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# Study of annealing-induced interdiffusion in $\text{In}_2\text{O}_3/\text{Ag}/\text{In}_2\text{O}_3$ structures by a combined X-ray reflectivity and grazing incidence X-ray fluorescence analysis



Bérenger Cabby<sup>a,\*</sup>, Fabio Brigidi<sup>d,g</sup>, Dieter Ingerle<sup>f</sup>, Emmanuel Nolot<sup>a</sup>, Giancarlo Pepponi<sup>d</sup>, Christina Strelj<sup>f</sup>, Luca Lutterotti<sup>e</sup>, Agathe André<sup>a</sup>, Guillaume Rodriguez<sup>a</sup>, Patrice Gergaud<sup>a</sup>, Magali Morales<sup>c</sup>, Daniel Chateigner<sup>b</sup>

<sup>a</sup> CEA, LETI, MINATEC Campus, 17 rue des Martyrs, 38054 Grenoble Cedex 9, France

<sup>b</sup> CRISMAT-ENSICAEN, IUT-Caen UCBN, 6 boulevard du Maréchal Juin, 14050 Caen Cedex 4, France

<sup>c</sup> CIMAP, 6 boulevard du Maréchal Juin, 14050 Caen Cedex 4, France

<sup>d</sup> MNF, CMM-Irst, Fondazione Bruno Kessler Via Sommarive 18, Povo, Trento 38123, Italy

<sup>e</sup> Department of Industrial Physics, University of Trento, via Sommarive 14, Povo, Trento 38123, Italy

<sup>f</sup> Atominstytut, Vienna University of Technology, Vienne 1020, Austria

<sup>g</sup> Department of Industrial Engineering, University of Trento, via Sommarive 9, Povo, Trento 38123, Italy

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## ABSTRACT

The combination of X-ray reflectivity and grazing incidence X-ray fluorescence has been applied to the characterization of an  $\text{In}_2\text{O}_3/\text{Ag}/\text{In}_2\text{O}_3$  stack for advanced photovoltaic applications. X-ray reflectivity is a well-known method for the characterization of multilayered structures by providing information on the thickness and the in-depth electronic density. Grazing incidence X-ray fluorescence provides information about the elemental depth distribution. As these techniques are based on similar measurement procedures and data evaluation approaches, their combination reduces the uncertainties of the individual techniques and provides an accurate depth-resolving analysis of multi-layers.

It has been shown that the combination of the techniques give insight into the material composition and the layers structure (thickness, density) as well as modifications induced by a thermal annealing.

As X-ray fluorescence signals have been acquired at different excitation energies, the influence of this parameter on the sensitivity of the measurements to the structural properties has been shown.

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## 1. Introduction

Transparent and conductive oxide (TCO) layers are essential components in several emerging photosensitive electronic products. Metal-oxides such as  $\text{SnO}_2$ ,  $\text{In}_2\text{O}_3$ , or  $\text{ZnO}$  have widely been used acting as transparent electrical contacts or electrodes in flat panel displays, touch screens, and thin film solar cells [1,2].

There is a wide range of requirements for these layers depending on the application. First, resistivity needs to be tailored to meet the electrical functionality. For solar cells, the sheet resistance ( $R_s$ ) must be in the 8–80  $\Omega/\text{sq}$  range [3]. The other key factor is the optical transmission ( $T$ ) of the layers. For photovoltaic applications, the material must be transparent in the visible spectral range.

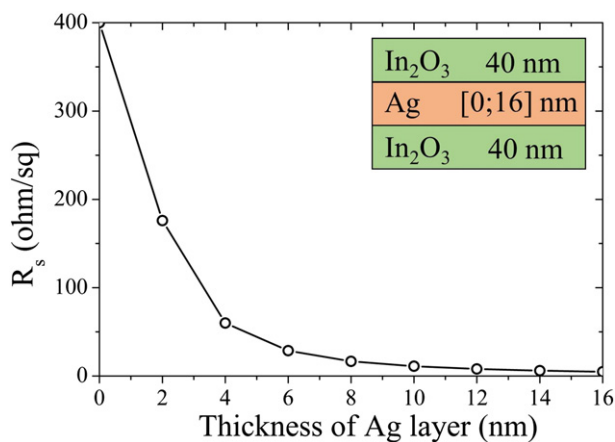
Due to a technological need for large-area photovoltaic devices with improved electrical and optical performances, the development of new structures is necessary. Recent investigations have shown that TCO/

metal/TCO multilayer systems are suitable as transparent conductive electrodes for photovoltaic devices [4,5].

