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N. Finck, M. L Schlegel, A. Bauer. Structural iron in dioctahedral and trioctahedral smectites: A polarized XAS study. *Physics and Chemistry of Minerals*, 2015, 42 (10), pp.847-859. 10.1007/s00269-015-0768-3 . cea-02383813

HAL Id: cea-02383813

<https://cea.hal.science/cea-02383813>

Submitted on 28 Nov 2019

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Structural iron in dioctahedral and trioctahedral smectites: A polarized XAS study

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Abstract

The chemical form of structural Fe in smectites influences many physicochemical properties of these clay minerals. Powder EXAFS data for structural Fe in smectites have been reported, however, the preferred orientation of clay platelets with respect to the X-ray beam may lead to erroneous conclusions on the local chemical environment. Dioctahedral montmorillonite and for the first time trioctahedral hectorite were prepared as textured samples and the Fe local environment was probed by analysis of the X-ray absorption pre-edge peaks at the magic angle and by polarized EXAFS (P-EXAFS) spectroscopy. Compared to powder measurements, overlapping contributions from shells with distinct orientations can be filtered more easily by P-EXAFS, thus decreasing uncertainties on structural parameters.

The pre-edge spectrum for montmorillonite is similar to spectra commonly reported for dioctahedral smectites. In contrast, the pre-edge spectrum of hectorite is notably distinct, and hints to either differences in the site symmetry and/or in covalence. In both smectites, Fe is surrounded by a first O shell at a distance consistent with sixfold coordinated Fe(III), suggesting that Fe(III) is located in the smectite octahedral sheet. This is corroborated by the distances and orientations of neighboring cationic shells, such as in-plane (Mg,Al) and out-of-plane Si shells. For montmorillonite, the results indicate Fe substitution for Al/Mg in the octahedral sheet, and a number of Fe neighbors consistent with random distribution in the octahedral sheet. For hectorite, results indicate a slight tendency for Fe atoms to form pairs in octahedral sheets; however, low numbers of neighboring cations were obtained, presumably a consequence of the presence of vacancies and/or Li in the vicinity of Fe, or of the coexistence of Fe and Mg neighbors with mutually cancelling EXAFS waves. Consistent with pre-edge data, the coordination numbers can also indicate some incoherency in Fe-cation interatomic distances in hectorite as a consequence of site distortion. These results suggest that Fe³⁺ in hectorite locally distorts the structure of the trioctahedral phyllosilicate, and tends to aggregate charge-deficient (i.e. vacant or Li⁺-containing) octahedral sites.

Keywords

Montmorillonite, Hectorite, Iron, polarized XAS

Introduction

Smectites are widespread clay minerals forming under surface and subsurface conditions (Meunier, 2005) and widely occur in weathering formations and sediments (Güven, 1988). These small, flat minerals possess a large surface area and permanent layer charge. As a consequence, they can absorb and retain a variety of ions or molecules and thus they dominate the physicochemical properties in the systems in which they are present. Their high reactivity has also spurred important industrial applications such as catalysis (e.g., Swarnakar et al., 1996; Güven, 2009) and health (Williams et al., 2009). Smectites are also constituents of bentonites and confer to these materials cost-efficient mechanical, hydraulic and chemical properties. Bentonites are extensively used as constituents of engineered barriers for the isolation of pollutants, and could be used as backfill in high-level nuclear waste repositories (Gates et al., 2009). Finally smectites such as hectorite can form as secondary minerals during the alteration of nuclear glass containing high-level waste (Thien et al., 2010). These newly formed minerals can further react with hazardous radionuclides released by alteration of nuclear glass. Thus, it is important to understand the smectite crystallochemical properties which ultimately control their reactivity.

Smectites can exhibit a large variation in chemical compositions as a result of their versatile structure and of cationic substitution in crystal sites. For example, montmorillonite, the common aluminum smectite in soils and sediments, is made by the condensation of one octahedral sheet sandwiched between two tetrahedral sheets (Figure 1), forming a TOT structure. Two-third of the octahedral sites are occupied mostly by Al(III), forming a dioctahedral framework in which Al(III) can be substituted by cations such as Mg(II) or Fe(II,III) (Stucki, 1988). In contrast, in trioctahedral hectorite all octahedral sites are filled, mostly by Mg(II), which can be substituted by Li(I) or trace amounts of Fe(II,III). Not only the extent, but also the distribution of chemical substitution can affect the chemical reactivity of smectite (Mering and Glaeser, 1954). For example, reduction or oxidation of redox sensitive cations will modify the layer charge and thus alter properties such as swelling and/or cation exchange capacity (Gates et al., 1996; Gates et al., 1998), which in turn may affect the retention capacities of bentonites.

Structural Fe in clay minerals can significantly impact the redox properties in the environments where they are present. The reduction of Fe(III) in smectite can be achieved by microbial (e.g. Pentráková et al., 2013) and by chemical treatments (e.g. Komadel et al., 2006; Stucki, 2006). For example, ferruginous smectite can be reduced by dithionite within hours and the reduction proceeds from the basal surfaces rather than from the particle edges (Komadel et al., 2006). Fe(II) in clay minerals is able to reduce heavy metals (Bishop et al., 2014) and radionuclides (Bishop et al., 2011; Yang et al., 2012), thereby modifying their mobility and availability. In contrast, Fe(III) may act as terminal electron acceptor for iron-reducing bacteria, which provides a mechanism of structural Fe(II) (re)generation (Ernstsen et al., 1998). The magnitude and reversibility of these reactions largely depend on the Fe local environment in the octahedral sheet, i.e., the formation of a solid solution or the occurrence of Fe-rich clusters (Drits and Manceau, 2000) which can be controlled in part by the smectite structure. For example, in dioctahedral smectites, two octahedral sites of distinct sizes can be identified, and Fe(III) would preferentially enter the largest one (Tsipursky and Drits, 1984), and thus be relatively dispersed. In contrast, preferential clustering of Fe was observed in trioctahedral micas, i.e., non-swelling phyllosilicates with TOT layers as in smectites (Manceau et al., 1990). This suggests that Fe distribution may differ depending on the

octahedral sheet framework (dioctahedral or trioctahedral) of smectites. Unfortunately, no study reported unambiguous evidence for random Fe distribution in trioctahedral smectites so far.

X-ray absorption spectroscopy (XAS), and more specifically extended X-ray absorption fine structure (EXAFS) spectroscopy, have proven to be reliable techniques to probe the molecular environment of transition metals. Powder EXAFS spectroscopy was used to investigate the Fe environment in montmorillonite (SWy-1) (Vantelon et al., 2003). No Fe-Fe pair was detected, suggesting that the structure has an ordered Fe distribution obeying an exclusion rule in the octahedral sheet. The spectroscopic analysis of such powder spectra is however complicated by textural effects which are difficult to avoid during sample preparation of minerals having a layered structure (Manceau, 1990b). Powder EXAFS spectra for such anisotropic samples should be recorded with the sample plane oriented at the magic angle (at which any angular dependence is extinguished), otherwise XAS analysis will provide an inaccurate description of the local environment. Unfortunately, information concerning sample preparation and orientation with respect to the in-coming X-ray beam is most of the time missing. Additionally, EXAFS analysis for such samples is complicated by interferences between EXAFS waves scattered by neighboring shells located at similar distances from the absorber. This complication can be overcome by polarized EXAFS (P-EXAFS) spectroscopy which can discriminate the contributions of absorber-backscatterer pairs with distinct crystallographic orientations. This is of high interest in investigations on clay minerals where cationic shells from the tetrahedral and octahedral sheets are separated by $< 0.25 \text{ \AA}$. The theoretical background and principle of P-EXAFS applied to smectite preparation is well documented (e.g. Manceau et al., 1998; Schlegel et al., 1999; Manceau et al., 2000a). Thus P-EXAFS can provide unique information on the nature, number, distance and orientation of neighboring shells. For Fe, this information can be complemented by the analysis of pre-edge peaks which are related to quadrupolar and (forbidden) dipolar transitions from the $1s$ to d orbitals. The position and amplitude of these peaks is known to be sensitive to the oxidation state and the site symmetry of Fe (Manceau and Gates, 1997; Wilke et al., 2001). The combination of pre-edge peaks and P-EXAFS spectroscopy was powerful in determining the Fe short-range environment in Fe-rich smectites such as nontronites (Manceau et al., 2000a; Manceau et al., 2000b; Gates et al., 2002).

