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INTEGRAL MEANS SPECTRUM OF WHOLE-PLANE SLE

DMITRY BELIAEV, BERTRAND DUPLANTIER, AND MICHEL ZINSMEISTER

ABSTRACT. We complete the mathematical analysis of the fine structure of harmonic measure on SLE curves that was initiated in Ref. [2], as described by the averaged integral means spectrum. For the unbounded version of whole-plane SLE as studied in Refs. [7, 16], a phase transition has been shown to occur for high enough moments from the bulk spectrum towards a novel spectrum related to the point at infinity. For the bounded version of whole-plane SLE of Ref. [2], a similar transition phenomenon, now associated with the SLE origin, is proved to exist for low enough moments, but we show that it is superseded by the earlier occurrence of the transition to the SLE tip spectrum.

1. INTRODUCTION

Harmonic measure is one of the fundamental objects in geometric function theory and its fine structure provides much information about the underlying geometry. We refer the reader to the survey by Makarov [17] or to the recent monograph [8]. In this article, we focus on the integral means spectrum of the harmonic measure (see definition below).

Schramm-Loewner Evolution (SLE) curves are (in part conjecturally) the conformally invariant scaling limits of interfaces in critical lattice models of statistical physics. Since its introduction by the late Oded Schramm [23] fifteen years ago, the SLE process has sparked intense interest both in mathematical and physical communities. (See Ref. [12] for a detailed study of SLE.) One direction of research is the study of geometrical properties of these curves, one approach to which is precisely to compute the multifractal spectra for the harmonic measure.

The first result in this direction is due to the second author who computed the SLE bulk spectrum by means of quantum gravity methods [3, 4], followed by Hastings who also computed a spectrum associated with the neighborhood of the SLE tip [10]. In Ref. [2], the first author and Smirnov provided a rigorous approach to the *average* integral means spectrum for whole-plane SLE, and showed that in a certain range of parameters (i.e., large enough negative moment orders), the tip spectrum dominates.

An *unbounded* version of whole-plane SLE was also studied in Refs. [6, 7] and [14, 15, 16] for which it was shown that for a large enough positive moment order, the (average) bulk integral mean spectrum undergoes a *phase transition* towards a novel form, which was argued to be related to the point at infinity. For the *bounded* version of whole-plane SLE as studied in Ref. [2], which is related by inversion to the unbounded one, a similar transition phenomenon may thus occur near the origin, which can be interpreted as the starting point of the random curve (see Definition 1.1 below). Indeed, it was observed in Refs. [16, 7] that the analysis provided in Ref. [2] is incomplete, and that the integral means spectrum could a priori be dominated by a novel spectrum, thought of as arising from the neighborhood of the starting point, and to be distinguished from that brought in by the vicinity of the tip.

The purpose of this work is to complete the analysis undertaken in Ref. [2], so as to rigorously establish, for the bounded version of whole-plane SLE as studied there, the

form of the averaged integral means spectrum for all moment orders. In particular, we show that for large enough negative moments, the new spectrum dominates the bulk one, but that both are still dominated by the tip spectrum in the same range of parameters. The results are summarized in Theorem 1.2. The existence of the new spectrum is established starting in Section 3.4. Following Ref. [7], the last Section 5.2 briefly discusses the relation of this spectrum to the derivative exponents of Ref. [13] for standard radial SLE, or, equivalently, to the non-standard tip exponents of Ref. [4]; it further heuristically suggests why the new spectrum should be associated with the ‘second tip’ of bounded whole-plane SLE, the image by inversion of the point at infinity in the unbounded version.

Before we proceed with the details of the analysis, let us also mention the work by Johansson Viklund and Lawler [11], who established the almost sure version of the SLE tip multifractal spectrum, that by Alberts, Binder and Viklund [1] on the almost sure dimension spectrum for SLE boundary collisions, as well as the recent preprint by Gwynne, Miller, and Sun [9], who used the so-called ‘‘Imaginary Geometry’’ of Miller and Sheffield [18, 19, 20, 21] to compute the a.s. value of the SLE bulk multifractal spectrum.

1.1. Definitions and Statements. Let $\Omega = \mathbb{C} \setminus K$ where K is a simply-connected compact set and let ϕ be a Riemann mapping from \mathbb{D}_- (i.e., the complement of the unit disk \mathbb{D}) onto Ω such that $\phi(\infty) = \infty$. The *integral means spectrum* of ϕ (or Ω) is defined as

$$\beta_\phi(t) = \beta_\Omega(t) = \limsup_{r \rightarrow 1^+} \frac{\log \int_0^{2\pi} |\phi'(re^{i\theta})|^t d\theta}{-\log(r-1)}.$$

For random fractals, it is natural to study *the average integral means spectrum*, which is defined as

$$\bar{\beta}(t) = \limsup_{r \rightarrow 1^+} \frac{\log \int_0^{2\pi} \mathbb{E} [|\phi'(re^{i\theta})|^t] d\theta}{-\log|r-1|}.$$

In this work, we are interested in the average integral means spectrum of whole-plane SLE curves.

Definition 1.1. Let $\xi(t) = \exp(i\sqrt{\kappa}B_t)$ be a two-sided Brownian motion on the unit circle with $t \in \mathbb{R}$ and $\kappa > 0$. The whole-plane SLE_κ is the family of conformal maps g_t satisfying

$$\partial_t g_t(z) = g_t(z) \frac{\xi(t) + g_t(z)}{\xi(t) - g_t(z)}, \quad (1.1)$$

with initial condition

$$\lim_{t \rightarrow -\infty} e^t g_t(z) = z, \quad z \in \mathbb{C} \setminus \{0\}.$$

This map g_t is a conformal map from $\mathbb{C} \setminus K_t$ onto \mathbb{D}_- , where the compact set K_t is the so-called *hull* of the SLE process, and it describes a family of hulls that grow from the origin towards infinity. This is the so-called *exterior* version of whole-plane SLE, as the map g_0 may be seen as the limit of a rescaled version of a radial SLE process growing from the unit disc towards infinity [2]. The integral means spectrum of whole-plane SLE is then defined as that of the *inverse map* $\phi = g_0^{-1}$. Another version describes hulls growing from infinity towards the origin, and is called *interior* whole-plane SLE. The integral means spectrum of this *unbounded* process is studied in Ref. [7].

Remark. The interior and exterior versions of whole-plane SLE are conjugate under the map $z \mapsto 1/z$, and their integral means spectra can be unified in a single formalism by considering *mixed moments*, involving powers of the moduli of the conformal map and of its derivative, as studied in Ref. [5].