To achieve the best balance between the optical and electrical properties of the structure, the metal layer has to be a low-resistivity material and Ag (1.6  $\mu\Omega \text{ cm}$  at 20 °C) is the most common interlayer used for TCO/metal/TCO electrodes [6]. Due to its high transmittance in the visible spectral range, indium oxide ( $\text{In}_2\text{O}_3$ ) also referenced as IO is used to embed the conductive layer. Moreover, the thickness of the layers has an important effect on the electrical conductivity and optical transmittance. As shown on Fig. 1 and by extrapolating from Fig. 2, a 6 nm thick Ag film embedded between two 40 nm thick IO films gives the best compromise between resistivity and transparency and hence the best performances for photovoltaic applications.

A combination of X-ray reflectivity (XRR) and grazing incidence X-ray fluorescence (GIXRF) measurements have been performed on such IO/Ag/IO structures. This approach has been carried out with the idea that the combined analysis would provide more detailed information than the sum of the individual techniques about the layer structure and its possible modifications due to a thermal annealing.

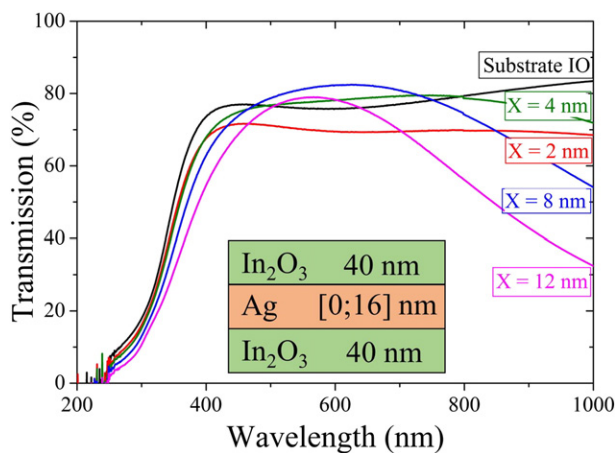
\* Corresponding author. Tel.: +33 438781906.  
E-mail address: [berenger.cabby@cea.fr](mailto:berenger.cabby@cea.fr) (B. Cabby).



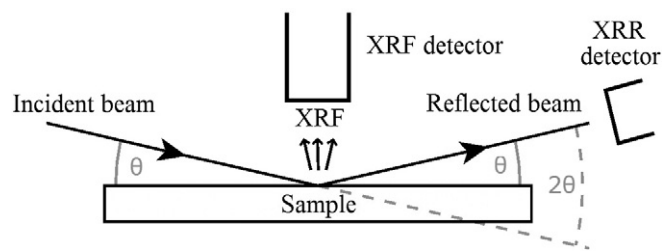
**Fig. 1.** Sheet resistance for IO/Ag/IO tri-layer samples with an Ag layer with varying thickness. 0 nm of Ag represents the sheet resistance of an 80 nm pure  $\text{In}_2\text{O}_3$  layer. For photovoltaic applications, acceptable sheet resistances have been obtained with Ag thickness superior to 4 nm.

GIXRF experiments have already been mainly employed for the characterization of thin films [7,8]. At incidence angles close to the critical angle of total external reflection, due to the interference between the incoming and the reflected beam, an X-ray standing wave field (XSW) is created above and inside the material [9]. By varying the incidence angle, the wave field intensity distribution changes and affects the X-ray fluorescence signal. This angular dependency of the GIXRF signals provides information on the elemental depth distribution as well as the thickness, density, and roughness of thin layers and interfaces.

XRR is a non-destructive, highly accurate method based on the reflection in the specular direction of X-rays at the surface and interfaces of the sample [10] and is used to determine thickness, in-depth electronic density, and roughness of thin layers. Both XRR and GIXRF methods can be operated simultaneously on a unique experimental setup using the same primary X-ray radiation (Fig. 3). The combination of the complementary information that these two techniques provide is used to reduce the uncertainties of the individual methods and allows a determination of the depth profile with a better accuracy [11]. The combined analysis of IO/Ag/IO stacking has been performed in order to obtain the material composition and the layers structure as well as their modifications induced by a thermal annealing.



