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NANO SAMPLING OF METALS WITH ULTRA SHORT LASER PULSES

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Cea

CEA Saclay micro-LIBS instrument for surface micro-imaging / elemental cartography



[*] D. Menut, P. Fichet, J.-L. Lacour, A. Rivoallan, and P. Mauchien, Micro-laser-induced breakdown spectroscopy technique: a powerful method for performing quantitative surface mapping on conductive and nonconductive samples, APPLIED OPTICS **42**, 2003, p. 6063





CEA Saclay near-field LA-ICP-MS installation



Nd-YAG laser (266 nm), pulse duration 4 ns. AFM: Atomic Force Microscope. ICP-MS: Inductively Coupled Plasma - Mass Spectrometry : double focusing sector field cea







[**] C. Jabbour, J.-L. Lacour, M. Tabarant, A. Semerok and F. Chartier, Development of a tip enhanced near-field laser ablation system for sub-micrometric analysis of solid samples, J. Anal. At. Spectrom., 2016, **31**, 1534–1541

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Aim of our theoretical studies



for
$$a \in [100 \,\mu\text{m} \rightarrow 10 \,\text{nm}]$$

for $\tau \in [4 \,\text{ns} \rightarrow 100 \,\text{fs}]$

$$Q(t, x, y, z) \Rightarrow T(t, x, y, z)$$

 $T_{max} = f$ (interaction parameters) $\Leftrightarrow T_{ablation}$

Laser pulses: 266 nm, 4 ns or 100 fs;

Samples: Si or Au

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ns-LASER HEATING : ONE-TEMPERATURE MODEL

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$$c(T)\frac{\partial T}{\partial t} = div(k(T)\nabla T) + Q(t,r,z)$$

$$Q(t,r,z) = I(t,r)A\alpha \exp(-\alpha z)$$
or
$$Q(t, x, y, z) = \alpha n \frac{c |\vec{E}|^2}{8\pi} \text{ from near-field calculations}$$

$$T(t = 0, r, z) = T_o$$
 and $\frac{\partial T(t, r, z)}{\partial z}\Big|_{z=0} = 0$

c (J m⁻³ K⁻¹) – thermal capacity; k (W m⁻¹ K⁻¹) – thermal conductivity; α (m⁻¹) – absorption coefficient of media; A – surface absorptivity at laser wavelength;

[***] A. Semerok, S. V. Fomichev, J.-M. Weulersse, F. Brygo, and P.-Y. Thro, Heating and ablation of tokamak graphite by pulsed nanosecond Nd-YAG lasers, JOURNAL OF APPLIED PHYSICS **101**, 084916, 2007.

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Si-bulk: c = 700 J/kg K, k = 149 W/m K, $\rho = 2330 \text{ kg/m}^3$, $\alpha = 200 \text{ µm}^{-1}$, A = 0.3. Laser: Gaussian F(r)= $F_o \exp -(r^2/a^2)$, $P(t) = P \sin^2(\pi t/2\tau)$ - pulse of $\tau = 4 \text{ ns}$, $\lambda = 266 \text{ nm}$, $F = 100 \text{ J/cm}^2$;



Heating temperature $\Delta T_{max} = f(a)$ for given $F = 100 \text{ J/cm}^2$, Si -bulk



Analytical solutions in linear approximation



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«Ablation» threshold *F_{th}* (J/cm²)





Normalized curves T(r) for different k (W/m K) (thermo conductivity coefficient)





Normalized curves T(r) for different laser pulse duration τ

Si, $\lambda = 266$ nm, $\alpha^{-1} = 5$ nm, $D_T = 0.8$ cm²/s, $L_T = 20$ nm ÷ 6.0 µm, a = 10 nm



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Two-temperature equations for fs laser heating of metals



$$c(T)\frac{\partial T}{\partial t} = div(k(T)\nabla T) + Q(t,r,z)$$



$$c_{e} \frac{\partial T_{e}}{\partial t} = div(k_{e} \nabla T_{e}) - \beta(T_{e} - T_{i}) + Q$$
$$c_{i} \frac{\partial T_{i}}{\partial t} = div(k_{i} \nabla T_{i}) + \beta(T_{e} - T_{i})$$

$$Q(t,r,z) = I(t,r)A\alpha \exp(-\alpha z)$$

or

Q – from near-field calculation

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Two-temperature "3D+t" heating model, Au

 $\Delta T_e (t, r = 0, z = 0) = f(t) \quad \text{and} \quad \Delta T_i (t, r = 0, z = 0) = f(t)$ $F = 100 \text{ mJ/cm}^2$



 $\Delta T_i max (t = \tau_{ei} \approx 100 \text{ ps}) \sim a (nm)$

Two-temperature "3D+t" heating model, Au

 $\Delta T_{e \text{ or } i} (t = t_n, r, z = 0) = f(r)$



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Two-temperature "3D+t" heating model, Au

 $\Delta T_{e \text{ or } i} (t = \overline{t_n, r} = 0, z) = f(z)$



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- 1) Nanometric laser ablation is possible with ns laser pulses when $L_T >> a$;
- 2) "Ablation threshold" F_{th} depends on laser spot diameter *a* for *a* < 1 μm : $F_{th} \sim 1/a$;
- 3) Low dependence of $T(r, z)/T_{max}$ on k (W/m K) for 3D heating regime;
- 4) Low dependence of $T(r, z)/T_{max}$ on $\tau(s)$ for 3D heating regime;
- **5)** T_{max} is for t $\approx \tau$ ei for fs heationf of metals:

[****] A. Semerok, S.V. Fomichev, C. Jabbour, J.-L. Lacour, M. Tabarant, F. Chartier, **Multi-parametric modeling of solid sample heating by nanosecond laser pulses in application for nano-ablation**, Oral presentation on LIBS'2016 conference, Chamonix; France; was sent in Spectrochimica Acta Part B.



Thank you for your attention



ABSTRACT

In sample microanalysis by laser ablation (LA), spatial resolution is determined by laser beam diffraction limits (of the order of a laser wavelength) and thermal diffusion of a deposited heating energy during laser pulse (proportional to the square root of pulse duration and matter diffusivity). Being limited by these laser beam features and those of heating energy, spatial resolution (a crater diameter) of \approx 1 µm was obtained with 4 ns laser pulses on 266 nm wavelength. To improve spatial resolution of microanalysis, the application of lasers with the pulses of shorter durations (ps and fs) may be advised. As another way to improve spatial resolution of microanalysis up to \approx 100 nm, one may advise LA with a highly localized laser field created by a tip near-field enhancement. The experiments with ns laser pulses were made along with multi-parametric theoretical studies based on one-temperature heating model.

In this work, the theoretical studies were extended on ultra-short laser pulses (ps or fs) to analyze the effect of pulse duration and matter properties (absorption coefficient, thermal conductivity and capacity) on the resulting temperature field spatial distribution T (t, x, y, z). A two-temperature model was applied for T (t, x, y, z) calculations. The results of these simulations are compared to temperature distributions for ns laser pulses. Discussion on advantages of ultra-short pulses application for LA with a tip near-field enhancement for consecutive chemical analysis with nanometric resolution will be presented.