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Deterministic model of PWR fast fluence for uncertainty propagations with the code APOLLO3[®]

Pietro Mosca¹, Laura Clouvel², Théophile Lacaze², Sébastien Lahaye², Jean Jacques Lautard² and Anne Marie Baudron²

ABSTRACT

The fast neutron fluence and the corresponding uncertainty are important parameters for reactor pressure vessel life time. This article presents one model, under development at CEA (Commissariat à l'Energie Atomique et aux Energies Alternatives), to carry out with the deterministic code APOLLO3[®] uncertainty calculations of the fast fluence for PWR irradiation surveillance. All calculations are made by MINARET, a 3D-SN solver of the APOLLO3[®] code which uses the discontinuous Galerkin finite elements approximation. The spatial mesh is unstructured and the transport calculations are parallelized with respect to the angular directions. In this numerical scheme, the multigroup cross-sections are sub-group self-shielded and collapsed over a dedicated energy mesh optimized by the Adaptive Energy Mesh Constructor (AEMC). Results from this model are encouraging with respect to the Monte Carlo reference TRIPOLI-4[®]. The integral of the flux over 1 MeV in the locations of interest (surveillance capsule and vessel) is calculated in less than 20 minutes with an error lower than 1%. Some examples of uncertainty calculations associated to design parameters in which the MINARET solver is coupled to the CEA uncertainty and sensitivity platform URANIE are also provided.

Keywords

RPV, NEUTRON FLUENCE, UNCERTAINTY, APOLLO3[®], URANIE, AEMC.

INTRODUCTION

The fast neutron fluence and the corresponding uncertainties are important parameters for the reactor pressure vessel life time. Usually for this kind of studies, reference tools are based on Monte Carlo codes. However, the CPU time required is still not reasonable for uncertainty and sensitivity analyses. In this perspective, this article presents the model, under development at CEA (Commissariat à l'Energie Atomique et aux Energies Alternatives), to carry out with the deterministic code APOLLO3[®] [1] uncertainty calculations of the fast fluence for PWR irradiation surveillance.

The new numerical scheme is based on the 3D-SN solver MINARET [2] of the APOLLO3[®] code which uses the discontinuous Galerkin finite elements approximation. The spatial mesh is unstructured and the transport calculations are parallelized with respect to the angular directions. In this scheme, the multigroup cross-sections are sub-group self-shielded and collapsed over a dedicated energy mesh optimized by the Adaptive Energy Mesh Constructor (AEMC) [3].

The structure of the paper is as follow. First the choice of the multigroup library and the

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optimization of the energy mesh to collapse cross-sections are described. Second the validation of the deterministic fast fluence scheme respecting the Monte Carlo reference TRIPOLI-4[®] [4] is shown on a three-dimensional PWR geometry. Then an example of uncertainty propagation coupling MINARET to the uncertainty and sensitivity platform URANIE [5] is presented and results are discussed.

VALIDATION OF THE MULTIGROUP LIBRARY

As a first step for the development of the deterministic fast fluence scheme, the choice of the multigroup library for the production of the core collapsed cross-section library is done. The considered libraries are mainly based on the JEFF3.1.1 evaluation and are declined over optimized 300, 600, 1200 group energy meshes by AEMC for fast reactor systems [3], [6]. A simplified slab model, which corresponds to a section cut of the reactor going from the core peripheral to the inner part of the vessel, is used to evaluate the accuracy of these different multigroup libraries. Vacuum boundary conditions are applied on the right end side and an external source distributed in energy according to the fissile spectra of the core is set on the left end side. The transport calculations are carried out with the APOLLO2 [7] discrete ordinates solver with a *S*32 angular discretization and a *P*5 approximation for the scattering cross sections. Cross-sections of all the light nuclides of the steel are subgroup self-shielded and differentiated per media of the slab. Table 1 shows the accuracy of the different deterministic solutions with respect to the reference TRIPOLI-4[®] using as discriminant parameter the total flux over 1 MeV, noted $\Phi_{1\text{MeV}}$ in this work. One can notice that passing from 300 to 600 group library the accuracy is almost increased of a factor two while a further energy mesh refining does not much more improve the accuracy of the solution. Fig. 1 shows that the corresponding multigroup fluxes in the vessel are within TRIPOLI-4[®] error bars in the energy range between 1 and 20 MeV. These results validate the use of the 600-group library for the production of the core collapsed cross-section library.

