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1 **The new international reference system for pure alpha-**
2 **and pure beta-emitting radionuclides**

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18 *Abstract*

19 The Bureau International des Poids et Mesures (BIPM) is developing with the support of
20 national metrology institutes an extension to the existing international reference system
21 of radionuclide metrology to address pure alpha-, pure beta-emitting and some electron
22 capture-decaying radionuclides. The new service is based on a liquid scintillation counter

23 and the triple to double coincidence (TDCR) technique. A dedicated facility to prepare
24 liquid scintillation sources and a TDCR system have been set up at the BIPM. The TDCR
25 system was conceived and commissioned in order to ensure the long-term reproducibility
26 of the reference it produces. The metrology is based on a dedicated approach which uses
27 the TDCR value as a parameter to stabilize the output quantities of the system against
28 changes in detection efficiency or asymmetry that could appear in a long-term period of
29 operation. The quality assurance based on periodic controls of toluene-based liquid
30 scintillation sources of ^3H and ^{14}C has permitted to demonstrate the reproducibility of the
31 system for a 20-month period.

32 **Keywords**

33 Radionuclide Metrology, Key comparisons, Liquid scintillation, TDCR, ESIR

34 **1. Introduction**

35 The Bureau International des Poids et des Mesures (BIPM) operates an international
36 reference system (the SIR) to compare standards of radioactivity realized by National
37 Metrology Institutes / Designated Institutes (NMI/DIs), to enable NMI/DIs to
38 demonstrate the equivalence of standards leading to a harmonized international
39 measurement system. The operation of the international measurement system is described
40 in more detail in a Mutual Recognition Arrangement, the CIPM MRA [1,2].

41 The SIR is a re-entrant ionization chamber; the output value is the ratio of the
42 ionization current measured with the source to be analyzed to the current produced by
43 long-lived sealed sources of ^{226}Ra [3–5]. In practice, the long-term reproducibility and the
44 precision achieved by this relative measurement allow to assess equivalences for most of
45 the gamma-emitting radionuclides. To cope with the increasing demand in nuclear
46 medicine in short-lived gamma-emitting radionuclides (e.g. $^{99\text{m}}\text{Tc}$, ^{18}F , ^{64}Cu , ^{11}C ...), a
47 transportable version called the SIRTI is also available [6]. The well-type NaI(Tl)
48 detector operates using as a reference: a long-lived sealed source of ^{94}Nb . It is linked to

49 the SIR and can travel to laboratories far from the BIPM to obtain equivalences for these
50 radionuclides. The main asset of the SIR/SIRTI services for NIM/DIs is to avoid
51 organizing large-scale comparison exercises implying send radioactive sources around
52 the world. The SIR is an on-demand service which permits the NIM/DIs to obtain
53 equivalence whenever they desire. Moreover, the efficiency in delivering comparison
54 results has been recently enhanced through a new database and automated reporting [7].

55 There are still some radionuclides that are not covered by the SIR/SIRTI services of
56 the BIPM: pure (or quasi-pure) alpha-particle emitters, pure (or quasi-pure) beta-particle
57 emitters, and some electron-capture decaying radionuclides (such as ^{55}Fe). The ionizing
58 particles emitted during the decay of these radionuclides cannot pass (or very weakly)
59 through the entry windows of the detectors and produce a sufficient signal. Meantime,
60 there is an increasing demand to evaluate equivalence of radioactive standards containing
61 these types of isotopes. New cancer treatments are emerging using radiopharmaceuticals
62 tagged by radionuclides decaying in a way to deposit very locally the dose [8–10]. Hence,
63 Targeted Beta-particle Therapy (TBT) and Targeted Alpha-particle Therapy (TAT) aim
64 to efficiently damage cancerous cells while sparing healthy tissues. Radionuclides that
65 are used (or studied to be used) in TBT are for instance: ^{32}P , ^{89}Sr , ^{90}Y , ^{169}Er , ^{177}Lu and the
66 ^{186}Re . Radionuclides that are used (or studied to be used) in TAT are for instance: ^{225}Ac ,
67 ^{211}At , ^{212}Bi , ^{212}Pb , ^{213}Bi , ^{227}Th and the ^{223}Ra . Some of these isotopes, such as the ^{177}Lu
68 used in the theragnostic perspective [11], emits enough x or gamma radiations to be
69 measured by the SIR/SIRTI [12–14]. However, it is not possible for most of them. In
70 addition to that, the need to assess the international equivalence of standards of pure
71 alpha- or beta-emitting radionuclides arises also for environmental issues where for
72 instance ^{89}Sr , ^{90}Y , ^{160}Tb , ^{161}Tb and ^{210}Po are of critical use for nuclear forensics [15–17].

73 Liquid scintillation counting appears to be the most suitable measurement techniques
74 to address the measurement issue associated with these radionuclides. This is why, the
75 extension of the SIR (called the Extended SIR (ESIR)) is under development since 2018
76 and based on a liquid scintillation counter and an approach derived from the primary
77 standardization technique called: Triple to Double Coincidence Ratio counting

78 (TDCR) [18,19]. On contrary to the SIR and SIRTl which achieve a long-term
79 reproducibility using long-lived radioactive sources and a relative measurement, the
80 ESIR implements another way to ensure reproducibility since the long-term stability of a
81 liquid scintillation source cannot be assumed [20] even for the nonaqueous ones (see
82 section 5). So, it has been decided to take benefit from the *TDCR* parameter which is
83 strongly correlated with the detection efficiency of the system to realize a self-
84 stabilization of the system. The output values delivered by the ESIR are comparison
85 indicators, I , stabilized using the *TDCR* value in different ways such as:

- 86 • I_0 is the non-stabilized comparison indicator,
- 87 • I_1 is the comparison indicator stabilized by a 2nd order polynomial function of
88 the *TDCR* value (built for robustness against efficiency changes),
- 89 • I_2 is the comparison indicator stabilized by the calculation of an equivalent
90 efficiency based on the TDCR model parametrized by fixed parameters (built
91 for robustness against both efficiency and asymmetry changes),
- 92 • I_3 is the comparison indicator obtain by making a monoenergetic assumption
93 in the TDCR model (it can be used for isotopes emitting close-to-
94 monoenergetic particles such as the ^{55}Fe).

95 Details on the calculation of these ESIR comparison indicators are given in [21–23].

96 In addition to this unique approach, it is also critical for this new service:

- 97 • to develop and maintain a source preparation facility with an accurate mass
98 measurement of active solutions,
- 99 • to develop and maintain a TDCR system suitable for long-term operations,

- 100 • to develop and maintain a proper associated quality assurance to monitor the long-
101 term reproducibility of the measurement system.

102

103 **2. The liquid scintillation source preparation facility**

104 To prepare liquid scintillation (LS) sources at the BIPM, aliquotes of an active solution
105 standardized by a laboratory are injected into LS vials containing LS cocktail.

106 *a. The liquid scintillation cocktail*

107 It was decided to use the Ultima-Gold (UG) cocktail from Perkin Elmer – (or equivalent).
108 It permits to properly dissolve the aqueous active solution while providing a good
109 scintillation yield. Ionisation quenching depends strongly on the properties of the LS
110 cocktail and it cannot be assumed a perfect reproducibility from batch to batch of UG
111 cocktails. However, as ionisation quenching effects are of relative minor importance for
112 medium and high energy depositions into the LS liquid, it can be assumed that UG (or
113 similar type of cocktails) will fit the purpose for the majority of radionuclides to address.

114 Nevertheless, the reproducibility of the ESIR when using UG cocktail could remains an
115 issue for low energy-deposited particles encountered for instance during decays of ^3H ,
116 ^{63}Ni , ^{55}Fe or ^{241}Pu . For these latter ones, the duration of key comparison exercises will be
117 limited to the validity period of a given batch of UG. This option imposes constraints on
118 laboratories as they cannot send their solutions when appropriate for them but only when
119 a comparison is planned and during a limited period. The second option to ensure on-
120 going comparisons as it is for the SIR, is to produce a specific LS cocktail at the BIPM
121 under reproducible and sustainable conditions. A preliminary study has permitted, through
122 a factorial design approach, to establish the preliminary composition of an LS cocktail
123 that can be reproducibly realised at the BIPM. The components are notably chosen for
124 not being impacted by the European REACH regulation. Thus, the diisopropylnaphthalen

125 entering as a solvent in UG composition is replaced by pseudocumen (1,2,4-
126 trimethylbenzene) and the anionic surfactant 4-NPE is replaced by bis(2-
127 ethylhexyl)hydrogene-phosphate. Also, to enhance the stability of the LS source in
128 regards with the variation of PH of active solutions, the primary fluorophore 2,5-diphenyl
129 oxazole (PPO) will be replaced by the more robust 1¹,2¹:2⁴:3¹-terphenyle (p-terphenyle).
130 Further studies will be realized to continue investigations on this option.

131 *b. The vials*

132 Glass vials will be the default type of vials used in the ESIR. However, as emphasis by
133 the recent CCRI(II).Fe-55 comparison [24], polyethylene or frosted glass vials are better
134 options to address ⁵⁵Fe and probably other radionuclides emitting low energy particles.
135 Offering a better diffusion of photons in conformance with the TDCR model,
136 polyethylene vials will be used for ⁵⁵Fe, ⁶³Ni, ³H and ²⁴¹Pu.

137 *c. The balance*

138 The balance is a Mettler Toledo XPE26C set on a marble table (see Fig. 1). It is
139 recalibrated each year by Mettler Toledo in compliance with the EURAMET calibration
140 guide N°18 [25]. The standard uncertainty for metallic mass in the range from 1 mg
141 to 100 mg is equal or better than 5 µg.



142

143

Fig. 1 Photography of the BIPM balance

144

The differential weighing technique [26] is used to measure the mass of solution put in the vials. The active solution is transferred to a pycnometer which is measured before and after having released some droplets into the vials. One difficulty observed when weighing pycnometers is the electrostatic perturbation from the dielectric material. To cope this issue, a electrostic discharger is set at the entrance of the balance to remove chages from the pycnomter. Also, prior to each source preparation, blank weighings (without droplet release) of the pycometer is realized to control that the reproducibility is contained within +/- 10 μg .

151

152

The facility operates for producing LS sources of beta-emitting radionuclides but canot adress alpha-emitting radionuclides. Indeed, the handling of the latter is only allowed within a fume hood whatever the source activity is. The setup of a new balance in a fume hood is in project to enable alpha-emitting radionuclides to be adressed by the ESIR.

155

156

d. The sources

157

10 LS sources are realized for each ESIR submission. The amount of standard solution injected into each vials is calculated so that the double coincidence count rate r_D is

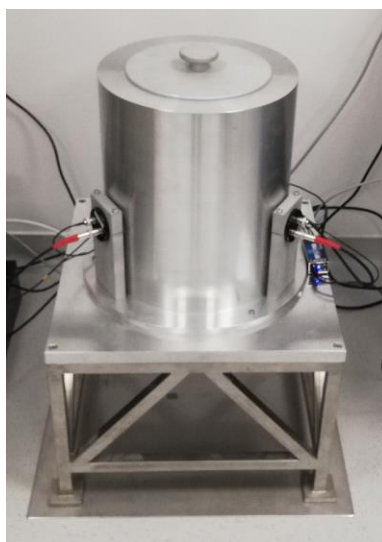
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159 between $5 \times 10^3 \text{ s}^{-1}$ and 10^4 s^{-1} . This constrain will guaranty to achieve a good counting
160 statistic and to operate the system always in the same range of counting dead time.

161 It is important to note that it is asumed that the LS sources are not stable. So, it is stated
162 in the procedure that the TDCR measurements must be done within the week just after
163 the source preparation.

164 **3. The TDCR system**

165 The TDCR system of the ESIR is presented in Fig. 2.



166

167 **Fig. 2** Photography of the ESIR TDCR system

168 As any TDCR system, it is composed of:

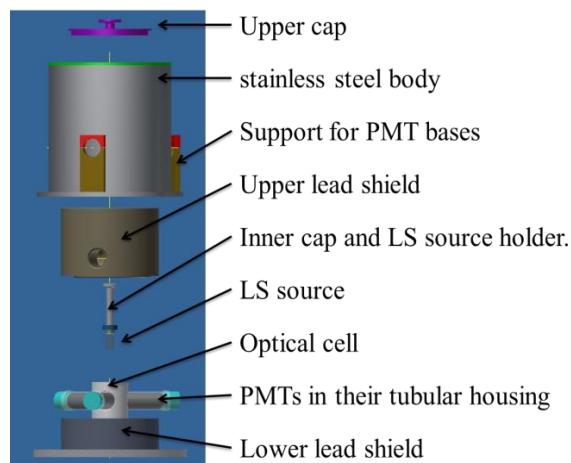
- 169 • a mechanical structure,
- 170 • an optical cell,
- 171 • a photo-conversion device using 3 photomultiplier tubes (PMTs),
- 172 • a specify signal processing electronics.

