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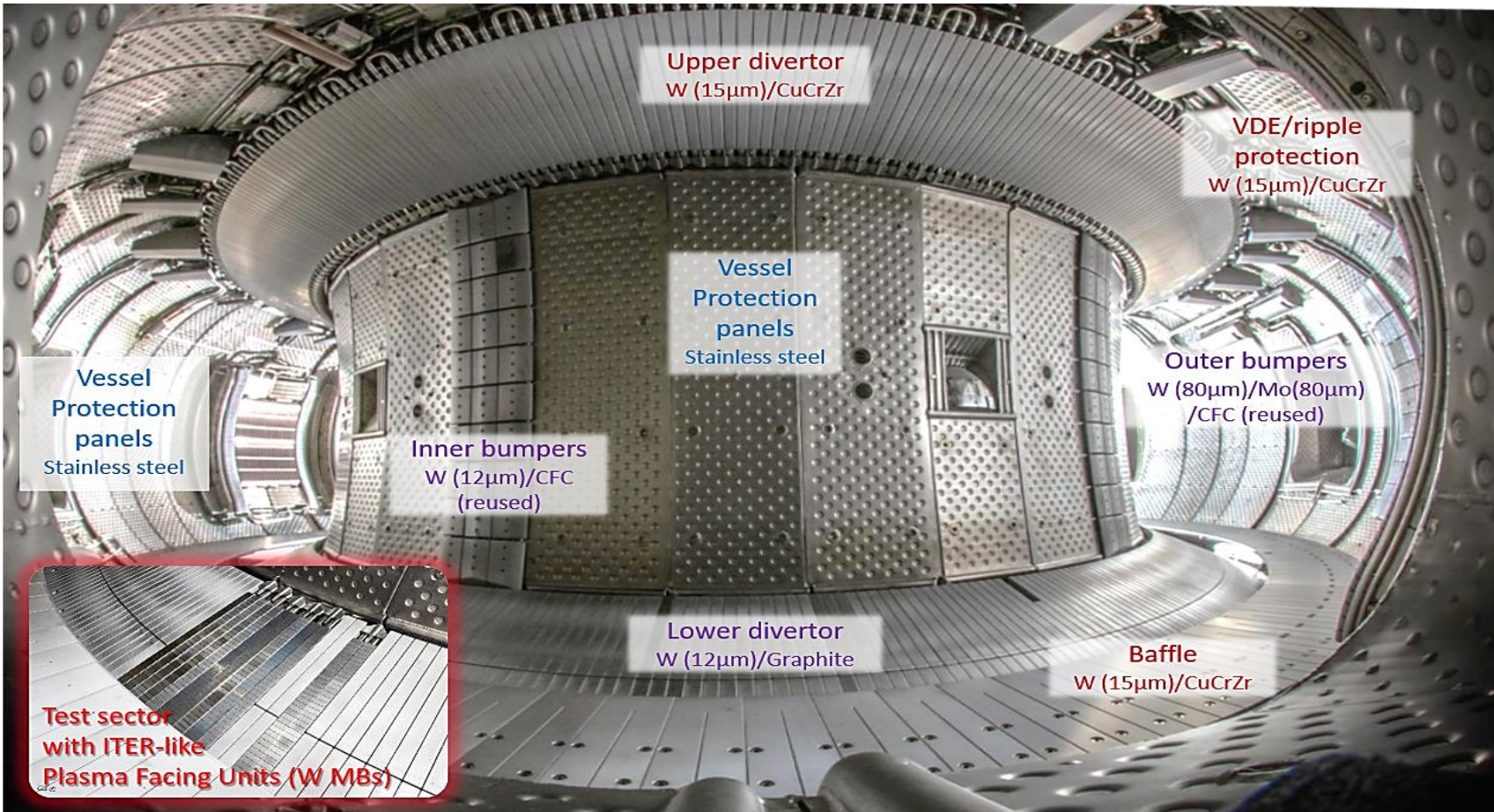


# Core radiative collapse in WEST LHCD plasmas: characterization and integrated modelling

EPS 30/06/2022

**V. Ostuni, J. Morales, J-F Artaud, C. Bourdelle, P. Manas, Y. Peysson, N. Fedorczak, R. Dumont, M. Goniche, P. Maget & WEST team** [west.cea.fr/WESTteam](http://west.cea.fr/WESTteam)

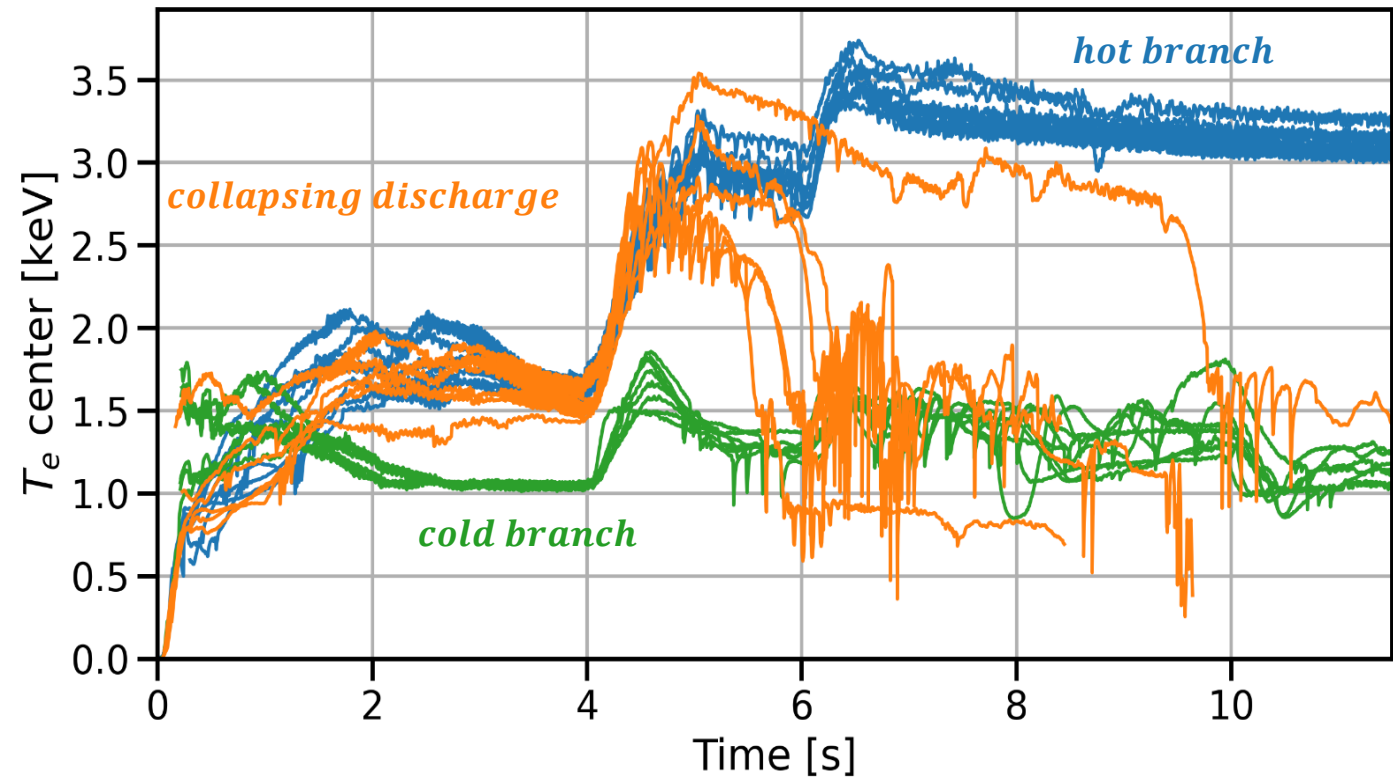
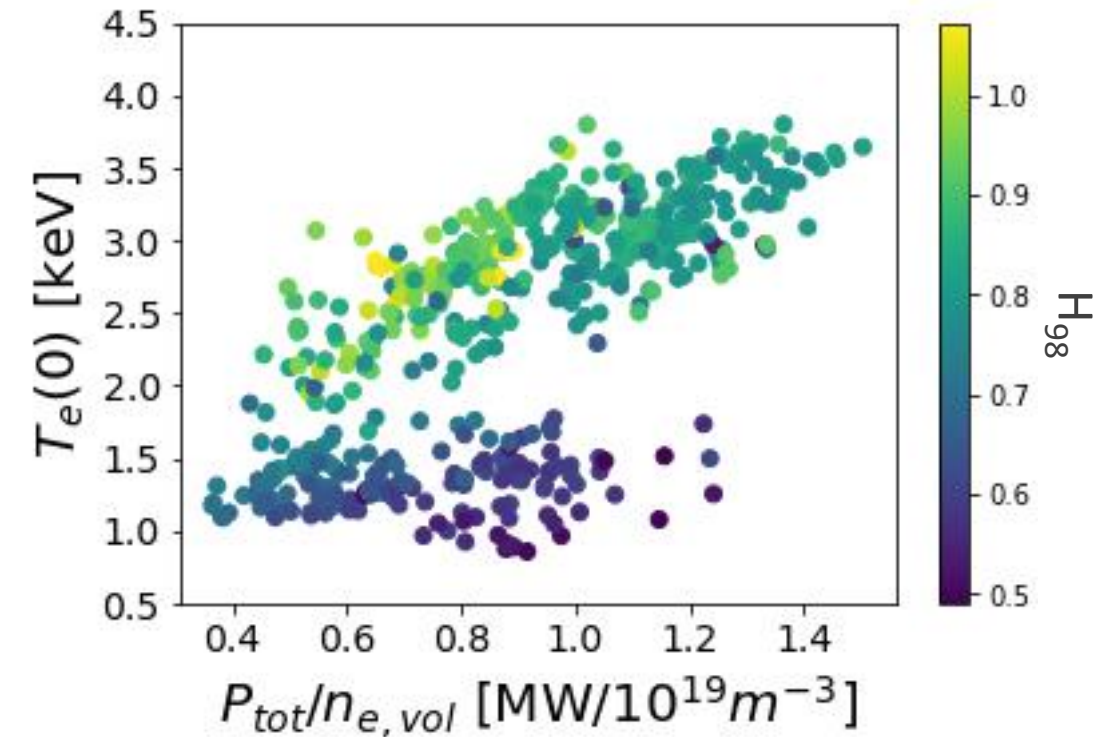




$B_T$	3.7 T
$R$	2.5 m
$a$	0.4/0.5 m
$P_{LH}$	2 launchers 7 MW
$P_{ICRH}$	3 antennas 9 MW

Dominant **electron heating**, no **external torque**.

In this talk: **L-mode** experiments, in **deuterium** with **LHCD** dominant heating.



Two questions arise from this analysis:

- Why are there two branches?