The main result of this paper is (Fig. 1)

Theorem 1.2. *i) The average integral means spectrum of the exterior whole-plane SLE $_{\kappa}$ is given by*

$$\begin{aligned}\beta_{\text{tip}}(t) &= -t - 1 + \frac{1}{4}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa t}), & t \leq t_2, \\ \beta_0(t) &= -t + \frac{4 + \kappa}{4\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa t}), & t_2 \leq t \leq t_3, \\ \beta_{\text{lin}}(t) &= t - \frac{(4 + \kappa)^2}{16\kappa}, & t_3 \leq t.\end{aligned}$$

ii) If in the definition of the spectrum we integrate over any set that excludes the neighbourhood of $\theta = 0$, which corresponds to excluding the influence of the tip of the curve, then the average spectrum is given by

$$\begin{aligned}\beta_1(t) &= -t - \frac{1}{2}(1 + \sqrt{1 - 2\kappa t}), & t \leq t_1, \\ \beta_0(t) &= -t + \frac{4 + \kappa}{4\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa t}), & t_1 \leq t \leq t_3, \\ \beta_{\text{lin}}(t) &= t - \frac{(4 + \kappa)^2}{16\kappa}, & t_3 \leq t.\end{aligned}$$

In the above, the transition values for the moment order t are given by

$$t_1 := -\frac{1}{128}(4 + \kappa)^2(8 + \kappa), \quad t_2 := -1 - \frac{3\kappa}{8}, \quad t_3 := \frac{3(4 + \kappa)^2}{32\kappa},$$

and such that $t_1 < t_2 < t_3$.

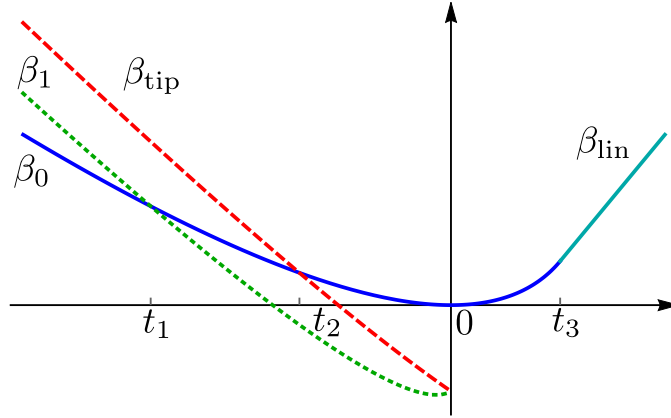


FIGURE 1. Plot showing the relative positions of the various forms taken by the integral means spectrum in Theorem 1.2: tip spectrum β_{tip} , novel spectrum β_1 , bulk spectrum β_0 and linear spectrum β_{lin} . The β_1 spectrum is the analogue for the exterior whole-plane SLE of the spectrum at infinity of Refs. [7, 16] for the interior case (see also Ref. [5]). Note that β_1 supersedes the bulk spectrum below t_1 , but stays below the tip spectrum.

Remark. Point *i)* in Theorem 1.2 agrees with Ref. [2], while Statement *ii)* pertains to this work.

2. DIFFERENTIAL OPERATOR

As seen above, the integral means spectrum is defined for a map from \mathbb{D}_- onto some domain, i.e., in terms of the inverse of the SLE map. The Loewner equation for g_t^{-1} is a PDE instead of an ODE and is much harder to work with. Instead, we study the *backward evolution*. It is well-known that for a standard radial [exterior] SLE process \hat{g}_t (i.e., obeying Eq. (1.1) for $t \geq 0$ with $\hat{g}_0(z) = z$), reversing time, i.e., changing sign in front of Eq. (1.1), leads to solutions $f_t := \hat{g}_{-t}$, $t \geq 0$, also called [backward] radial SLE, that have, up to conjugation by rotation, the same distribution as \hat{g}_t^{-1} (see Ref. [2], Lemma 1, which is an analog of Lemma 3.1 in Ref. [22]).

To compute the average integral means spectrum, let us then introduce the function $\tilde{F}(z; \tau) := \mathbb{E}[|f'_\tau(z)|^t]$, where $f_\tau(z)$, $\tau \geq 0$, is such a backward radial SLE $_\kappa$ process. (Actually, this function depends also on the moment order t and on κ , but these are fixed throughout the proof and we will no longer mention the dependence thereof.) The function \tilde{F} is $C^{\infty,1}$ in z, τ , since $f'_\tau(z)$ is by Loewner theory, and this remains true under expectation by standard dominated convergence and distortion theorems. As was shown in Ref. [2], it satisfies a parabolic PDE:

Lemma 2.1. *The function $\tilde{F}(z; \tau)$ is a solution to the PDE in variables r, θ, τ ,*

$$\begin{aligned} t \frac{r^4 + 4r^2(1 - r \cos \theta) - 1}{(r^2 - 2r \cos \theta + 1)^2} \tilde{F} + \frac{r(r^2 - 1)}{r^2 - 2r \cos \theta + 1} \tilde{F}_r \\ - \frac{2r \sin \theta}{r^2 - 2r \cos \theta + 1} \tilde{F}_\theta + \frac{\kappa}{2} \tilde{F}_{\theta, \theta} - \tilde{F}_\tau = 0, \end{aligned} \quad (2.1)$$

where $z = re^{i\theta}$.

Lemma 3 in Ref. [2] further shows that there exists a limit to $e^{-\tau} f_\tau(z)$ as $\tau \rightarrow +\infty$, which has the same distribution as $g_0^{-1}(z)$, where g_t is a whole-plane SLE process as in Definition 1.1; hence we introduce

$$F(z) := \lim_{\tau \rightarrow +\infty} e^{-\tau t} \tilde{F}(z; \tau),$$

where limit and expectation commute by the same arguments as above. The average integral means spectrum of the exterior whole-plane SLE is thus associated with the singular behavior of $\int_0^{2\pi} F(re^{i\theta}) d\theta$ as $r \rightarrow 1^+$.

Multiplying Eq. (2.1) by $e^{-\tau t}$, and passing to the limit, shows that F is a solution to the PDE in r, θ ,

$$\begin{aligned} t \left(\frac{r^4 + 4r^2(1 - r \cos \theta) - 1}{(r^2 - 2r \cos \theta + 1)^2} - 1 \right) F + \frac{r(r^2 - 1)}{r^2 - 2r \cos \theta + 1} F_r \\ - \frac{2r \sin \theta}{r^2 - 2r \cos \theta + 1} F_\theta + \frac{\kappa}{2} F_{\theta, \theta} = 0. \end{aligned} \quad (2.2)$$

The exchange above, of the $\tau \rightarrow +\infty$ limit and of partial derivatation of $\tilde{F}(z; \tau)$ with respect to r and θ , is justified by the fact that the τ -family, $e^{-\tau} f'_\tau(z)$, and all its z -derivatives are normal, i.e., uniformly bounded in any compact of \mathbb{D}_- , so that the spatial derivatives of $e^{-\tau t} \tilde{F}(z; \tau)$ form an equicontinuous family. A further requirement is that $\lim_{\tau \rightarrow +\infty} \frac{\partial}{\partial \tau} (e^{-\tau t} \tilde{F}) = 0$. Use of the Schramm-Loewner equation (1.1) for f_τ shows that

$$\frac{\partial}{\partial \tau} |e^{-\tau} f'_\tau(z)|^t = 2t |e^{-\tau} f'_\tau(z)|^t \operatorname{Re} \frac{\xi(\tau)^2}{(f_\tau(z) - \xi(\tau))^2}.$$

Classical Koebe distortion theorems then show that the right-hand side is bounded by $C(z)e^{-2\tau}$, with C defined on \mathbb{D}_- ; this insures both the validity of the exchange of expectation and τ -derivation, and the vanishing limit above.