173 Technological options for these elements must be specified to address long-term
174 reproducibility and maintainability. To this end, an effort is made to maximize the
175 detection efficiency (close to 1) to minimize the part of the method in the long-term
176 reproducibility strategy. This property is particularly of importance for isotopes emitting
177 low energy particles such as ^3H , ^{63}Ni , ^{55}Fe , ^{241}Pu because the model from which I_1 , I_2 and
178 I_3 were conceived comes close to its limits to describe the underlying physics due notably
179 to low photoconversion statistics, significance of the ionization quenching, delayed
180 fluorescence, etc. Also, materials and electronics must be chosen in a way to guaranty the
181 sustainability of the reference produced by the system even when parts of the equipment
182 are damaged or come to obsolescence.

183 *a. The mechanical structure*

184 Unlike many TDCR systems, it was chosen not to implement a fourth gamma-rays
185 channel or an embedded source of ^{241}Am for Compton coincidence feature. These options
186 gives the possibility of performing $4\pi(\text{LS})\beta\text{-}\gamma$ counting [27] or Compton source
187 efficiency tracing [28]. Another TDCR system will be dedicated for these options
188 because it is preferable for the ESIR to simplify the design as much as possible to meet
189 our specific need exclusively. Thus, the absence of a fourth gamma-rays channel permit
190 to setup a 10 cm-thick gamma-rays shielding with a greater spherical coverage and
191 therefore a lower background noise. Also, the implementation of an ^{241}Am sealed source
192 for Compton coincidence requires coring a collimator hole through the optical cell. The
193 absence of this feature optimizes the symmetry (120° angular spacing between PMTs) of
194 the system. This optimizes the relevance of the I_1 indicator (not implementing asymmetry
195 correction) and minimizes the stabilization effort carried out by the I_2 comparison
196 indicator to maintain reproducible output values under possible asymmetry.

197 The assembly is placed on a reinforced support and incorporated into a cylindrical steel
198 cover ensuring light tightness. All this mechanical structure was designed and machined
199 at the BIPM workshop. Fig. 3 shows the exploded view of the BIPM TDCR system.



200

201

Fig. 3 Exploded view of the BIPM TDCR system

202

b. The optical cell

203

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205

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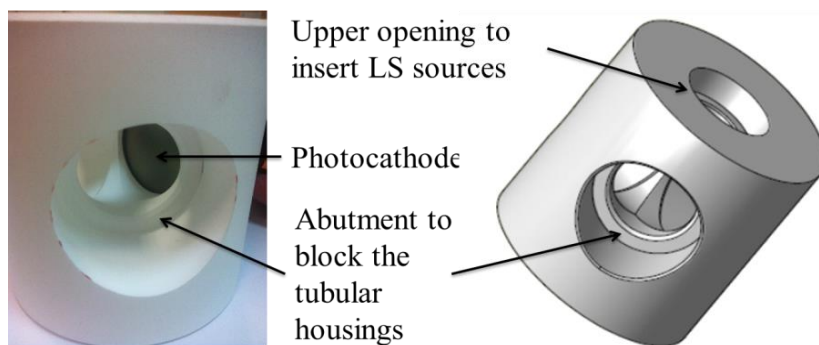
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210

211

It was decided to machine the optical cell in low density PTFE (Gigahertz-Optik ODM98-B [29]). This material, used by the PTB in one of their TDCR systems, has a very good reflectivity (98%) for visible photons. Moreover, this material is also known for its great immunity to chemical attacks and to radiation damages [30]. Therefore, it gives the TDCR system a greater long-term stability of the optical properties than traditional painting of titanium oxide, which requires to be periodically renewed. The optical cell (presented in Fig. 4) has been machined at the BIPM workshop. Like the photocathodes, the internal surfaces of the optical cell are frosted to maximize the stochastic scattering of photons in accordance with the TDCR model.

212



213 **Fig. 4** Photography and 3-D drawing of the optical cell

214 *c. The photoconversion device*

215 The insertion of PMTs through the shield to reach the optical cell is a delicate point
216 where light leaks can occur. To cope with this problem, a tube adjusted to the diameter of
217 the tubular housings of the PMTs has been conceived (see Fig. 5a). This tube comes to
218 lock in abutment of the optical cell so that the photocathode penetrates therein. The PMT
219 is fixed in the tubular housing thanks to an O-ring (see Fig. 5b).



220

221

(a)

(b)

222 **Fig. 5** Tubular housing of PMT tubes. (a) with a PMT inserted; (b) empty housing
223 showing the O-ring

224 The PMTs are Hamamatsu R331-05 selected to have an identical quantum efficiency of
225 +/- 1 % for an excitation at the wavelength of 420 nm. As can be seen in Fig. 5a, the
226 bialkali photocathodes are curved and frosted in order to respectively maximize the
227 optical efficiency and the stochastic diffusion of fluorescence photons. The curvature
228 allows the photocathode to approach the LS source borders within 2 mm. Furthermore,
229 the window material is low K content to minimize background noise. Before being
230 inserted into its tubular housing, the PMT is covered with a mu-metal foil to limit the
231 impact of possible electromagnetic disturbances on electron multiplication. The cathode
232 is polarized with a positive voltage. And, in accordance with our strategy to avoid

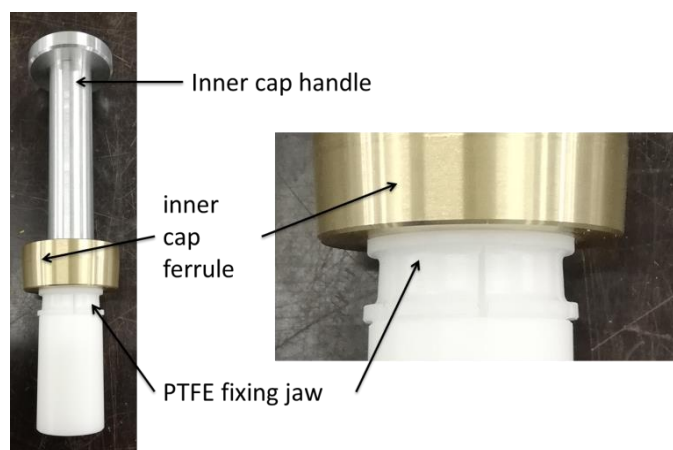
233 dispensable features which could be sources of disturbance for the reference to be
234 produced, no potentiometric defocusing has been implemented.

235 *d. The LS source positioning system*

236 The previous BIPM TDCR system used an extruded base at the bottom of the cell to
237 position the source. This approach turns out to have the following two disadvantages:

- 238 • having a LS source off-center from the axis of the photocathodes that is not
239 optimal to maximize the detection efficiency,
- 240 • achieving in practice mediocre reproducibility of the positioning of the LS
241 source.

242 For the new system, a special plug has been designed to hold the LS source in a
243 reproducible position maximizing the detection efficiency. The source volume is
244 maintained in the axis of the photocathodes using a jaw build in PTFE which grab the
245 caps of the LS vials (see Fig. 6). This approach is implemented within POLATOM's
246 TDCR systems and gives hope for an almost perfect reproducibility of the positioning of
247 the source.



248

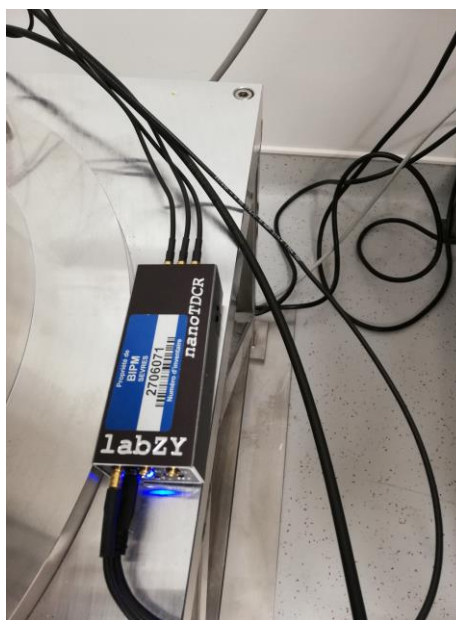
249 **Fig. 6** Photography of the inner cap of the BIPM TDCR system handling a LS source

250

e. The electronics

251 The PMTs are polarized at a voltage equal to +2500 V using NHQ103M HV units. This
252 voltage value is ideal to enable the system to detect properly single photoelectron events.
253 The anode signals from the three PMTs are processed by the nanoTDCR module [31,32]
254 (see **Fig. 5**). It embeds in its FPGA the same extended dead time processing as in the
255 reference MAC3 unit [33] which is unfortunately no more maintained. The nanoTDCR
256 provides the following quantities:

- 257
- the triple coincidence count rate r_T ,
 - 258 • the logical sum of the double coincidence count rate r_D ,
 - 259 • the double coincidence count rates r_{AB} , r_{BC} , r_{AC} of each pair of the three channels
260 A, B and C.
 - 261 • the count rates r_A , r_B , r_C from channels A, B and C.



262

Fig. 7 Photography of the nanoTDCR module

263

264 The key comparison indicators of the ESIR are built from these quantities such as:

- 265
- I_0 only uses r_D ,
 - 266 • I_1 uses r_D , the *TDCR* value, R_T/R_D , plus two preset parameters α_1 and α_2 ,

267 • I_2 uses r_D and the TDCR values, R_T/R_{AB} , R_T/R_{BC} , R_T/R_{AC} , plus preset parameters
268 for the TDCR model: the Birks constant k_B , the stopping power $dE/dx(E)$ and the
269 energy spectrum $S(E)$,

270 • I_3 uses r_T , r_{AB} , r_{BC} , r_{AC} .

271 The count rates r_A , r_B , r_C are used to correct the measurement from accidental
272 coincidences [34]. Details are available in [35].

273 To enhance the long-term maintainability of the ESIR, a redundant electronic based on
274 the CEAN digitizer DT5730S is planned to be setup. The list mode data acquisition will
275 permit to have more control on the algorithmic implemented for the TDCR signal
276 processing.

277 **4. Commissioning of the TDCR system**

278 The ESIR aims to produce a fixed reference. So it is of importance that parameters used
279 in the signal processing are defined once and for all. The extended dead time algorithm
280 requires the following parameters:

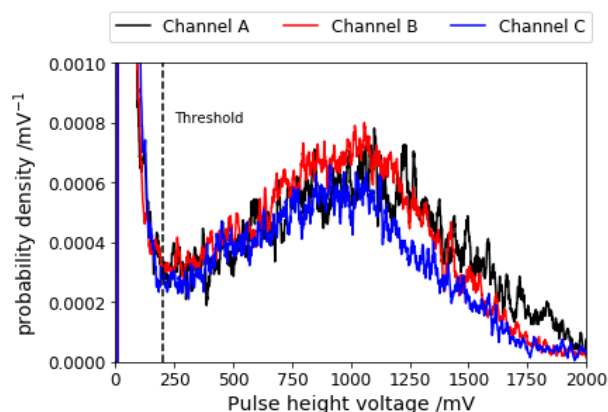
- 281 • The discrimination thresholds on each channel,
- 282 • The extended dead time period,
- 283 • The coincidence resolving time.

284 One interesting feature of the nanoTDCR is to process the signal with two different
285 extended dead times and two different resolving times. The setup of these signal
286 processing parameters is described below.

287 *a. Setup of the discrimination thresholds*

288 Fig. 8 presents the pulse high spectra recorded without LS sources in the TDCR system.
289 The left component comes from the thermal electronic noise and the spread peak centred

290 on about 1 V is induced by the dark noise related to the stochastic thermionic emission
291 from photocathodes. The three spectra as almost superimposed which means that PMTs
292 displayed equivalent thermal noise magnitude and voltage gain. The first set step is to set
293 the discrimination threshold in the valley of the spectra in order to maximize the signal to
294 noise ratio, i.e., minimizing the thermal noise impact while maximizing the detection
295 probability of single electron pulses. According to the shape observed in Fig. 6,
296 thresholds are therefore set at -200 mV for all the 3 channels A, B and C.