In this study, the Fe chemical state and distribution in dioctahedral montmorillonite (SWy-1) and trioctahedral hectorite (SHCa-1) were investigated by XAS. These smectites are taken as representative of dioctahedral and trioctahedral clay minerals. In addition, these minerals are also model systems for material present in the backfill of a nuclear waste repository site and as a secondary phase frequently detected in alteration experiments of nuclear glass, respectively. As such, thorough determination of their crystallochemical properties is useful to improve the robustness of long-term models of waste confinement. In this study, oriented self-supporting films were prepared in order to probe the Fe K -edge by P-EXAFS spectroscopy for SWy-1 and, for the first time, SHCa-1. Also, P-XAS data collected and analyzed to understand how Fe can affect the local structural properties of dioctahedral or trioctahedral smectites.

Experimental

Samples

Montmorillonite (SWy-1) and hectorite (SHCa-1) were purchased from the Source Clays Repository of the Source Clay Mineral Project (SCMP). Montmorillonite was purified according to a standard procedure

described in (Golubev et al., 2006) and the fraction $< 1 \mu\text{m}$ was isolated. For hectorite the fraction $< 2 \mu\text{m}$ was isolated by sedimentation and purified following the procedure described in (Schlegel et al., 1999). The purified montmorillonite and hectorite fractions were saturated with Na. The elemental compositions of the purified fractioned samples were determined by acid digestion and ICP-MS (Thermo X-Series II) analysis of the resulting solutions. The results are presented in Table 1. Besides the elements listed in Table 1, hectorite also contains large amounts of fluorine (2.7 wt% for the SCMP hectorite, Thomas et al., 1977), but this element was not quantified because the analysis involved acid digestion using concentrated HF.

To remove possible X-ray amorphous ferric oxides during the purification, hectorite underwent a treatment with dithionite-citrate-bicarbonate that imposes strongly reducing conditions, followed by treatment with hydrogen peroxide that imposes strongly oxidizing conditions. With dithionite, a complete reduction of structural Fe from a trivalent to a divalent oxidation state can be achieved in smectites, such as reported for the Fe-rich SWa-1 (Gorski et al., 2012) and for the montmorillonite SWy-2 (Neumann et al., 2011), and subsequent re-oxidation can be achieved with hydrogen peroxide (Gorski et al., 2013). Furthermore, reported spectroscopic data showed that structural changes in the Fe coordination environment and in the silicate lattice can be caused by the Fe reduction and that these are not fully reversible (e.g., SWy-2 and SWa-1, (Neumann et al., 2011; Gorski et al., 2013)). Similarly to the reported montmorillonite data, Fe in hectorite may also have underwent a reduction–oxidation cycle, resulting in partly irreversible structural modifications. However, our infrared spectra do not reveal any strong perturbation of the hectorite lattice, certainly due to the low Fe content.

Purified Na-saturated smectites were first characterized by powder X-ray diffraction and Fourier-transform infrared (FTIR) spectroscopy. The X-ray diffractograms were collected on oriented samples with a D8 Advance (Bruker) diffractometer (Cu K_{α} radiation) equipped with an energy dispersive detector (Sol-X). The FTIR spectra were obtained with a Bruker IFS 55 spectrometer, equipped with an ATR (Attenuated Total Reflectance) cell and an MCT (Mercury Cadmium Telluride) detector. Self-standing films were prepared by slow filtration of the clay suspensions on a $0.025 \mu\text{m}$ pore size filter (Millipore) (Schlegel et al., 1999). This protocol readily provided highly textured self-supporting films to perform polarized XAS experiments, with a and b layer axes of the crystallites randomly oriented in the plane of the film.

X-ray absorption spectroscopy (XAS)

Polarized Fe K -edge XAS data of the oriented clay samples were collected at the FAME beamline (Proux et al., 2005) at the ESRF (Grenoble, France) with a ring energy of 6 GeV and a ring current of 50-90 mA. The data were collected in fluorescence-yield detection mode using a 30-element Ge solid state detector (Canberra) at angles (α) between the electric field of the X-ray beam and the clay layer plane of 10, 35, 55 and 80°. The energy calibration was done by setting the first inflection point of the Fe K -edge of an iron foil at 7112.0 eV.

The pre-edge spectra were modeled by pseudo-Voigt line shapes (Westre et al., 1997) after subtraction of the baseline with an exponential function. The energy position, the full width at half-maximum and the peak height were fit using a fixed 50:50 Gaussian:Lorentzian ratio for the peaks. The pre-peak intensity was calculated as the integrated areas of the pseudo-Voigt functions used for peak modeling.

Analysis of the EXAFS data was performed following standard procedures by using the Athena and Artemis interfaces to the Iffit software (Ravel and Newville, 2005). The spectra were extracted from the raw data and the Fourier transforms (FTs) were obtained from the $k^3 \times \chi(k)$ functions (3.2 – 10.2 Å⁻¹ for both samples at all angles). The spectral data were fit in R space (1.3 – 3.8 Å for both samples at all angles) using phase and amplitude functions calculated with feff8.4 (Ankudinov et al., 1998). The amplitude reduction factor (S_0^2) was set to 0.70, a value well suited to properly fit the data of α -Fe (data not shown). In order to generate theoretical paths, the published structures of montmorillonite (Tsipursky and Drits, 1984) and hectorite (Breu et al., 2003) were used with the absorber (Fe) located at octahedral position. Since the differences in backscattering amplitude between neighboring Mg, Al and Si are limited, the XAS fitting models were simplified. In montmorillonite, octahedral and tetrahedral contributions were fit by using only Al and Si backscatterers, respectively. In hectorite, only Mg (octahedral) and Si (tetrahedral) contributions were considered. Similarly, O and F only differ by $Z \pm 1$ so that F neighbors in hectorite can be fitted with EXAFS phase and amplitude functions for O with only minimal discrepancy. Thus the first coordination shell of Fe in hectorite can be considered to be made of O atoms only for fitting purposes.

For a given film, the data were fit simultaneously at all angles using a single value of ΔE , and for a given shell a common bond length and mean square displacement (“Debye-Waller” factor). The Fe:Al (montmorillonite) and Fe:Mg (hectorite) ratios were kept equal at all angles. The uncertainties are typically ± 0.02 Å on the distances for well-resolved atomic shells and ± 20 % on the coordination numbers. The experimental uncertainty on α is estimated to $\pm 1^\circ$. The fit quality was quantified by the R_f factor representing the absolute misfit between theory and data (Ravel, 2000).

Results and discussions

X-ray diffraction and infrared spectroscopy

X-ray diffractograms for the two smectites are presented in Figure 2. The basal spacing for montmorillonite ($d(001) = 12.2(2)$ Å) and hectorite ($d(001) = 14.8(2)$ Å) are typical of smectite with one or two interlayer water molecules, respectively (Meunier, 2005). Quartz (peak at $26.6^\circ 2\theta$) is also present in small quantities in montmorillonite, but could not be detected in hectorite. Quartz does not contain Fe and thus will not affect XAS measurements. No other crystalline phase could be detected and the flat backgrounds of the diffractograms indicate the absence of additional X-ray amorphous phase.

The IR spectrum of montmorillonite (Figure 2) matches reported data (Madejova and Komadel, 2001), with an –OH stretching band of structural hydroxyl groups at 3623 cm^{-1} and Si–O stretching bands at 989 cm^{-1} and 1116 cm^{-1} (longitudinal mode). The spectrum also contains AlAl–OH, AlFe(III)–OH and AlMg–OH deformation bands at 914 cm^{-1} , 882 cm^{-1} and 848 cm^{-1} , respectively. The spectrum of hectorite also matches reported data (Madejova and Komadel, 2001), with –OH stretching bands of structural hydroxyl groups (3674 cm^{-1}) and bonded water (3621 cm^{-1}). The presence of adsorbed water is further indicated by two broad bands at 3380 cm^{-1} (–OH stretching) and $\sim 3265\text{ cm}^{-1}$ (overtone of –OH deformation of water). This finding indicates that the interlayer of hectorite contains more water than that of montmorillonite for which these bands are absent, what is consistent with 2-layer and 1-layer hydrate state, respectively, determined by XRD. In hectorite, the bands at 960 cm^{-1} and 689 cm^{-1} correspond to Si–O stretching and the band at 651 cm^{-1} to Mg₃–OH deformation.

The trioctahedral character of this smectite is revealed by the $\text{Mg}_3\text{-OH}$ band and the position of the structural hydroxyl group. The Si–O stretching bands at 780 cm^{-1} and 797 cm^{-1} can be related to the presence of quartz in both samples (Madejova and Komadel, 2001). The two bands are of similar intensity in hectorite, but not in montmorillonite, possibly because the band at 797 cm^{-1} in montmorillonite can also contain contribution from the Fe(III)Mg–OH bending mode (Gates, 2008). No other phases could be detected.