It is easy to see that after a change of variables, this equation is still of *parabolic* type, where θ plays the rôle of a spatial variable, and $r \rightarrow 1^+$ corresponds to time going to infinity [2].

Instead of polar coordinates, it turns out to be more convenient to work with (z, \bar{z}) coordinates, where $F = F(z, \bar{z})$ now formally depends on both complex variables, and where Eq. (2.2) becomes,

$$\begin{aligned} \Lambda F(z, \bar{z}) &= 0, \\ \Lambda &:= -\frac{\kappa}{2} (z\partial - \bar{z}\bar{\partial})^2 + \frac{z+1}{z-1} z\partial + \frac{\bar{z}+1}{\bar{z}-1} \bar{z}\bar{\partial} - t \left(\frac{1}{(z-1)^2} + \frac{1}{(\bar{z}-1)^2} \right). \end{aligned} \quad (2.3)$$

(See Ref. [7] for details.)

We wish to study how the solutions to Eq. (2.3) behave when z approaches the unit circle. In contradistinction to the interior case [7], it seems difficult to construct explicit solutions, so we are left with constructing sub- and super-solutions with same boundary behaviors [2].

Following Refs. [2, 7], we consider the action of Λ on functions of the peculiar form,

$$\psi(z, \bar{z}) := (z\bar{z} - 1)^{-\beta} g(u) = (|z|^2 - 1)^{-\beta} g(u), \quad (2.4)$$

where g is a C^2 function of $u := (1-z)(1-\bar{z}) = |1-z|^2$. By looking at the leading terms in Eq. (2.3) for ψ as $r \rightarrow 1^+$, one obtains a ‘boundary equation’ for g (see Refs [2, 7] for details),

$$(t(2-u) - 2\beta) g(u) + \left(\frac{\kappa}{2}(2-u) - (4-u) \right) u g'(u) + \frac{\kappa}{2}(4-u) u^2 g''(u) = 0; \quad (2.5)$$

from now on, we assume that g in (2.4) satisfies this equation.

Let us then consider the action of the differential operator (2.3) on ψ (2.4) in \mathbb{D}_- and follow Ref. [7, Section 4.2.]. Making use of Eq. (2.5) to eliminate the second derivative of g , one obtains after some computation,

$$\begin{aligned} \frac{\Lambda\psi}{\psi} &= (z\bar{z} - 1) \left[\frac{1}{u}(t - \beta) + \frac{1}{4-u}(\beta + t) - \left(\frac{\kappa}{2} - 1 - \frac{2\kappa}{4-u} \right) \frac{g'(u)}{g(u)} \right] \\ &+ \frac{(z\bar{z} - 1)^2}{u^2} \left[\frac{1}{4-u} \left(-2t - 2\beta - \kappa u \frac{g'(u)}{g(u)} \right) + \left(\frac{\kappa}{2} + 1 \right) \left(u \frac{g'(u)}{g(u)} \right) \right]. \end{aligned} \quad (2.6)$$

We are also interested in the action of Λ on ψ functions with logarithmic corrections [2]. Let us introduce

$$\ell_\delta = \ell_\delta(z, \bar{z}) := (-\log(z\bar{z} - 1))^\delta.$$

It is then easy to see that,

$$\frac{\Lambda(\psi\ell_\delta)}{\psi\ell_\delta} = \frac{\Lambda(\psi)}{\psi} - \frac{2\delta z\bar{z}}{u(-\log(z\bar{z} - 1))}. \quad (2.7)$$

For certain choices of exponent β and of g , some leading terms in (2.6) may cancel out so that the second term on the r.h.s. of Eq. (2.7) dominates. The latter has a sign opposite to the arbitrary sign of δ , which means that if ψ is positive, we shall be able construct sub- and super-solutions $\psi\ell_\delta$ with growth rate arbitrary close to that of ψ .

3. BOUNDARY SOLUTIONS AND THEIR EXTENSION TO \mathbb{D}_-

3.1. Hypergeometric equation. We shall be interested in functions g of the form $g(u) = u^\gamma g_0(u)$, $\gamma \in \mathbb{R}$, and such that Eq. (2.5) may reduce to a hypergeometric-type equation on g_0 [2]. Upon substituting this into (2.5) and factoring out u^γ , we obtain,

$$(2\beta(\gamma) - 2\beta + A(\gamma)u) g_0(u) + \left(\frac{\kappa}{2}(2-u) + (\kappa\gamma - 1)(4-u)\right) u g_0'(u) + \frac{\kappa}{2}(4-u) u^2 g_0''(u) = 0, \quad (3.1)$$

in terms of the quadratic polynomials,

$$\beta(\gamma) := \kappa\gamma^2 - \left(\frac{\kappa}{2} + 2\right) \gamma + t, \quad (3.2)$$

$$A(\gamma) := -\frac{\kappa}{2}\gamma^2 + \gamma - t. \quad (3.3)$$

For the choice of parameter $\beta = \beta(\gamma)$, Eq. (3.1) then reduces to an hypergeometric equation,

$$A(\gamma)g_0(u) + \left[\frac{\kappa}{2}(2-u) + (\kappa\gamma - 1)(4-u)\right] g_0'(u) + \frac{\kappa}{2}(4-u)u g_0''(u) = 0, \quad (3.4)$$

which is the same as Eqs. [7, (184)] or [2, (17)]. Following either of these papers, we see that the general solution is,

$$g_0(u) = C_0 {}_2F_1(a, b, c, u/4) - C'_0 (u/4)^{1/2-a-b} {}_2F_1(a', b', c', u/4), \quad (3.5)$$

where

$$\begin{aligned} a = a(\gamma) &= \gamma - \gamma_+, & b = b(\gamma) &= \gamma - \gamma_-, & c &= \frac{1}{2} + a + b, \\ a' &= \frac{1}{2} - a, & b' &= \frac{1}{2} - b, & c' &= \frac{1}{2} + a' + b', \end{aligned} \quad (3.6)$$

and where $\gamma_\pm := (1 \pm \sqrt{1 - 2\kappa t})/\kappa$ are the two roots in $A(\gamma) = -\frac{\kappa}{2}(\gamma - \gamma_+)(\gamma - \gamma_-)$.