**Fig. 2.** Optical transmittance for IO/Ag/IO tri-layer samples with an Ag layer with varying thickness. For photovoltaic applications, the transparency in the visible is necessary and sample with an Ag thickness of 12 nm are not adapted.



**Fig. 3.** Experimental setup for combined XRR and GIXRF analysis.

**Table 1**

Samples analyzed using combined XRR-GIXRF.

Sample	Top layer	Annealing
A	40 nm	
B	40 nm	200 °C for 1 h
C	15 nm	
D	15 nm	200 °C for 1 h

## 2. Experiments

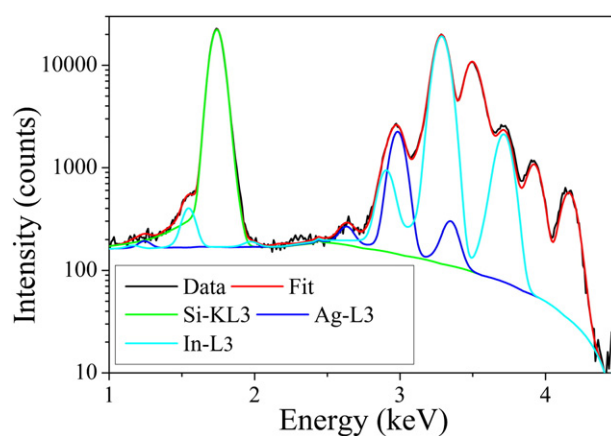
### 2.1. Sample preparation

The structure of the samples studied by combined XRR-GIXRF analysis consists in a 6 nm Ag layer embedded between two 40 nm  $\text{In}_2\text{O}_3$  layers deposited on a 500 nm  $\text{SiO}_2$  / $\text{Si}(001)$  substrate. Each layer has been deposited via PVD (physical vapor deposition) with a low-temperature deposition process.

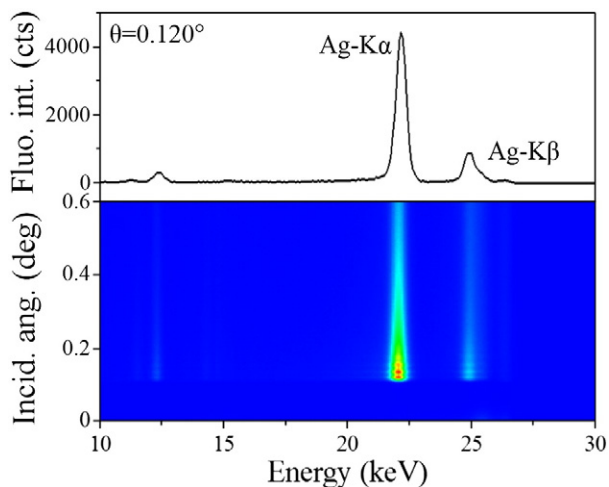
In order to simulate a realistic technological integration, a thermal budget has been applied to duplicate samples. The annealing has been made in a Tempress furnace system at 200 °C during 1 h. To improve the sensitivity on the  $\text{In}_2\text{O}_3$ /Ag interface, the thickness of the top IO layer has been reduced via CMP (chemical mechanical polishing) to a thickness of approximately 15 nm. The same approach was used by Giubertoni et al. to gain sensitivity on buried layers when doing grazing incidence absorption studies [12]. The list of analyzed samples is reported in Table 1.

### 2.2. Acquisition

The first measurement campaign was performed at the Atominstitut (Vienna, Austria) where both XRR and GIXRF measurements have



**Fig. 4.** Fluorescence intensity extraction from a measurement at 1.2° with a Cu-tube. Due to the correlation of Ag-L3 and In-L3 fluorescence peaks, the use of a third-party software like PyMCA [14] is necessary.