The XRD and IR data indicate that the smectite preparations were free from Fe-containing ancillary phases, attesting to the success of the purification procedures. Also, –OH bands where the hydroxyl groups are connected to Fe octahedra were detected in the IR spectrum of montmorillonite, but not in hectorite, probably due to the low Fe content of this mineral. FeFe–OH deformation bands were absent in the two spectra, excluding Fe clustering and thus hinting at a random Fe distribution, at least in montmorillonite.

Pre-edge spectroscopy

The Fe *K*-edge XAS spectra have a weak pre-edge feature $\sim 15\text{ eV}$ below the main absorption edge arising predominantly from $1s \rightarrow 3d$ (quadrupolar) transitions. The $1s \rightarrow 3d$ transition is electric dipole-forbidden by parity considerations in a centrosymmetric octahedral environment (e.g., O_h symmetry) (Westre et al., 1997). However, a very weak pre-edge feature, split into t_{2g} - and e_g -like components, is still experimentally observed, the intensity depending on the local symmetry and on the cation electronic properties. The pre-peak intensity can thus be used to evaluate the Fe oxidation state and coordination geometry (Manceau and Gates, 1997; Manceau et al., 2000a; Wilke et al., 2001). For example, a higher pre-edge amplitude is observed for fourfold coordinated Fe compared to sixfold coordinated Fe, a consequence of mixing $3d$ and $4p$ atomic orbitals in T_d symmetry (Westre et al., 1997).

The pre-edge is located at around 7114.5 eV for both montmorillonite and hectorite (Figure 3), indicating the Fe is in an oxidation state of +III in both clays. The intensity of the pre-edges is weak ($\sim 1.5\%$ of the main edge), suggesting that Fe atoms are predominantly sixfold coordinated. The pre-edge features were modeled with pseudo-Voigt line shapes (Westre et al., 1997) to gain additional information on their amplitude and position (Figure 4). At $\alpha = 35^\circ$ (polarized and powder XAS data are identical at this angle), the difference in energy (Table 2) of the components for montmorillonite (1.49 eV) compared to hectorite (1.45 eV) is below the uncertainty associated to the position of the features ($\sim 0.1\text{ eV}$). The total pre-peak intensities (i.e., total integrated areas, Table 2) equal 0.097 and 0.094 for montmorillonite and hectorite, respectively. These low values compare with those obtained for sixfold coordinated Fe(III) (Wilke et al., 2001). Also, the total pre-peak intensities are similar in both smectites. Interestingly, the ratio between the intensity of the contributions at lower to higher energies equals 1.31 for montmorillonite and 1.00 for hectorite. The intensity ratio of montmorillonite is close to expectation for octahedral high-spin ferric Fe (3:2) (Westre et al., 1997). The lower value for hectorite may best be explained by site distortion for Fe substituting for Mg/Li in the trioctahedral smectite. However, the low pre-peak intensity compared to the main edge suggested that the site is centrosymmetric and that such site distortion is actually limited. Alternatively, the observed difference may be explained by distinct covalence of the chemical bonds between Fe and the surrounding ligands (Westre et al., 1997). The e_g set of the $3d$ orbitals is more covalent than the t_{2g} set due to σ bonding with the ligands. Hectorite contains significant amount of fluorine substituting for OH groups, and because F has an electronegativity higher than O, the observed

differences may also be attributed to Fe binding to F in hectorite. Regardless of the correct interpretation, these dissimilarities in pre-peak intensities indicate slight differences in Fe local environments.

The pre-peak intensity ratios do not exhibit a monotonic trend with α , consistent with previous observations on Fe-rich phyllosilicate {Dyar, 2001 #74}. For montmorillonite, the ratio increases from 10° to 55° and then decreases whereas for hectorite it increases from 10° to 35° and then decreases. Since a dipole transition would be dichroic in nature, this complex angular behavior suggests that the transition is predominantly quadrupole (Hahn et al., 1982). The weakness of the dipole contribution corroborates the predominantly octahedral nature of the coordination polyhedra {Munoz, 2013 #75}.

EXAFS spectroscopy

EXAFS spectra. Both sets of P-EXAFS spectra exhibit a significant angular dependence over the whole k range and well-defined isosbestic points are observed, attesting to the high degree of particle orientations in the self-supporting films (Figure 5). The variation in amplitude and position of the maxima at $\sim 4.2, 5.3, 6.5, 8.0$ and 9.1 \AA^{-1} are consistent with the presence of atomic shells with distinct orientations, e.g., in-plane Mg/Al/Fe and out-of-plane Si atoms. The large dependence on polarization is also consistent with the successful preparation of highly oriented clay films.

The powder EXAFS spectra recorded at $\alpha = 35^\circ$ for montmorillonite and hectorite show subtle differences in the position and amplitude of oscillation maxima, e.g., at $k \sim 5.0, 6.3$ and 8.0 \AA^{-1} (Figure 3). The presence of distinct frequencies suggests that Fe has distinct molecular environments in terms of nature and number of backscattering atoms, consistent with the dioctahedral nature of montmorillonite and the trioctahedral nature of hectorite. This is also clearly seen on the corresponding Fourier transforms (FT) where the contributions from the cationic shells differ.

Fourier transforms. The FTs contain several peaks (Figure 5). Peak A ($R + \Delta R \sim 1.6 \text{ \AA}$) corresponds to the first oxygen shell (O1) and its magnitude is higher in montmorillonite than in hectorite. Best fits (Figures 6 and 7) were obtained for Fe–O1 interatomic distances of $R_{\text{Fe-O1}} = 2.01(1) \text{ \AA}$ in montmorillonite and $2.00(1) \text{ \AA}$ in hectorite (Tables 3 and 4). These bond lengths are characteristic of sixfold-coordinated Fe(III) and match reported values for Fe located at clay octahedral sites (Manceau et al., 1998; Manceau et al., 2000a; Vantelon et al., 2003). The greater amplitude of peak A for montmorillonite correlates with a greater number of detected O1 neighbors, e.g., at $\alpha = 35^\circ$, $N_{\text{O1}} = 5.8(5)$ for montmorillonite and $5.1(5)$ for hectorite (Tables 3 and 4). The slightly lower N_{O1} in hectorite can be related to a larger structural disorder resulting in a broader distribution of Fe–(O,F) bond lengths (Manceau et al., 2000a) around a mean value, which is consistent with the pre-edge results.

The amplitude of peak A decreases for both smectites, and the decrease is more pronounced in montmorillonite than in hectorite (Figure 5). Fits (Figures 6 and 7) indicate that N_{O1} decreases from $6.3(5)$ at $\alpha = 10^\circ$ to $4.9(4)$ at $\alpha = 80^\circ$ for montmorillonite and from $5.4(5)$ to $4.6(4)$ for hectorite (Tables 3 and 4). Using the relationship between N as a function of α and the angle β between the X-ray absorber-backscatterer pair and the normal to the film plane (Schlegel et al., 1999), the orientation of the O1 shell can be estimated. The calculated values of β_{O1} equal $58.1(1.5)^\circ$ and $56.9(7)^\circ$ for montmorillonite and hectorite, respectively. Though the value

obtained for montmorillonite deviates significantly from the value expected for fully symmetric octahedra (54.7°), they are consistent with the octahedral flattening observed in most hydrous TOT phyllosilicates (Güven, 1988).

Peak B ($R + \Delta R \sim 2.75 \text{ \AA}$) corresponds to the sum of contributions from the nearest octahedral (Mg,Al,Fe) shell (montmorillonite) or (Mg,Li) shell (hectorite), the nearest tetrahedral Si shell and the next-nearest O shell. The position, amplitude and angular dependences of this peak depend on the interferences between the EXAFS contributions from the octahedral and tetrahedral sheets. For example, peak B for montmorillonite displays a complex polarization dependence: its amplitude decreases from $\alpha = 10^\circ$ to $\alpha = 35^\circ$ and then increases with α , and the position and imaginary parts shift to higher $R + \Delta R$ values with α (Figure 5). In smectites, this angular behavior is diagnostic of in-plane predominant contributions from (Mg,Al) cations in the octahedral sheet, and out-of-plane contributions from tetrahedral Si (Schlegel et al., 1999). This peak was fit assuming Fe, Al and Si contributions at $R_{\text{Fe-Fe1}} = 3.02(1) \text{ \AA}$, $R_{\text{Fe-Al1}} = 3.04(1) \text{ \AA}$ and $R_{\text{Fe-Si1}} = 3.22(1) \text{ \AA}$, respectively (Table 3). Best-fit bond distances match values obtained from diffraction data (Tsipursky and Drits, 1984). For increasing α values, the coordination numbers decrease from 0.5(1) to 0.1(1) for the Fe shell, from 3.4(3) to 0.6(1) for the Al shell, and increase from 1.4(7) to 6.6(7) for the Si shell (Table 3). These angular dependences are consistent with Al and Fe shells oriented predominantly in the in-plane dimension while the Si shell is oriented in the out-of-plane direction. The values number of neighboring atoms obtained at $\alpha = 35^\circ$ ($N_{\text{Fe1}} = 0.4(1)$, $N_{\text{Al1}} = 2.6(3)$ and $N_{\text{Si1}} = 3.4(3)$) are in reasonable agreement with coordination numbers for montmorillonite (3 (Al,Fe) neighbors in-plane and 4 (Si,Al) neighbors out-of-plane). The fit results indicate a small but significant contribution from neighboring Fe and thus the presence of Fe-Fe pairs. This finding is based on the ability of P-EXAFS to filter out overlapping contributions from shells with distinct orientations, thereby decreasing uncertainties on structural parameters. This can explain, at least partly, why our results slightly differ from earlier conclusions based on the analysis of powder EXAFS spectra (Vantelon et al., 2003).