Hypergeometric functions are singular at $u = 4$, but the solution should be smooth at $z = -1$, which means that the coefficients C_0 and C'_0 should be chosen in such a way that the singular parts cancel out [2]. This precisely happens, up to a constant factor, for

$$g_0(0) = C_0 = \frac{\Gamma(3/2 - a - b)}{\Gamma(1/2 - a)\Gamma(1/2 - b)}, \quad C'_0 = \frac{\Gamma(c)}{\Gamma(a)\Gamma(b)}, \quad (3.7)$$

such that near $u = 4$ (see Ref. [2]),

$$g_0(u) = \frac{1}{\sqrt{\pi}} \left(\frac{1}{2} - a - b\right) + O(4 - u). \quad (3.8)$$

3.2. Action of Λ on trial functions ψ . Following Ref. [7, Section 4.2.3], let us now return to the action (2.6), this time for $\psi(z, \bar{z}) = (z\bar{z} - 1)^{-\beta} u^\gamma g_0(u)$, and for the choice $\beta = \beta(\gamma)$ (3.2). Using $u g'/g = u g'_0/g_0 + \gamma$, we finally obtain,

$$\begin{aligned} \frac{\Lambda\psi}{\psi} &= (z\bar{z} - 1) \left[\frac{1}{u} [t + \gamma - \beta(\gamma)] - \frac{2A(\gamma)}{4-u} - \left(\frac{\kappa}{2} - 1 - \frac{2\kappa}{4-u}\right) \frac{g'_0(u)}{g_0(u)} \right] \\ &+ \frac{(z\bar{z} - 1)^2}{u^2} \left[\frac{1}{4-u} \left(4A(\gamma) - \kappa u \frac{g'_0(u)}{g_0(u)}\right) + \left(1 + \frac{\kappa}{2}\right) \left(\gamma + u \frac{g'_0(u)}{g_0(u)}\right) \right]. \end{aligned} \quad (3.9)$$

In Eq. (3.9), there seems to be an unexpected singularity at $u = 4$ ($z = -1$), but in fact the choice of constants C_0 and C'_0 in Eq. (3.7), made to ensure that $g_0(u)$ is regular at $u = 4$, further yields

$$\frac{g'_0(u)}{g_0(u)} = -\frac{1}{2}ab + O(4-u) = \frac{1}{\kappa}A(\gamma) + O(4-u),$$

so that the singularity at $u = 4$ cancels out in the action (3.9).

In the $u \rightarrow 0$ limit, Eq. (3.5) shows that (up to a non-vanishing coefficient)

$$u \frac{g'_0(u)}{g_0(u)} \sim u^\alpha, \quad \alpha := \min \left\{ \frac{1}{2} - a - b, 1 \right\}, \quad (3.10)$$

Remark. In the case where C_0 vanishes, $\alpha = 1/2 - a - b$.

The second line in Eq. (3.9) is in this limit,

$$\frac{(z\bar{z} - 1)^2}{u^2} [C(\gamma) + O(u^\alpha)], \quad (3.11)$$

with

$$C(\gamma) := A(\gamma) + \left(1 + \frac{\kappa}{2}\right) \gamma = -\frac{\kappa}{2}\gamma^2 + \left(2 + \frac{\kappa}{2}\right) \gamma - t. \quad (3.12)$$

3.3. Beliaev-Smirnov solution. The first pair of exponents β and γ that we are interested in have been introduced in Ref. [2], so as to cancel the leading singularity in (3.11). For a given t , there are two solutions to $C(\gamma_0) = 0$, and we consider the particular values,

$$\begin{cases} \gamma_0 = \frac{1}{2\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa t}), \\ \beta_0 = \beta(\gamma_0) = \frac{\kappa}{2}\gamma_0^2 = -t + (2 + \frac{\kappa}{2})\gamma_0 = -t + \frac{4+\kappa}{4\kappa}(4 + \kappa - \sqrt{(4 + \kappa)^2 - 8\kappa t}). \end{cases} \quad (3.13)$$

The first condition is for the boundary solution g_0 (3.5) to be bounded (i.e., u^{γ_0} should be the only singular term in g), which means that $1/2 - a - b \geq 0$, where now $a = a(\gamma_0)$, $b = b(\gamma_0)$ as in Eqs. (3.6). Simple algebra shows that this is equivalent to $t \leq t_3$.

The second condition is that $g_0(u)$ should be positive for the whole range $u \in [0, 4]$ when z describes the unit circle. It was observed in Ref. [2] that this happens when $1/2 - b > 0$. It was erroneously stated there that this is always true, while in fact this holds only for $t_1 < t$. (See Proposition 5.1 below.)

For $t \in (t_1, t_3)$ we thus have that g_0 is bounded and positive, and the arguments in Ref. [2], which we detail and refine here, stay valid. To study the spectrum, one has to analyze the behavior of $\Lambda\psi_0$, where $\psi_0(z, \bar{z}) = (z\bar{z} - 1)^{-\beta_0} u^{\gamma_0} g_0(u)$ for values β_0, γ_0 as in Eq. (3.13), and the associated hypergeometric function g_0 (3.5). This is provided by Eq. (3.9), where (3.12) and (3.13) give the explicit coefficients, $t + \gamma_0 - \beta(\gamma_0) = 2t - (1 + \frac{\kappa}{2})\gamma_0$, and $A(\gamma_0) = -(1 + \frac{\kappa}{2})\gamma_0$.

As a preparation also for a complete analysis below, let us detail various radial limits as $|z| \rightarrow 1$ in action (3.9), while recalling the geometrical constraint, $|z| - 1 \leq |z - 1| = u^{1/2}$. *Generic case:* $z\bar{z} - 1 \rightarrow 0$, but u is bounded away from 0. In this case, $\Lambda\psi_0/\psi_0 = O(z\bar{z} - 1)$.