- Why do the **25%** of the **discharges** in the hot branch **collapse**?

Phenomena observed in **JET** during  $I_p$  ramp up and **FTU**

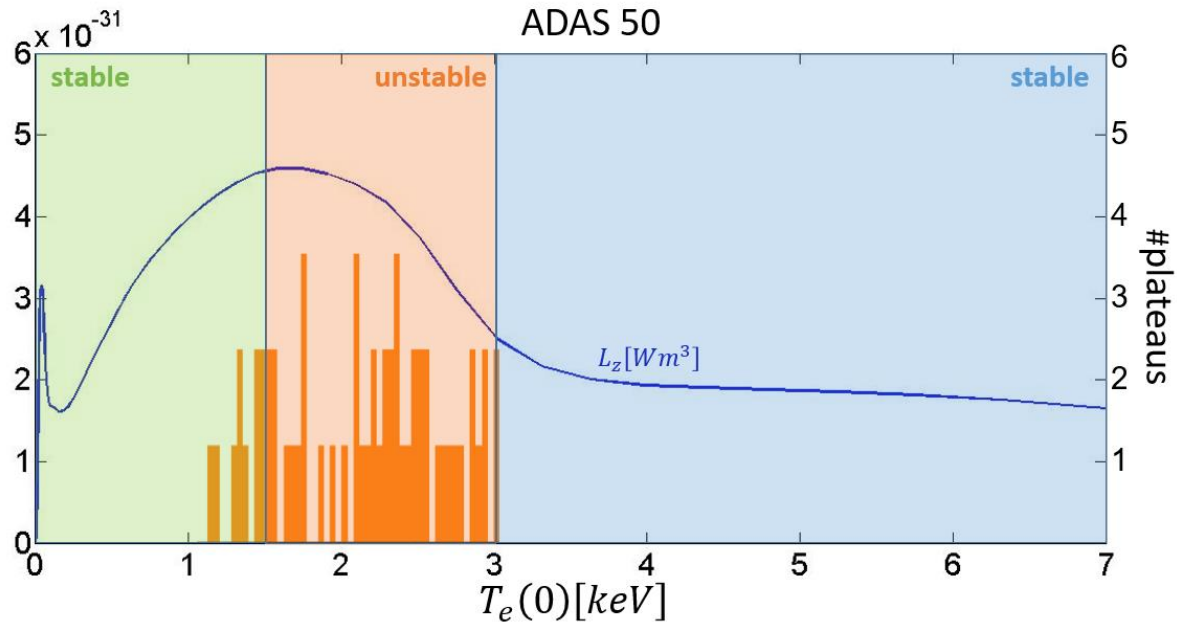
[C.D. Challis NF 2020]

[P. Buratti PPCF 1997]

- **Characterize** and **understand** the dynamics of the **radiative collapse**;
- **Reproduce** the rapid **collapse** of **central temperature** using an integrated modelling framework;
- Conclusions.

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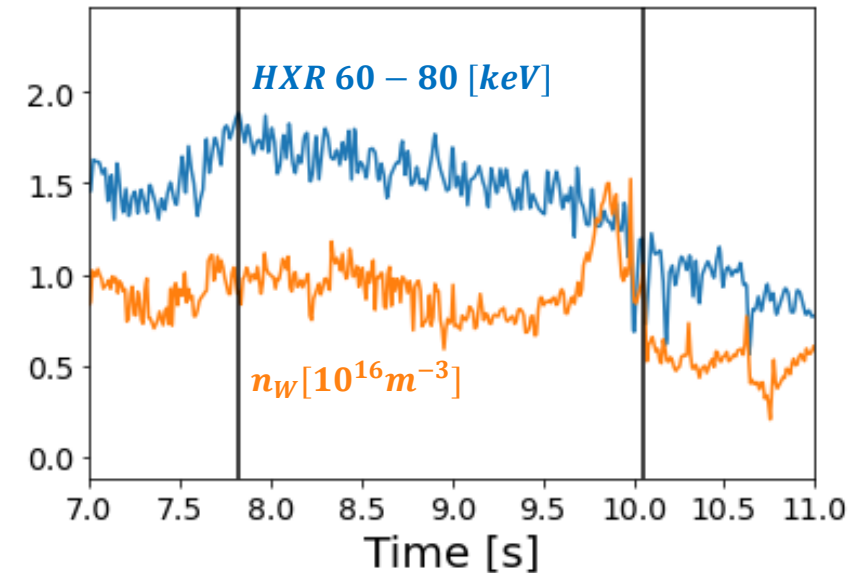
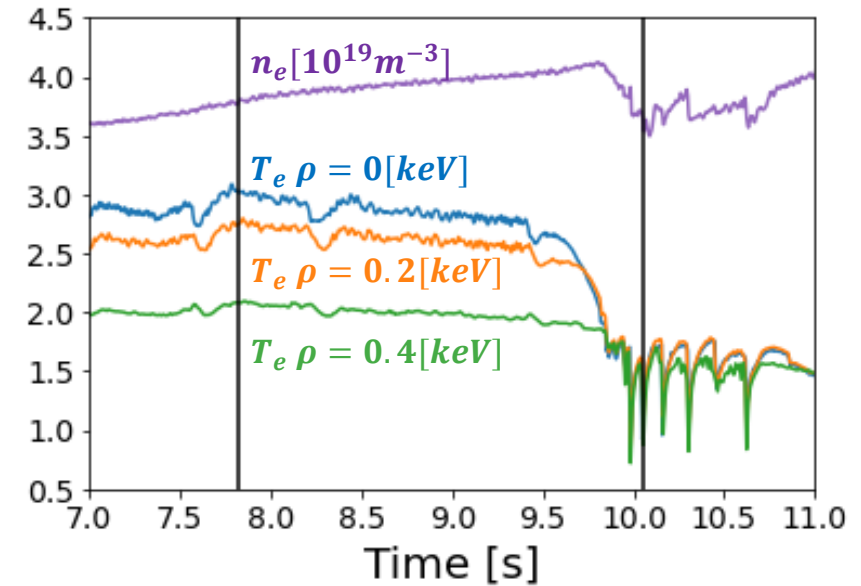




Prior to the collapse:

- a slow **increase** of **electron density**;
- a slow **decrease** of **electron temperature**;
- a **constant W** density;
- a slow **decrease** of central **HXR signal**: signature of a lower fast electron production by **LHCD**.

During the collapse : **peaking of the tungsten profile**.



- Characterize and understand the dynamics of the radiative collapse;
- **Reproduce the rapid collapse of central temperature using an integrated modelling framework;**
- Conclusions.

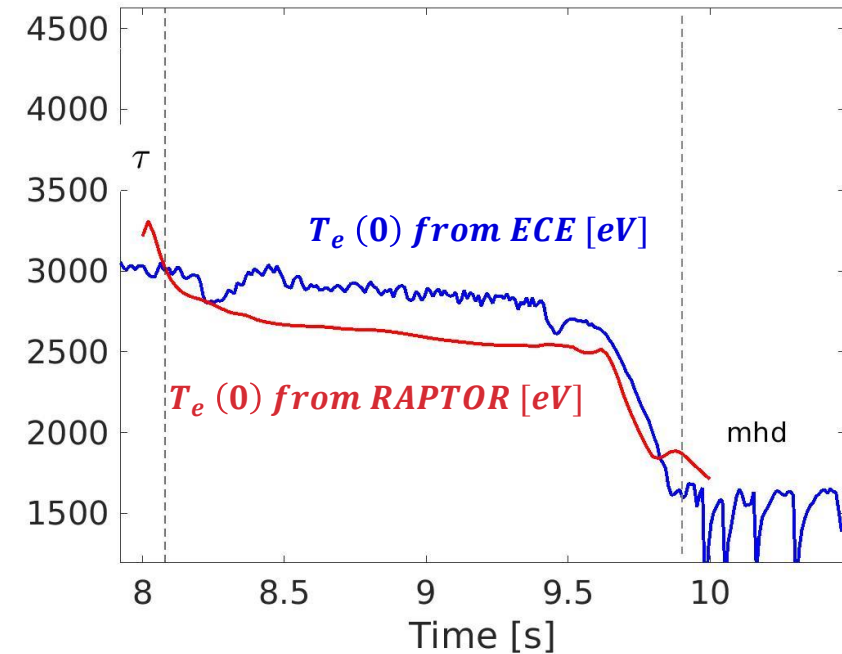
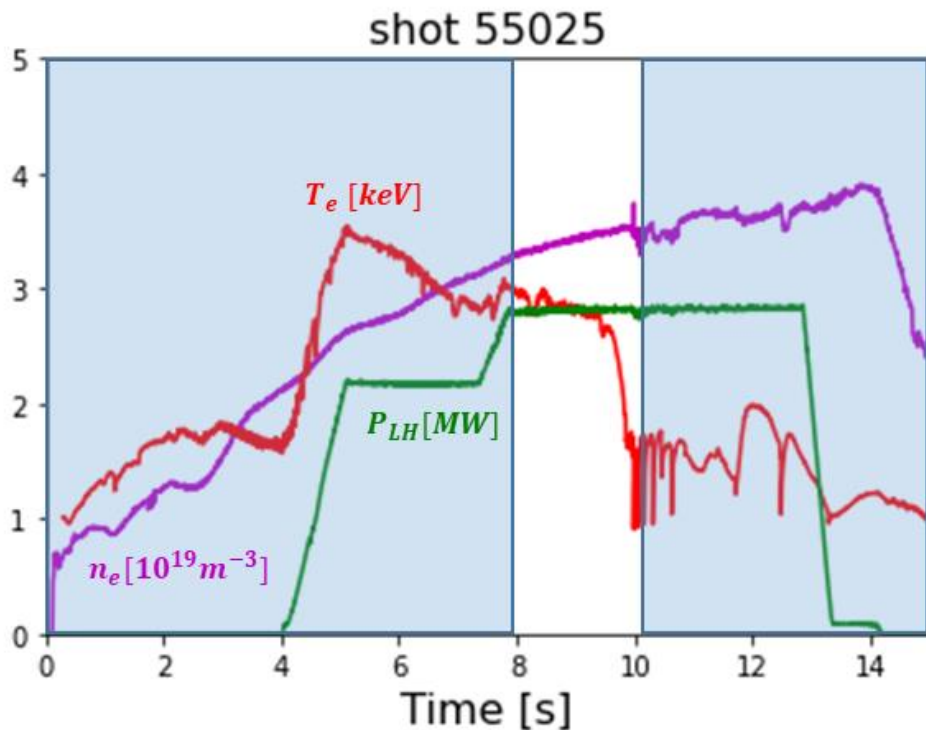