297

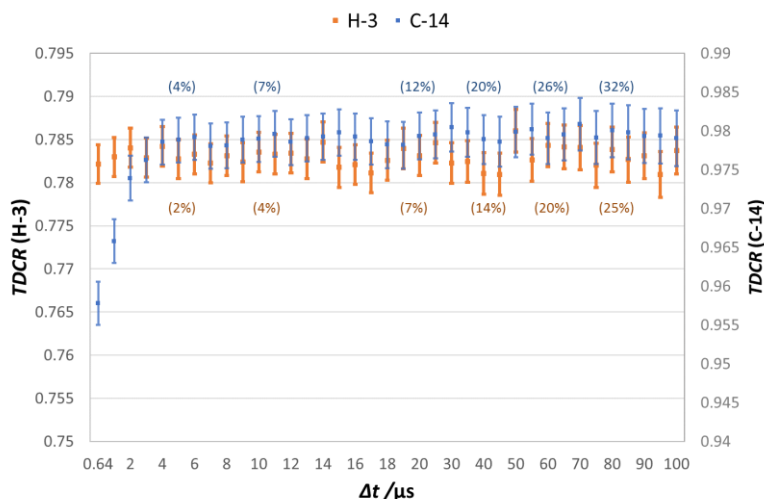
298 **Fig. 8** Pulse high spectra of single photoelectron pulses from the three channels of the
299 TDCR system

300 *b. Setup of the extended dead time*

301 After-pulses induce a metrological challenge in every measurement systems
302 implementing PMTs [36]. To limit after-pulses impact, the extended dead time method
303 paralyzes the signal processing during an imposed period Δt after each triggering
304 occurring on one of the three channels.

305 Considering a resolving time set to 50 ns, the TDCR value was estimated as a function of
306 the extended dead time period Δt for 2 toluene-based standard sources of ^3H (3.2 kBq)
307 and ^{14}C (2.0 kBq) and presented in Fig. 9. The influence of after-pulses on the
308 measurement seems to be a function of the energy of the beta particles and therefore of
309 the amplitude of the primary pulse. No significant impact is observed in the case of
310 tritium and a Δt of 5 μs seems sufficient for ^{14}C sources to filter this type of disturbance.

311 To maximize the statistical precision of the counting and the accuracy of the live-time
 312 estimation, the dead time of the measurement must be minimized. Thus, by taking a
 313 reasonable safety margin, a first Δt equal to 10 μs will ensure after-pulse filtering while
 314 minimizing the dead time. The second Δt is set equal to a larger value, 50 μs , and will
 315 permit to compare measurement results obtained with these two paralyzing times to
 316 assess the impact of after pulses.



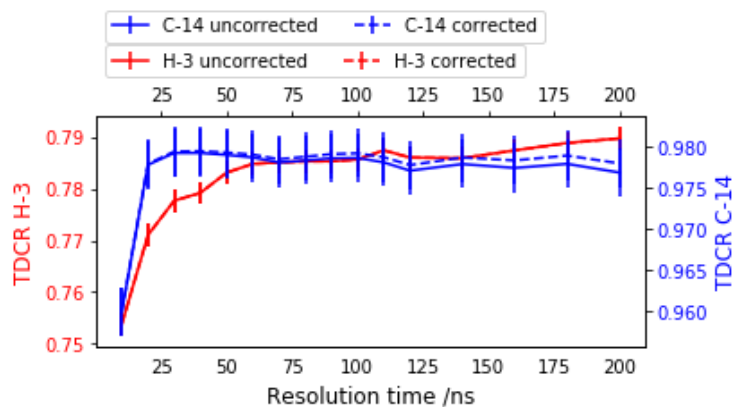
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318 **Fig. 9** *TDCR* value as a function of the extended dead time for the two sources of ^3H and
 319 ^{14}C . The measurement dead time is displayed in brackets.

320 *c. Setup of the coincidence resolving time*

321 The resolution time $\Delta\tau$ is the time window during which any triggering between the two
 322 or three channels will be considered as a double or triple coincident event. Too short a $\Delta\tau$
 323 leads to a risk of non-detection. Conversely, too long a resolution time increases the risk
 324 of accidental coincidences. Keeping the same experimental configuration as before, the
 325 evolution of the *TDCR* value as a function of the resolving time is studied. It is observed
 326 in Fig. 10 a disparity of behaviour between ^3H and ^{14}C with the resolving time. While the
 327 *TDCR* value of ^{14}C reaches a constant (or slightly decreasing) value from a resolution
 328 time greater than 30 ns, the *TDCR* value associated with ^3H increases monotonically.
 329 This can be explained by the fact that the stopping power of very low energy electrons is

330 greater than for higher energy electrons. The decays of tritium nuclei will therefore have
 331 a greater tendency to produce triplet states of excitement leading to delayed fluorescence
 332 or phosphorescence [37]. Anyway, this phenomenon depending on the scintillating
 333 liquid, concerns almost exclusively low energy emitting isotopes such as ^{55}Fe or tritium.
 334 Although it can potentially cause a bias in the context of primary measurements, these
 335 biases are constant for a given scintillator and can therefore be assumed in ESIR
 336 operation. So, a value of $\Delta\tau$ equal to 50 ns will be considered whatever the radionuclide
 337 is. This choice guarantees a margin compared to the minimum of 30 ns observed here and
 338 greatly minimizes the risk of accidental coincidence. A measurement with a resolution
 339 time of 100 ns will also be simultaneously recorded in order to control the stability of the
 340 scintillating liquid with respect to the phenomenon described above.



341

342 **Fig. 10** TDCR value as a function of the coincidence resolving time for the two sources
 343 of ^3H and ^{14}C .

344 **5. The quality assurance**

345 The quality assurance will play a critical role to control the long-term reproducibility of
 346 the reference values produced by the ESIR. A set of procedures were established and
 347 specific attention was carried on the automation of the measurement analysis by a Python
 348 program and by establishing a Git-controlled database in the same framework as it was
 349 recently established for the SIR [7].

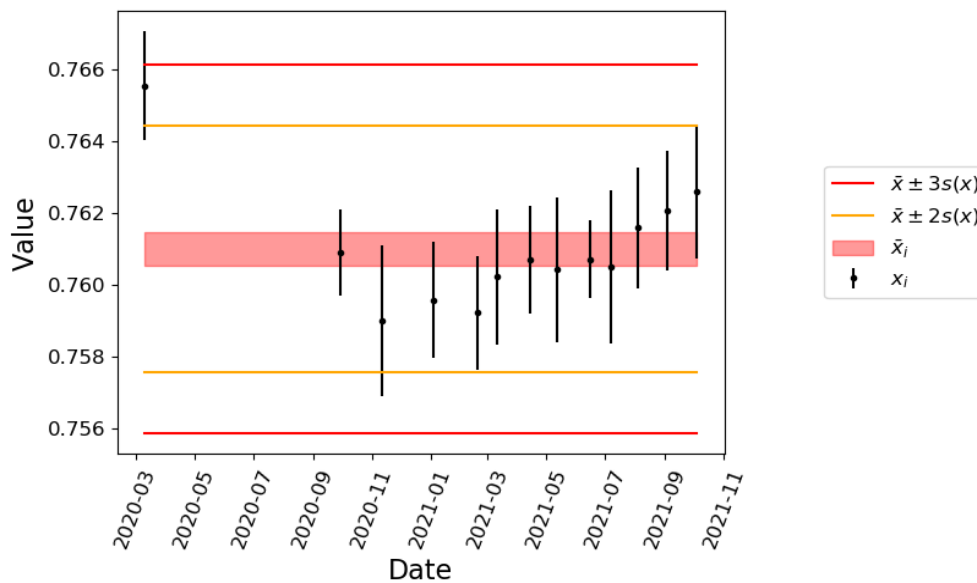
350 Besides, a monthly periodic control of the system was setup since its commissioning in
351 August 2020. The check sources are 6 toluene-based LS sources (^{14}C sources and ^3H
352 sources) known for their good stability over a long-term period. The reproducibility of
353 the ESIR is assessed by performing statistical analysis on these quality measurements. As
354 recommended by the ISO17025:2017, control charts (x-chart, R-chart, CUSUM chart and
355 EWMA chart) are applied to detect outliers or drifts signaling that the system is not under
356 statistical control. Also, the consistency of the data set with the uncertainty evaluation of
357 each point is tested using a Chi-squared test.

358 The results for one of the tritium sources for which a first measurement point was
359 obtained in March 2020 are displayed in Figs. 11-14. The error bars are standard
360 uncertainty associated with the measurement points, x_i . The lower and upper control
361 limits are the mean of all the measurement points ± 2 times (yellow lines) or 3 times
362 (red lines) standard deviation (type A evaluation). The central red line represents the
363 mean with a thickness equal to ± 1 standard deviation of the mean value.

364 The raw comparison indicator I_0 and the $TDCR$ values (Figs. 11 and 12) show clearly that
365 the system is not by itself stable. Even if the raw measurement data I_0 passes tests from
366 control charts, it fails the Chi-squared test with a p-value equal to 0.002 (corresponding to
367 a Birge ratio equal to 1.76). So, the hypothesis of consistency of the measurements
368 according to their individual uncertainty evaluations can be rejected with a 95 %
369 confidence level. This means that it exists slight (about 0.2 %) but significant drifts in the
370 detection efficiency that are not taken into account in the uncertainty budget.

371 For the same data set, key comparison indicators I_1 and the I_2 are displayed
372 in Fig. 13 and Fig. 14. They both pass tests from control charts and the Chi-squared test
373 with p-values close to 1. This result stresses the robustness provided by the approach
374 developed for the ESIR. It demonstrates that the ESIR is capable to provide reproducible
375 values I_1 and I_2 during a period of 20 months.

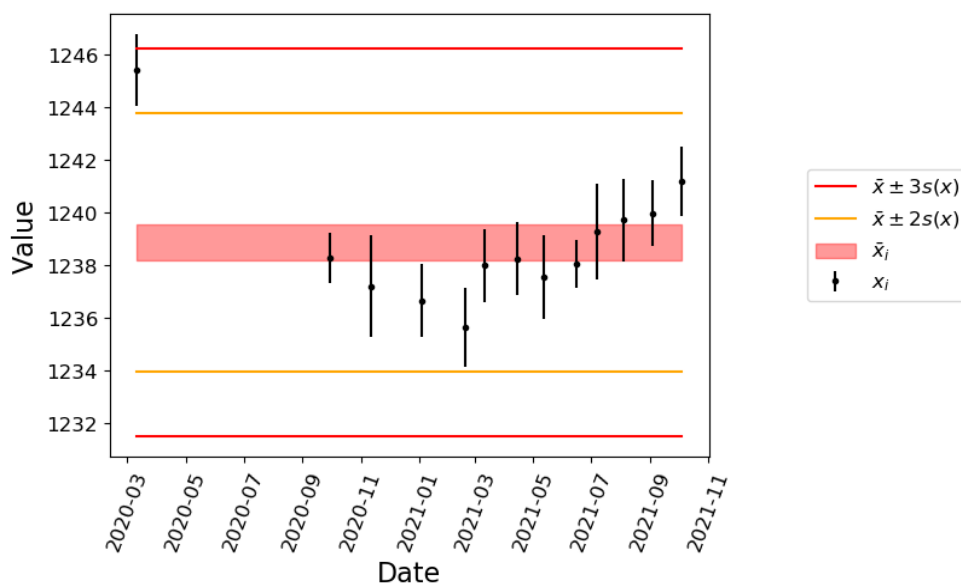
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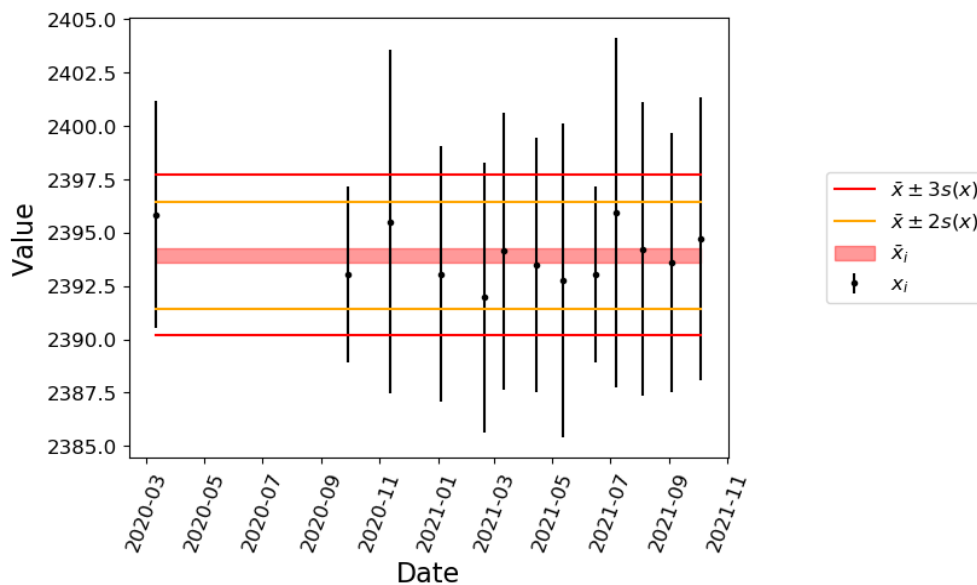
Fig. 11 TDCR values obtained for the ^3H source HDCO since March 2020