Special case: z approaches 1, $u \rightarrow 0$, so that $z\bar{z} - 1 = O(u^{1/2})$. In this case $g'_0/g_0 = O(u^\alpha/u)$ in (3.10) is dominated by $1/u$, as we have shown before that $1/2 - a - b > 0$ for $t < t_3$, so that $\alpha > 0$. Thus the first line in (3.9) is of order $O(u^{-1/2})$. The second line, as given by Eq. (3.11), seems to be of order $O(u^{-1})$, but because of the very choice

of $\gamma = \gamma_0$ such that $C(\gamma_0) = 0$, it is of order $O(u^{\alpha-1})$, with $\alpha > 0$ for $t < t_3$. The results of this discussion can simply be recast as,

$$\frac{\Lambda\psi_0}{\psi_0} = \frac{z\bar{z} - 1}{u}O(u^0) + \frac{(z\bar{z} - 1)^2}{u^2}O(u^\alpha), \quad |z| - 1 \leq u^{1/2}. \quad (3.14)$$

To complete the proof of Theorem 1.2 in the range $t \in (t_1, t_3)$, one then considers as in Ref. [2], the set of logarithmic modifications of ψ_0 for all δ ,

$$\ell_\delta\psi_0(z, \bar{z}) = (-\log(z\bar{z} - 1))^\delta (z\bar{z} - 1)^{-\beta_0} u^{\gamma_0} g_0(u).$$

Recalling the action (2.7), one sees that for $|z|$ close enough to 1 and for all u , both terms in $\Lambda\psi_0/\psi_0$ (3.14) are dominated by the logarithmic second term on the r.h.s. of (2.7):

$$\left| \frac{\Lambda\psi_0}{\psi_0} \right| \leq \frac{2|\delta|z\bar{z}}{u(-\log(z\bar{z} - 1))}. \quad (3.15)$$

The positive function $\psi_0\ell_\delta$ is thus a sub-solution for $\delta > 0$ or a super-solution for $\delta < 0$, so that the integral means spectrum is $\beta_0(t)$ for $t \in [t_2, t_3)$, or $\beta_{\text{tip}}(t) = \beta_0(t) - 2\gamma_0(t) - 1$, for $t \in (t_1, t_2)$ where $2\gamma_0(t) + 1 < 0$.

By Hölder's inequality, the spectrum is convex, and by standard distortion theorems, bounded by t for $t > 0$. This, together with the fact that $\partial_t\beta_0(t_3) = 1$, establishes Theorem 1.2 for $t \geq t_3$.

3.4. Below t_1 . At t_1 , $b = 1/2$, and $g_0(0) = C_0 = 0$ in Eq. (3.7). Actually, this vanishing also happens for higher half-integer values $b = n + 1/2$. Proposition 5.2 below shows that there indeed exist a finite set of integers \mathcal{J}_κ , and a finite discrete set \mathcal{T}_κ of moment orders,

$$\mathcal{J}_\kappa := \{n \in \mathbb{N}, 0 \leq n \leq \lfloor \kappa^{-1} \rfloor\} \quad (3.16)$$

$$\mathcal{T}_\kappa := \{t_{1-n}, n \in \mathcal{J}_\kappa\} \quad (3.17)$$

$$t_{1-n} = t_{1-n}(\kappa) := -\frac{(1+2n)(8+\kappa-2n\kappa)(4+\kappa+2n\kappa)(4+\kappa-2n\kappa)}{128(1-n\kappa)^2}, \quad (3.18)$$

such that,

$$b = b(\gamma_0(t_{1-n})) = n + \frac{1}{2}, \quad n \in \mathcal{J}_\kappa, \quad g_0(0) = C_0 = 0, \quad t \in \mathcal{T}_\kappa. \quad (3.19)$$

The $n = 0$ case corresponds, for any value of κ , to the point t_1 as above; note also that for $\kappa > 1$, $\mathcal{T}_\kappa = \{t_1\}$, whereas strictly positive values of n exist in \mathcal{J}_κ only for $0 < \kappa \leq 1$.

In Eq. (3.7), observe now that $1/2 - a > 0$, because $a = a(\gamma_0) < 0$ in Eq. (3.6), and recall that $1/2 - a - b > 0$ for $t < t_3$. Therefore, the sign of $g_0(0)$ (3.7) is given by that of $\Gamma(1/2 - b)$, and by the very property of analytical continuation of the Γ -function, is thus alternating in the successive intervals $t \in (t_{-n}, t_{1-n})$, being positive or negative for n odd or even, respectively. For later convenience, let us then introduce,

$$\sigma = \sigma(t) := \text{sgn } g_0(0) = (-1)^{n-1}, \quad t \in (t_{-n}, t_{1-n}), \quad n \in \mathcal{J}_\kappa. \quad (3.20)$$

Notice that owing to Eq. (3.8), $g_0(4)$ is always *positive* for $t < t_3$. Then, in the interval of moment orders, $t \in (t_{-n}, t_{1-n})$, with $n \in \mathcal{J}_\kappa$, the graph of $g_0(u)$ possesses exactly $n + 1$ simple zeroes over the interval $u \in (0, 4)$.

3.5. Power-law solution. Note that the hypergeometric equation (3.4) becomes degenerate when $A(\gamma) = 0$, with g_0 a constant solution and g of the form u^γ . As before, there are two solutions, γ_\pm , to this system, which correspond to the degenerate cases $a = 0$ or $b = 0$ in Eqs. (3.5) and (3.6), and to $C_0 = 0$ in Eq. (3.7). As in the case of interior whole-plane SLE [7], we are especially interested in the pair,

$$\begin{cases} \gamma_1 := \gamma_+ = \frac{1}{\kappa}(1 + \sqrt{1 - 2\kappa t}), \\ \beta_1 = \beta(\gamma_1) = -t - \frac{\kappa}{2}\gamma_1 = -t - \frac{1}{2}(1 + \sqrt{1 - 2\kappa t}). \end{cases} \quad (3.21)$$

Remark. Both β_0 (3.13) and β_1 (3.21) are given by the same quadratic function $\beta(\gamma)$ (3.2), in terms of γ_0 and γ_1 , respectively.

When plugging $\psi_1 = \psi_1(z, \bar{z}) := (z\bar{z} - 1)^{-\beta_1} u^{\gamma_1}$ into Eq. (3.9), because $A(\gamma_1) = 0$ and $g'_0 = 0$, many terms disappear, and the result is simply,

$$\frac{\Lambda\psi_1}{\psi_1} = \frac{z\bar{z} - 1}{u} \left(2t + \left(1 + \frac{\kappa}{2} \right) \gamma_1 \right) + \frac{(z\bar{z} - 1)^2}{u^2} \left(1 + \frac{\kappa}{2} \right) \gamma_1. \quad (3.22)$$

As before, we distinguish two cases: $|z| \rightarrow 1$, but u is bounded away from 0; $u \rightarrow 0$, and $|z| - 1 = O(u^{1/2})$. In the first case, $\Lambda\psi_1/\psi_1$ is $O(|z| - 1)$; in the second case, it is $O(u^{-1})$.

4. MIXING THE TWO SOLUTIONS

For $t < t_1$, $g_0(u)$ changes sign at least once over the interval $(0, 4)$, invalidating the proof of Section 3.3. Recall that at the origin its sign alternates, as described in Eq. (3.20). The idea is thus to try and combine the two functions ψ_0 and ψ_1 into

$$\psi := \sigma\psi_0 + \psi_1 = \sigma g_0(u) u^{\gamma_0} (z\bar{z} - 1)^{-\beta_0} + u^{\gamma_1} (z\bar{z} - 1)^{-\beta_1}, \quad (4.1)$$

where σ is as in Eq. (3.20), so as to restore overall positivity for ψ . Then, in the action $\Lambda(\psi\ell_\delta)$, the differential operator Λ will act differently on ψ_0 and ψ_1 , still maintaining the possibility to build sub- and super-solutions in this way. In this section, we are mostly interested in $t < t_1$, but some arguments are independent of that assumption, provided one assumes that, e.g., $t < 0$, so that both ψ_0 and ψ_1 are defined.