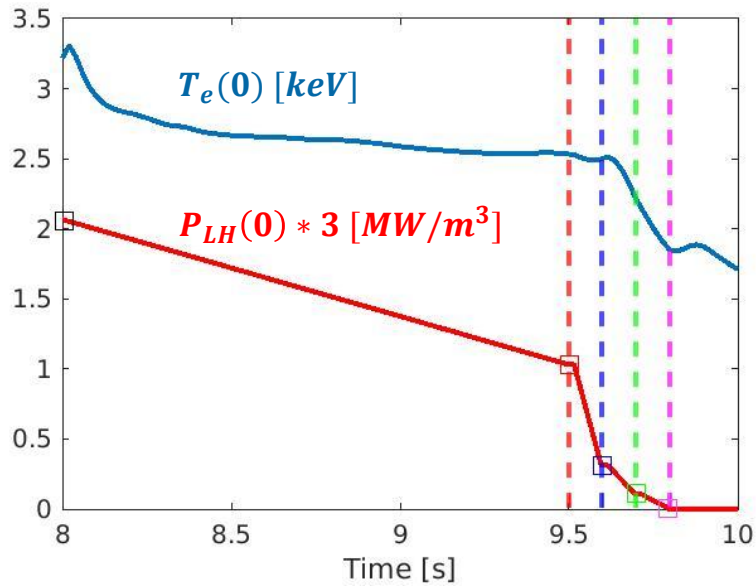
1D transport code: **RAPTOR** [F. Felici NF 2018]

**Turbulent transport: 10D Neural Network version** trained on **QuaLiKiz**.

[K. van de Plassche PP 2020]

$T_e$ and $T_i$	Predicted using QLKNN-10D
$P_{LHCD}$	LUKE
$n_e$	interpretative
$n_w$	Interpretative, based on bolometry inversion
$P_{rad}$	Using ADAS database

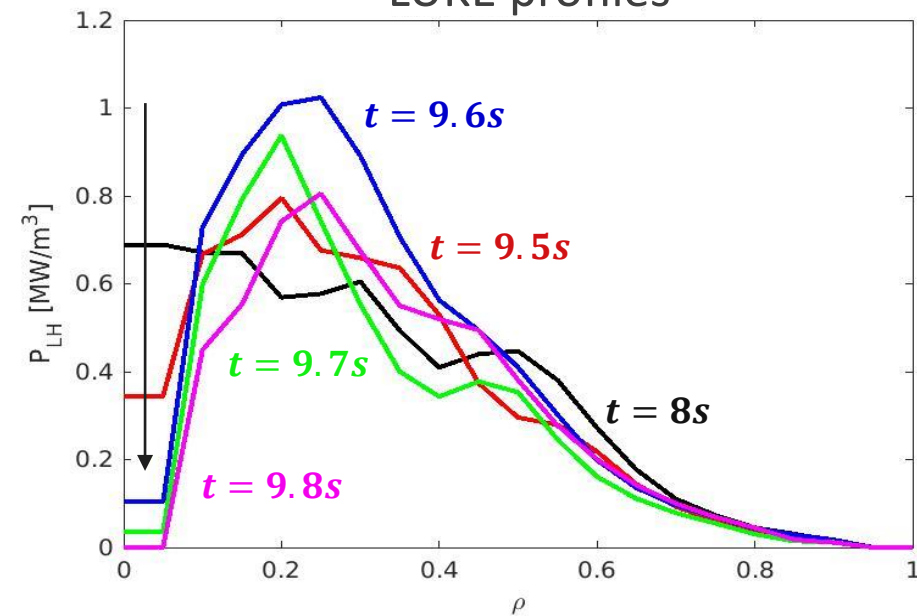




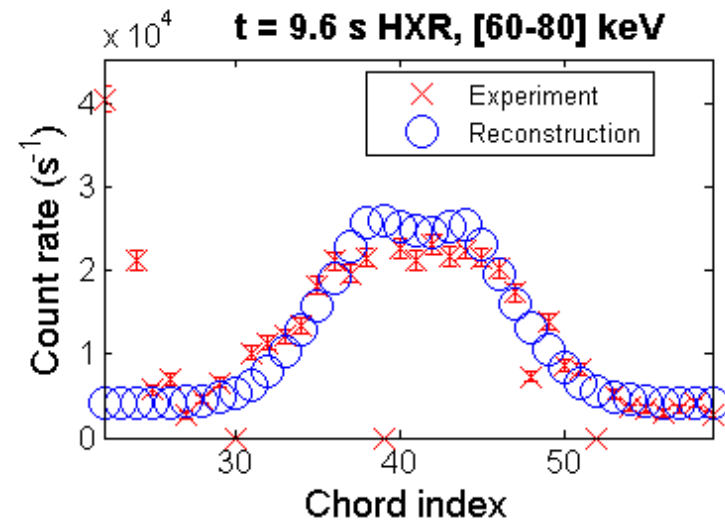
$d T_i$	$P_{LHCD}$	$n_e$
Predicted using QLKNN-10D	LUKE	interpretat

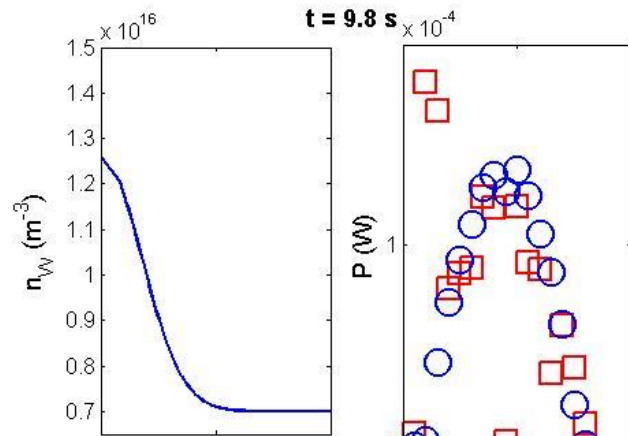
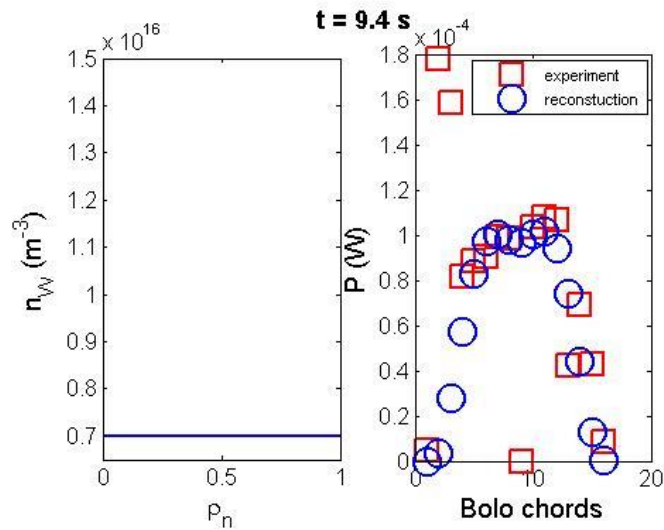
Predicted using QLKNN-10D

LUKE profiles



**LHCD** multipass absorption computed by **LUKE** code [Peysson FST 2014] **challenging** for core **absorption** computation.

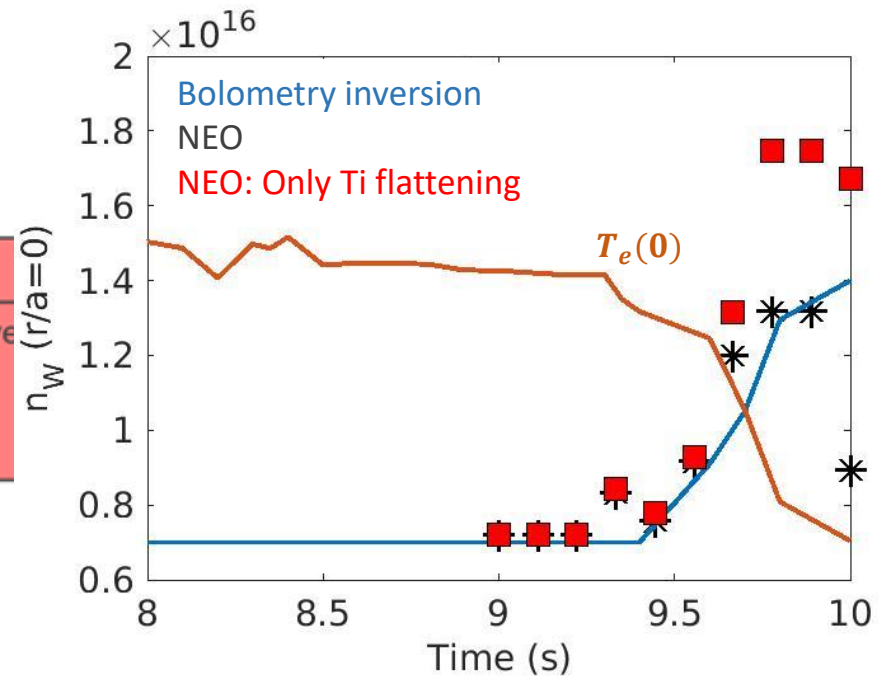


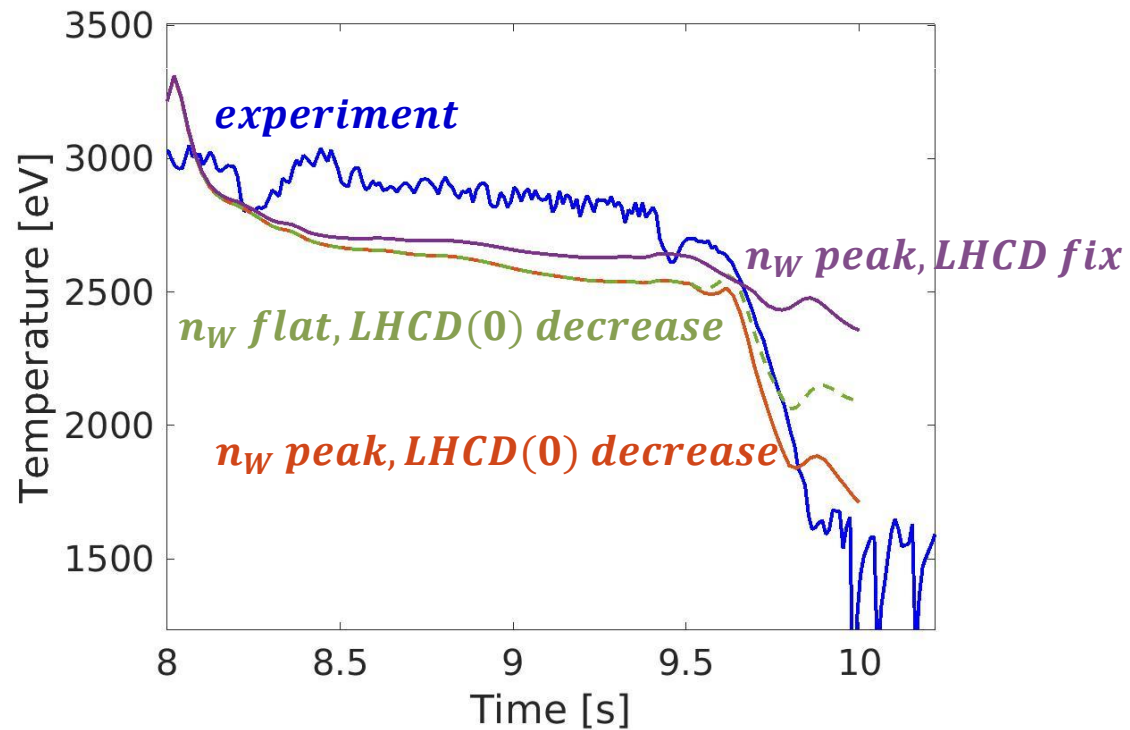


$T_e$ and $T_i$	$P_{\text{LHCD}}$	$n_e$	$n_W$
Predicted using QLKNN-10D	LUKE	interpretative	Interpretative based on bolometry inversion

$n_W$  approximated by a **Gaussian symmetric** around the center **on top of a flat profile**.