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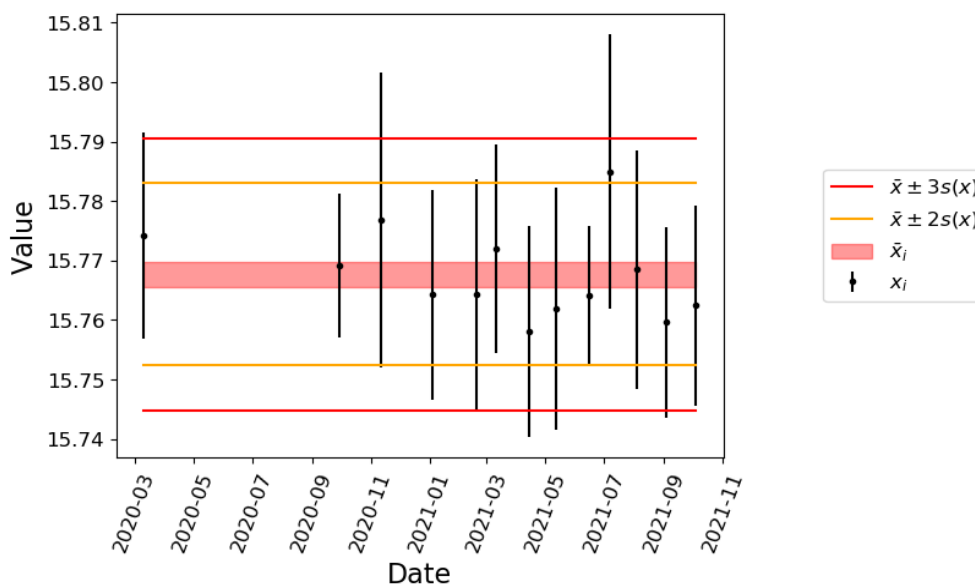
Fig. 12 I_0 values obtained for the ^3H source HDCO since March 2020



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Fig. 13 I_1 values obtained for the ^3H source HDCO since March 2020



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Fig. 14 I_2 values obtained for the ^3H source HDCO since March 2020

385

Results for the other check standards are available as supplementary information.

386

387 **Conclusions**

388 A liquid scintillation source facility and a new TDCR system have been set up at the
389 BIPM to extend the SIR to pure alpha-, pure beta-emitting and some electron capture-
390 decaying radionuclides. The objective is to improve the traceability chain for
391 radionuclides of interest for new cancer treatments and some environmental issues.

392 The service was commissioned in August 2020. Since this data, the ESIR operates,
393 through a set of procedures and fixed defined parameters, to produce highly reproducible
394 measurement quantities. These quantities called key comparison indicators will permit to
395 establish reference reference values (KCRV) for these radionuclides and degrees of
396 equivalence for the participating NMI/DIs. The ESIR aims to be an on-demand service
397 for at least isotope emitting medium and high energy particles. A CCRI(II) pilot study on
398 Co-60 key comparison is underway to validate the ESIR against the SIR in the context of
399 an international key comparison exercise. This pilot study implies 12 NMI/DIs and is
400 planned to be completed in 2022.

401 Meantime, the periodic quality control using toluene-based liquid scintillation sources of
402 ^3H and ^{14}C have been launched. They have permitted to demonstrate the capability of the
403 ESIR to produce key comparison indicators with a good reproducibility for a 20-month
404 period. This preliminary result is encouraging on the stability attainable by the unique
405 method of the ESIR based on the TDCR technique.

406 Many challenges and perspectives are under study with ESIR such as,

- 407 • the adaptation of the source preparation facility to address alpha-emitting
408 radionuclides (e.g., ^{211}At , ^{212}Bi , ^{212}Pb , ^{213}Bi , ^{227}Th),
- 409 • the set-up of a redundant digital signal processing,
- 410 • the development of an LS cocktail prepared at the BIPM to allow on demand
411 comparison for radionuclides emitting low energy particles (e.g., ^{55}Fe , ^{63}Ni , ^3H
412 and ^{241}Pu),

- 413 • the development of a transportable ESIR (analog to the SIRTI) using the
414 transportable TDCR system developed by the LNE/LNHB [38].

415 **References**

- 416 [1] CIPM 2003 Mutual recognition of national measurement standards and of
417 calibration and measurement certificates issued by national metrology institutes
- 418 [2] Karam L R 2007 Application of the CIPM MRA to radionuclide metrology
419 *Metrologia* **44** 4–9
- 420 [3] Rytz A 1983 The international reference system for activity measurements of γ -
421 ray emitting nuclides *Int. J. Appl. Radiat. Isot.* **34** 1047–56
- 422 [4] Ratel G 2007 The Système International de Référence and its application in key
423 comparisons *Metrologia* **44** 7–16
- 424 [5] Karam L, Judge S and Louw W 2019 The international measurement system for
425 radionuclide metrology: A strategy for the future *Appl. Radiat. Isot.* **154** 108838
- 426 [6] Michotte C, Nonis M, Bobin C, Altzizoglou T and Sibbens G 2013 *The SIRTI: a*
427 *new tool developed at the BIPM for comparing activity measurements of short-*
428 *lived radionuclides world-wide*
- 429 [7] Coulon R, Courte S, Judge S, Michotte C and Nonis M 2021 Digitalization of the
430 reporting of key comparisons for radionuclide metrology *Meas. Sci. Technol.*
431 **Submitted** 113015
- 432 [8] Sgouros G, Bodei L, McDevitt M R and Nedrow J R 2020 Radiopharmaceutical
433 therapy in cancer: clinical advances and challenges *Nat. Rev. Drug Discov.* **19**
434 589–608
- 435 [9] Parker C, Lewington V, Shore N, Kratochwil C, Levy M, Lindén O, Noordzij W,
436 Park J and Saad F 2018 Targeted Alpha Therapy, an Emerging Class of Cancer

- 437 Agents *JAMA Oncol.* **4** 1765
- 438 [10] Duchemin C, Ramos J P, Stora T, Ahmed E, Aubert E, Audouin N, Barbero E,
439 Barozier V, Bernardes A-P, Bertreix P, Boscher A, Bruchertseifer F, Catherall R,
440 Chevally E, Christodoulou P, Chrysalidis K, Cocolios T E, Comte J, Crepieux B,
441 Deschamps M, Dockx K, Dorsival A, Fedosseev V N, Fernier P, Formento-
442 Cavaier R, El Idrissi S, Ivanov P, Gadelshin V M, Gilardoni S, Grenard J-L,
443 Haddad F, Heinke R, Juif B, Khalid U, Khan M, Köster U, Lambert L, Lilli G,
444 Lunghi G, Marsh B A, Palenzuela Y M, Martins R, Marzari S, Menaa N, Michel
445 N, Munos M, Pozzi F, Riccardi F, Riegert J, Riggaz N, Rinchet J-Y, Rothe S,
446 Russell B, Saury C, Schneider T, Stegemann S, Talip Z, Theis C, Thiboud J, van
447 der Meulen N P, van Stenis M, Vincke H, Vollaire J, Vuong N-T, Webster B,
448 Wendt K and Wilkins S G 2021 CERN-MEDICIS: A Review Since
449 Commissioning in 2017 *Front. Med.* **8**
- 450 [11] Iravani A, Violet J, Azad A and Hofman M S 2020 Lutetium-177 prostate-specific
451 membrane antigen (PSMA) theranostics: practical nuances and intricacies *Prostate*
452 *Cancer Prostatic Dis.* **23** 38–52
- 453 [12] Michotte C, Ratel G, Courte S, Johansson L, Keightley J, Arinc A, Bakhshandear
454 E, Pommé S, Altitzoglou T, Paepen J and Van Ammel R 2014 BIPM comparison
455 BIPM.RI(II)-K1.Lu-177 of activity measurements of the radionuclide 177 Lu for
456 the NPL (UK) and the IRMM (EU), with linked results for the comparison
457 CCRI(II)-K2.Lu-177 *Metrologia* **51** 06002–06002
- 458 [13] Michotte C, Courte S, Nonis M, Coulon R, Judge S, Kossert K and Nähle O 2021
459 Final report of the new BIPM comparison BIPM.RI(II)-K1.Ac-225 of activity
460 measurements of the radionuclide 225 Ac to include the 2019 result of the PTB
461 (Germany) *Metrologia* **58** 06018
- 462 [14] Michotte C, Courte S, Nonis M, Coulon R, Judge S, Ratel G, Fréchou C, Cassette
463 P, Keightley J, Kossert K and Nähle O 2021 Final report of the new BIPM
464 comparison BIPM.RI(II)-K1.Ra-223 of activity measurements of the radionuclide

- 465 223 Ra including the 2014 result of the NPL (United Kingdom), the 2014 result of
466 the PTB (Germany) and the 2018 result of the LNE-LNHB (France) *Metrologia* **58**
467 06007
- 468 [15] Wallenius M, Lützenkirchen K, Mayer K and Varga Z 2018 Actinides: Nuclear
469 Forensics *Encyclopedia of Inorganic and Bioinorganic Chemistry* (Chichester,
470 UK: John Wiley & Sons, Ltd) pp 1–17
- 471 [16] Jiang J, Davies A, Thorne K and Gilligan C 2017 Rapid analysis of ⁸⁹Sr and ⁹⁰Sr
472 in nuclear forensics samples *J. Radioanal. Nucl. Chem.* **311** 1417–25
- 473 [17] Jiang J, Davies A V. and Britton R E 2017 Measurement of ¹⁶⁰Tb and ¹⁶¹Tb in
474 nuclear forensics samples *J. Radioanal. Nucl. Chem.* **314** 727–36
- 475 [18] Pochwalski K, Broda R and Radoszewski T 1988 Standardization of pure beta
476 emitters by liquid-scintillation counting *Int. J. Radiat. Appl. Instrumentation. Part*
477 **39** 165–72
- 478 [19] Broda R 2003 A review of the triple-to-double coincidence ratio (TDCR) method
479 for standardizing radionuclides *Appl. Radiat. Isot.* **58** 585–94
- 480 [20] Nedjadi Y, Duc P F, Bochud F and Bailat C 2016 On the stability of ³H and ⁶³Ni
481 Ultima Gold liquid scintillation sources *Appl. Radiat. Isot.* **118** 25–31
- 482 [21] Coulon R, Broda R, Cassette P, Courte S, Jerome S, Judge S, Kossert K, Liu H,
483 Michotte C and Nonis M 2020 The international reference system for pure β-
484 particle emitting radionuclides: an investigation of the reproducibility of the results
485 *Metrologia* **57** 035009
- 486 [22] Coulon R, Judge S, Liu H and Michotte C 2021 The international reference
487 system for pure beta-particle emitting radionuclides: an evaluation of the
488 measurement uncertainties *Metrologia* **58** 025007
- 489 [23] Kossert K, Sabot B, Cassette P, Coulon R and Liu H 2020 On the photomultiplier-
490 tube asymmetry in TDCR systems *Appl. Radiat. Isot.* **163** 109223