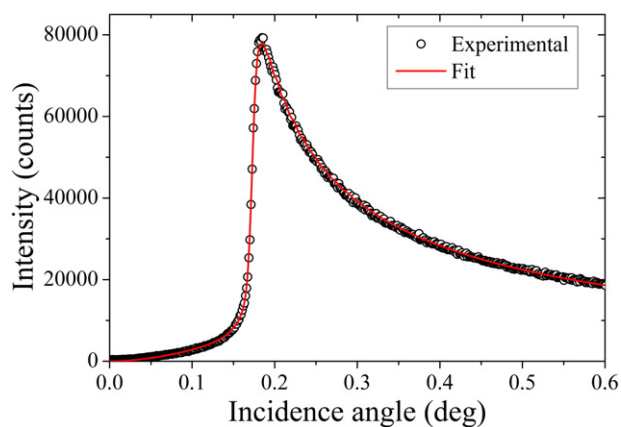


**Fig. 5.** Fluorescence intensity from a measurement at 26,400 eV with a synchrotron radiation at an angle of incidence  $\theta = 0.12^\circ$ . With this excitation energy, the fluorescence peaks of interest are not overlapping.

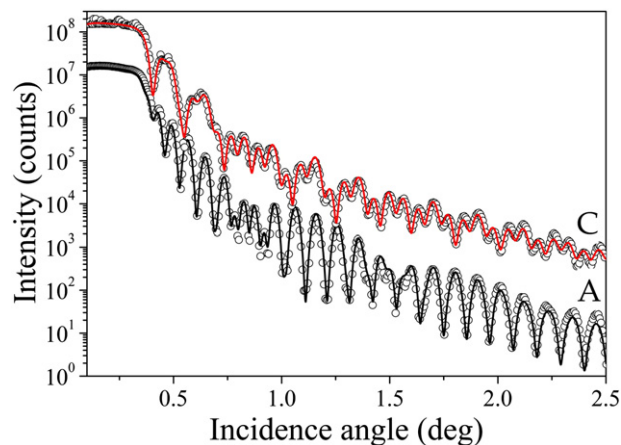
been carried out on a self-developed system with a 2 kW fine focus Cu anode tube and a multilayer monochromator selecting the Cu-K $\alpha$  line (8050 eV), both mounted on a theta-2theta goniometer under vacuum [13]. An Amptek Silicon Drift Detector (SDD) was used to measure the reflected beam. A Zr filter was inserted in front of this detector at low angles to reduce the primary beam intensity. The fluorescence signal was measured with a Vortex SDD placed at  $\Psi_d = 90^\circ$  above the sample. Spectra were acquired from  $\theta = 0^\circ$  to  $3^\circ$  with an angular step of  $0.05^\circ$  for a total count time of approximately 1 h.

As shown in Fig. 4, In-L and Ag-L fluorescence lines as well as Si pile-up peak overlap in the 3 keV range. At low incidence angles (typically  $< 0.4^\circ$ ), due to the absorption of the incident beam in the top IO layer, the fluorescence signal of the Ag-L $\alpha$  line is low and completely covered by the In-L fluorescence lines. At this excitation energy, working under vacuum is necessary in order to get rid of Ar-K $\alpha$  lines (2957 eV) and avoid further overlaps with the fluorescence lines of interest.

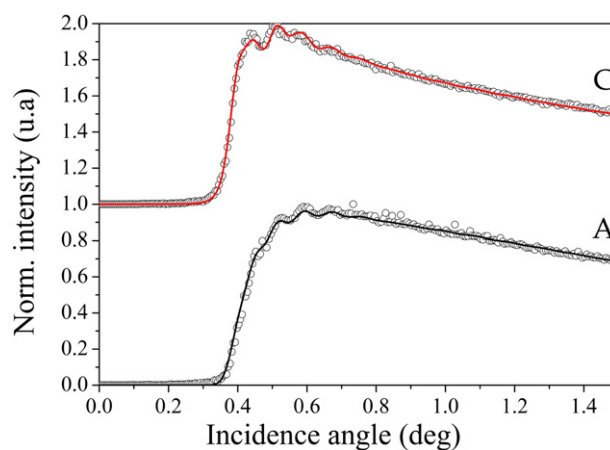
In order to overcome signal extraction difficulties due to overlaps and analyze more intense fluorescence K-lines, additional GIXRF measurements have been carried out on the Gilda (BM08) line at ESRF-Grenoble. To improve the sensitivity, the measurements have been realized at 26,400 eV above the Ag-K but below the In-K edge. To avoid sample edge effects during the measurements, the horizontal size of the primary beam was reduced by a 50  $\mu\text{m}$  slit. A 13-element



**Fig. 6.** GIXRF fit on JGIXA of a 50 nm Ni sample measured at 26,400 eV at ESRF. Fitting of reference sample allows the determination of geometrical factors such as the divergence of the beam and the size of the illuminated area.