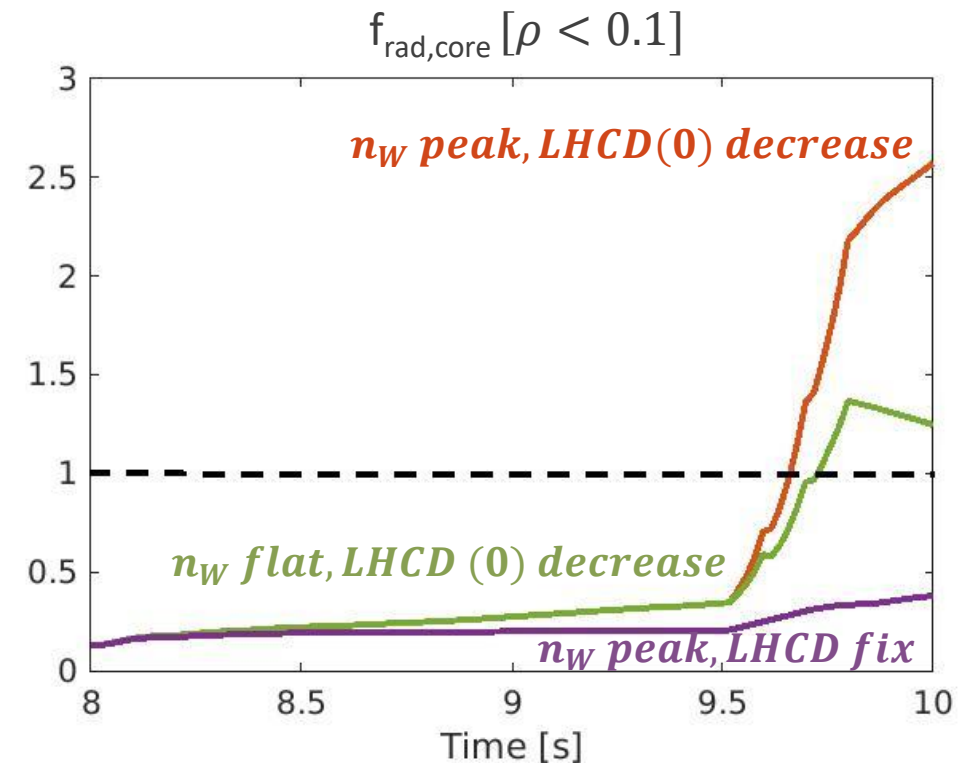
Neoclassical transport computed with NEO [Belli 2008]





Only with both the contribution of the **tungsten accumulation** and the **decrease of the central LHCD absorption** the **speed of collapse is reproduced**.

The **collapse** is not reproduced if the **radiated power** inside  $\rho=0.1$  does not **overtake the central electron heating**.





- Characterize and understand the dynamics of the radiative collapse;
- Reproduce the rapid collapse of central temperature using an integrated model framework;
- **Conclusions.**

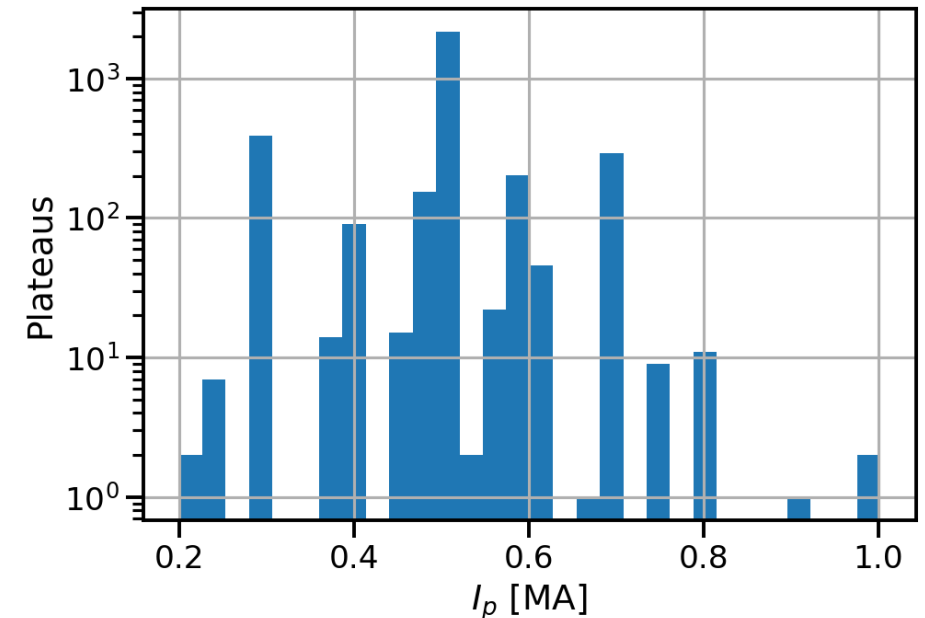
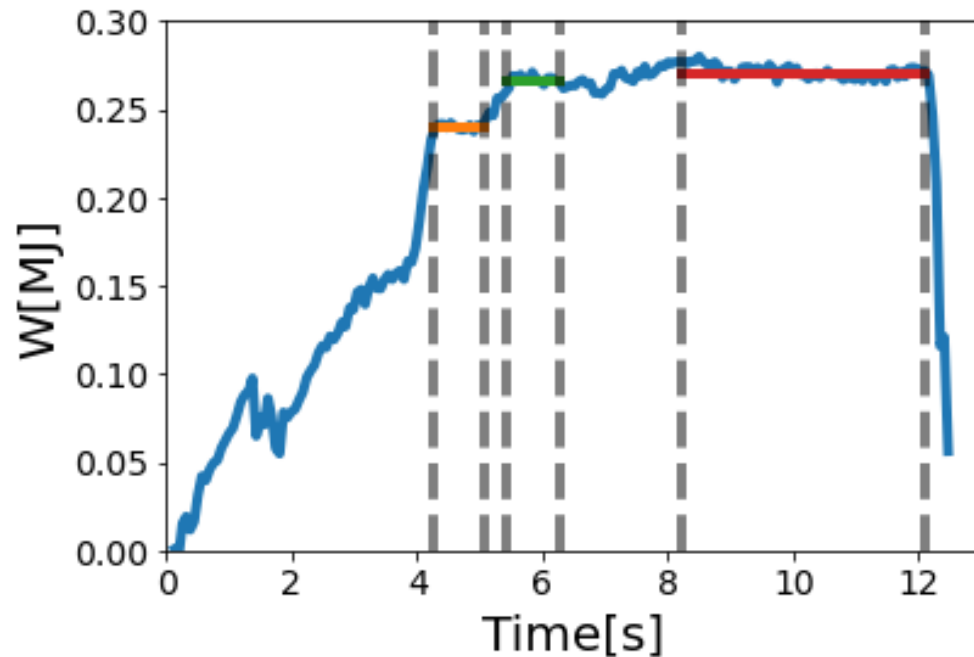
- In the database, L mode **plateaus** where  $T_e(0)$  rapidly **collapse** are identified in the range of **1.5 keV** and **3 keV**.
- Time sequence of  $T_e(0)$  collapse acceleration:
  - The **slight density rise** leads to **less on-axis LHCD power deposition** enhancing the central  $T_e$  **reduction**;
  - **$\nabla T$  reduction** in core leads to a **reduction of W core temperature screening** hence **to core W accumulation**.
- even in absence of external torque and particle source, **core electron heating is required** to burn-through W cooling factor.
  - Core **ICRH** absorption **optimization** ongoing;
  - **3MW of ECRH** from 2023.



## BACK UP SLIDES

DE LA RECHERCHE À L'INDUSTRIE

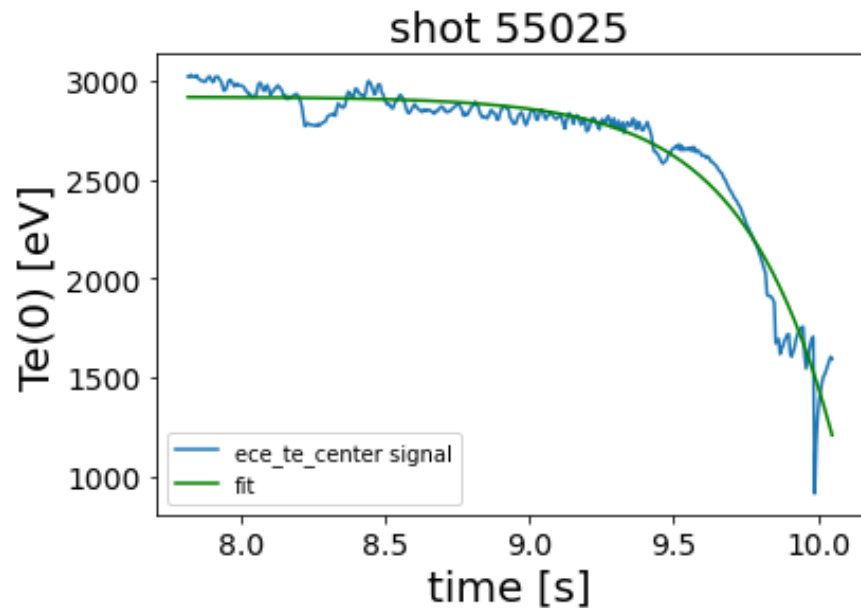
- The database contains the **mean values** and the **standard deviations** of different diagnostic measurements at each **plateau** (quasi-steady-state).
- The plateaus are identified intersecting the **total power plateaus** and the **plasma current plateaus**.
- There are **285** discharges with **732** plateaus in the database.



