *Phys. Rev. Lett.*, 95:170602, 2005.
- [6] L. Bieberbach. Über die Koeffizienten derjenigen Potenzreihen, welche eine schlichte Abbildung des Einheitskreises vermitteln. *S-B. Preuss. Akad. Wiss.*, 1:940–955, 1916.

- [7] I. Binder and B. Duplantier. Multifractal properties of harmonic measure and rotation for SLE, 2017. In preparation.
- [8] L. de Branges. A proof of the Bieberbach conjecture. *Acta Math.*, 154:137–152, 1985.
- [9] B. Duplantier. Conformally Invariant Fractals and Potential Theory. *Phys. Rev. Lett.*, 84:1363–1367, 2000.
- [10] B. Duplantier. Conformal fractal geometry & boundary quantum gravity. In M. L. Lapidus and M. van Frankenhuysen, editors, *Fractal geometry and applications: a jubilee of Benoît Mandelbrot, Part 2*, volume 72 of *Proc. Sympos. Pure Math.*, pages 365–482. Amer. Math. Soc., Providence, RI, 2004.
- [11] B. Duplantier and I. A. Binder. Harmonic Measure and Winding of Conformally Invariant Curves. *Phys. Rev. Lett.*, 89:264101, 2002.
- [12] B. Duplantier and I. A. Binder. Harmonic measure and winding of random conformal paths: A Coulomb gas perspective. *Nucl. Phys. B [FS]*, 802:494–513, 2008.
- [13] B. Duplantier, Nguyen Thi Phuong Chi, Nguyen Thi Thuy Nga, and M. Zinsmeister. Coefficient estimates for whole-plane SLE processes. July 2011. <http://hal.inria.fr/hal-00609774>.
- [14] Bertrand Duplantier, Chi Nguyen, Nga Nguyen, and Michel Zinsmeister. The Coefficient Problem and Multifractality of Whole-Plane SLE and LLE. *Ann. Henri Poincaré*, 16(6):1311–1395, 2014. <http://arxiv.org/pdf/1211.2451v2.pdf>.
- [15] J. Feng and T. H. MacGregor. Estimates on the integral means of the derivatives of univalent functions. *J. Anal. Math.*, 29:203–231, 1976.
- [16] John B. Garnett and Donald E. Marshall. *Harmonic Measure*. Cambridge University Press, 2005.
- [17] H. Grunsky. Koeffizienten Bedingungen für schlicht abbildende meromorphe Funktionen. *Math. Z.*, 45:29–61, 1939.
- [18] E. Gwynne, J. Miller, and X. Sun. Almost sure multifractal spectrum of SLE. December 2014. arXiv:1412.8764.
- [19] M. B. Hastings. Exact Multifractal Spectra for Arbitrary Laplacian Random Walks. *Phys. Rev. Lett.*, 88:055506, 2002.
- [20] F. Johansson Viklund and G.F. Lawler. Almost sure multifractal spectrum for the tip of an SLE curve. *Acta Math.*, 209(2):265–322, 2012.
- [21] A. Kemppainen. Stationarity of SLE. *J. Stat. Phys.*, 139:108–121, 2010.
- [22] K. Kytölä and A. Kemppainen. SLE local martingales, reversibility and duality. *J. Phys. A: Math. Gen.*, 39:L657–L666, 2006.
- [23] G. F. Lawler, O. Schramm, and W. Werner. Values of Brownian intersection exponents. II. Plane exponents. *Acta Math.*, 187(2):275–308, 2001.
- [24] G. F. Lawler, O. Schramm, and W. Werner. Conformal invariance of planar loop-erased random walks and uniform spanning trees. *Ann. Probab.*, 32(1B):939–995, 2004.
- [25] Thanh Binh Le. Around Milin’s conjecture and SLE maps, 2010. Mémoire de M2, Université d’Orléans.
- [26] N. A. Lebedev and I. M. Milin. On the coefficients of certain classes of univalent functions. *Mat. Sb.*, 28:359–400, 1951. (In Russian).
- [27] I. Loutsenko. SLE $_{\kappa}$: correlation functions in the coefficient problem. *J. Phys. A Math. Theor.*, 45(26):265001, 2012.
- [28] I. Loutsenko and O. Yermolayeva. On exact multi-fractal spectrum of the whole-plane SLE. 2012. arXiv:1203.2756.
- [29] I. Loutsenko and O. Yermolayeva. Average harmonic spectrum of the whole-plane SLE. *J. Stat. Mech.*, page P04007, 2013.
- [30] I. Loutsenko and O. Yermolayeva. New exact results in spectra of stochastic Loewner evolution. *J. Phys. A Math. Theor.*, 47(16):165202, 2014.

- [31] K. Löwner. Untersuchungen über schlichte konforme Abbildungendes Einheitskreises. *Math. Annalen*, 89:103–121, 1923.
- [32] N. G. Makarov. Fine structure of harmonic measure. *Rossiiskaya Akademiya Nauk. Algebra i Analiz*, 10:1–62, 1998. English translation in *St. Petersburg Math. J.* 10: 217–268 (1999).
- [33] I. M. Milin. Estimation of coefficients of univalent functions. *Dokl. Akad. Nauk SSSR*, 160:196–198, 1965.
- [34] Ch. Pommerenke. *Boundary Behaviour of Conformal Maps*. Grundlehren der mathematischen Wissenschaften, Vol. 299. Springer, Berlin, 1992.
- [35] M. S. Robertson. On the theory of univalent functions. *Ann. of Math*, 37:374–408, 1936.
- [36] S. Rohde and O. Schramm. Basic Properties of SLE. *Ann. of Math.*, 161:883–924, 2005.
- [37] I. Rushkin, E. Bettelheim, I. A. Gruzberg, and P. Wiegmann. Critical curves in conformally invariant statistical systems. *J. Phys. A: Math. Gen.*, 40:2165–2195, 2007.
- [38] O. Schramm. Scaling limits of loop-erased random walks and uniform spanning trees. *Israel J. Math*, 118:221–288, 2000.

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