In contrast to montmorillonite, the amplitude of peak B in hectorite increases with α and its position varies only slightly (Figure 5). Peak B was modelled considering Mg, Fe and Si shells at $R_{\text{Fe-Mg1}} = 3.03(1) \text{ \AA}$, $R_{\text{Fe-Fe1}} = 3.03(3) \text{ \AA}$ and $R_{\text{Fe-Si1}} = 3.24(1) \text{ \AA}$, respectively (Table 4). The bond distances are consistent with values obtained from diffraction data for Mg and Si shells (Breu et al., 2003) and $R_{\text{Fe-Fe1}}$ is typical of edge-sharing Fe octahedra. The coordination numbers of octahedral Mg and Fe decrease from 2.8(2) to 0.6(1) and from 0.3(1) to 0.1(1), respectively, and increase for Si from 2.1(2) to 6.9(7), for increasing α values (Table 4). The angular dependences are consistent with Mg and Fe shells located in-plane and Si located out-of-plane, but the in-plane number of neighboring Mg+Fe are much lower than expected from the smectite composition. For a trioctahedral structure, six octahedral and four tetrahedral cationic neighbors would be expected. The detected number of tetrahedral neighbors equals 3.4(6), reasonably close to the theoretical value of four, but the number of octahedral neighbors equals ~ 2 , much lower than the expected value of six (Table 4). Several explanations could possibly explain the deviation from the ideal values. First, hectorite contains Li substituting for Mg (Table 1) and Li is too light to be detected by EXAFS spectroscopy. The impact of this substitution on N_{Mg} can be estimated from the Li:Mg ratio of $\sim 1:10$ in hectorite. Assuming random distributions of Mg and Li in the octahedral sheet, the detected number of Mg neighbors should be close to 5.4, a value still significantly higher than the experimental value. Second, as Fe(III) substitutes for Mg(II), the increase in local charge may be balanced by vacancies or, more likely by Li(I) preferential insertion in adjacent octahedral sites, thereby lowering the number

of detected octahedral neighbors. A third hypothesis is that electronic waves backscattered by octahedral Fe and Mg/Al are out of phase (Manceau, 1990a) leading to destructive interferences as can be seen in Figures 6 and 7. Owing to this interference, the apparent numbers of Fe and Mg neighbors are mutually lowered and so actual number of neighboring Fe (and Mg) cations. Please note that such interference would also occur in powder spectra, and would be further compounded by the EXAFS contribution of out-of-plane Si. Peak B for hectorite also contains contribution from the next nearest oxygen shell ($N_{O2} = 0.6(1.1)$) located on the basal plane ($R_{Fe-O2} = 3.48(1) \text{ \AA}$). This shell is not detected in the dioctahedral smectite because of corrugation of the basal oxygen planes that increases the average spread of the Fe–O2 distances (Manceau et al., 2000b). From the increase of N_{O2} with α (Table 4) this shell is clearly located out-of-plane, matching expectation for a shell located at the basal plane.

Two additional O shells were also used to improve the fit of peak B at higher distances in both smectites (Tables 3 and 4). The O3 shell originates from O atoms of adjacent octahedra ($R_{Fe-O3} = 3.78\text{-}3.79 \text{ \AA}$) and the O4 shell from atoms of the basal planes of the silicate sheet ($R_{Fe-O4} = 4.01(1) \text{ \AA}$), and the bond distances are close to reported values (Tsipursky and Drits, 1984; Breu et al., 2003). For montmorillonite, the number of detected backscatterers at $\alpha = 35^\circ$ ($N_{O3} = 5.8(1.3)$ and $N_{O4} = 4.2(1.9)$) match coordination numbers expected for smectites ($N_{O3} = 6$ and $N_{O4} = 4$), whereas they are much lower than expected for hectorite ($N_{O3} = 2.8(1.5)$ and $N_{O4} = 1.0(1.3)$). This again suggests that the phyllosilicate local structure around Fe(III) is distorted in hectorite.

For montmorillonite and hectorite, N_{O4} increases from 2.4(2.4) to 8.0(1.3) and from 0.1(1) to 5.5(5), respectively, for α increasing from 10° to 80° . These variations match expectations for the O4 shell located at the smectite basal planes, i.e., in out-of-plane direction. The O3 coordination number is invariant within uncertainty both for montmorillonite ($N_{O3} = 5.4(1.6)$ to $6.1(9)$ for $\alpha = 10$ to 80° , respectively) and hectorite ($N_{O3} = 2.4(1.3)$ to $4.1(1.9)$ for $\alpha = 10$ to 80° , respectively). Thus no clear angular dependence can be obtained from these values but they are compatible with the known orientation of the O3 shell inclined toward the octahedral sheet {Schlegel, 2001 #76}.

Discussion

X-ray diffraction and infrared data showed that both purified smectites are free from Fe-containing impurities, implying that Fe can only be located in the smectite structure. Infrared data are also consistent with a dioctahedral structure for montmorillonite and a trioctahedral structure for hectorite. The pre-edge and EXAFS spectroscopic data indicate that Fe is in trivalent oxidation state and located in the octahedral sheet but that the Fe local environments in the two clays are dissimilar. The bond lengths and the angular dependences of the neighboring atomic shells confirm this structural assignment and indicate that the amount of Fe in the tetrahedral sheet is marginal.

In montmorillonite, the number of cationic neighbors detected in-plane ($N_{Fe1} + N_{Al1} = 3$) is in agreement with a dioctahedral structure ($N = 3$). Fits to the data are consistent with Fe substituting for Al/Mg in the octahedral layer. Fe-Fe pairs were detected, but with a low number of backscatterers ($N_{Fe1} = 0.4(1)$). The elemental composition (Table 1) indicates a moderate Fe content (2.3 wt%) and the corresponding (Al+Mg):Fe molar ratio equals 9.4. Considering a fully random Fe distribution with all Al located in the octahedral sheet, about 0.3 neighboring Fe atom would be expected. This calculated value is similar to the fit results within

uncertainties. This finding excludes Fe clustering and suggests that Fe is randomly distributed in montmorillonite. Earlier powder EXAFS analysis (Vantelon et al., 2003) also indicated an ordered Fe distribution in this montmorillonite but obeying an exclusion rule. However, powder XAS applied to sheet silicates is affected by larger uncertainties due to the interferences between waves backscattered by shells located at similar distances {Manceau, 1990 #55}. This limitation was overcome in the present study by using the improved sensitivity of P-EXAFS spectroscopy, and low number of Fe backscatterers could be unambiguously detected. This improved sensitivity of P-EXAFS spectroscopy was subsequently applied to investigate in detail and for the first time, the Fe distribution in the trioctahedral hectorite.