- 491 [24] Broda R, Bonková I, Capogni M, Carconi P, Cassette P, Coulon R, Courte S,
492 Felice P De, Dziel T, Fazio A, Frechou C, Galea R, García-Toraño E, Kołakowska
493 E, Kossert K, Krivošík M, Lech E, Lee K B, Liang J, Listkowska A, Liu H,
494 Navarro N, Nähle O J, Nowicka M, Rooy M van, Sabot B, Saganowski P, Sato Y,
495 Tymiąski Z, Yunoki A, Zhang M and Ziemek T 2021 The CCRI(II)-K2.Fe-
496 55.2019 key comparison of activity concentration measurements of a 55 Fe
497 solution *Metrologia* **58** 06010
- 498 [25] EURAMET 2015 *Guidelines on the calibration of non-automatic weighing*
499 *instruments*
- 500 [26] Lourenço V and Bobin C 2015 Weighing uncertainties in quantitative source
501 preparation for radionuclide metrology *Metrologia* **52** S18–29
- 502 [27] Bobin C and Bouchard J 2006 A $4\pi(\text{LS})\beta\text{-}\gamma$ coincidence system using a TDCR
503 apparatus in the β -channel *Appl. Radiat. Isot.* **64** 124–30
- 504 [28] Cassette P and Do P 2008 The Compton source efficiency tracing method in
505 liquid scintillation counting: A new standardization method using a TDCR counter
506 with a Compton spectrometer *Appl. Radiat. Isot.* **66** 1026–32
- 507 [29] Anon ODM98-B
- 508 [30] Xiaoguang Zhong, Li Yu, Jiazhen Sun and Yuefang Zhang 1993 Radiation
509 stability of PTFE irradiated under various conditions *Polym. Degrad. Stab.* **39**
510 187–91
- 511 [31] Mitev K, Cassette P, Jordanov V, Liu H R and Dutsov C 2017 Design and
512 performance of a miniature TDCR counting system *J. Radioanal. Nucl. Chem.* **314**
513 583–9
- 514 [32] Jordanov V, Cassette P, Dutsov C and Mitev K 2020 Development and
515 applications of a miniature TDCR acquisition system for in-situ radionuclide
516 metrology *Nucl. Instruments Methods Phys. Res. Sect. A Accel. Spectrometers,*
517 *Detect. Assoc. Equip.* **954** 161202

- 518 [33] Bouchard J and Cassette P 2000 MAC3: An electronic module for the processing
519 of pulses delivered by a three photomultiplier liquid scintillation counting system
520 *Appl. Radiat. Isot.* **52** 669–72
- 521 [34] Dutsov C, Cassette P, Sabot B and Mitev K 2020 Evaluation of the accidental
522 coincidence counting rates in TDCR counting *Nucl. Instruments Methods Phys.*
523 *Res. Sect. A Accel. Spectrometers, Detect. Assoc. Equip.* **977** 164292
- 524 [35] Michotte C, Courte S, Nonis M, Coulon R, Judge S, Nähle O and Takács M P
525 2021 Update of the BIPM comparison BIPM.RI(II)-K1.Sr-85 of activity
526 measurements of the radionuclide 85 Sr to include the 2018 result of the PTB
527 (Germany) *Metrologia* **58** 06006
- 528 [36] Akgun U, Ayan A S, Aydın G, Duru F, Olson J and Onel Y 2008 Afterpulse
529 timing and rate investigation of three different Hamamatsu Photomultiplier Tubes
530 *J. Instrum.* **3** T01001–T01001
- 531 [37] Bobin C, Thiam C, Chauvenet B and Bouchard J 2012 On the stochastic
532 dependence between photomultipliers in the TDCR method *Appl. Radiat. Isot.* **70**
533 770–80
- 534 [38] Anon Discussion with Benoit Sabot, 2021
535
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537

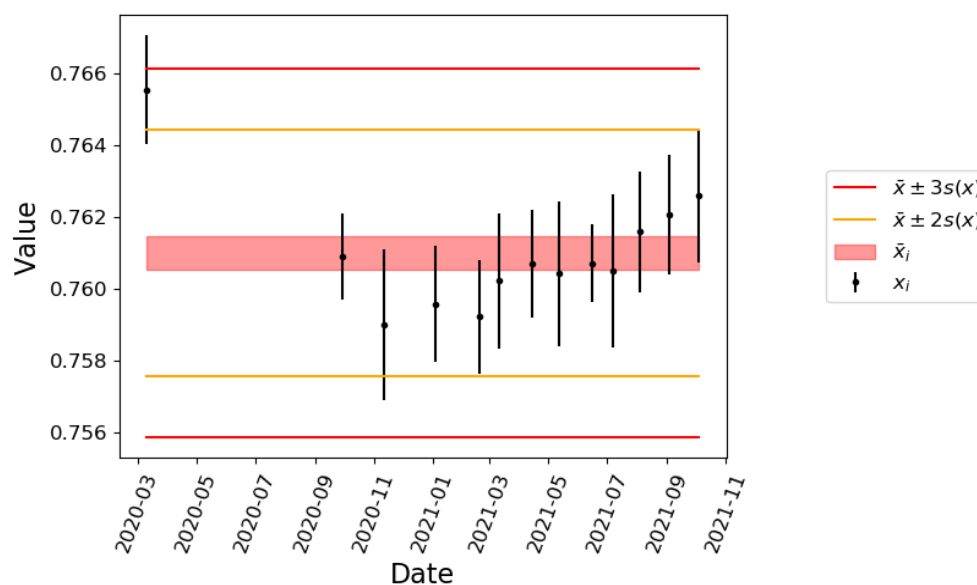
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Supplementary information

539 The full analysis of the period quality control is reported in this appendix. The error bars
540 are standard uncertainty associate with the measurement points, x_i . The lower end upper
541 control limits are the mean of all the measurement points ± 2 times (yellow lines) or 3
542 times (red lines) standard deviation (type A evaluation). The central red line represents
543 the mean with a thickness equal to ± 1 standard deviation of the mean value.

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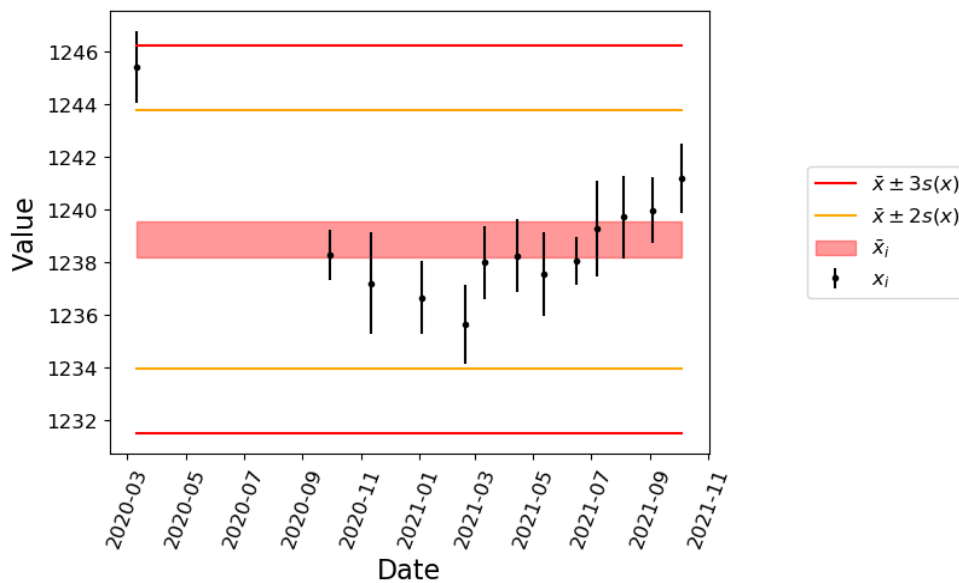
a. Toluene-based tritium standard HDCO



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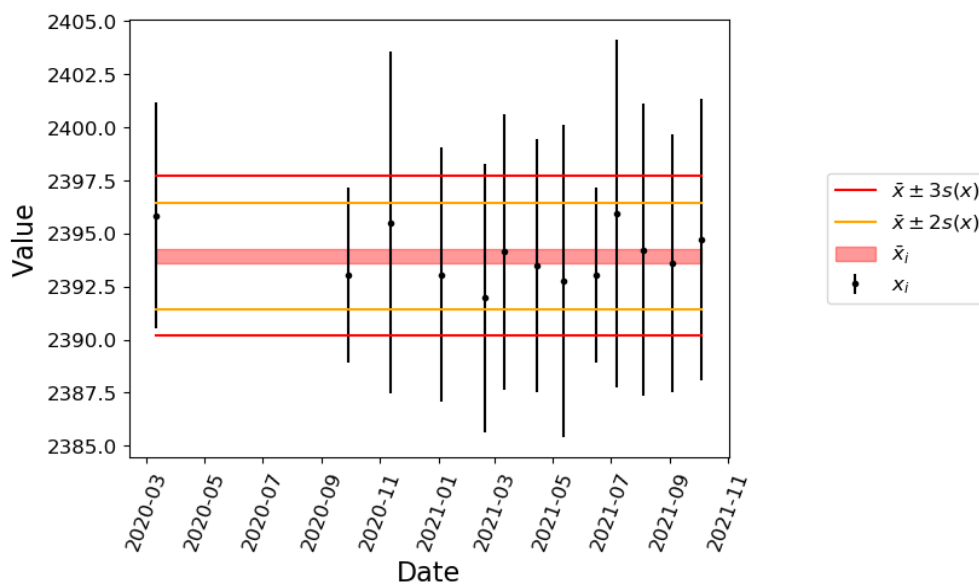
Fig. 1 TDCR values obtained for the ^3H source HDCO since March 2020



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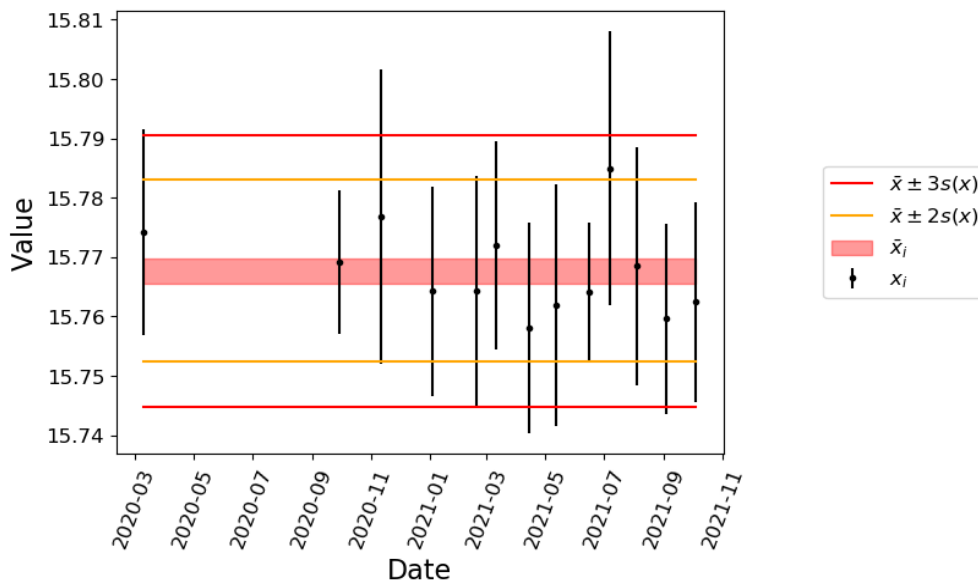
Fig. 3 I_0 values obtained for the ^3H source HDCO since March 2020



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Fig. 2 I_1 values obtained for the ^3H source HDCO since March 2020



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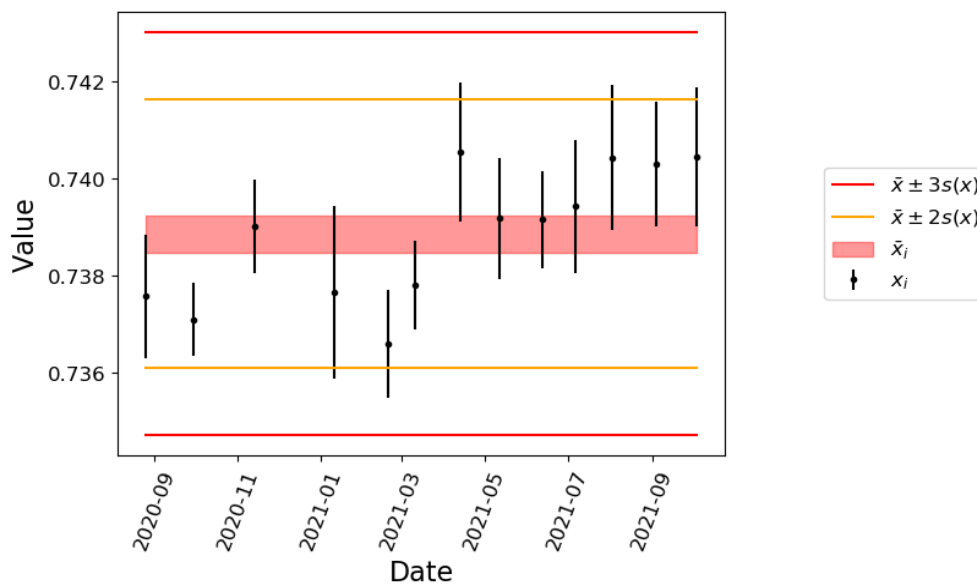
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Fig. 2 I_2 values obtained for the ^3H source HDCO since March 2020