**Fig. 7.** XRR measurements (open circles) and fits (solid lines) of the as-deposited samples A and C at 8.05 keV.



**Fig. 8.** Ag-L $\alpha$  GIXRF measurements (open circles) and fits (solid lines) of the as-deposited samples A and C at 8.05 keV.

high-purity Ge-SSD (solid state detector) with a maximum count rate of 80 k counts per second was used to measure the fluorescence. Energy dispersive X-ray spectra were acquired from  $\theta = 0^\circ$  to  $0.6^\circ$  with an angular step of  $0.001^\circ$  for a total acquisition time of 10 minutes. The energy tunability of the synchrotron source allowed the study of the in-depth Ag-K $\alpha$  line isolated from the indium contributions (Fig. 5).

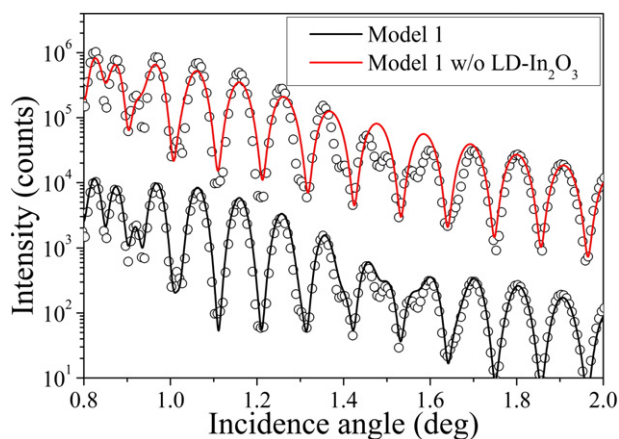
### 2.3. Data reduction and numerical models

Extraction of fluorescence intensities from data measured with the Cu-tube has been done using the software PyMCA [14]. The background

**Table 2**

Model 1—Thickness and densities of the as-deposited samples obtained by comparison of JGIXA fitting and experimental results measured with Cu-K $\alpha$  energy.

Sample A	Thickness (nm)	Density ( $\text{g}\cdot\text{cm}^{-3}$ )
LD-In $_2$ O $_3$	1.2 nm	4.04
In $_2$ O $_3$	40.1 nm	6.91
Ag	6.0 nm	9.83
In $_2$ O $_3$	40.0 nm	6.95
Sample C	Thickness (nm)	Density ( $\text{g}\cdot\text{cm}^{-3}$ )
LD-In $_2$ O $_3$	1.4 nm	4.1
In $_2$ O $_3$	15.0 nm	6.61
Ag	6.0 nm	9.87
In $_2$ O $_3$	40.0 nm	6.88



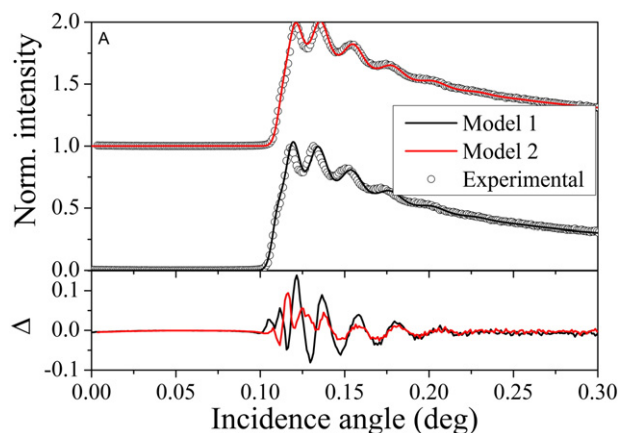
**Fig. 9.** XRR measurements (open circles) and fits (solid lines) of the as-deposited samples A at 8.05 keV with and without the introduction of the low-density  $\text{In}_2\text{O}_3$  top layer in the model. It can easily be seen that in order to reproduce the experimental data around  $1.5^\circ$ , the introduction of a low-density top layer is necessary.

has been simulated by the strip background model implemented in the software. For the measurements done at ESRF, as the peaks are not overlapped and the signal-to-noise ratio is high, the integral of the fluorescence counts in a region of interest (ROI) has been used.