In hectorite all octahedral sites are filled by various cations such as Li, Mg, Al and Fe, and in addition some F substitutes for OH groups. This complex chemical composition of the octahedral sheet could somewhat complicate the powder EXAFS investigation of the Fe environment, thus requiring the use of P-EXAFS spectroscopy. The pre-edge is particularly sensitive to the electronic and geometric structure of the Fe site. The analysis of the hectorite pre-edge suggests differences in the binding environment compared to montmorillonite that may be attributed to some F substitution for OH groups. This substitution possibly distorted the site geometry and/or modified the ligand field. Although OH and F have similar ionic radii ($r(\text{OH}) = 1.32 \text{ \AA}$, $r(\text{F}) = 1.28 \text{ \AA}$ in II-coordination; Shannon, 1976), F is more electronegative than O, which can decrease the length of the Fe–F bond compared to the Fe–O one. This is confirmed by reported $R_{\text{Mg-O}}$ and $R_{\text{Mg-F}}$ distances in hectorite (Breu et al., 2003) differing by $\sim 0.06 \text{ \AA}$. The Fe–(O,F)₆ coordination octahedron is thus probably slightly distorted, as corroborated by the slightly larger mean square displacement of the O1 shell in hectorite ($\sigma^2 = 0.006 \text{ \AA}^2$) compared to montmorillonite ($\sigma^2 = 0.005 \text{ \AA}^2$). The distribution of interatomic distances around a mean value dampens the EXAFS oscillatory contribution from the (O,F)₁ shell, thus leading to a slightly lower number of neighboring (O,F) atoms. Furthermore, the mean square displacement is also larger for octahedral neighbors in hectorite (Mg, Fe) compared to montmorillonite (Al, Fe). However, previous NMR data have demonstrated Fe–F avoidance in phyllosilicate, i.e., F is usually not associated to Fe octahedra {Sanz, 1983 #79}. Also, NMR investigation on hectorite failed to report evidence for Fe–F association, although this could be due to the minute amount of structural Fe (0.2 %wt) compared to the high F content (up to 5.18 wt% in some preparations) {Thomas, 1977 #4008}. In addition, any distortion due to the presence of F is unlikely to result in a large incoherency of the distances between Fe and the neighboring cations. Indeed, the limited nature of this disorder is attested by the ability of conventional diffraction methods to derive a structural model for hectorite without the need of complex corrections {Oberlin, 1966 #77; Breu, 2003 #7}. The benign nature of this distortion is corroborated by the detection of similar numbers of Si atoms at very similar distances and comparable mean square displacement in both smectites.

Hectorite has a low Fe content (0.2 wt%) and the corresponding (Al+Mg):Fe molar ratio equals 130. Considering a fully random distribution of Fe in the octahedral sheet, the number of neighboring Fe detected by EXAFS would be about 0.06, slightly smaller than the detected value of 0.2(1), values close to the lower limit of detection of EXAFS. This would suggest that Fe atoms within the octahedral sheet may preferentially form some pairs, as reported for trioctahedral micas (Manceau et al., 1990). However, the extended formation of Fe(III) clusters within the octahedral sheet is ruled out because it would have led to much greater N_{Fe} values.

Fe is also surrounded by Mg neighbors, but the number of detected Mg is much smaller than the theoretical number for random distribution in hectorite ($N_{\text{Mg}} = 5.4$), which cannot be explained by structural disorder alone. In hectorite, almost all octahedral sites are filled with cations of similar sizes when sixfold coordinated by oxygen atoms ($r^{\text{VI}}(\text{Mg(II)}) = 0.72 \text{ \AA}$ and $r^{\text{VI}}(\text{Li(I)}) = 0.76 \text{ \AA}$ (Shannon, 1976)), but of different charge. The sixfold coordinated Fe(III) is slightly smaller ($r^{\text{VI}}(\text{Fe(III)}) = 0.65 \text{ \AA}$) than Mg or Li so that the substitution would not result in significant lattice strain. However, Fe(III) insertion in octahedral sites may be locally charge-balanced by octahedral vacancies and/or by Li(I) filling neighboring octahedra. Because Li cannot be detected by EXAFS spectroscopy, these two hypotheses cannot be discriminated, but both of them would result in a low number of detected octahedral neighbors. In fact, the presence of vacancies surrounding Fe(III) would correspond to the local formation of a dioctahedral-like environment around Fe(III), a little similar to the Fe(III) local environment in montmorillonite. However, it would probably lead to a large deficit of local charge. Interestingly, in dioctahedral smectites, deficit of local charge can be decreased by thermal diffusion of Li(I) cations in vacant octahedra {Komadel, 2005 #80}. In the case of hectorite, such a diffusion would not be driven by heat, though, but would readily occur during crystallization. Other charge-balance mechanisms include the deprotonation of hydroxyl groups of the octahedral sheet close to Fe(III), but this would probably lead to severe distortion of the Fe coordination sphere. Still another mechanism of local charge compensation would be the substitution of Si(IV) by Al(III) in the tetrahedral sheet. From our EXAFS results, however, we surmise that local charge compensation by Li .

Conclusions

The local chemical environment around Fe in the dioctahedral smectite montmorillonite and, for the first time, in the trioctahedral smectite hectorite, was characterized by P-EXAFS spectroscopy. Iron(III) in these smectites is surrounded by in-plane (Mg,Al,Fe) atoms and out-of-plane Si atoms at distances in agreement with previous studies. Furthermore, Fe-Fe pairs in limited number could be unambiguously detected, owing to the ability of P-EXAFS to filter out overlapping contributions from shells with distinct orientations. Low numbers of neighboring Fe were detected, and no extended clustering could be observed in either of the minerals, indicating a next-to-random distribution of Fe in the octahedral sheet. Furthermore, the XAS data indicate that the substitution of Fe for Al in the dioctahedral montmorillonite does only marginally affect the local site symmetry. In contrast, the substitution of Mg(II) by Fe(III) creates an excess of charge, which is balanced by a deficit of charge in the neighboring octahedra, either in the form of vacancies or as Mg(II) substitutions by Li(I). If confirmed, this result would provide a nice example on how impurities in clay minerals can drive heterogeneities in the charge distributions at clay mineral surfaces. The additional amount of lattice flexibility offered by the presence of vacant sites would also explain the ability of hectorite to incorporate large cations such as lanthanides {Finck, 2009 #81}, and, possibly, actinides.

Acknowledgements

The European Synchrotron Radiation Facility is acknowledged for provision of synchrotron radiation beamtime. We thank O. Proux for assistance during XAS data collection at the FAME beamline.

Figure captions

Figure 1. Top: Structure of smectite. T: tetrahedral sheet, O: octahedral sheet. Middle: projection down c^* of trioctahedral framework. Bottom: projection down c^* of dioctahedral framework. Blue triangles are $(\text{Si,Al})\text{O}_4$ tetrahedra, one tetrahedral sheet is not shown. M1 denotes *trans* sites, and M2 *cis* sites.

Figure 2. X-ray diffractogram (basal spacing or $d(001)$ is indexed) and ATR-FTIR spectrum of the clay fractions investigated in this study. Top: Montmorillonite. Bottom: Hectorite.

Figure 3. Comparison of experimental Fe *K*-edge XANES (left), and EXAFS spectra (middle) recorded at $\alpha = 35^\circ$ for montmorillonite and hectorite with the corresponding Fourier transforms (right).

Figure 4. Normalized pre-edge spectra (solid black line) and best fits (dashed red lines) with individual components (blue and green thin lines) for montmorillonite (left) and hectorite (right).

Figure 5. Experimental P-EXAFS spectra (left) with the corresponding Fourier transforms (right) for montmorillonite and hectorite.

Figure 6. Experimental (solid black lines) and modeled (dashed red lines) polarized Fourier transforms for the montmorillonite data. The right panel indicates the contributions at $\alpha = 35^\circ$ of the single atomic shells.

Figure 7. Experimental (solid black lines) and modeled (dashed red lines) polarized Fourier transforms for the hectorite data. The right panel indicates the contributions at $\alpha = 35^\circ$ of the single atomic shells.

Table captions

Table 1. Main cations content of montmorillonite (SWy-1) and hectorite (SHCa-1) determined by chemical analysis (Other includes oxygen).

Table 2. Results from the pre-edge fitting (uncertainties on the energies are estimated to 0.1 eV and 10 % on the intensities).