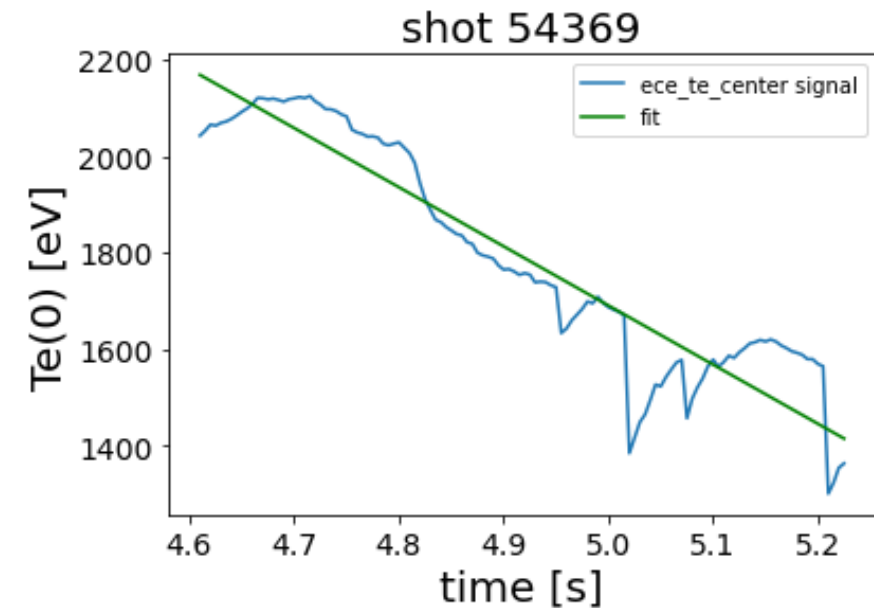


Automatic identification of radiative collapses within plateaus:

*exponential fit:  $T_e(0) = -e^{(a+bt)} + c$*   
 with  $a < 0$  for the concavity,

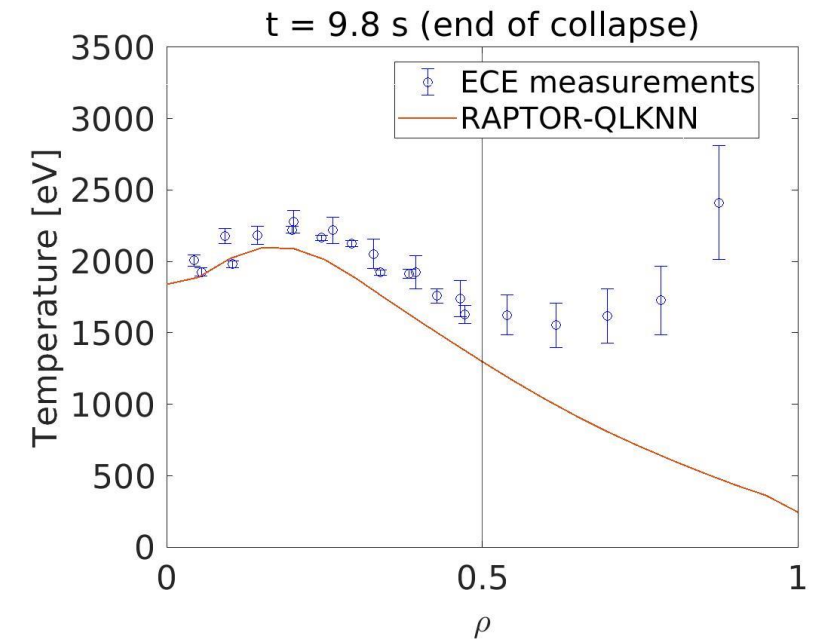
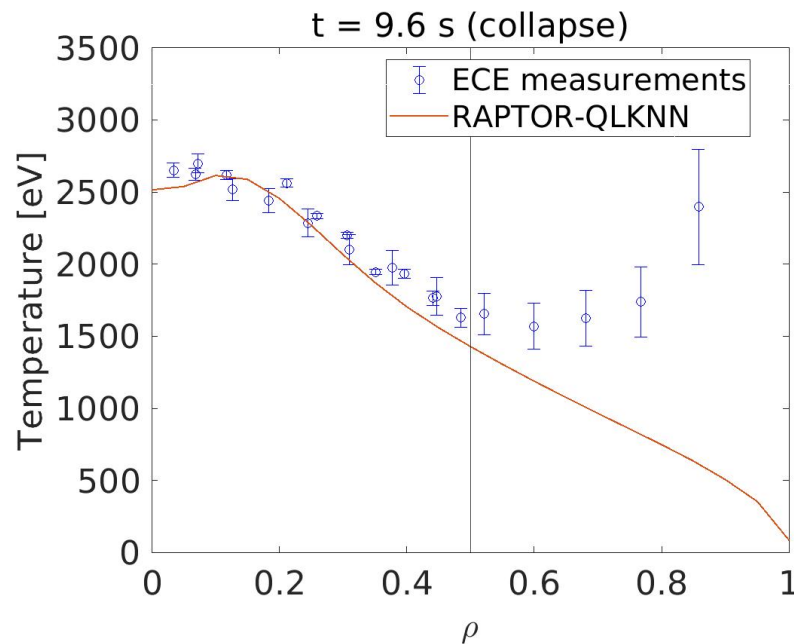
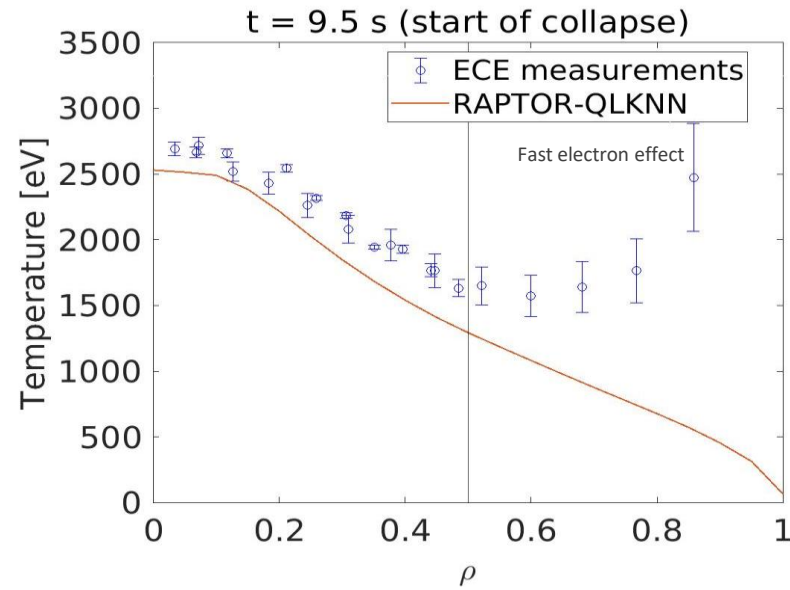


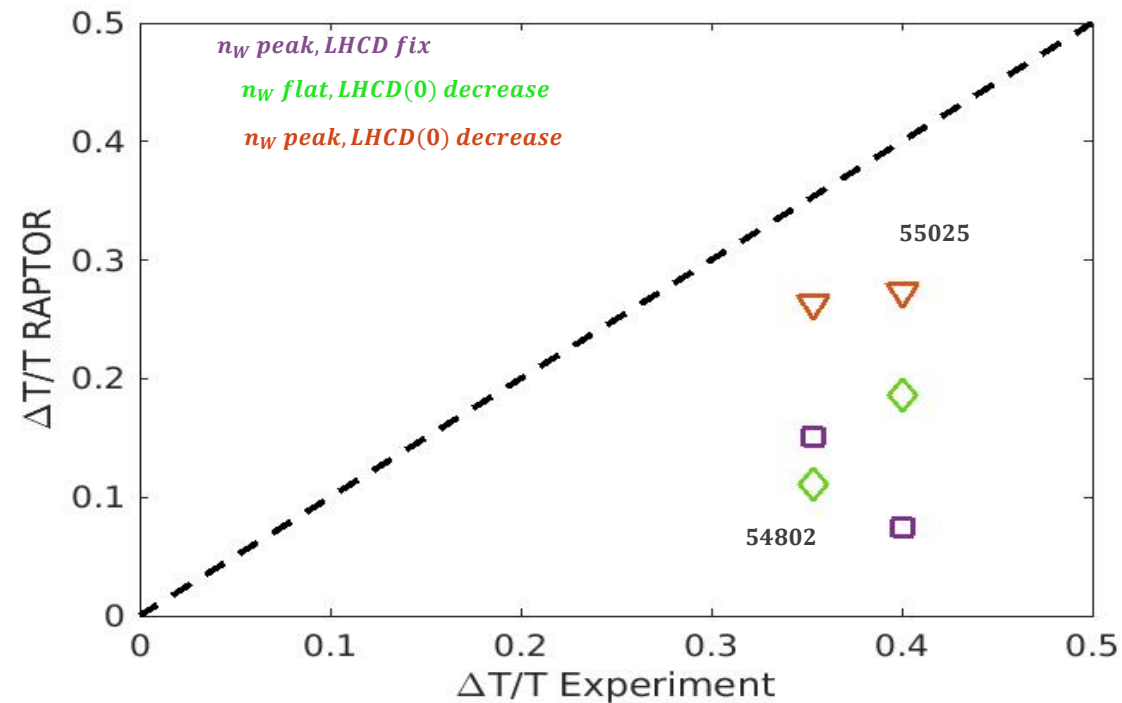
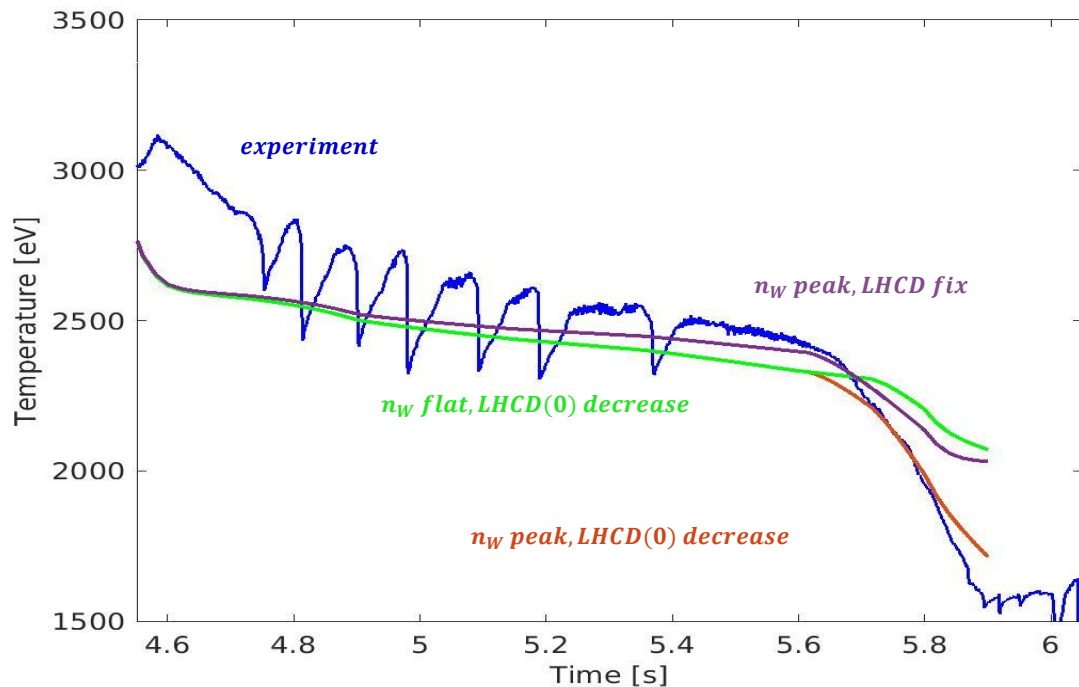
*linear function fit*  
 with *slope*  $< -830$



Three categories identified:

**unstable plateaus** where a collapse takes place, **cold stable plateaus** and remaining plateaus called **stable hot**.