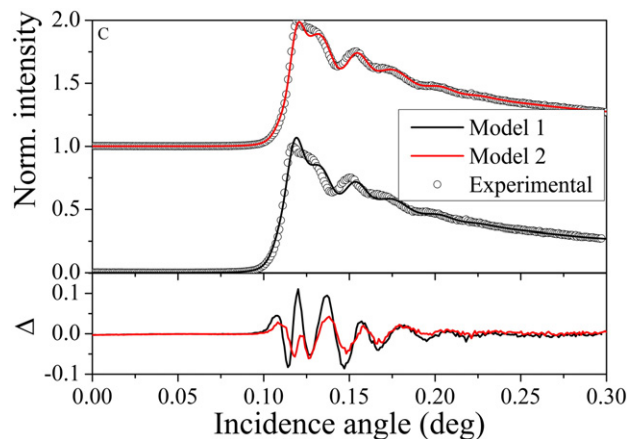
The modeling of reflectivity and fluorescence measurements has been carried out using the JGIXA software [15] with two refinement configurations. As GIXRF have been acquired at different excitation energies, a first model consists in fitting XRR and GIXRF spectra acquired at the same excitation energy, while in the second model, the excitation energies of XRR and GIXRF are different.

For the optimization procedure, two sets of parameters, *i.e.* for the sample and the experimental setup, are defined as starting values. The sample parameters to be evaluated for each layer of the stack are the composition, thickness, mass density, and roughness. Each can be independently refined. Due to existing local minima, global optimization algorithms are required to find the global minimum. These algorithms reduce the dependency on initial values of the parameters and hence are less relying on a priori knowledge. The fitting of the measured data has been optimized with a differential evolution (DE) algorithm included in the software [16].

The angular profile of GIXRF also depends on experimental setup parameters such as the divergence of the beam, the size of the beam, and



**Fig. 10.** Ag-K $\alpha$  GIXRF Gilda measurements and fits of sample A. The difference curve is represented at the bottom of the figure. Model 2 has been obtained by combining sets of data acquired at different excitation energies. It offers smaller differences between experimental and simulated data.



**Fig. 11.** Ag-K $\alpha$  GIXRF Gilda measurements and fits of sample C. The difference curve is represented at the bottom of the figure. Model 2 has been obtained by combining sets of data acquired at different excitation energies. It offers smaller differences between experimental and simulated data.

the inspected area. They have been determined by the measurement on the same experimental setups and evaluation of already known samples. As seen in Fig. 6, a 50 nm Ni layer deposited on a 300 nm  $\text{SiO}_2/\text{Si}$  substrate has been used as reference sample for the determination of the geometrical factors. These experimental parameters are then fixed during the simulations of the TCO/metal/TCO samples.

### 3. Results and discussion

#### 3.1. As-deposited samples

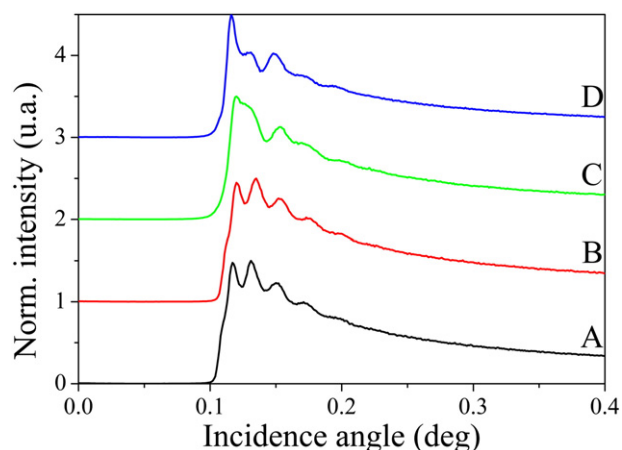
Data measured with the standard laboratory Cu-tube (open circles) and JGIXA simulations (solid lines) for the as-deposited samples (sample A and sample C) are shown in Figs. 7 and 8. The agreement between the experimental and simulated points is very good.