Table 3. Quantitative EXAFS analysis for montmorillonite^a

Table 4. Quantitative EXAFS analysis for hectorite^a

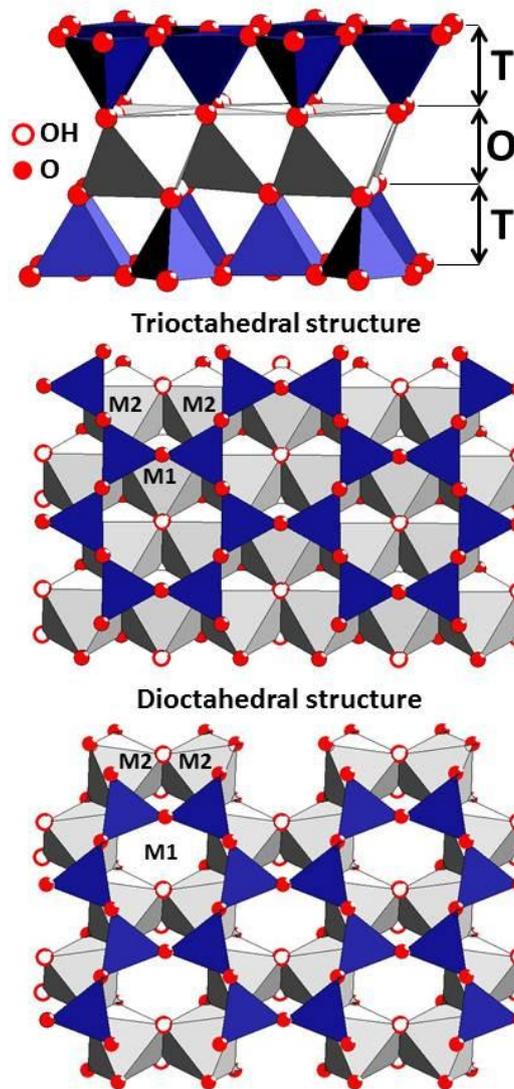


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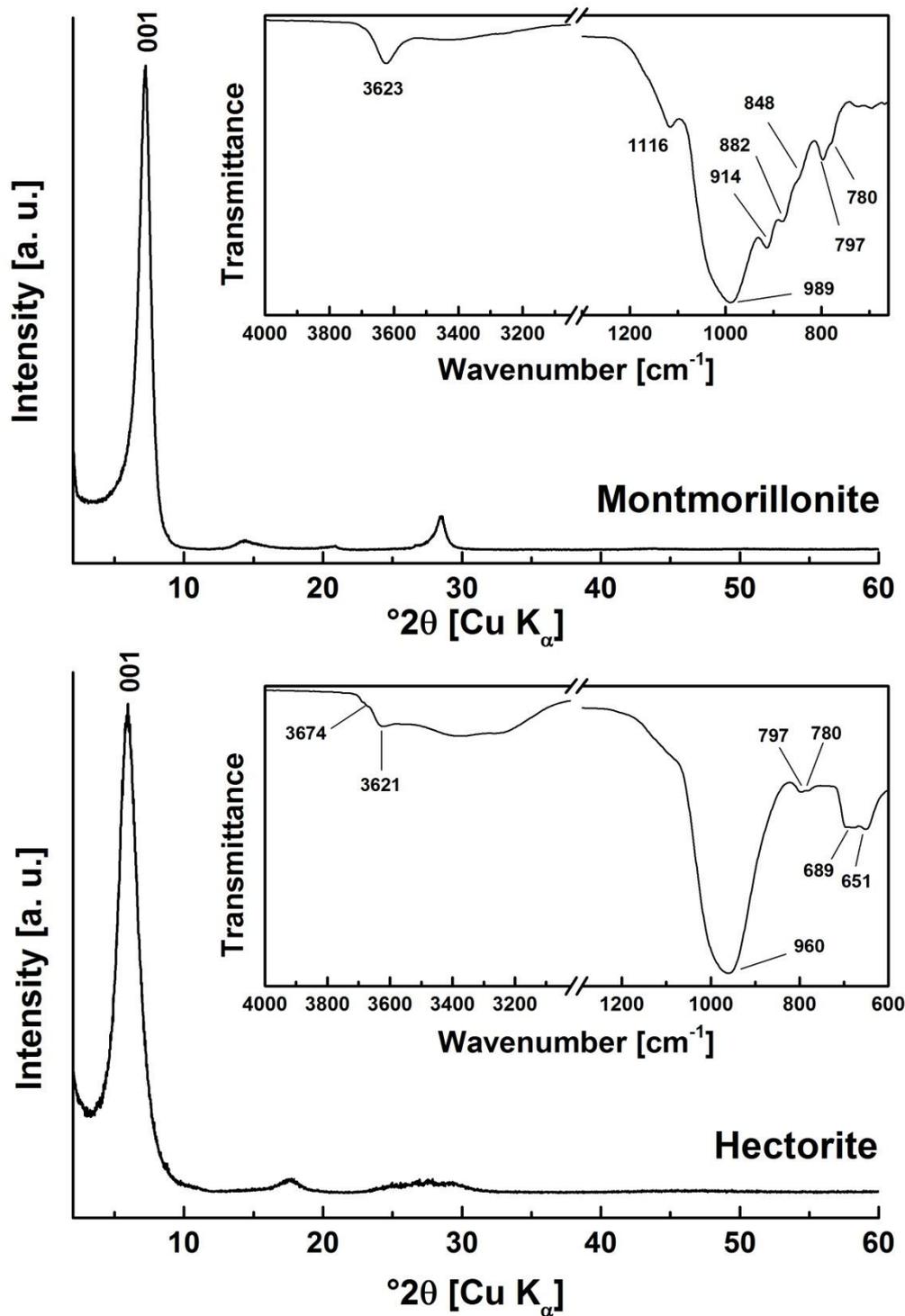


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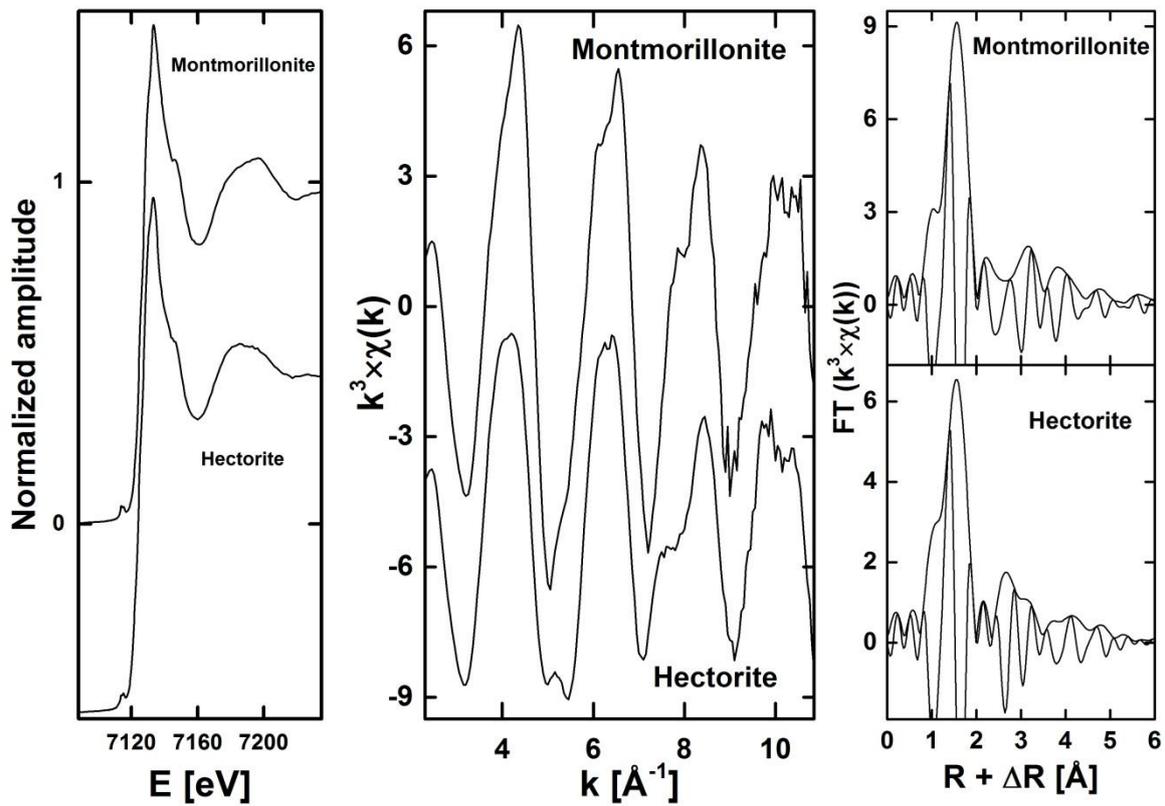