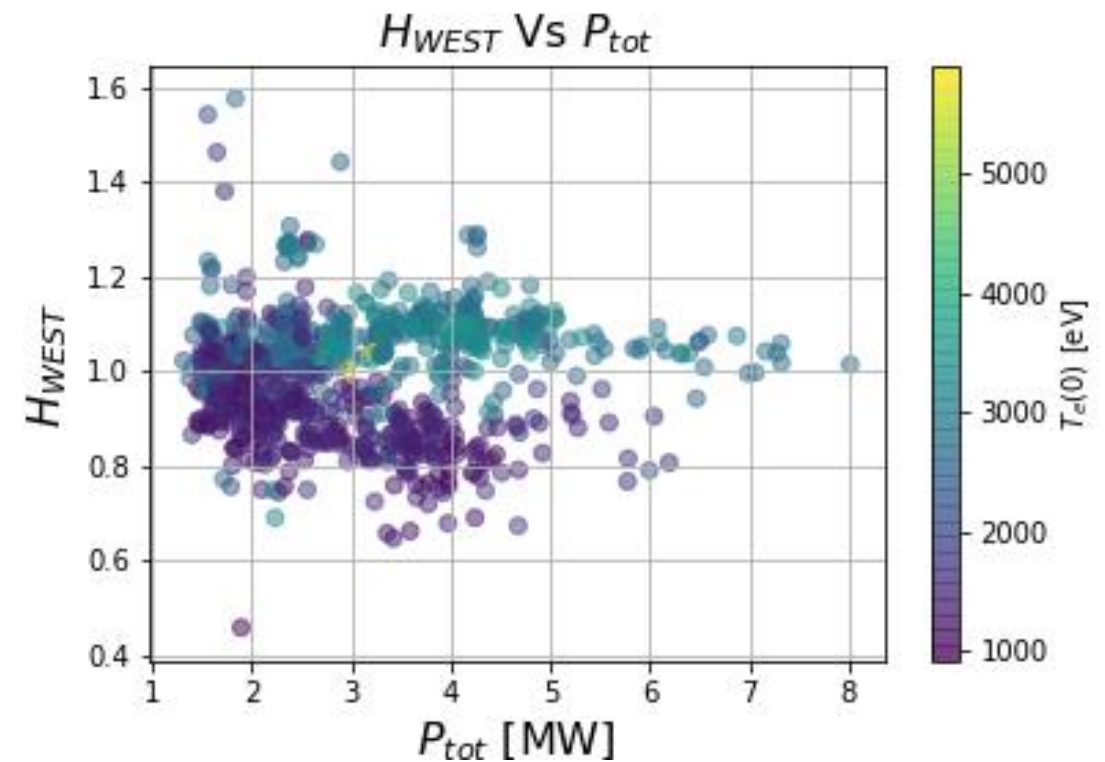
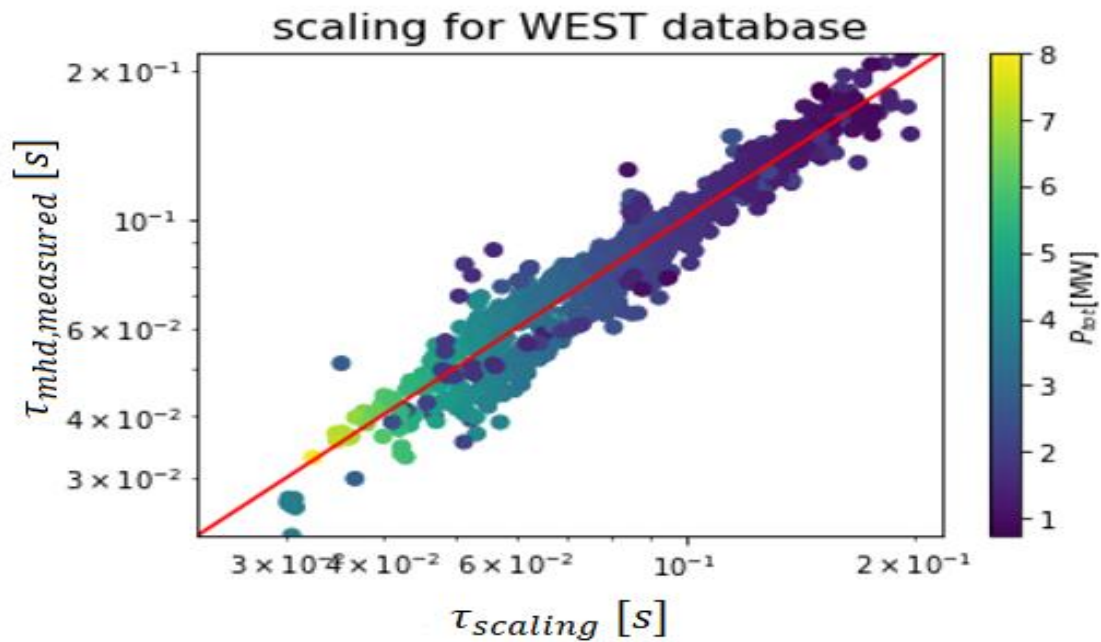


The same procedure was carried out for the **54802 discharge**.

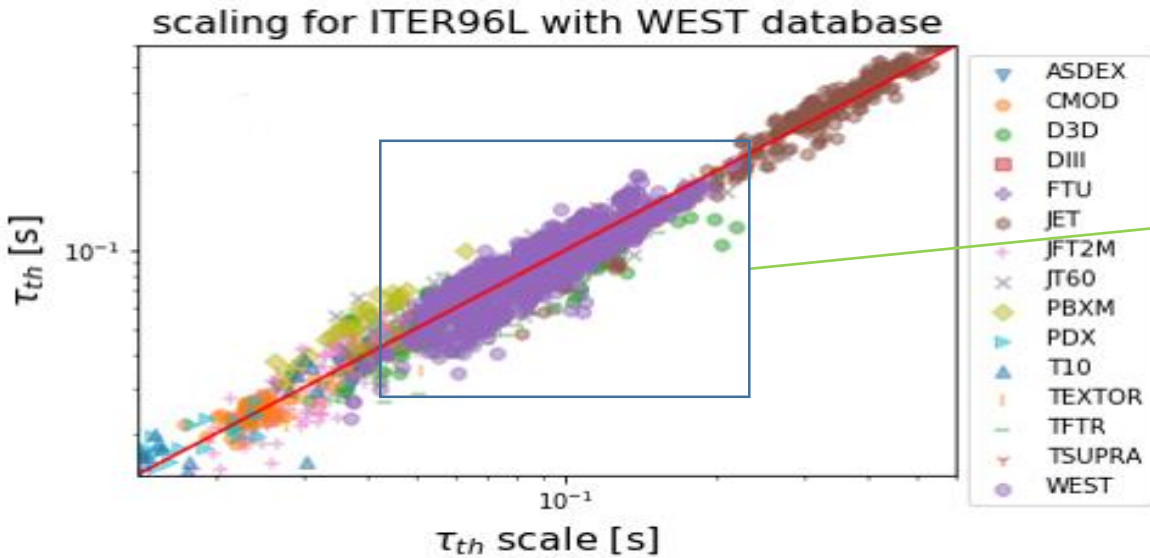
The temperature collapses are associated to **variable contributions** of **LHCD** core heating depletion and **tungsten peaking**.

$$\tau_{mhd,measured} = \frac{W_{mhd}}{P_{tot}} = \frac{\frac{3}{2} \int_V P dV}{P_{ohmic} + P_{aux}}$$

$$H_{WEST} = \frac{\tau_{mhd,measured}}{\tau_{scaling}}$$



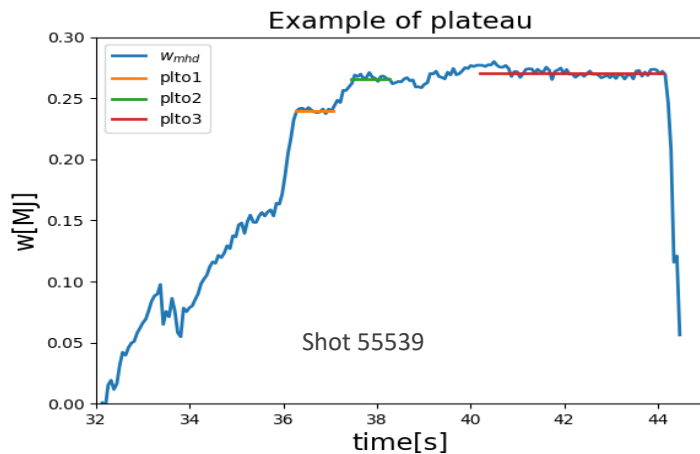




Filter:

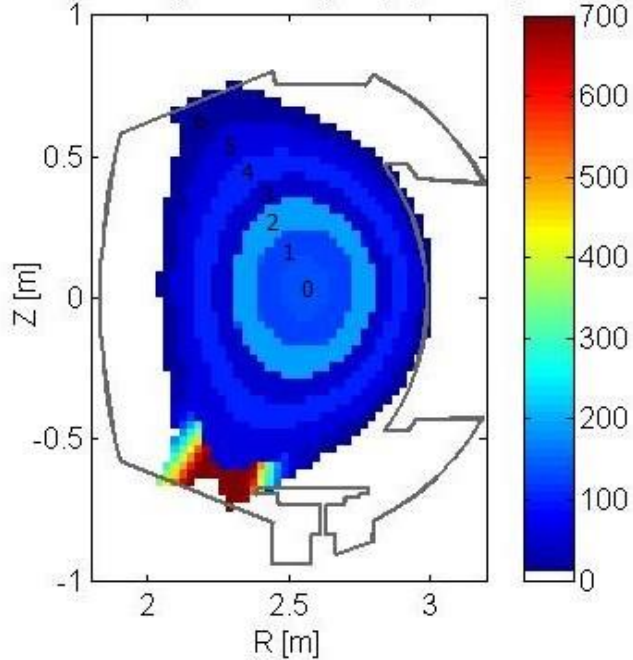
- Small reconstruction errors;
- Diverted plasma;
- L-mode plateaus;
- LSN
- Only deuterium shots

$$\tau = C I_p^{\alpha_{I_p}} B^{\alpha_B} P_{tot}^{\alpha_{P_{tot}}} n_e^{\alpha_{n_e}} M^{\alpha_M} R^{\alpha_R} \epsilon^{\alpha_\epsilon} k^{\alpha_k}$$