Lemma 4.1. *There is $r_0 > 1$ such that $\psi = \sigma\psi_0 + \psi_1 > 0$ for all z such that $1 < |z| < r_0$.*

Proof. First of all, note that both ψ_0 and ψ_1 are continuous in the complement of the unit disc, and that $\psi_1 > 0$ everywhere. For $t > t_1$ we have $g_0 > 0$, hence both terms in ψ are positive. For $t = t_1$, $g_0(u) > 0$ for $u > 0$, so that $\psi > 0$. For $t < t_1$, this is no longer true, but $\sigma g_0(u) \geq 0$ in some neighborhood $u \in [0, u_0]$ of $z = 1$, hence $\psi > 0$ there. Outside of this neighborhood, u is uniformly bounded away from 0, and also bounded above by $(|z| + 1)^2$; therefore, for any $r_0 > 1$, there exist positive constants c_0, c_1 such that for all z such that $1 < |z| < r_0$ and $u = |z - 1|^2 > u_0$,

$$\begin{aligned} 0 \leq |\psi_0| &\leq c_0 (z\bar{z} - 1)^{-\beta_0}, \\ \frac{1}{c_1} (z\bar{z} - 1)^{-\beta_1} &< \psi_1 < c_1 (z\bar{z} - 1)^{-\beta_1}. \end{aligned}$$

For $t < t_1$, we have that $\beta_1 > \beta_0$ (See Prop. 5.3 below and Fig. 1), so that for $|z|$ sufficiently close to 1, we have $\psi_1 > |\psi_0|$, hence $\psi > 0$. \square

Lemma 4.2. *For $t < t_1$ and $t \notin \mathcal{T}_\kappa$, there is $r_0 > 1$ such that $\Lambda(\psi\ell_\delta)$ for ψ (4.1) has a constant sign in the annulus $1 < |z| < r_0$, which depends only on that of δ .*

Proof. As shown by Eq. (2.7), multiplying ψ by a logarithmic factor ℓ_δ results in an additional term in the action of the differential operator, whose sign depends on that of δ only. We shall show that near the boundary of the unit disc, this additional term is the main one, hence $\Lambda(\psi\ell_\delta)$ has constant sign there. Since ψ is positive, this implies that $\psi\ell_\delta$ is a sub- or super-solution for $\delta > 0$ or $\delta < 0$, respectively.

As was shown in (2.7), (3.14), (3.15) and (3.22), $\Lambda(\psi\ell_\delta)/\ell_\delta$ can be written (up to smaller order terms) as the sum of four terms

$$\frac{\Lambda(\psi\ell_\delta)}{\ell_\delta} = \sigma\psi_0 \frac{-2\delta r^2}{u(-\log(r-1))} + \psi_1 \frac{r-1}{u} + \psi_1 \frac{(r-1)^2}{u^2} + \psi_1 \frac{-2\delta r^2}{u(-\log(r-1))}, \quad (4.2)$$

where $r = |z|$. We denote these terms by I – IV. Here, we omitted the positive constants in front of II, III, but terms I and IV are written in their *complete* form. For functions A and B of r , we shall use the short-hand notations, $A \gtrsim B$ for $A \geq cB$ with c some positive constant, and $A \approx B$ when both $A \gtrsim B$ and $A \lesssim B$ hold.

Below, we consider three cases, and show that in each case one of the logarithmic terms (either I or IV) is the leading one when $r \rightarrow 1^+$.

Case 1: u is bounded away from zero. It is obvious that $IV \gtrsim II, III$ as $(r-1) \rightarrow 0$. For u bounded away from zero, we have $\psi_0 \approx (r-1)^{-\beta_0}$ and $\psi_1 \approx (r-1)^{-\beta_1}$. Since $\beta_1 > \beta_0$ for $t < t_1$ (Proposition 5.3), we have $IV \gtrsim I$. Also notice that ψ_1 is positive, hence for sufficiently small $r-1$, the sign of $\Lambda(\psi\ell_\delta)$ is opposite to that of δ .

Case 2: We assume that $(r-1)^{2-\varepsilon} < u < u_0$, where $\varepsilon > 0$ and where, as above, $u_0 > 0$ is chosen such that $\sigma g_0(u) > 0$ for $0 < u < u_0$. Then,

$$\begin{aligned} II &= \psi_1 \frac{r-1}{u}, \\ III &\lesssim \psi_1 \frac{(r-1)^2}{u} \frac{1}{(r-1)^{2-\varepsilon}} = \psi_1 \frac{(r-1)^\varepsilon}{u}, \\ IV &\approx \psi_1 \frac{-\delta}{u(-\log(r-1))}, \end{aligned}$$

and obviously IV dominates II and III. Since $\sigma\psi_0$ is positive for our choice of u_0 , we have that both main terms I and IV in (4.2) have a sign opposite to that of δ .

Case 3: We assume that $r-1 > u^{1/2+\varepsilon}$ for some $\varepsilon > 0$ that will be determined later. We also recall that $r-1 \leq u^{1/2}$. In this case $-\log(r-1) \approx -\log u$. First, we notice that

$$\begin{aligned} I &\approx \psi_0 \frac{-\delta}{u(-\log u)}, \\ II &\lesssim \psi_1 \frac{u^{1/2}}{u}, \\ III &\lesssim \psi_1 \frac{1}{u}, \\ IV &\approx \psi_1 \frac{-\delta}{u(-\log u)}, \end{aligned}$$

so that both III and IV are dominated by $III' := \psi_1/u$. We would like to show that $III' \lesssim I$, and this requests comparing ψ_1 to ψ_0 . Proposition 5.4 below precisely gives that for $t < t_1$, one has $\beta_{\text{tip}} = \beta_0 - 2\gamma_0 - 1 > \beta_1$, hence

$$(r-1)^{-\beta_1} < (r-1)^{-\beta_0} (r-1)^{2\gamma_0+1},$$

and

$$\psi_1 \approx (r-1)^{-\beta_1} u^{\gamma_1} < (r-1)^{-\beta_0} u^{\gamma_0} u^{\gamma_1-\gamma_0} (r-1)^{2\gamma_0+1} \approx \sigma \psi_0 u^{\gamma_1-\gamma_0} (r-1)^{2\gamma_0+1}.$$

Note that the last estimate requires that $g_0(0) \neq 0$, hence the condition $t \notin \mathcal{T}_\kappa$ in Lemma 4.2. We expect $u^{\gamma_1-\gamma_0} (r-1)^{2\gamma_0+1}$ to be bounded by some positive power of u . As was shown in Ref. [2], the definition of the threshold t_2 for the tip relevance is that $2\gamma_0+1 < 0$ for $t < t_2$, and we have here $t < t_1 < t_2$. Hence, for $r-1 > u^{1/2+\varepsilon}$,

$$u^{\gamma_1-\gamma_0} (r-1)^{2\gamma_0+1} < u^{\gamma_1-\gamma_0+(1/2+\varepsilon)(2\gamma_0+1)} = u^{\gamma_1+1/2+\varepsilon(2\gamma_0+1)}.$$

Since $\gamma_1 > 0$, for sufficiently small ε , the latter is bounded by some positive power of u , e.g., by $u^{1/2}$. Hence for this ε , we have that

$$\psi_1 \lesssim \sigma \psi_0 u^{1/2}. \quad (4.3)$$

Thus we see that $\text{III}' \lesssim \text{I}$, and that the sign of $\Lambda(\psi l_\delta)$ is given by that of I in Eq. (4.2), which is opposite to the sign of δ .