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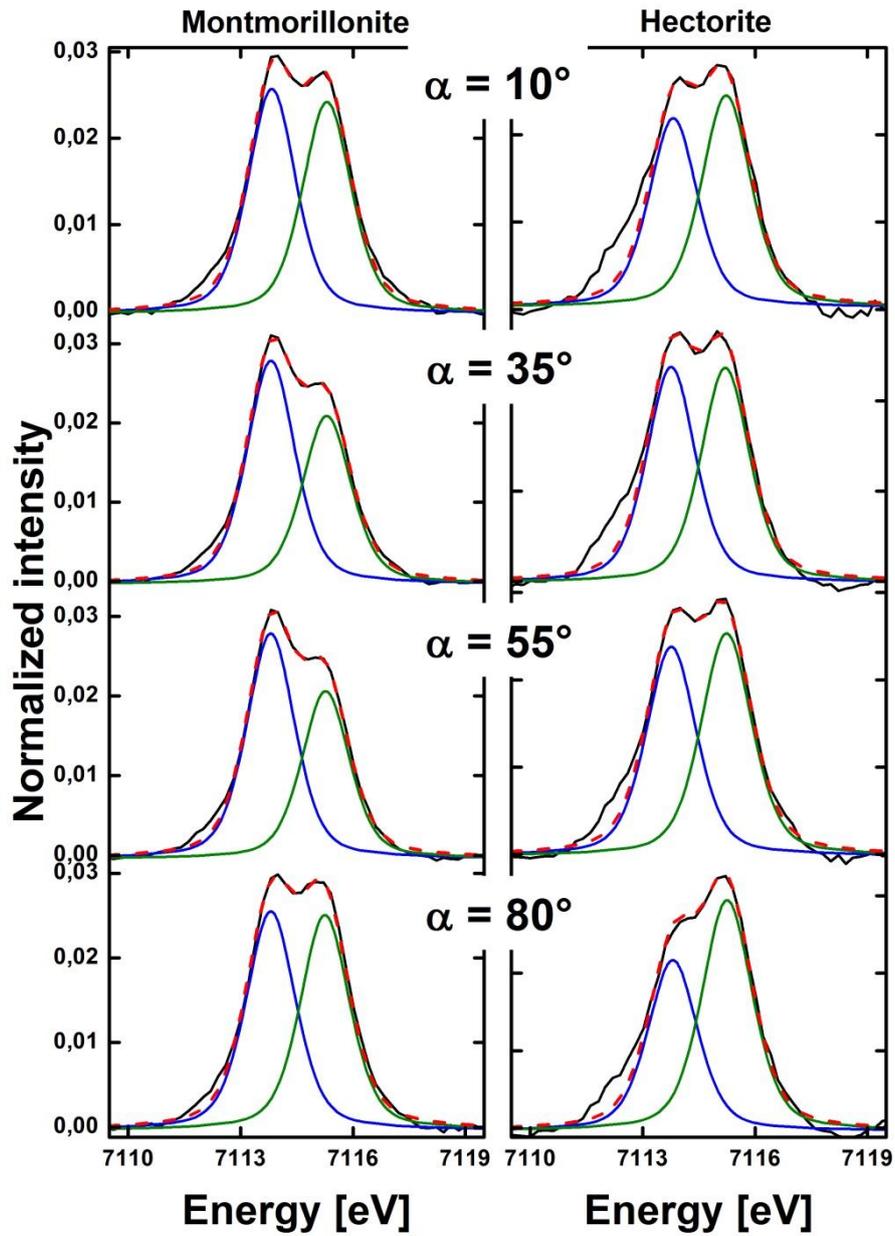


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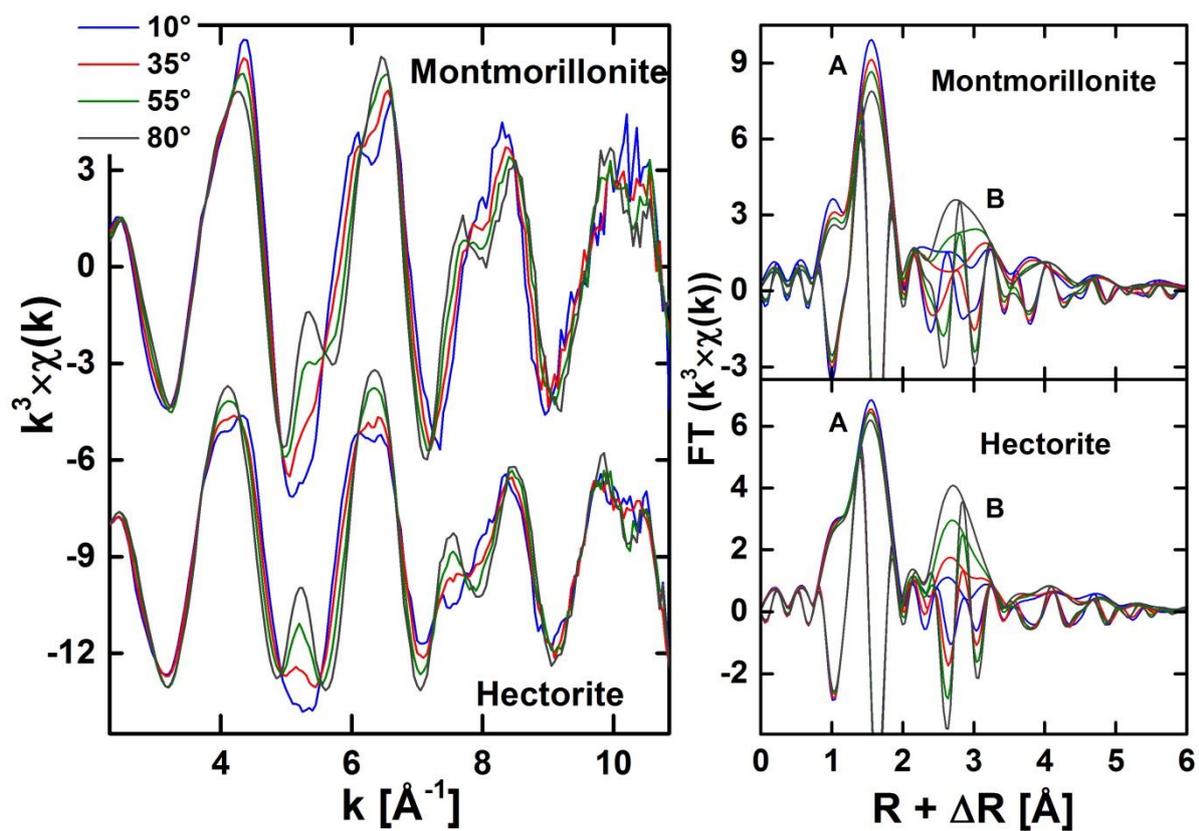


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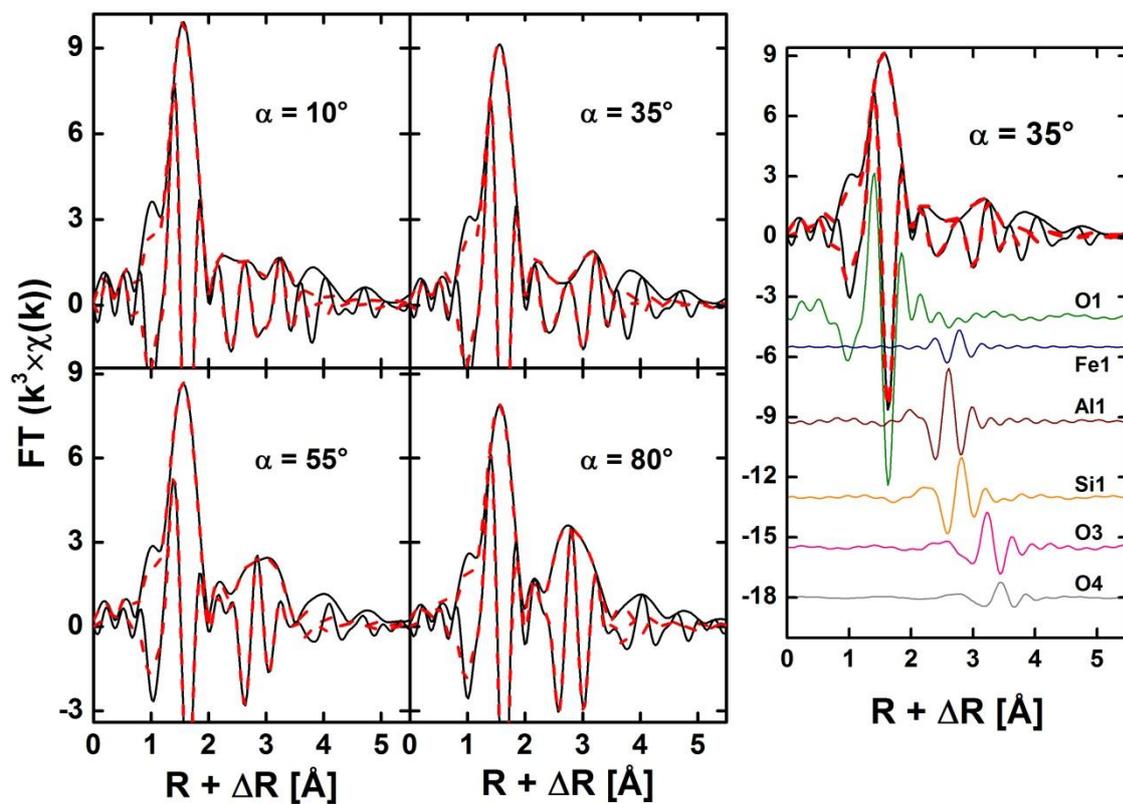