Table 1 : Total flux over 1 MeV in a slab obtained running APOLLO2 with different multigroup libraries and compared to TRIPOLI-4[®] reference. A2 and T4 stand for APOLLO2 and TRIPOLI-4[®] respectively.

$1 - \Phi_{1\text{MeV}}^{\text{A2}} / \Phi_{1\text{MeV}}^{\text{T4}}$	300 groups	600 groups	1200 groups	3 σ (T4)
Capsule	-1.48%	-0.81%	-0.69%	0.38%
Vessel	-1.25%	-0.69%	-0.44%	1.20%

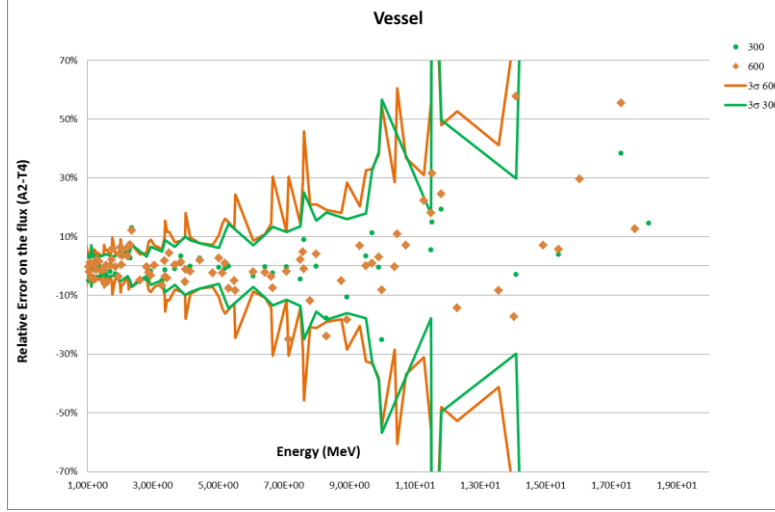


Figure 1 Multigroup flux over 1 MeV in the vessel medium of the slab obtained running APOLLO2 with 300 and 600-group libraries and compared to TRIPOLI-4[®] reference. A2 and T4 stand for APOLLO2 and TRIPOLI-4[®] respectively.

OPTIMIZATION OF THE ENERGY MESH FOR CROSS-SECTION COLLAPSING

Since the memory required by the 600 group library is too huge to be used in a three-dimensional fluence core calculation, a cross-section collapsing is mandatory. To address such a problem, the approach described in [8] has been adapted to the fluence context. The optimization is accomplished on the slab geometry of the previous section, considering a pure uranium and plutonium fission source. The energy mesh bounds are adapted by AEMC [3] in order to minimize the errors between two successive SN transports. The reference SN calculation is carried out over the fine 600-group energy mesh and the approximate one is obtained from collapsed cross-sections by flux homogenization over a coarser energy mesh.

To accomplish the optimization, it is defined the following functional that depends on the coarse mesh M_c of NG groups:

$$F(M_c) = \sum_{n=1}^N \alpha_n [F_n(M_c)]^2 + \sum_{k=1}^{NG} [G_k(M_c)]^2, \quad (1)$$

where $N = NI \times NM \times NR \times NG \times NC$ is the total number of components, NI is the total number of isotopes, NM is the total number of media, NR represents the total number of reactions ($NR = 2$ corresponding to absorption and scattering) and NC is the total number of problems to be considered in the optimization ($NC=2$ for the two types of sources). The $F_n(M_c)$ components represent the discrepancy between the reference isotopic reaction rate and its approximate value $T_{\rho,x,i}^{r,G}$, $T_{\rho,x,i}^{a,G}$ renormalized to the corresponding total isotopic reaction rates $T_{\rho,x,i}^r(M_c)$, $T_{\rho,x,i}^a(M_c)$:

$$F_n(M_c) = \left[\frac{T_{\rho,x,i}^{r,G}(M_c)}{T_{\rho,x,i}^r(M_c)} - \frac{T_{\rho,x,i}^{a,G}(M_c)}{T_{\rho,x,i}^a(M_c)} \right], \quad \alpha_n = \frac{T_{\rho,x,i}^r(M_c)}{T_{\rho,i}(M_c)}, \quad T_{\rho,i}^{r,a}(M_c) = \sum_x T_{\rho,x,i}^{r,a}(M_c), \quad (2)$$

$$1 \leq \rho \leq NR, 1 \leq G \leq NG, 1 \leq i \leq NI, 1 \leq x \leq NM \times NC.$$

The use of both the weights α_n and the renormalization is justified in [8]. The $G_k(M_c)$ components measure the discrepancy between the reference and the approximate leakage terms:

$$G_n(Mc) = \left[\frac{J^{r,G}(Mc)}{J^r(Mc)} - \frac{J^{a,G}(Mc)}{J^a(Mc)} \right]. \quad (3)$$

The 600-group library has 96 groups above 1 MeV, the remaining groups beneath, which are useless with regard to the problem being treated, are collapsed in a unique group. Indeed, the flux responsible for damaging the vessel is currently defined as the total neutron flux greater than 1 MeV. AEMC has therefore determined a coarse mesh consisting of a single group between 10^{-11} MeV and 1 MeV and (NG - 1) energy groups greater than 1 MeV (with NG < 96 the available number of groups).

Fig. 2 shows the dependency of the functional $F_n(Mc)$ with respect to the number of groups. One can notice that the accuracy does not improve itself over 19 groups. The same tendency has been verified comparing $\Phi_{1\text{MeV}}$, which are obtained with different optimized energy meshes, and the Monte Carlo reference.

These results lead to the use of the 19-group optimized energy mesh in the core fluence calculations. This energy mesh contains 1 group between 10^{-11} and 1 MeV and 18 groups between 1 and 20 MeV.

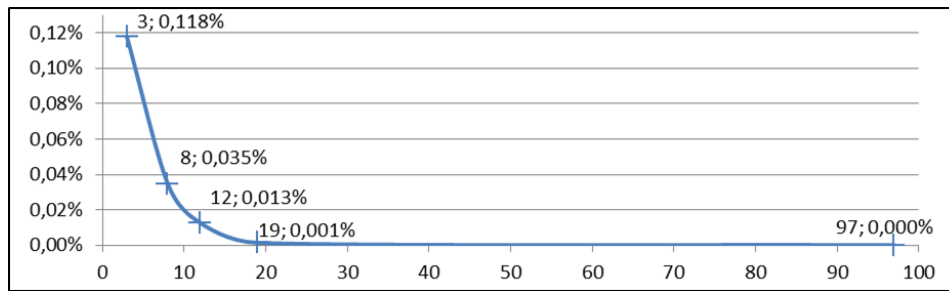


Figure 2 Functional value of the energy mesh optimization versus the number of groups

VALIDATION OF THE DETERMINISTIC FLUENCE MODEL

The validation of the deterministic fluence model is carried out modeling one eighth of a 900 MWe PWR core. The reflector geometry is obtained by axial extrusion of five radial basic geometries that model the baffles, the surveillance capsules, the stiffeners, the cold and the hot water part at the bottom and the top of the core (Fig. 3). The core is composed of 26 assemblies and each assembly is a standard 17x17 UO₂ homogeneous pin cells. The homogenization is justified by the fact that neutron mean free path in UO₂ fuel and in borated hot water is higher than the size of the cells at energies beyond 1 MeV. Only the four last peripheral assemblies of the core are taken into account in the transport calculation (Fig. 3). This simplification reduces of one third the CPU time preserving the accuracy of the solution because 99% of the fluence in the vessel and in the surveillance capsules is due to neutrons from the peripheral assemblies.

The power distribution in the core is an input datum which is directly converted into a neutron source radially homogenized on the fuel cells and discretized into 56 axial planes.