A first structure (model 1, Table 2) resulting from refinement of these data has been obtained. The incorporation of a top indium oxide layer with a low density ( $\text{LD-In}_2\text{O}_3$ ) is necessary in order to reproduce all the fine details of the reflectivity curves and improve the refinement agreement (Fig. 9). For both samples, the surface and interfacial roughness ranges from 5 to 12 Å. The accuracy of XRR for the determination of thickness and roughness is high with uncertainties lower than 5 Å and 1 Å, respectively [17]. Thanks to the decoupling of the sample parameters, the determination of the electronic density is reinforced by the use of the XRR and GIXRF curves for a joint fit. This aspect remains sensitive despite weak signal-to-noise ratio and rather smoothed fluorescence fringes.

Compared to the laboratory GIXRF measurements, higher-energy experiments carried out at the Gilda ESRF beam line have both the advantage of non-overlapped fluorescence lines and a much stronger

**Table 3**  
Model 2—Thickness and densities of the as-deposited samples obtained from correlating JGIXA fits of the experimental data measured at two different energies (respectively 26.4 and 8.05 keV).

Sample A	Thickness (nm)	Density ( $\text{g}\cdot\text{cm}^{-3}$ )
$\text{LD-In}_2\text{O}_3$	0.9	4.44
$\text{In}_2\text{O}_3$	40.1	7.28
Ag	6	9.8
$\text{In}_2\text{O}_3$	40	7.1
Sample C	Thickness (nm)	Density ( $\text{g}\cdot\text{cm}^{-3}$ )
$\text{LD-In}_2\text{O}_3$	1.4	4.31
$\text{In}_2\text{O}_3$	15	6.95
Ag	6	10.4
$\text{In}_2\text{O}_3$	40	6.99



**Fig. 12.** Ag-K $\alpha$  GIXRF measurement of the IO/Ag/IO samples at 26,400 eV. Thinning of the top IO layer has greatly improved the sensitivity of the GIXRF technique. Differences in the GIXRF profile can be noted between the C (as deposited) and the D (annealed) sample.

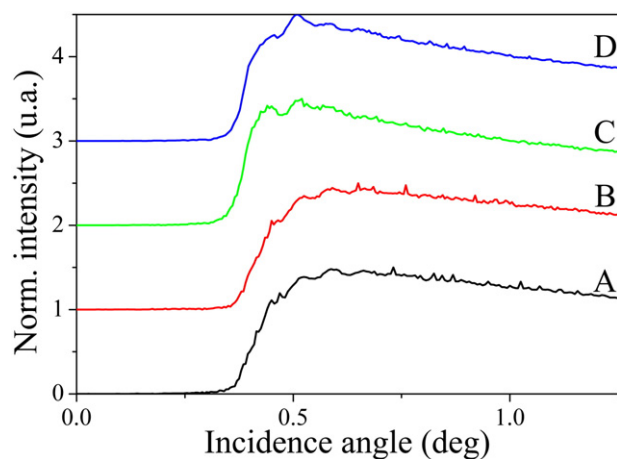
dependence of the fluorescence intensities with the incidence angle. New simulations have been done with this set of data (Figs. 10 and 11).

By refinement of parameters, an improved model (model 2, Table 3) has been obtained by combining the reflectivity measured at Cu-K $\alpha$  energy and the fluorescence measured at the excitation energy of 26,400 eV. The difference curve between calculated and experimental data for the two models compares favorably for model 2, both for XRR and GIXRF experiments via the correction of the densities values. These density variations (around 5% for each layer) do not have any effects on the reflectivity calculations.

The sensitivity of the combined analysis has been greatly improved by using X-ray measurement done at adapted primary energies. The joint fit of standard laboratory XRR and synchrotron-based GIXRF has provided more accurate information about the in-depth profile.