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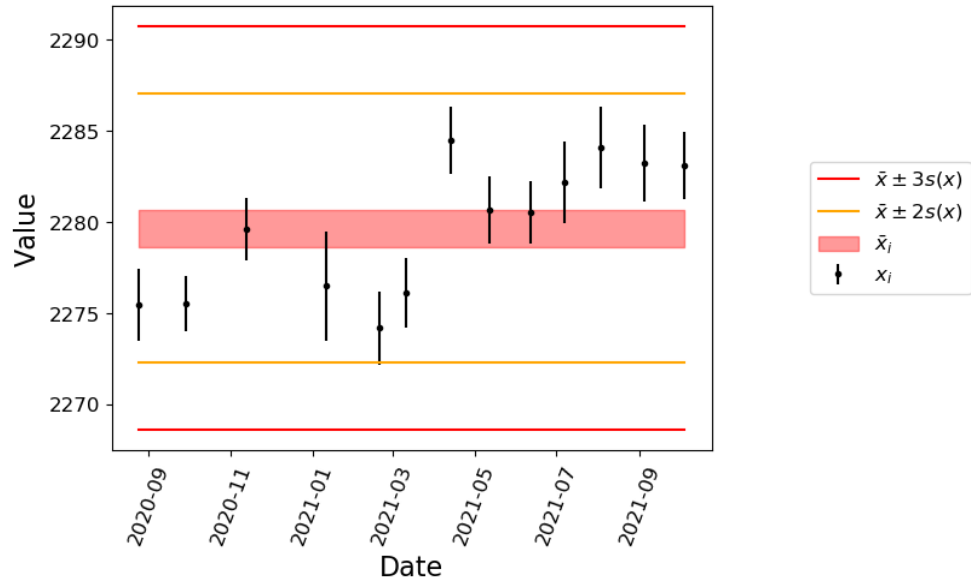
b. Toluene-based tritium standard HMI



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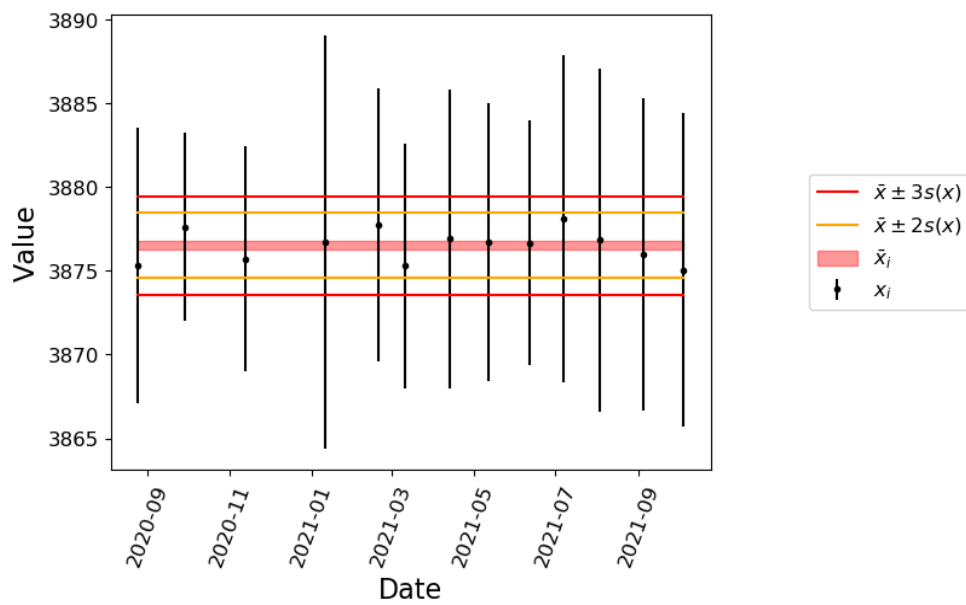
Fig. 1 TDCR values obtained for the ^3H source HMI since August 2020



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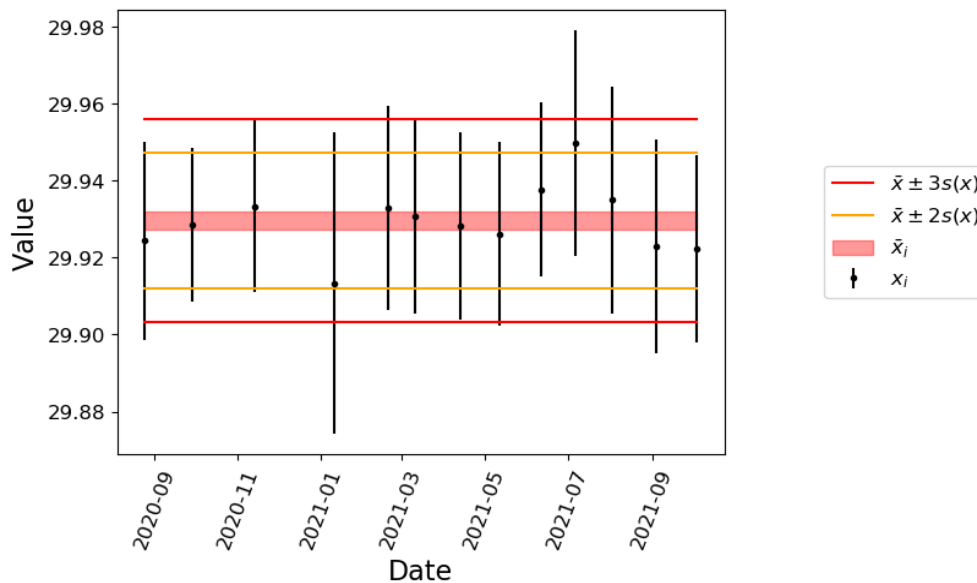
Fig. 2 I_0 values obtained for the ^3H source HMI since August 2020



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Fig. 2 I_1 values obtained for the ^3H source HMI since August 2020



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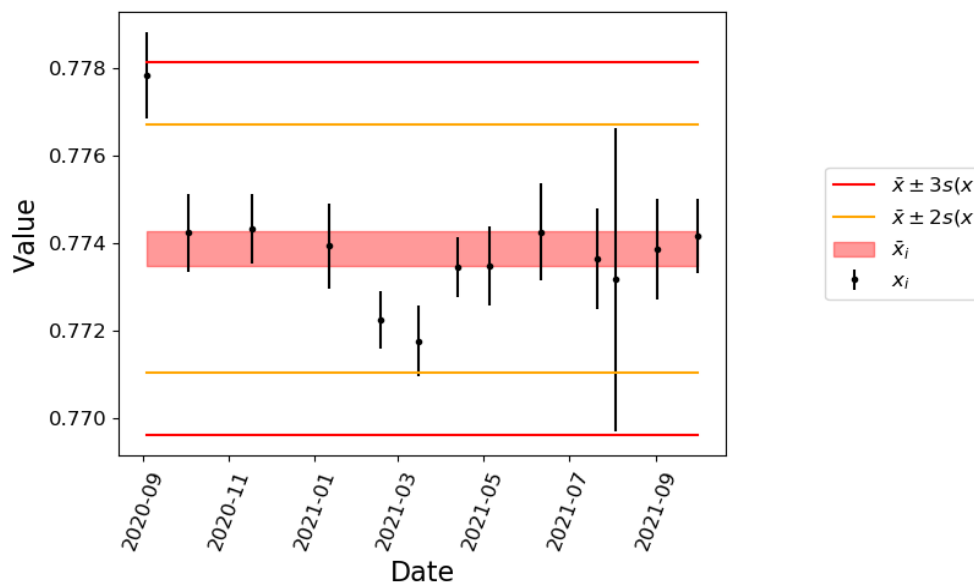
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Fig. 2 I_2 values obtained for the ^3H source HMI since August 2020

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c. Toluene-based tritium standard H3-117TOL

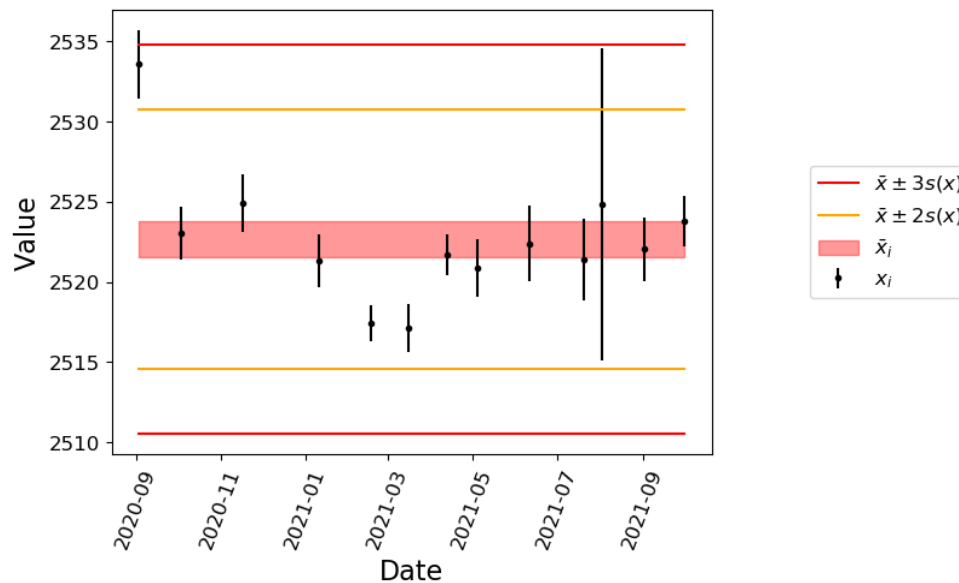


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Fig. 1 TDCR values obtained for the ^3H source H3-117TOL since August 2020

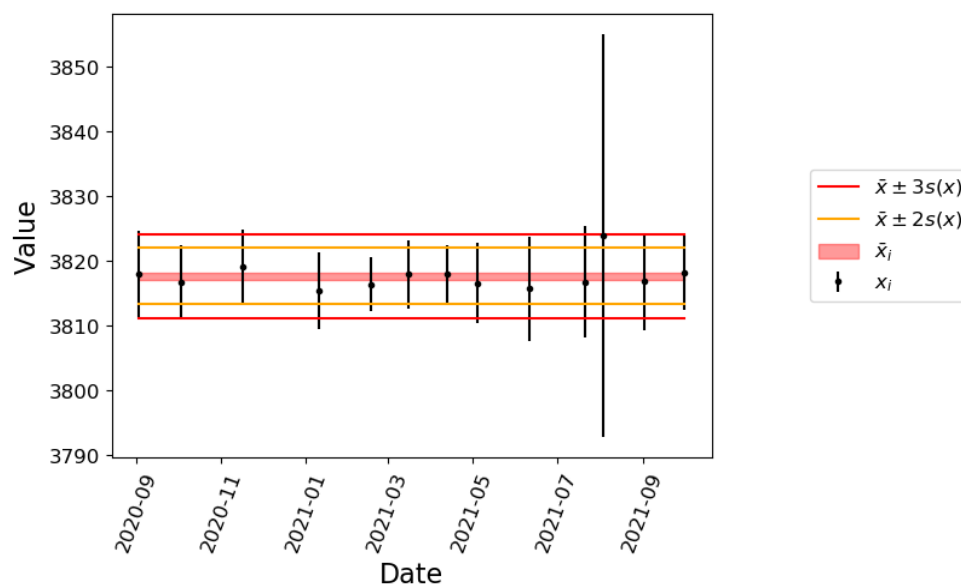
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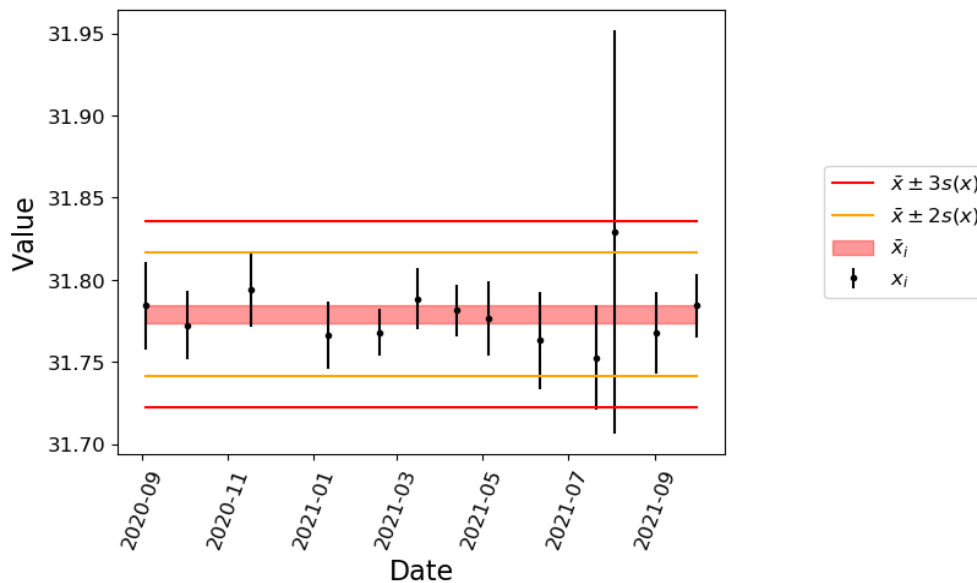
Fig. 2 I_0 values obtained for the ^3H source H3-117TOL since August 2020



