

# From Fick's second law and Kirchhoff's current law in a subgranular sintering model: densification rate



S. Martin<sup>1,2</sup>, J. Léchelle<sup>1\*</sup>, R. Boyer<sup>3</sup>, K. Saikouk<sup>3</sup>, K. Saleh<sup>2</sup>, M. Guessasma<sup>4</sup>, J. Fortin<sup>4</sup>

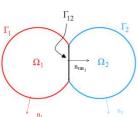
CEA, DEN, DEC, SPUA, LMPC, F-13108 Saint-Paul Lez Durance, France (jacques.lechelle@cea.fr; + 33 (0) 4 42 25 73 61)
Équipe d'Accueil Transformations Intégrées de la Matière Renouvelable (EA 4297), Université de Technologie de Compiègne, 60200 Compiègne, France
ALTP, CMI, Aix-Marseille Université, Technopôle Château-Gombert, 39, rue F. Joliot Curie, F-13453 Marseille cedex 13, France
Université de Picardie Jules Verne (UPJV), Laboratoire des Technologies Innovantes, 48 rue Ostende 02100 St Quentin, France

# Context

> Sinterability of powders for nuclear fuel pellets is of major importance for fuel fabrication, to this aim a sub-granular model is under development

> In the framework of the modelling of sintering on a sub-granular scale, a densification law has been derived in the case of two grains. Predictions from these laws are

# **Model description**



1. Elastic behaviour at each time step: Displacement field computation with Von Neumann Boundary conditions: Laplace law

For  $\Gamma_{\rm l}$  and  $\ \Gamma_{\rm 2}$  free surfaces:  $\overline{\overline{\sigma}}\left(\vec{n}\right)\cdot\vec{n}=-p_{gaz}-2\gamma_{SV}k_{SV}$ with  $\sigma$  the stress tensor,  $\vec{n}$  the outer normal,  $\gamma_{SV}$  the Solid/Vapor surface energy and  $\boldsymbol{k}_{SV}$  the local mean curvature

For  $\Gamma_{12}$  grain boundary:  $\overline{\overline{\sigma_1}}(\overrightarrow{n_1}) \cdot \overrightarrow{n_1} = \overline{\overline{\sigma_2}}(\overrightarrow{n_1}) \cdot \overrightarrow{n_1} - 2\gamma_b k_b$ 

## 2. Chemical potential µ with a mechanical origin:

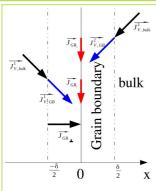
$$\mu = -\frac{1}{3}tr(\sigma) - \frac{1}{2}\Theta\left(\sigma:s:\sigma\right) \ \ \text{with s the compliance matrix}$$

### 3. Onsager first law: mass flow

Bulk, free surface and grain boundary fluxes:  $\vec{J} = -L\vec{\nabla} \left(\frac{\mu}{T}\right)$ 

## 4. Onsager second law: interface movements

Normal speed V<sub>int</sub> of interfaces → irreversible mass flow  $\iint \rho \overrightarrow{V_{int}} \cdot \vec{n} dS = M \iint \left( \left( \overrightarrow{j_{\text{GB}_{\perp}}} + \overrightarrow{j_{\text{GB}_{\perp}}} \right) \cdot \vec{n} + \text{div}_{S} \left( \overrightarrow{j_{GB}} - \overrightarrow{j_{GB}} \right) + \left( \delta \left( \overrightarrow{j_{\text{GB}_{\perp}}} - \overrightarrow{j_{\text{GB}_{\perp}}} \right) \cdot \overrightarrow{t} \right) \right) dS$ 

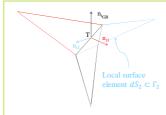


It has been shown that if  $\dot{V}_1$  is the total mass of matter (or quasi-incompressible volume) produced along the grain boundary by grain 1 per time unit in

$$\dot{V}_1 = -\Omega \int \int_{\Sigma_{\text{\tiny CBL}}} (\overrightarrow{j_{\text{\tiny GBL}}^1} \, \overrightarrow{n}^1 + \text{div}_{\text{S}} (\overrightarrow{j_{\text{\tiny GB}}^1} + (\delta \overrightarrow{j_{\text{\tiny GBL}}^1} \, \overrightarrow{t}) \overrightarrow{t}) dS$$

 $\dot{V}_1$  total mass of matter (or quasi-incompressible volume) produced along the grain boundary by grain 1 per time unit in direction of grain 2

$$\mathbf{X} \quad \dot{V}_2 = -\Omega \int \int_{\Sigma_{GB}} (\vec{j}_{GB_{\perp}}^2 \vec{n}^2 + \operatorname{div}_S (\vec{j}_{GB}^{2} + (\delta \vec{j}_{GB_{\perp}}^2 \vec{t}) \vec{t}) dS$$



At triple point T:  $\overrightarrow{J_{S_1}} \cdot \overrightarrow{n_{t_1}} = -\overrightarrow{J_{GB_1}} \cdot \overrightarrow{n_{t_{GB}}}$  with  $\overrightarrow{n_{t_1}}$  the normal to the triple edge tangent with surface 1 Une normal to the triple edge tangent with surface So that if  $\overrightarrow{J_{V_1}}$  can be neglected and Ostrogradski relation is used:  $\overrightarrow{V_1} = \Omega \oint_{\Gamma_{12}} \overrightarrow{J_{S_1}} \cdot \overrightarrow{n_{t_1}} \, dt$ 

 $\overrightarrow{J_{S_1}} = -L\overrightarrow{\nabla_S}\left(\frac{\mu}{\tau}\right)$  with  $\mu$  the chemical potential

The shrinkage of two grains reads:  $\dot{h} = \frac{1}{2A_{GR}} (\dot{V}_1 + \dot{V}_2)$  with  $A_{GB}$  the area of the grain boundary

# Model predictions for densification

### ➤ A consequence of Kirchoff Current Law

The shrinkage rate depends only on the shape of the free surfaces in the vicinity of the triple line,

For rigid grains ( $\mu_{bulk} = \mu_0$ ), the chemical potential at a surface can be expressed without solving any mechanical Partial Differential Equations:

$$\mu = \mu_{bulk} \, + \, \Omega \left\{ \gamma_S \left[ \frac{1}{R_1} + \frac{1}{R_2} \right] + \frac{1}{R_1} \frac{\partial^2 \gamma_S}{\partial^2_{n_1}} + \frac{1}{R_2} \frac{\partial^2 \gamma_S}{\partial^2_{n_2}} \right\}$$

### > A consequence of the surface chemical potential expression

For spherical rigid grains (R<sub>1</sub>=R<sub>2</sub>=sphere radius)  $\overline{V}_{SR} = \vec{0}$  diffusion along free surfaces can only occur if the material is non-isotropic, i.e.:  $\frac{\partial^2 Y_S}{\partial^2 n_t} \neq 0$ 

For rigid grains, absence of diffusion perpendicular to the grain boundary ( $\overrightarrow{j_{GB^{\perp}}} = \overrightarrow{0}$ ) i.e. with only grain boundary and surface diffusion the shrinkage rate would be zero:  $\dot{h} = 0$ 

The latter conclusions are changed when grains are no more considered as rigid but for instance as elastic at a given time (quasi-static step) during the sintering process  $(\mu_{bulk} = f(\sigma) \neq \mu_0)$ .

# Conclusions

- > The mechanical model developed can be a useful tool for the study of impact of initial grain morphology upon sintering
- > Non obvious consequences of the mechanical part of the chemical potential upon grain shrinkage rate can be discussed