The axial temperature profiles of the water in the core and in the reflector correspond to standard operating conditions, elsewhere only average operating temperature values are considered. Since Doppler effect is negligible on light nuclei and also on heavy nuclei above the resonance energy range (energy higher than 300 keV), the water density is the only considered physical parameter to be modified by the temperature. All the cross-sections are calculated at the average operating

temperatures.

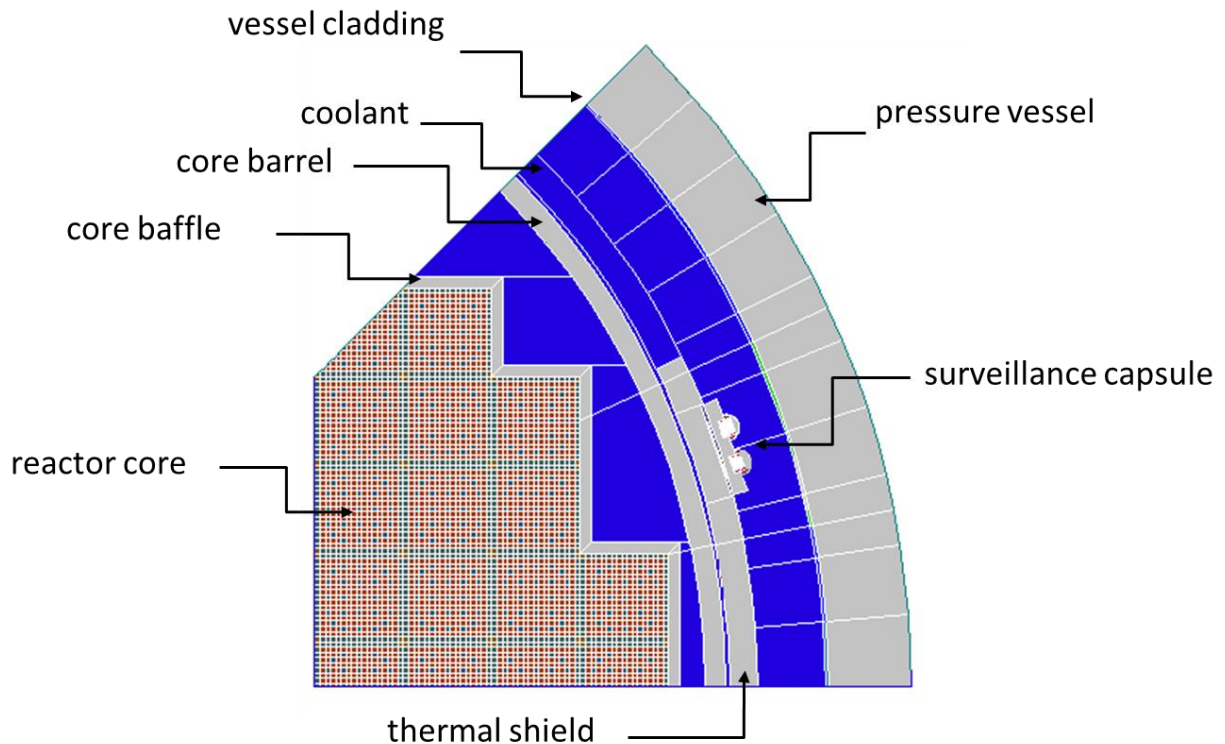


Figure 3: Two-dimensional geometry of the PWR core

All the transport calculations are carried out by the MINARET [2] 3D SN-solver of the APOLLO3[®] code. This solver uses the discontinuous Galerkin finite elements approximation, the spatial mesh is unstructured and the transport calculations are parallelized with respect to the angular directions. The solution in each cell can be radially and axially expanded up to a polynomial of order 1. The solution method is iterative and accelerated by DSA (Diffusion Synthetic Acceleration). The total FME (Finite Element Matrix) is solved by transport sweeps without spatial decomposition keeping only in memory the matrices of the DSA.

The multigroup cross-sections are obtained collapsing the sub-group self-shielded cross-sections of the 600-group library with the flux of the slab model over the optimized 19 group-energy mesh.