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Fig. 2 I_1 values obtained for the ^3H source H3-117TOL since August 2020



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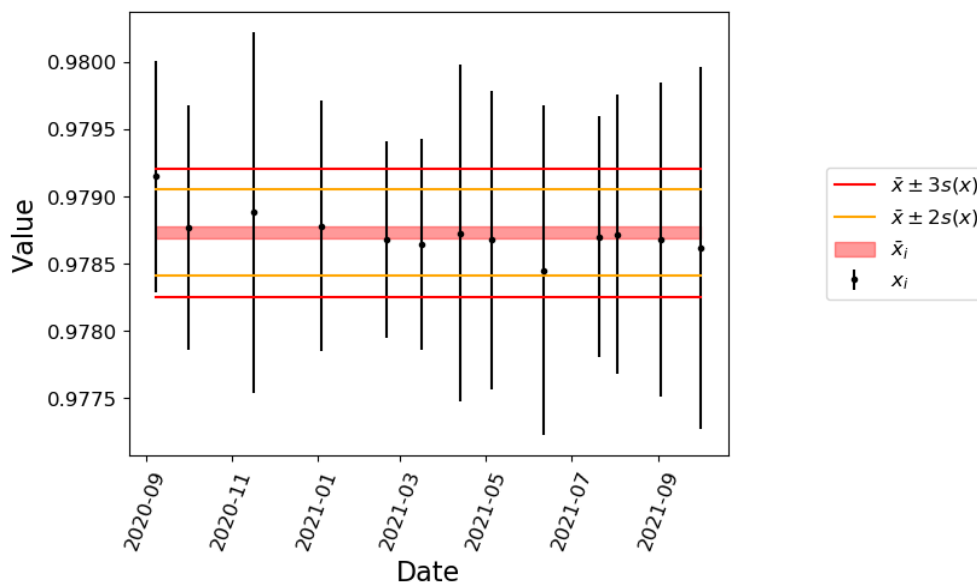
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Fig. 2 I_2 values obtained for the ^3H source H3-117TOL since August 2020

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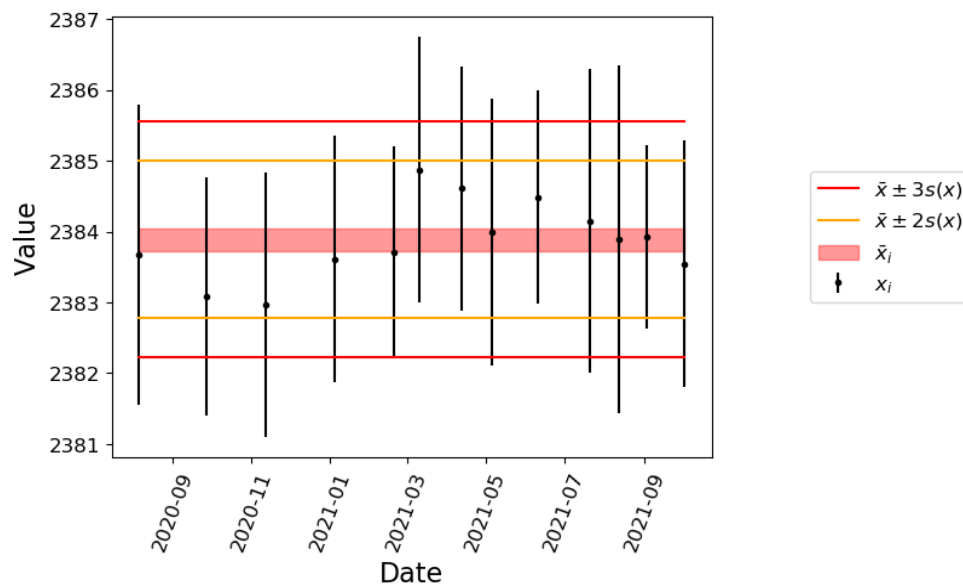
d. Toluene-based ^{14}C standard CCT



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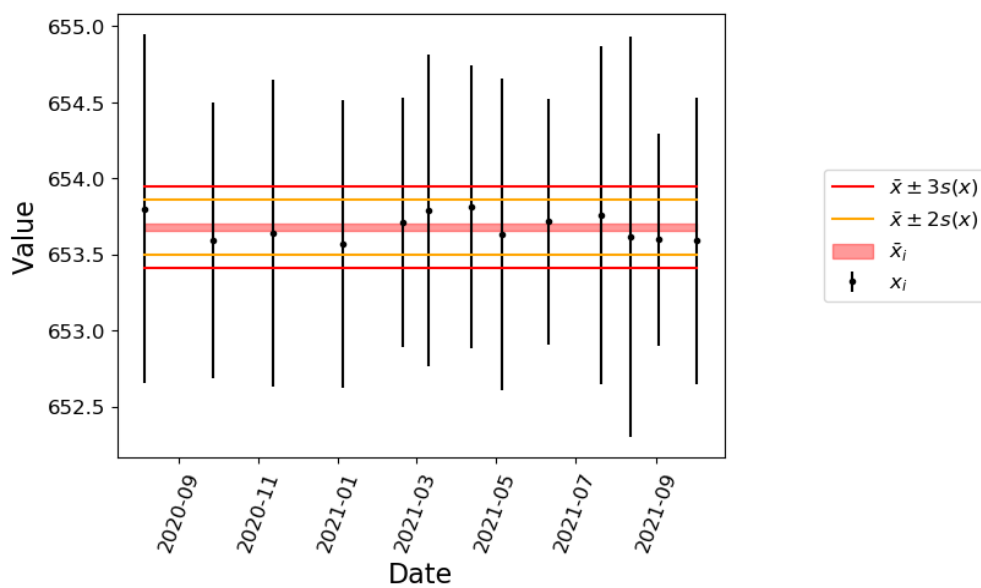
Fig. 1 TDCR values obtained for the ^{14}C source CCT since August 2020



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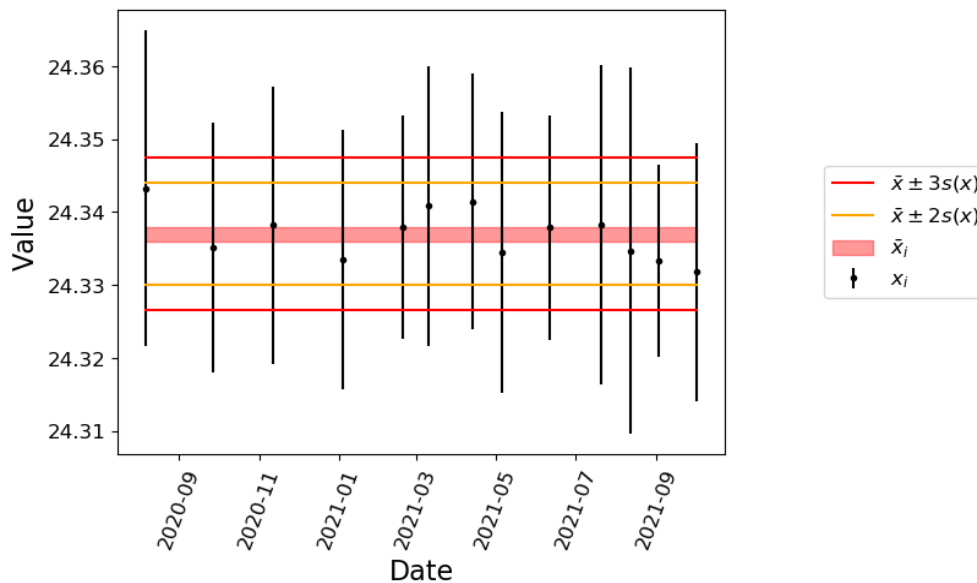
Fig. 2 I_0 values obtained for the ^{14}C source CCT since August 2020



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Fig. 2 I_1 values obtained for the ^{14}C source CCT since August 2020



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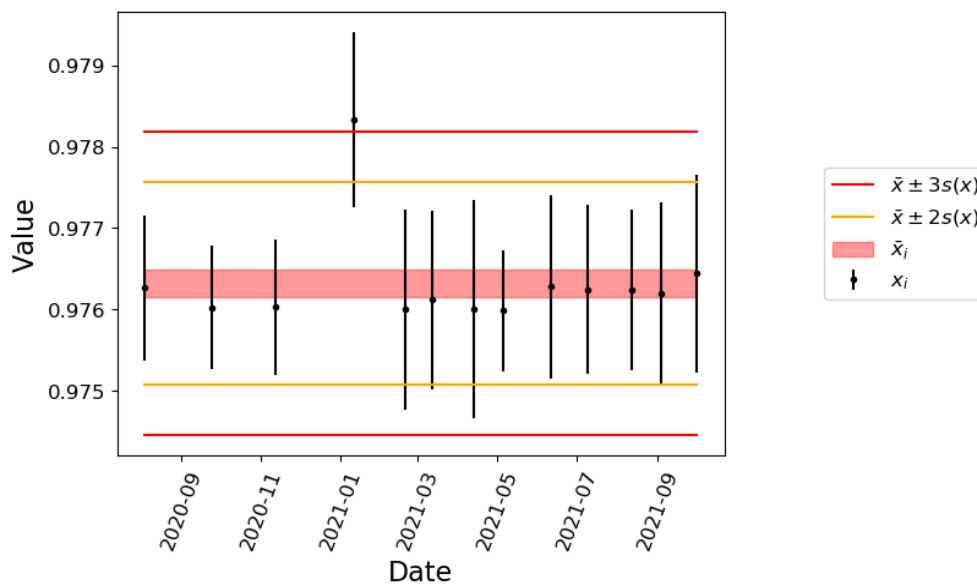
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Fig. 2 I_2 values obtained for the ^{14}C source CCT since August 2020

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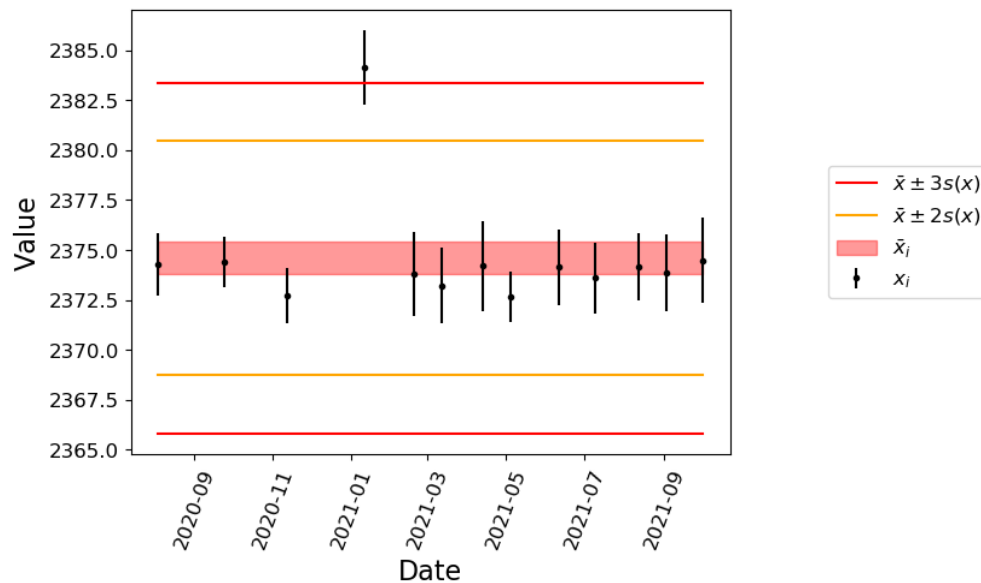
e. Toluene-based ^{14}C standard CMF