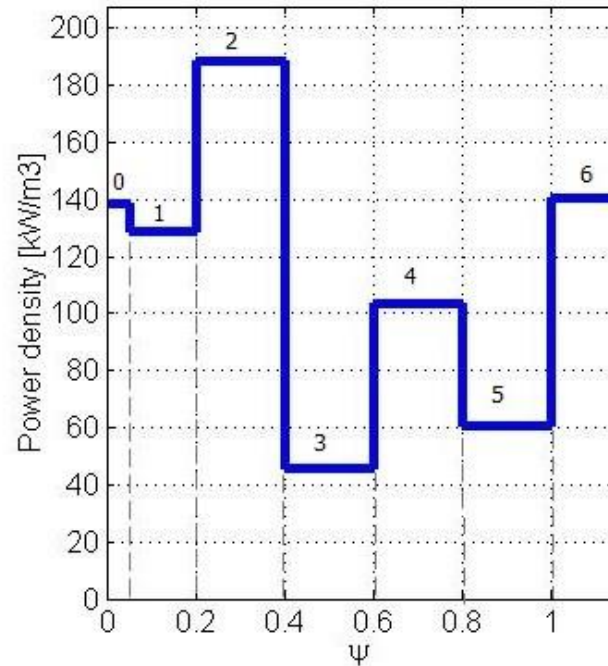


$$\begin{aligned} \alpha_{I_p} &= 1.35 \\ \alpha_{n_e} &= -0.16 \\ \alpha_{P_{tot}} &= -0.75 \end{aligned}$$

$\tau_{scaling}$

55025 bolometry emissivity map [kW/m<sup>3</sup>] 8<t<8.5

55025 inversion profile 8&lt;t&lt;8.5



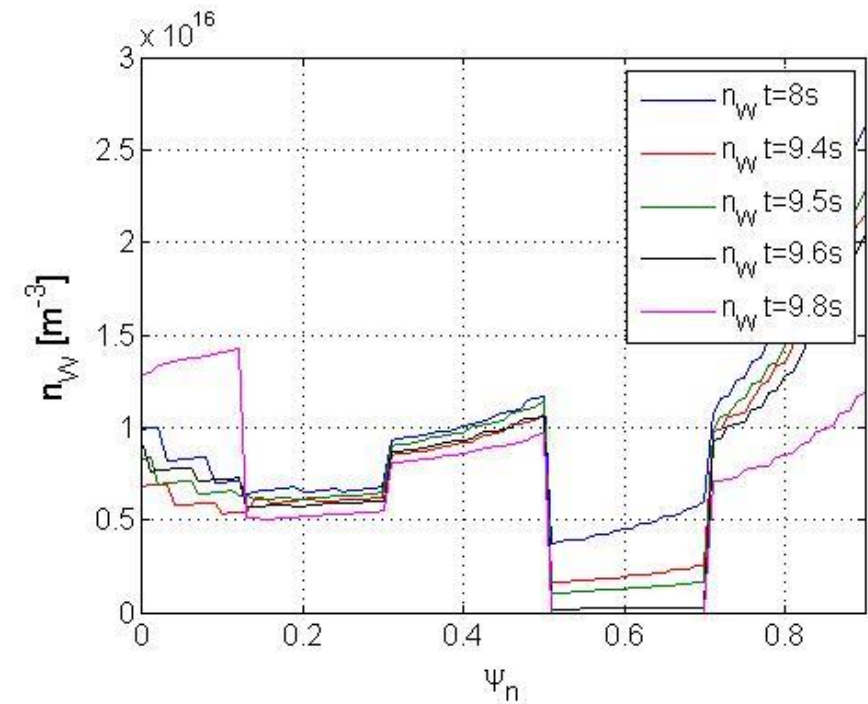
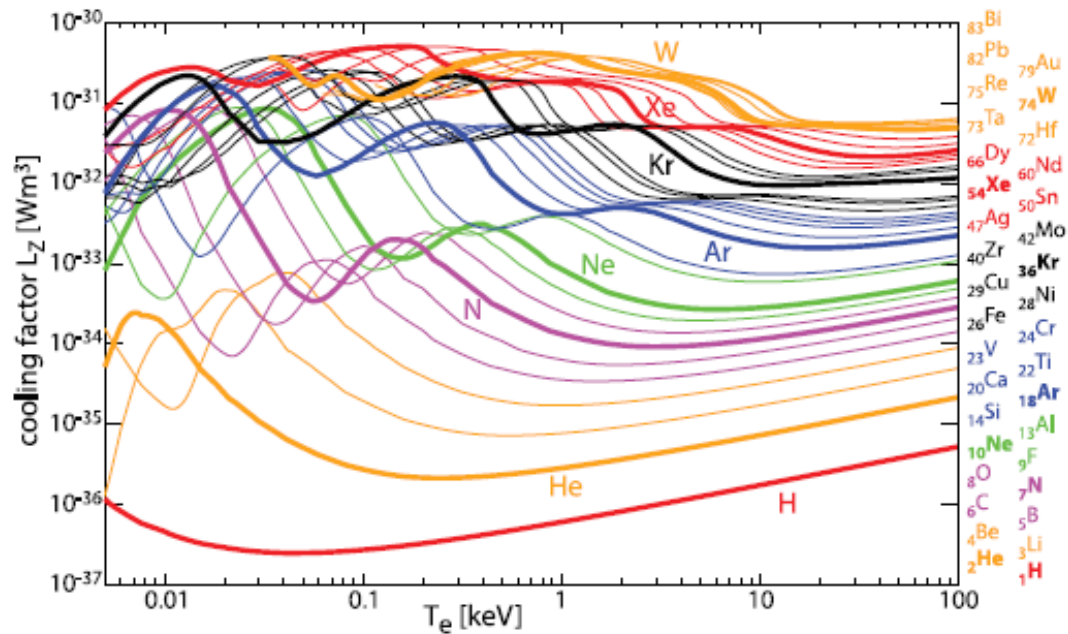
A **bolometer** measures a line-integrated value of the **local radiative emissivity** along a viewing line of sight.

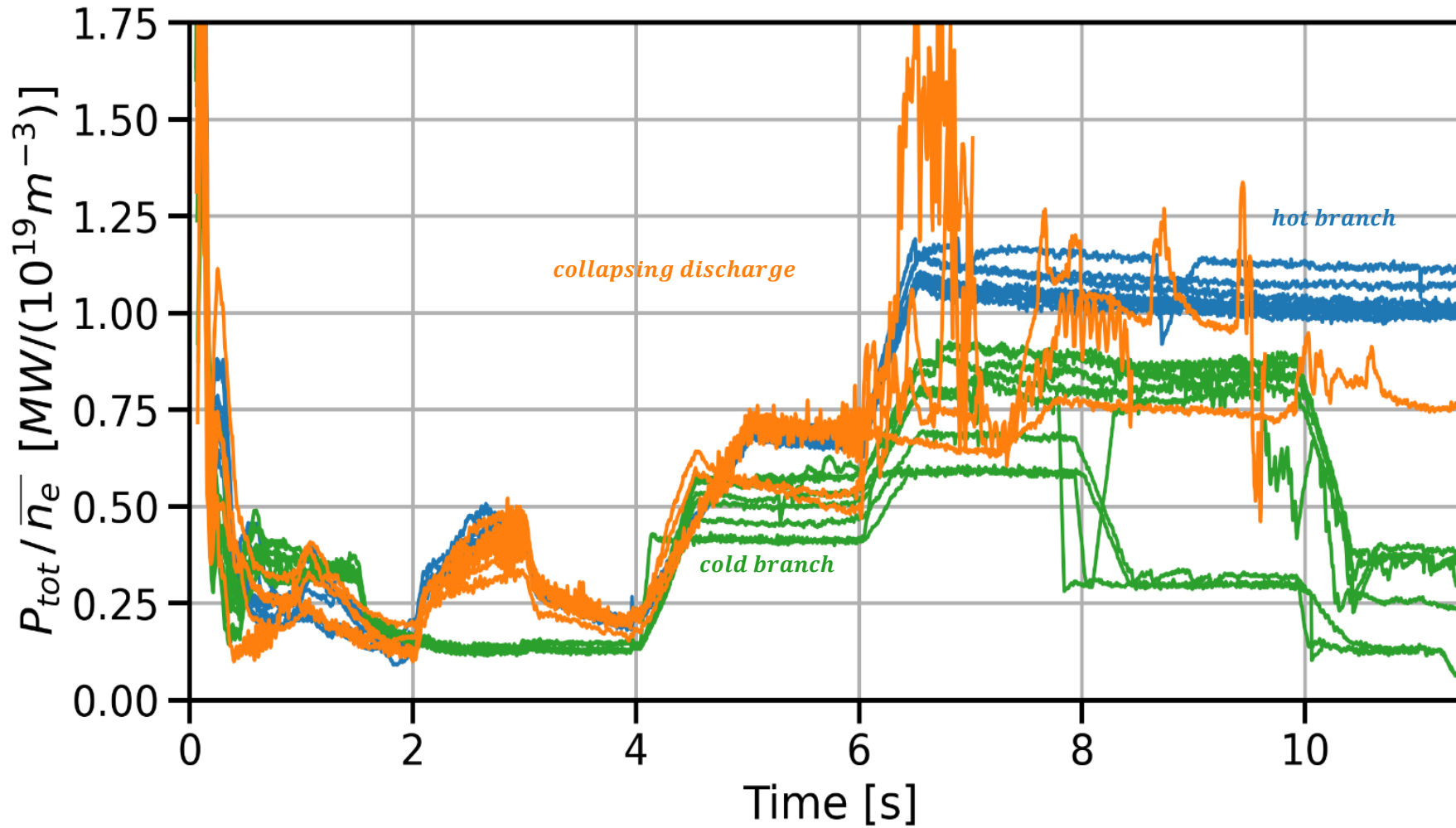
A **concentric layer decomposition** of local plasma emissivity is assumed together **with asymmetric factors** in the SOL and edge. The tomography inversion is computed.