The MINARET calculations are carried out using a S32 angular discretization, a *P3* approximation for the scattering cross sections, a radial mesh size of 2 cm, an axial mesh size 10 cm. The radial and the axial polynomial order are set equal to 1 and 0 respectively and the relative error per group on the flux is imposed lower than 10^{-4} . These parameters are issued of a preliminary convergence analysis of the solution.

The reference calculation is carried out with the TRIPOLI-4[®] Monte Carlo code using a geometrical description without homogenization of the core. Multigroup and pointwise cross-sections are issued from the processing of the JEFF3.1.1 evaluation.

From Table 2 one can observe that relative error on the total flux over 1 MeV is lower than 1%. Fig. 4 shows that the multigroup flux discrepancy in the vessel is lower than 5% between 1 and

20 MeV. These results prove that the higher complexity of the 3D core test does not impact the accuracy of the solution because it is comparable to the slab test (cf. Fig. 1).

The APOLLO3[®] transport calculation took 20 minutes and required a total memory of 113 Gb with a MPI (Message Passing Interface) parallel computing on a LINUX CenOS6.4 cluster of 4 processors with 20 cores each. The corresponding TRIPOLI-4[®] reference calculations required 6 hours with the same demand of processor resources.

These results support the use of MINARET to perform uncertainty studies of neutron fluence using a Monte Carlo approach.

Table 2 : Total flux over 1 MeV for the 900 MWe PWR obtained with APOLLO3[®] and compared to the reference TRIPOLI-4[®]. A3 and T4 stand for APOLLO3[®] and TRIPOLI-4[®] respectively.

Medium	$1 - \Phi_{1\text{MeV}}^{A3} / \Phi_{1\text{MeV}}^{T4}$	3σ (T4)
Vessel	0.61 %	0.24%
Capsule 20°	0.86 %	0.09%
Capsule 17°	0.30 %	0.09%

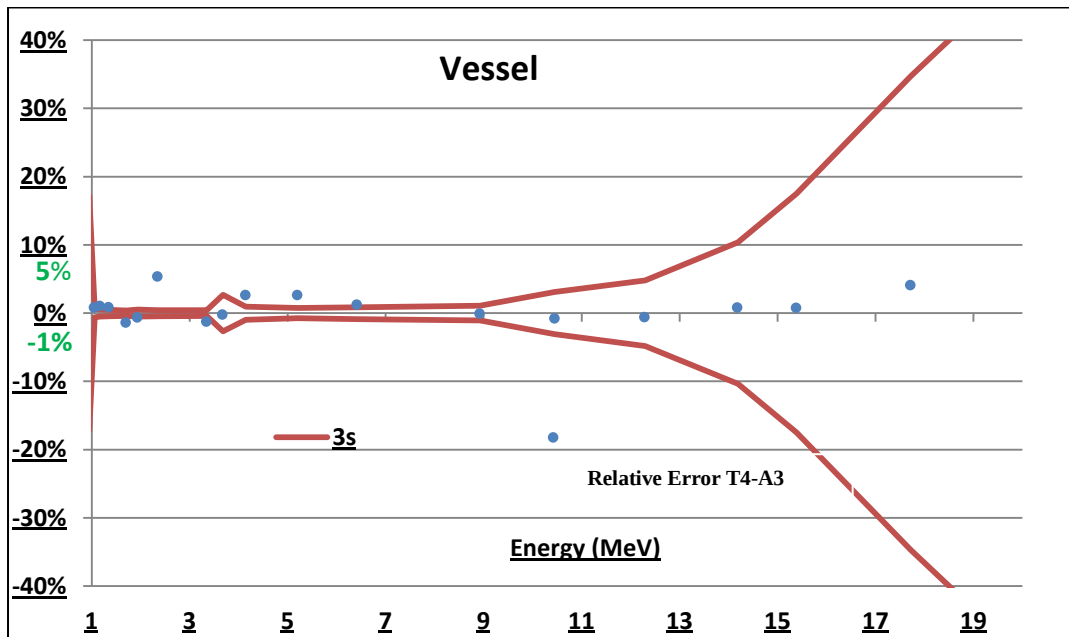


Figure 4 : Multigroup flux over 1 MeV in the vessel of 900 MWe PWR obtained with APOLLO3[®] and compared to the reference TRIPOLI-4[®]. A3 and T4 stand for APOLLO3[®] and TRIPOLI-4[®] respectively.