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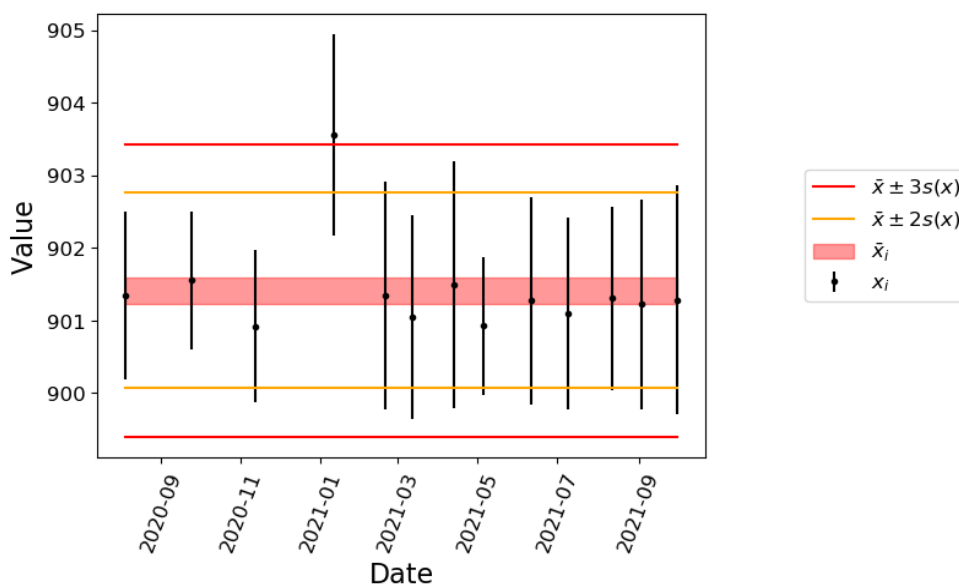
Fig. 1 TDCR values obtained for the ^{14}C source CMF since August 2020



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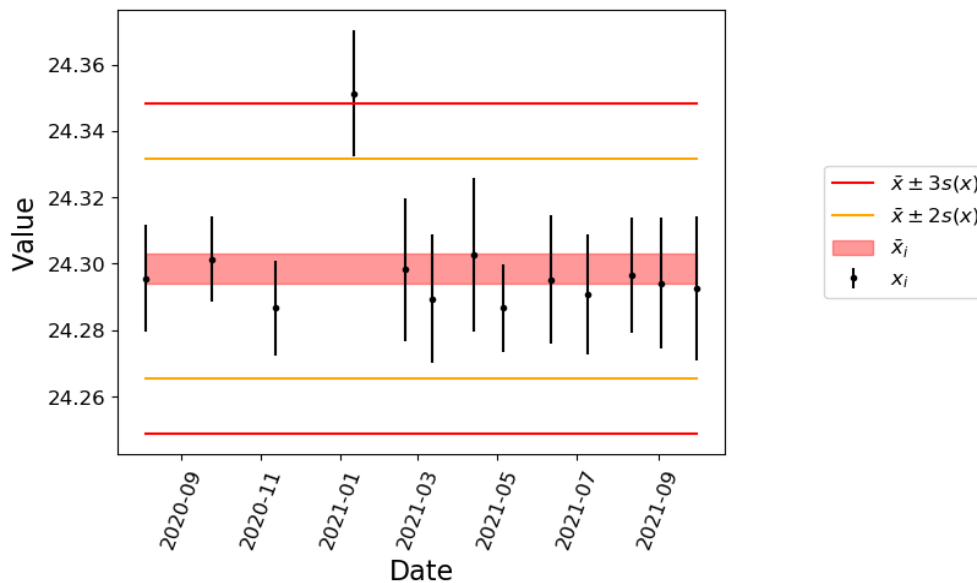
Fig. 2 I_0 values obtained for the ^{14}C source CMF since August 2020



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Fig. 2 I_1 values obtained for the ^{14}C source CMF since August 2020



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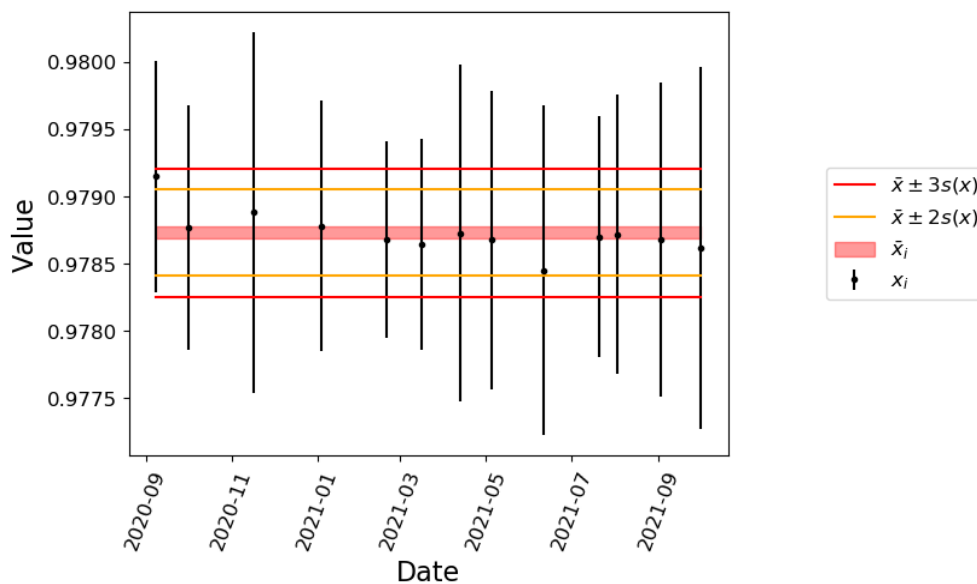
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Fig. 2 I_2 values obtained for the ^{14}C source CMF since August 2020

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f. Toluene-based ^{14}C standard C14-117TOL

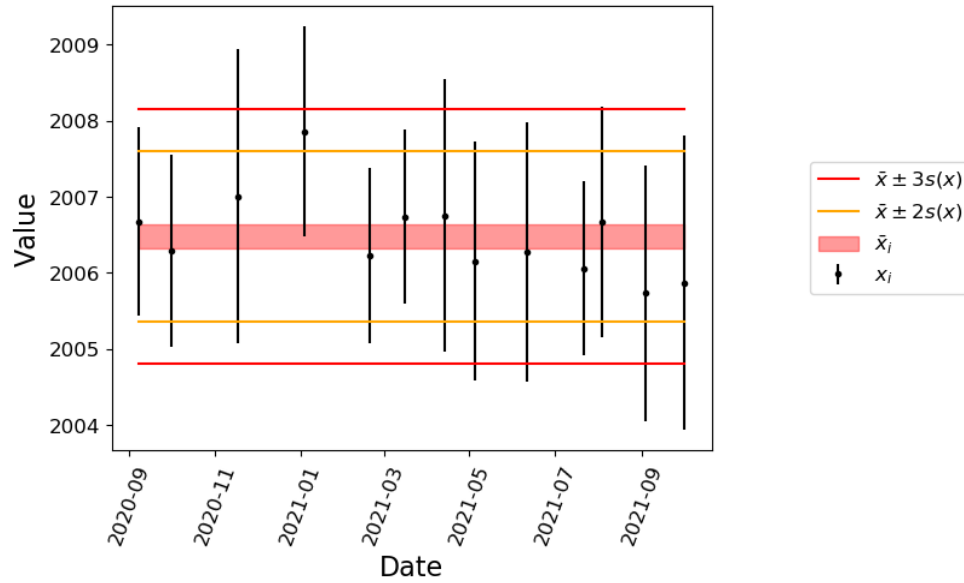


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Fig. 1 TDCR values obtained for the ^{14}C source C14-117TOL since August 2020

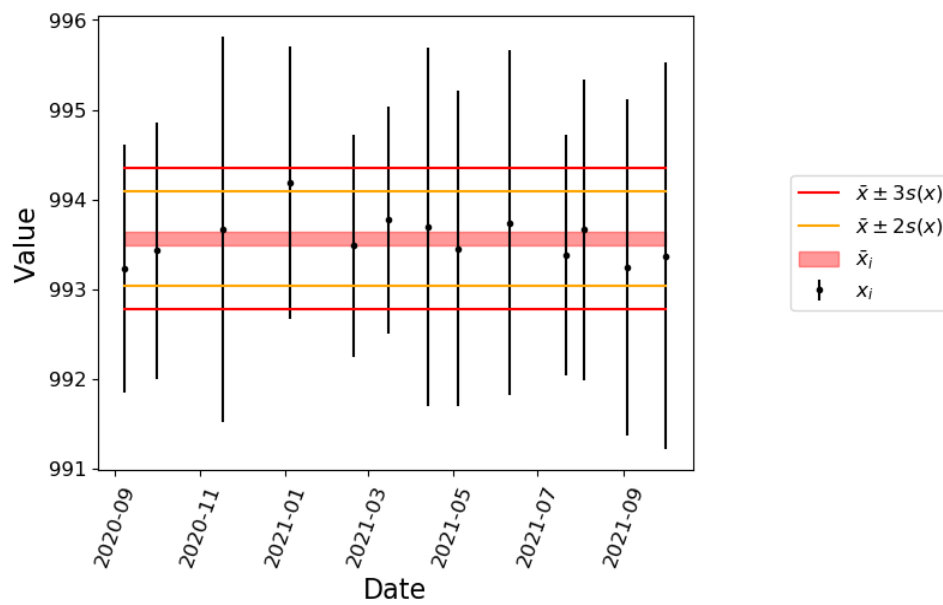
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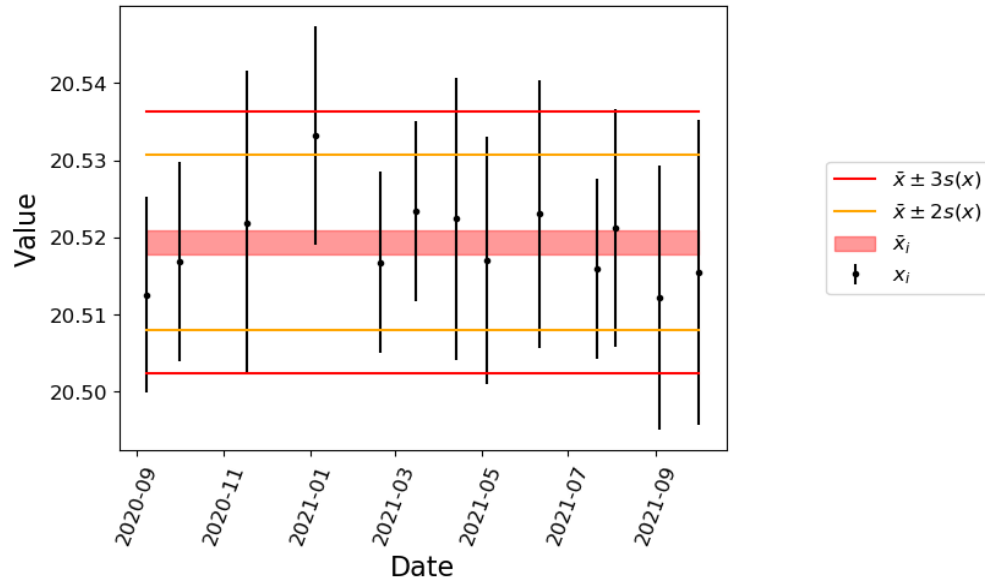
Fig. 2 I_0 values obtained for the ^{14}C source C14-117TOL since August 2020



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Fig. 2 I_1 values obtained for the ^{14}C source C14-117TOL since August 2020



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Fig. 2 I_2 values obtained for the ^{14}C source C14-117TOL since August 2020

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