The **power emission** in  $\frac{W}{m^3}$  **for each layer** are estimated.

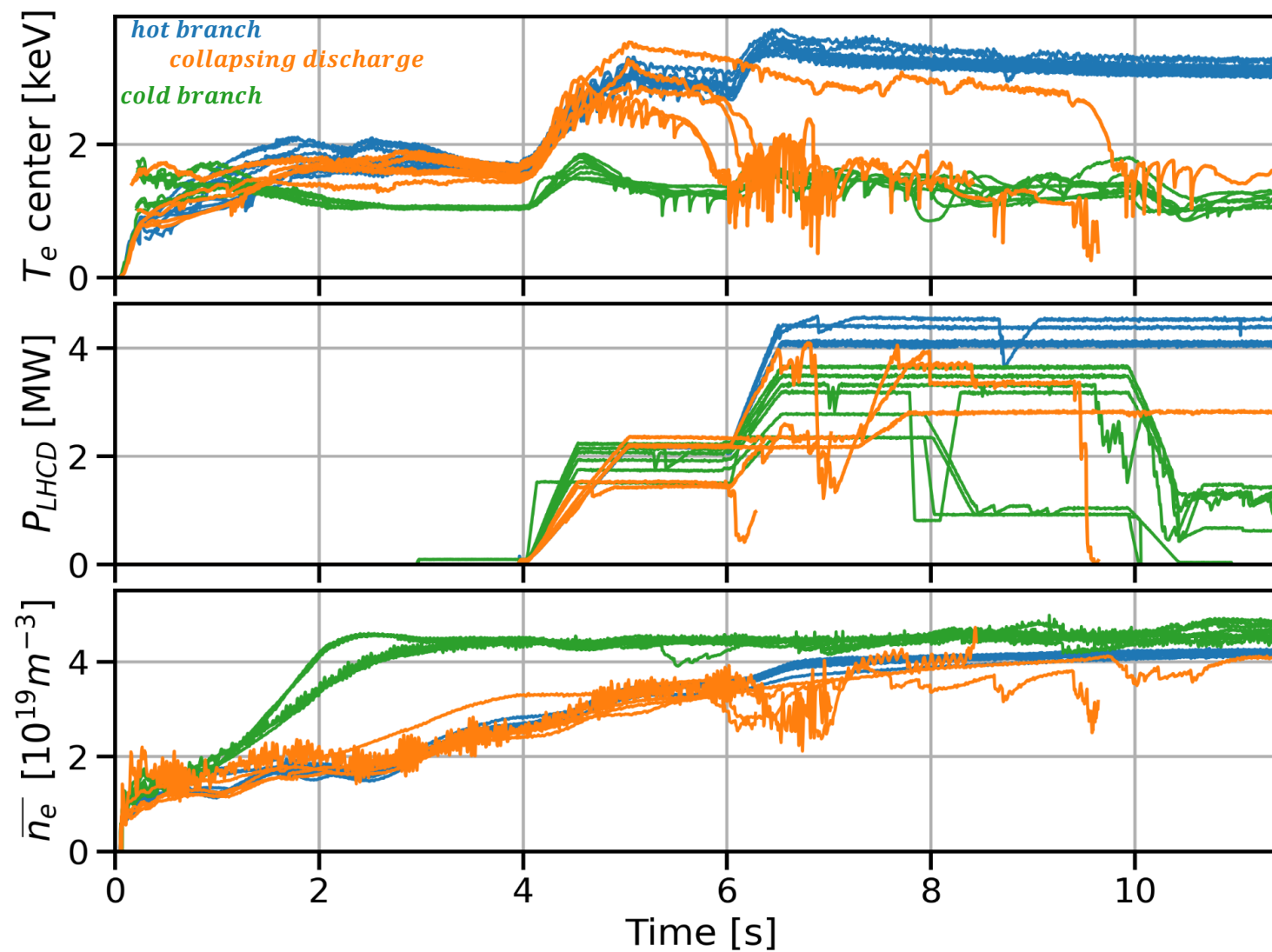
Above **1KeV** the **main radiator** is the **tungsten**.

Assuming that all the **emission radiation** comes from the **tungsten**, it is possible to use the inverse of bolometry to compute **its density** at each layer.





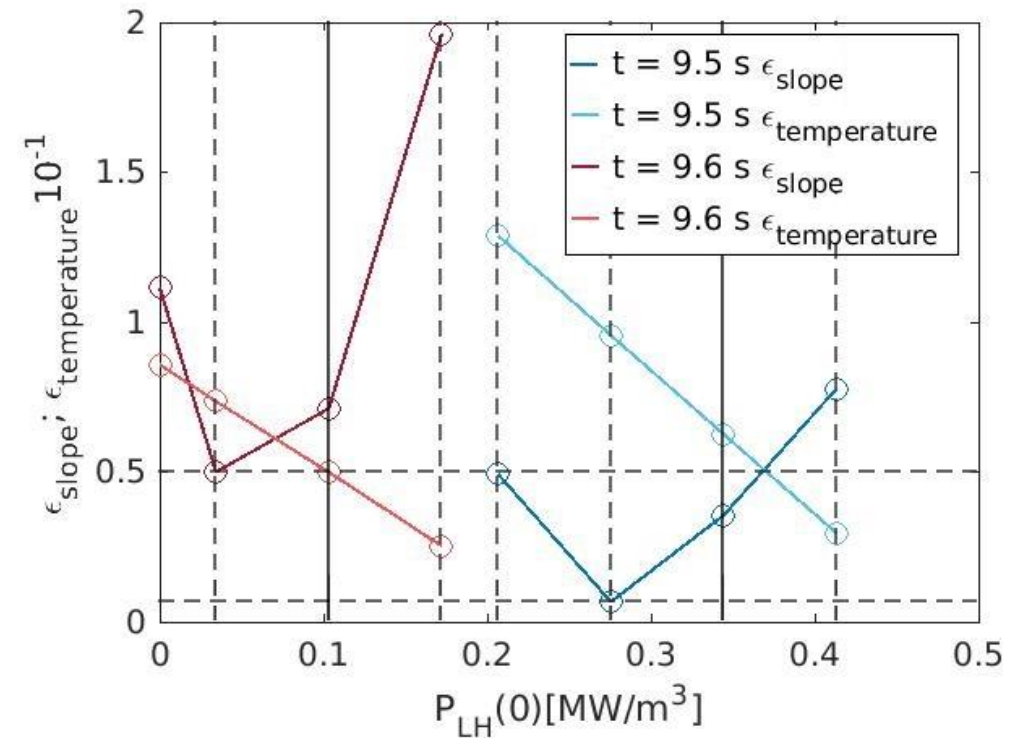




The central value is adjusted to find the optimum between  $\epsilon_{slope}$  and  $\epsilon_{temperature}$ .

$$\epsilon_{slope} = \frac{|(T_{t-1}^{exp} - T_t^{exp}) - (T_{t-1}^{simu} - T_t^{simu})|}{(T_{t-1}^{exp} - T_t^{exp})}$$

$$\epsilon_{temperature} = \frac{|T^{exp} - T^{simu}|}{T^{exp}}$$



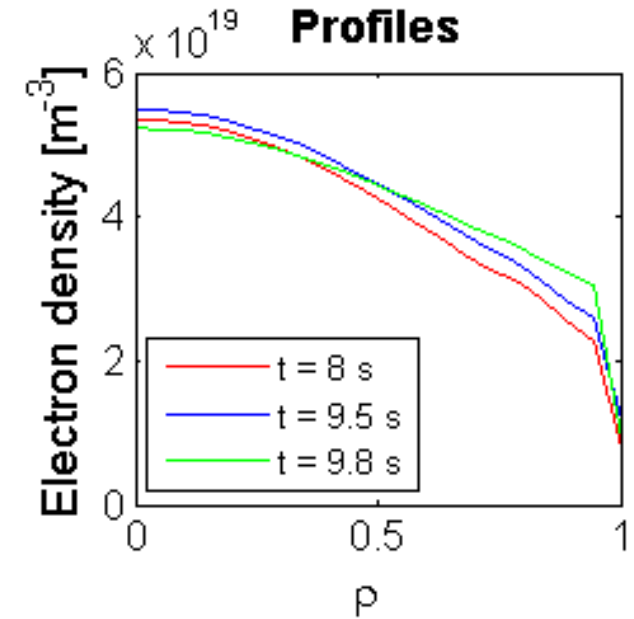
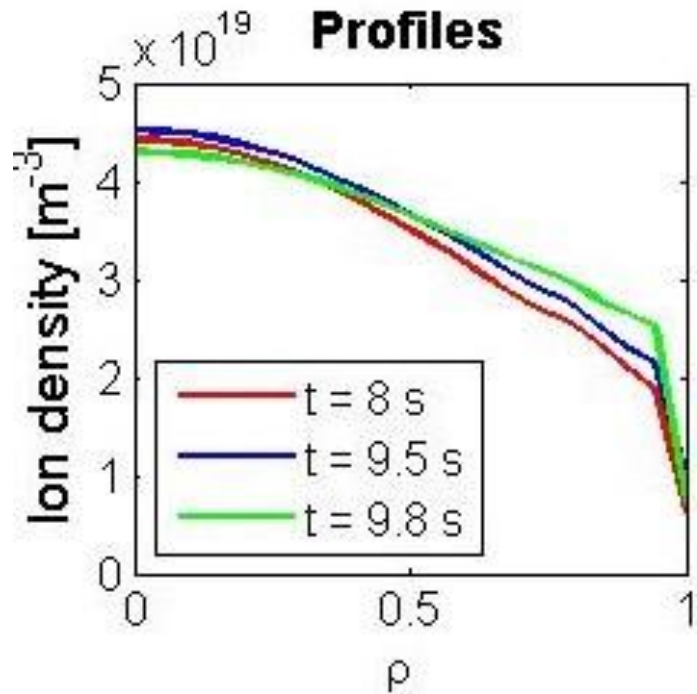
W neoclassical transport (in absence of poloidal asymmetries)

$$\Gamma_W = -Zq^2 D n_W \left[ \frac{1}{Z} \frac{\nabla n_W}{n_W} - \frac{\nabla n_i}{n_i} + \frac{1}{2} \frac{\nabla T_i}{T_i} \right]$$

Diffusion

Convection due to ion  
density peaking  
(inward)

Convection due to ion  
thermal screening,  
**outward**



The plateaus in H mode in C4:

- In USN 4 plateaus
- In LSN 10 plateaus

Max duration of H mode on WEST 4 s.

The transitions are observed in the hot and cold branches.

The transitions are unstable because  $P_{\text{rad}}$  increases when the pedestal is formed and so the power that crosses the separator is reduced and we go back to L mode.

In C6 we will raise the density to high power because we think we were on the low density branch of the L-H transition.