### 3.2. Annealed samples

After annealing at 200 °C, we tested the efficiency of the combined XRR and GIXRF approach to detect temperature-induced modifications of the stack. In Figs. 12 and 13, the Ag-K $\alpha$  fluorescence angular intensities for the four samples are presented. Thinning of the top IO layer has improved the sensitivity of GIXRF to interfaces and reveals differences much more illustrated between samples C and D, than between A and B on which only a slight shift of the whole spectra is being visible.



**Fig. 13.** Ag-L $\alpha$  GIXRF measurements of the IO/Ag/IO samples at 8050 eV. Thinning of the top IO layer has greatly improved the sensitivity of the GIXRF technique. Differences in the GIXRF profile can be noted between the C (as deposited) and the D (annealed) sample.

**Table 4**

Thickness and densities of the sample D obtained from correlating JGIXA fits of the experimental data measured at two different energies (respectively 26.4 and 8.05 keV).

Sample D	Thickness (nm)	Density (g.cm <sup>-3</sup> )
LD-In <sub>2</sub> O <sub>3</sub>	2	6.12
In <sub>2</sub> O <sub>3</sub>	16.3	7.18
In <sub>2</sub> O <sub>3</sub> ; Ag	1.1	8.36
Ag	5.7	9.9
In <sub>2</sub> O <sub>3</sub> ; Ag	0.6	7.47
In <sub>2</sub> O <sub>3</sub>	40	7

Such behavior was also observed on the Ag-L $\alpha$  GIXRF profiles measured with laboratory instrument.

One could interpret this difference in the GIXRF profile as an annealed-induced interdiffusion between the Ag and the In<sub>2</sub>O<sub>3</sub> layers. We tested a rough model for such an interdiffusion by incorporating intermediate density 1 nm thick layers around the Ag/In<sub>2</sub>O<sub>3</sub> interfaces (Table 4). The results of the combined fit of the lab-based XRR and synchrotron-based GIXRF are presented in the Fig. 14 and Table 3. The simulation highlights the diffusion of the Ag as an intermixing of elements is obtained in the transition layers. The surface and interfacial roughness ranges from 4 to 10 Å.

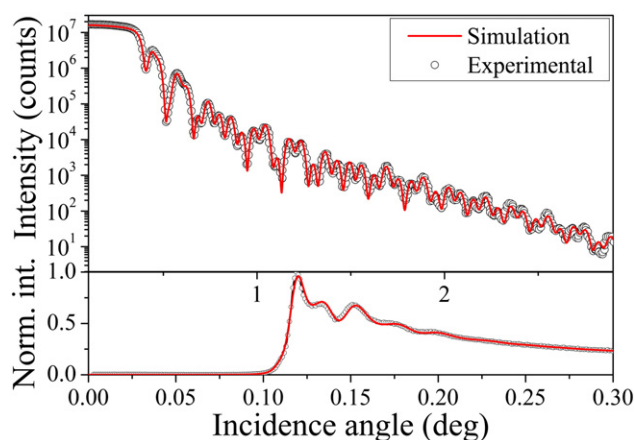
The combination of XRR and GIXRF allows detecting and characterizing a thin diffusion profile at the interface of two layers. Even on non-optimized experimental setup (*i.e.* laboratory Cu-tube), the diffusion induced by thermal annealing is noticeable.

## 4. Conclusions and outlook

The combined evaluation of GIXRF and XRR resulted in a non-destructive and precise characterization of the In<sub>2</sub>O<sub>3</sub>/Ag/In<sub>2</sub>O<sub>3</sub> multi-layered structure. It has been shown that the sensitivity of the joint approach can be optimized with the choice of the primary energy.

A thin interdiffusion profile induced by annealing has been observed and characterized. Sensitivity to structural changes in the depth of 0.5–1 nm was achieved for measurements performed at a synchrotron facility but also for the ones acquired with the laboratory experimental setup.

The effect of the annealing on material properties such as the crystallographic phase, the size of the crystallites, and the residual stress have been observed by X-ray diffraction measurements (XRD). In order to obtain a full characterization of these photovoltaic devices, the combination of XRD, XRR, and GIXRF in respect with the models presented in this work will be done.



**Fig. 14.** XRR and GIXRF measurements and fitting of the sample D with reflectivity measured at 8.05 keV and fluorescence at 26.4 keV.

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