UNCERTAINTY PROPAGATION

The deterministic fluence scheme is used in order to lead a simplified uncertainty propagation to evaluate the impact of technological parameters (manufacturing tolerance of in-vessel components and water temperature). This study is carried out coupling the MINARET solver to the CEA uncertainty and sensitivity platform URANIE [5] which pilots the design of experiment (DOE).

Table 3 shows the perturbed parameters and the corresponding variations in the study. In our case, we consider that the variation of the uncertain variables follows a uniform law centered on the nominal value; that is, the realization of the geometric values in the finite set, defined by the uncertainties, is equally probable. For temperature sampling, the perturbation is the same at each axial step of the core, since total power of the reactor is not an uncertain variable of the study. The Latin hypercube sampling method is used by URANIE for the sampling. The DOE is composed of 2000 samples for each uncertainty variables in the study. The choice of the uniform distribution law is justified by the fact that only the extreme values for each uncertain variable are provided and any other laws would give supplementary and arbitrary information.

Table 3 : Uncertainty variables

Parameter	Perturbation
Azimuth	$\alpha^{\pm 0.33^\circ}$
Internal Vessel Radius	$R_{intves}^{\pm 5.8mm}$
Internal Envelope Radius	$R_{intenv}^{\pm 4mm}$
External Envelope Radius	$R_{extEnv}^{\pm 1mm}$
Internal Shield Radius	$R_{intsh}^{\pm 4.8mm}$
Capsule thickness	$E_{capthick}^{\pm 0.38mm}$
Water gap	$E_{watgap}^{\pm 0.5mm}$
Capsule position	$E_{cappos}^{\pm 0.2mm}$
Water temperature	$T_{wat}^{\pm 1^\circ}$

With URANIE a sensitivity analysis is also carried out, in order to evaluate if the variations of Φ_{1MeV} can be approached by a linear, additive or nonlinear form of the uncertainty variables, and also to determine the most influential ones in the point of interest.

Fig. 5 illustrates the impact of the nine variables on Φ_{1MeV} in the surveillance capsules and the vessel. One can observe that the relation between Φ_{1MeV} and the uncertainty variables is well fitted by a linear relation because the coefficient of determination R^2 is higher than 0.99. The sensitivities are thus simply proportional to the coefficient of the linear regression. Fig. 6 classes the sensitivities of the uncertainty variables in the surveillance capsules. In the 17° capsule, the angular position, the water temperature, the internal radius of the shield and of the envelope have more important impact on Φ_{1MeV} than the other parameters because of much higher range of variation. A similar analysis shows that the vessel is very sensitive to its internal radius, and much less to the internal radii of the envelope and the shield. It is almost unaffected by other uncertain variables because of its distance from the core and the shielding effect produced by the interposed water.

	Capsule 17°	Capsule 20°	Vessel Cladding	Vessel
Temperature				
Azmut				
Vessel Radius				
Int Envelope Radius				
Ext Envelope Radius				
Int Shield Radius				
Capsule thickness				
Water gap				
Capsule position				
R ²	0,99711	0,994609	0,999351	0,999507
Mean	1,82E+11	1,56E+11	2,50E+10	9,73E+09
Sigma	3,43E+09	1,84E+09	1,18E+09	3,44E+08
Sigma/Mean (%)	1,89	1,18	4,71	3,54
Nominal Value	1,82E+11	1,56E+11	2,50E+10	9,72E+09
Ecart (%)	-0,13	-0,11	-0,15	-0,11

Figure 5 : DOE for the total flux over 1 MeV in surveillance capsules, the vessel cladding and the vessel