Altogether, these three cases show that in some neighborhood of the unit circle, the sign of $\Lambda(\psi l_\delta)$ is constant and opposite to that of δ . \square

Lemmas 4.1 and 4.2 together show that ψl_δ is a sub- or super-solution, depending on the sign of δ . Following Ref. [2], one obtains that for F such that $\Lambda F = 0$, there exist positive constants c_2 and c_3 , such that in some annulus adherent to \mathbb{D} , $c_2 \psi l_{-\delta} < F < c_3 \psi l_\delta$, with $\delta > 0$. We conclude that F behaves like ψ (4.1), up to arbitrary small logarithmic correction.¹ This completes the proof of Theorem 1.2 for $t < t_1$ and $t \notin \mathcal{T}_\kappa$.

Lastly, when t belongs to the discrete set \mathcal{T}_κ (3.17), because of (3.19), $g_0(u)$ (3.5) vanishes too fast at $u = 0$, and neither an upper bound like (4.3) holds, nor Lemma 4.2. But having established Theorem 1.2 for $t \notin \mathcal{T}_\kappa$ suffices to prove it for all t , by simply invoking the *convexity* property of the integral means spectrum [17].

5. PHASE TRANSITIONS

5.1. Loci of various spectra. In this section, we prove that Figure 1 gives an accurate description of the phase transitions between different parts of the integral means spectrum. Since all β 's and γ 's are given by simple algebraic equations, all inequalities are in principle elementary, but could be a bit fiddly if not addressed in the right way.

First, notice that γ_0 is increasing and negative for $t < 0$, whereas γ_1 is decreasing and positive. Next, we shall need to study $\gamma_0 + \gamma_1$. Differentiating with respect to t yields

$$\frac{\partial}{\partial t}(\gamma_0 + \gamma_1) = \frac{1}{\sqrt{(4+\kappa)^2/4 - 2t\kappa}} - \frac{1}{\sqrt{1-2t\kappa}} < 0,$$

so that $\gamma_0 + \gamma_1$ is decreasing; since $\gamma_0(0) = 0$ and $\gamma_1(0) = 2/\kappa$, this gives that $\gamma_0 + \gamma_1 > 2/\kappa$ iff $t < 0$.

Proposition 5.1. *Let $b = b(\gamma_0)$ be as in (3.6), then $1/2 - b > 0$ if and only if $t > t_1$.*

Proof. Owing to (3.6), $b = \gamma_0 + \gamma_1 - 2/\kappa$. Solving the equation $b = 1/2$ yields $t = t_1$. Since b is decreasing with t , we obtain that $1/2 - b > 0$ if and only if $t > t_1$. \square

Proposition 5.2. *The set of equations, $b = n + \frac{1}{2}$, $n \in \mathbb{N}$, is realized at the finite set of points $\mathcal{T}_\kappa := \{t_{1-n}, n \in \mathcal{J}_\kappa\}$, where $\mathcal{J}_\kappa := \{n \in \mathbb{N}, 0 \leq n \leq \lfloor \kappa^{-1} \rfloor\}$ and where t_{1-n} is given by Eq. (3.18).*

¹We believe that the logarithmic correction is not really there and that F/ψ is bounded.

Proof. For $t \in (-\infty, 0]$, we have that $b \in [0, \frac{1}{2} + \frac{1}{\kappa})$, which defines the possible range of values $n \in \mathcal{J}_\kappa$ where $b = n + \frac{1}{2}$. Solving the latter equation yields (3.18). \square

Proposition 5.3. *We have that $\beta_1(t) > \beta_0(t)$ if and only if $t < t_1$.*

Proof. First of all, recall that all exponents β 's are given by the same quadratic function $\beta(\gamma)$ (3.2). Since $\gamma_0 \neq \gamma_1$, $\beta(\gamma_1) = \beta(\gamma_0)$ if and only if $\gamma_1 + \gamma_0 = (4 + \kappa)/2\kappa$. We know that this happens at $t = t_1$ only. Computing at $t = 0$, we have $\beta_0(0) = 0$ and $\beta_1(0) = -1$, hence $\beta_1(t) > \beta_0(t)$ if and only if $t < t_1$. \square

Note that Propositions 5.1 and 5.3 imply that the point where the construction in Ref. [2] breaks down and the point where β_1 exceeds β_0 coincide.

Proposition 5.4. *We have that $\beta_1 < \beta_{\text{tip}} = \beta_0 - 2\gamma_0 - 1$ if and only if $t < 0$.*

Proof. From Eqs. (3.13) and (3.21) we have,

$$\beta_0 = -t + \frac{4 + \kappa}{2}\gamma_0, \quad \beta_1 = -t - \frac{\kappa}{2}\gamma_1,$$

so that

$$\beta_{\text{tip}} - \beta_1 = \frac{\kappa}{2}(\gamma_0 + \gamma_1) - 1 = \frac{\kappa}{2}b(\gamma_0),$$

which is positive iff $t < 0$. \square

5.2. Second tip and derivative exponents. Of particular interest here is the *packing spectrum* associated with the β_1 spectrum of whole-plane SLE $_\kappa$ (see Ref. [17] for a detailed discussion of the different spectra of the harmonic measure and their relations),

$$s(t) := \beta_1(t) - t + 1 = -2t + \frac{1}{2} - \frac{1}{2}\sqrt{1 - 2\kappa t}. \quad (5.1)$$

In the domain $t \leq t_1 < 0$, $s(t)$ is decreasing; its inverse is

$$t = -\nu(s), \quad \nu(s) := \frac{s}{2} + \frac{1}{16} \left(\kappa - 4 + \sqrt{(4 - \kappa)^2 + 16\kappa s} \right), \quad (5.2)$$

where, as remarked in Ref. [7, Section 4.4], $\nu(s)$ coincides with the non-standard multifractal tip exponents as obtained in Ref. [4, Eq. (12.19)], or with the so-called *derivative exponents* as obtained in Ref. [13, Eq. (3.1)] for standard (interior or exterior) radial SLE. For \hat{g}_τ such a radial SLE $_\kappa$ of hull K_τ , the exponent ν describes the exponential decay $e^{-\nu\tau}$ in time $\tau > 0$, of the moment of the boundary derivative modulus, $\mathbb{E}[|\hat{g}'_\tau(z)|^s]$, for $z \in \partial\mathbb{D} \setminus K_\tau$. It also governs the same exponential decay of the moment of order s , $\mathbb{E}[L_t^s]$, of the harmonic measure L_t of $\partial\mathbb{D} \setminus K_\tau$ in $\mathbb{D} \setminus K_\tau$, as seen from the origin in the interior case, or from infinity in the exterior case.