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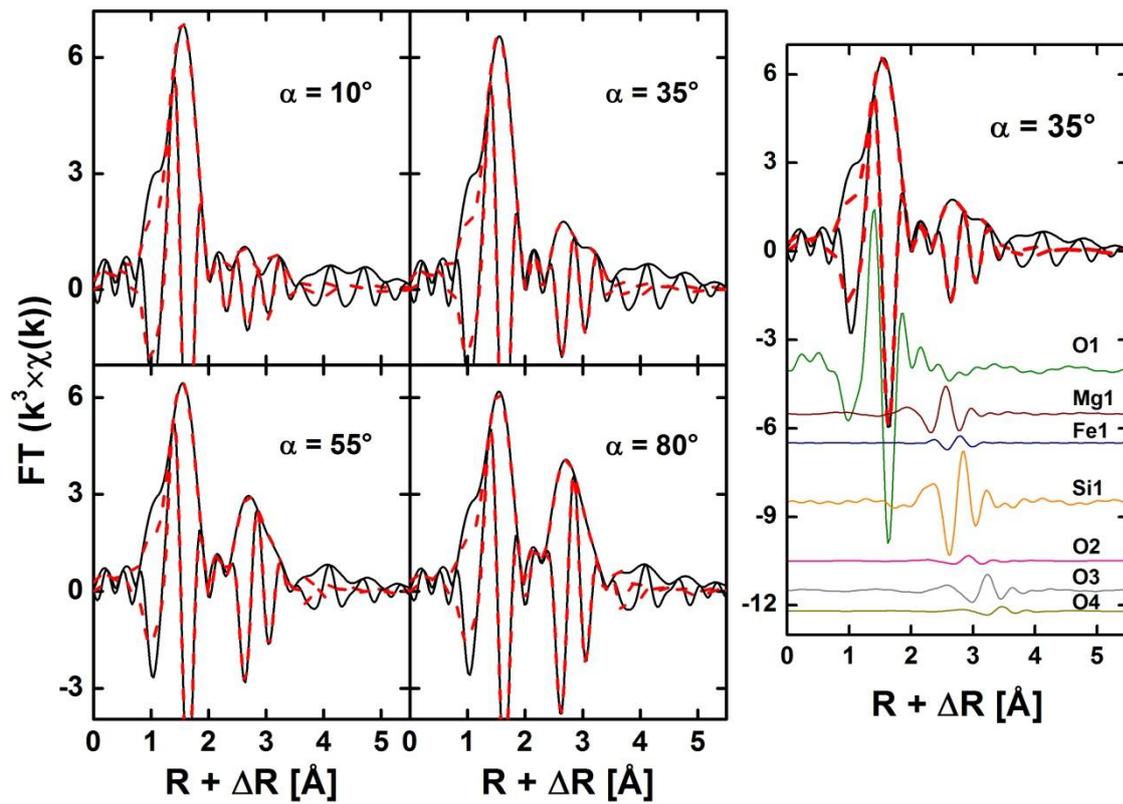


Figure 7. Experimental (solid black lines) and modeled (dashed red lines) polarized Fourier transforms for the hectorite data. The right panel indicates the contributions at $\alpha = 35^\circ$ of the single atomic shells.

Table 1. Main cations content of montmorillonite (SWy-1) and hectorite (SHCa-1) determined by chemical analysis (Other includes oxygen).

(wt %)	Li	Mg	Al	Si	Fe	Other	Total
SWy-1	<0.01	1.36 (±0.01)	8.93 (±0.07)	25.84 (±3.30)	2.30 (±0.14)	61.57	100.00
SHCa-1	0.39 (0.01)	12.15 (±0.09)	0.40 (±0.01)	21.08 (±2.69)	0.24 (±0.01)	65.74	100.00

Table 2. Results from the pre-edge fitting (uncertainties on the energies are estimated to 0.1 eV and 10 % on the intensities).

Sample	α	Pre-edge peak energy [eV]	Pre-edge peak intensity ¹	Peak energy difference ²	Peak intensity ratio ³
Montmorillonite	10°	7113.83	0.050	1.48	1.04
		7115.31	0.048		
	35°	7113.81	0.055	1.49	1.31
		7115.30	0.042		
	55°	7113.81	0.055	1.46	1.38
		7115.27	0.040		
	80°	7113.81	0.051	1.44	1.02
		7115.25	0.050		
Hectorite	10°	7113.82	0.042	1.41	0.88
		7115.23	0.048		
	35°	7113.76	0.047	1.45	1.00
		7115.21	0.047		
	55°	7113.77	0.048	1.47	0.92
		7115.24	0.052		
	80°	7113.81	0.044	1.45	0.72
		7115.26	0.061		

¹peak intensities correspond to integrated areas. ²peak energy differences correspond to energy differences between the lowest energy and the highest energy pre-edge features. ³ratio of the lowest energy peak intensity to the highest energy peak intensity.

Table 3. Quantitative EXAFS analysis for montmorillonite^a

α	Fe \leftrightarrow O1			Fe \leftrightarrow Fe1			Fe \leftrightarrow Al1		
	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]
10°	6.3(5)			0.5(1)			3.4(3)		
35°	5.8(5)	2.01(1)	0.005	0.4(1)	3.02(1)	0.004	2.6(3)	3.04(1)	0.004
55°	5.5(5)			0.2(1)			1.4(5)		
80°	4.9(4)			0.1(1)			0.6(1)		

α	Fe \leftrightarrow Si1			Fe \leftrightarrow O3			Fe \leftrightarrow O4			ΔE_0^b	R_f^c
	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]	<i>N</i>	<i>R</i> [Å]	σ^2 [Å ²]	[eV]	($\times 10^3$)
10°	1.4(7)			5.4(1.6)			2.4(2.4)				5.6
35°	3.4(3)	3.22(1)	0.007	5.8(1.3)	3.79(1)	0.004	4.2(1.9)	4.01(2)	0.007	6.8(4)	5.4
55°	5.2(5)			6.0(1.1)			6.3(1.6)				5.6
80°	6.6(7)			6.1(9)			8.0(1.3)				8.3

^aThe data were fit over the entire range considering the data from both tables: *N* is the coordination number, *R* is the interatomic distance, σ^2 is the mean squared displacement. The number in parentheses indicates the uncertainty. ^bShift in ionization energy, threshold energy E_0 taken as zero crossing of the second derivative. ^cFigure of merit of the fit (Ravel, 2000).

Table 4. Quantitative EXAFS analysis for hectorite^a

α	Fe \leftrightarrow O1			Fe \leftrightarrow Mg1			Fe \leftrightarrow Fe1			Fe \leftrightarrow Si1		
	N	R [Å]	σ^2 [Å ²]	N	R [Å]	σ^2 [Å ²]	N	R [Å]	σ^2 [Å ²]	N	R [Å]	σ^2 [Å ²]
10°	5.4(5)			2.8(2)			0.3(1)			2.1(2)		
35°	5.1(5)	2.00(1)	0.006	1.7(6)	3.03(1)	0.008	0.2(1)	3.03(3)	0.008	3.4(6)	3.24(1)	0.007
55°	4.9(4)			1.1(6)			0.1(1)			5.1(6)		
80°	4.6(4)			0.6(1)			0.1(1)			6.9(7)		

α	Fe \leftrightarrow O2			Fe \leftrightarrow O3			Fe \leftrightarrow O4			ΔE_0^b	R_f^c
	N	R [Å]	σ^2 [Å ²]	N	R [Å]	σ^2 [Å ²]	N	R [Å]	σ^2 [Å ²]	[eV]	($\times 10^3$)
10°	0.1(1)			2.4(1.3)			0.1(1)				4.7
35°	0.6(1.1)	3.48(1)	0.006	2.8(1.5)	3.78(1)	0.008	1.0(1.3)	4.01(1)	0.008	6.6(2)	3.6
55°	1.4(1.1)			3.9(1.6)			3.7(1.4)				2.9
80°	1.4(6)			4.1(1.9)			5.5(5)				4.4

^aThe data were fit over the entire range considering the data from both tables: N is the coordination number, R is the interatomic distance, σ^2 is the mean squared displacement. The number in parentheses indicates the uncertainty. ^bShift in ionization energy, threshold energy E_0 taken as zero crossing of the second derivative. ^cFigure of merit of the fit (Ravel, 2000).

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