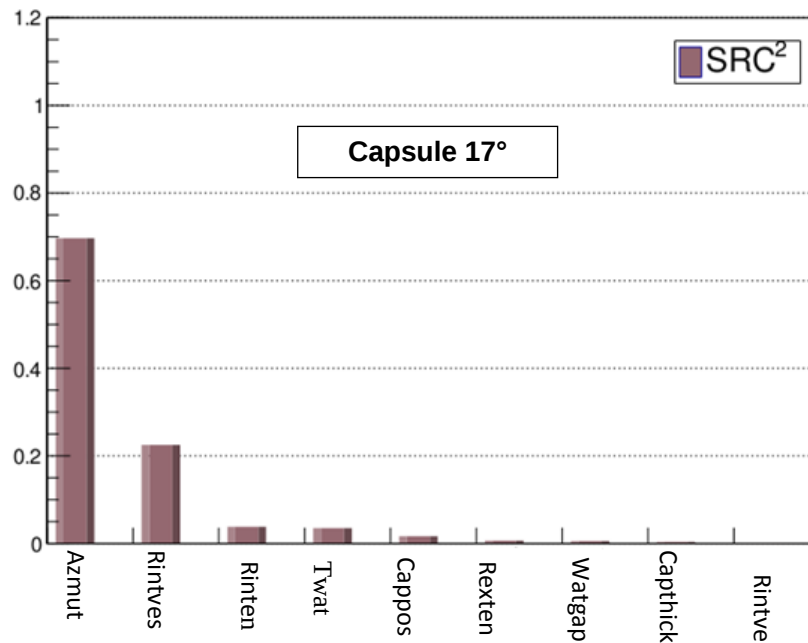


Figure 6 : Sensitivities of the total flux over 1 MeV in the surveillance capsule 17°

Fig. 7 shows the impact of the uncertain variables on the multigroup flux in the surveillance capsules. The angular position is anti-correlated with the multigroup flux because the capsule gets closer to the vessel point of highest flux (“hot spot”) when the angle decreases. The most impacted energetic range is between 1 and 3 MeV because by approaching the peripheral assemblies the weight of the fission spectrum increases (energy range between 1 and 3 MeV) and the probability is higher that neutrons with lower energy reach the capsule. The shield internal radius and the multigroup flux are positively correlated in the range between 2 and 20 MeV, because the slowing-down power of the steel decreases due to the reduction of its corresponding thickness. On the other hand, the anti-correlation below 3 MeV is due to a high neutron transfer below 1 MeV. This effect is produced by the increase of the water thickness and by the high cross-section value of the water in this energy range (Fig. 8).

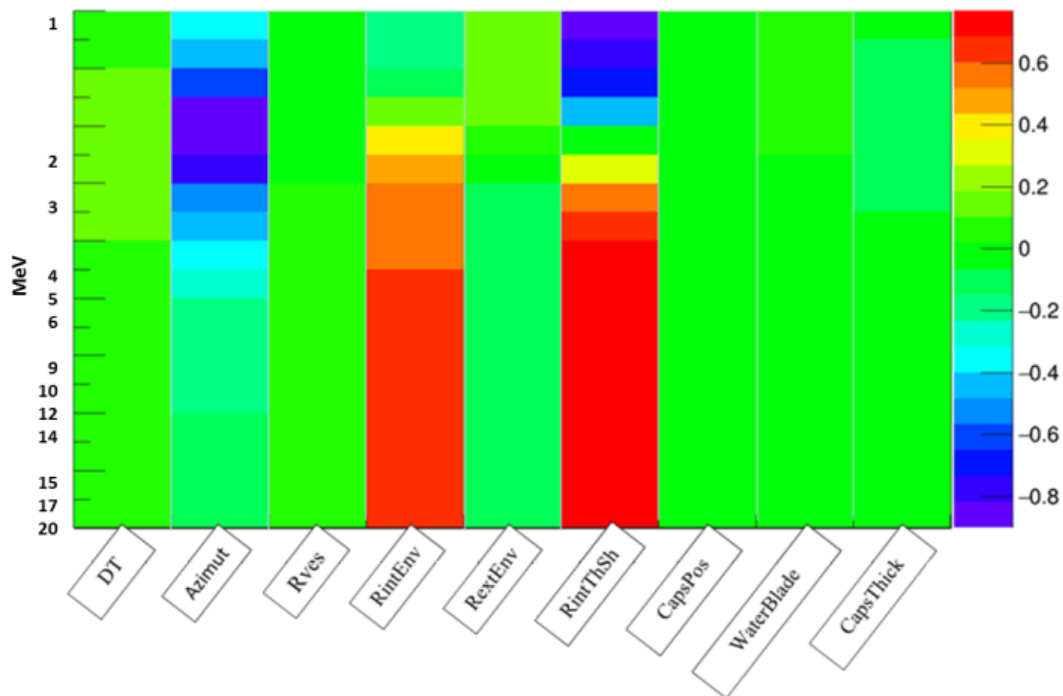


Figure 7: Correlation between variables and multigroup flux in the surveillance capsule.