In the case of the interior whole-plane SLE, of map f_0 from \mathbb{D} to a slit domain [7], one has $f_0 = \lim_{\tau \rightarrow +\infty} e^\tau \hat{g}_{-\tau}$, which is in law the same as $\lim_{\tau \rightarrow +\infty} e^\tau \hat{g}_\tau^{-1}$, with \hat{g}_τ standard interior radial SLE; in this limit, the unit circle is pushed back to infinity as $e^\tau \partial\mathbb{D}$. Ref. [7, Section 4.4, Figure 8] then provides a heuristic explanation of the inverse relation between (5.1) and (5.2) as due, in the integral means $\int_{r\partial\mathbb{D}} |f'_0(z)|^t |dz|$ where $r \rightarrow 1^-$, to the local boundary contribution of the image under \hat{g}_τ of $\partial\mathbb{D} \setminus K_\tau$, i.e., in the limit $\tau \rightarrow +\infty$, of the pre-image under f_0 of the point at infinity.

The two whole-plane maps, interior f_0 and exterior g_0^{-1} , are naturally conjugate under the inversion map $z \mapsto 1/z$, as are the interior and exterior versions of radial SLE \hat{g}_τ ; in the exterior case, the unit circle shrinks as $e^{-\tau} \partial\mathbb{D}$ to a vanishingly small circle around the origin, to which the whole-plane SLE curve is anchored via this ‘second tip’. This strongly

suggests that for bounded whole-plane SLE, the β_1 -spectrum is due to the presence of the second tip, image of the point at infinity in the unbounded case.

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REFERENCES

- [1] T. Alberts, I. Binder, and F. Viklund. A dimension spectrum for SLE boundary collisions. *Commun. Math. Phys.*, 343:273–298, 2016.
- [2] D. Beliaev and S. Smirnov. Harmonic measure and SLE. *Commun. Math. Phys.*, 290(2):577–595, 2009.
- [3] B. Duplantier. Conformally invariant fractals and potential theory. *Phys. Rev. Lett.*, 84(7):1363–1367, 2000.
- [4] B. Duplantier. Conformal fractal geometry & boundary quantum gravity. In M. L. Lapidus and M. van Frankenhuysen, editors, *Fractal geometry and applications: a jubilee of Benoît Mandelbrot, Part 2*, volume 72 of *Proc. Sympos. Pure Math.*, pages 365–482. Amer. Math. Soc., Providence, RI, 2004.
- [5] B. Duplantier, H. Ho, B. Le, and M. Zinsmeister. Logarithmic Coefficients and Generalized Multifractality of Whole-Plane SLE. Apr. 2015. arXiv:1504.05570.
- [6] B. Duplantier, C. Nguyen, N. Nguyen, and M. Zinsmeister. Coefficient estimates for whole-plane SLE processes. July 2011. <https://hal.archives-ouvertes.fr/hal-00609774v1>.
- [7] B. Duplantier, C. Nguyen, N. Nguyen, and M. Zinsmeister. The Coefficient Problem and Multifractality of Whole-Plane SLE and LLE. *Ann. Henri Poincaré*, 16(6):1311–1395, 2014. arXiv:1211.2451.
- [8] J. B. Garnett and D. E. Marshall. *Harmonic measure*, volume 2 of *New Mathematical Monographs*. Cambridge University Press, Cambridge, 2007. Reprint of the 2005 original.
- [9] E. Gwynne, J. Miller, and X. Sun. Almost sure multifractal spectrum of SLE. (2014). arXiv:1412.8764.
- [10] M. B. Hastings. Exact multifractal spectra for arbitrary Laplacian random walks. *Phys. Rev. Lett.*, 88(5):055506, 2002.
- [11] F. Johansson Viklund and G. Lawler. Almost sure multifractal spectrum for the tip of an SLE curve. *Acta Math.*, 209(2):265–322, 2012.
- [12] G. F. Lawler. *Conformally Invariant Processes in the Plane*, volume 114 of *Mathematical Surveys and Monographs*. American Mathematical Society, Providence, RI, 2005.
- [13] G. F. Lawler, O. Schramm, and W. Werner. Values of Brownian intersection exponents. II. Plane exponents. *Acta Math.*, 187(2):275–308, 2001.
- [14] I. Loutsenko. SLE $_{\kappa}$: correlation functions in the coefficient problem. *J. Phys. A Math. Gen.*, 45(26):265001, 2012.
- [15] I. Loutsenko and O. Yermolayeva. On exact multi-fractal spectrum of the whole-plane SLE. Mar. 2012. arXiv:1203.2756.
- [16] I. Loutsenko and O. Yermolayeva. Average harmonic spectrum of the whole-plane SLE. *J. Stat. Mech.*, 2013(4):P04007, 2013.

- [17] N. G. Makarov. Fine structure of harmonic measure. *St. Petersburg Math. J.*, 10(2):217–268, 1999.
- [18] J. Miller and S. Sheffield. Imaginary geometry I: interacting SLEs. *Probab. Theory Relat. Fields*, 164(3):553–705, 2016. arXiv:1201.1496.
- [19] J. Miller and S. Sheffield. Imaginary geometry II: reversibility of $SLE_{\kappa}(\rho_1; \rho_2)$ for $\kappa \in (0, 4)$. *Ann. Probab.*, 44(3):1647–1722, 2016. arXiv:1201.1497.
- [20] J. Miller and S. Sheffield. Imaginary geometry III: reversibility of SLE_{κ} for $\kappa \in (4, 8)$. *Ann. Math.*, 184:1–32, 2016. arXiv:1201.1498.
- [21] J. Miller and S. Sheffield. Imaginary geometry IV: interior rays, whole-plane reversibility, and space-filling trees. *Probab. Theor. Relat. Fields*, 2017. To appear. arXiv:1302.4738.
- [22] S. Rohde and O. Schramm. Basic properties of SLE. *Ann. Math. (2)*, 161(2):883–924, 2005.
- [23] O. Schramm. Scaling limits of loop-erased random walks and uniform spanning trees. *Israel J. Math.*, 118:221–288, 2000.

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