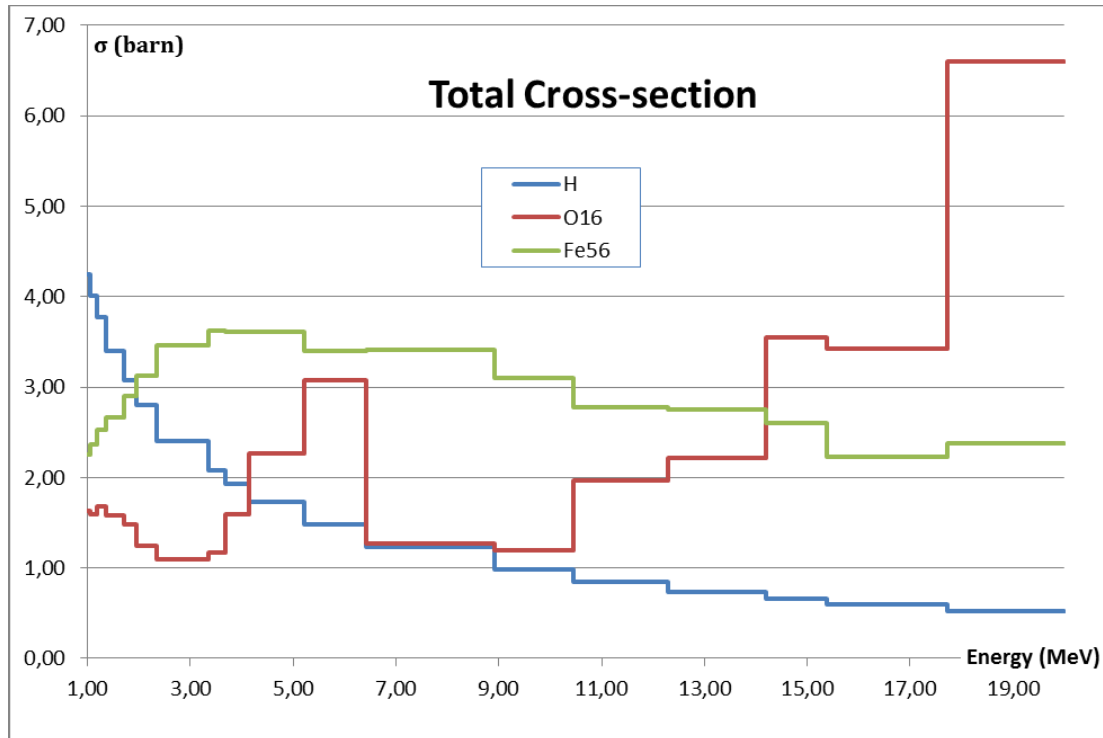


Figure 8: Multigroup total cross-sections of different isotopes

CONCLUSIONS

This paper shows the new deterministic fast fluence scheme with the APOLLO3[®] code developed at CEA.

The use of sub-group self-shielded cross-section collapsed over an optimized collapsing energy mesh by AEMC allows to carry out very precise deterministic calculations of neutron flux in the vessel. Results of this deterministic model are encouraging, the integral of the flux over 1 MeV in the location of interest is calculated in less than 20 minutes with an error lower than 1% with respect to the TRIPOLI-4[®] Monte Carlo reference.

An example of uncertainty propagation with the CEA uncertainty and sensitivity platform URANIE is also provided. This example offers an alternative approach to the classic first order sensitivity analysis to perform uncertainty studies like in [9]. It will be pursued in order to estimate the impact of nuclear data (cross-sections, fission spectra) and of other technological parameters (power distribution, power intensity ...).

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