

DE LA RECHERCHE À L'INDUSTRIE



www.cea.fr

On the effect of buoyancy on lateral migration of bubbles in turbulent flows: insights from Direct Numerical Simulations

Guillaume Bois¹, A. du Cluzeau¹, A. Toutant²

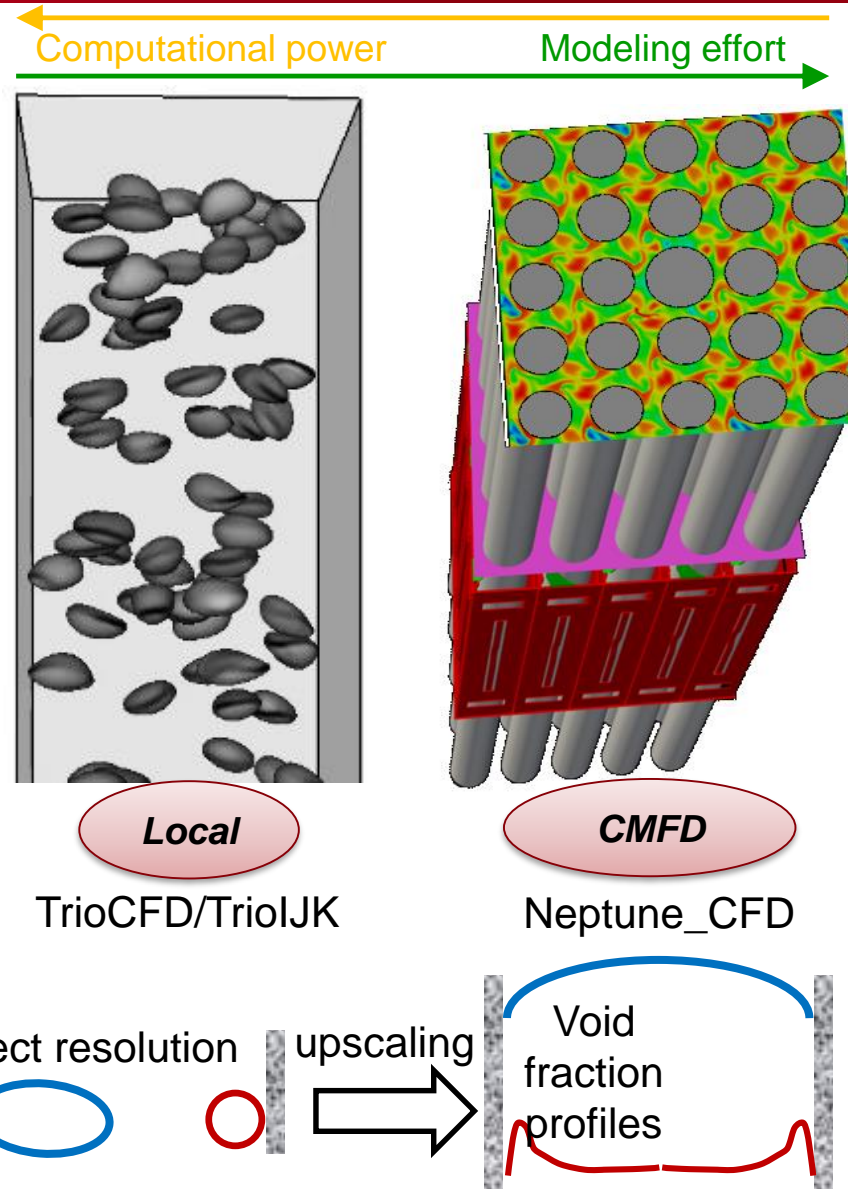
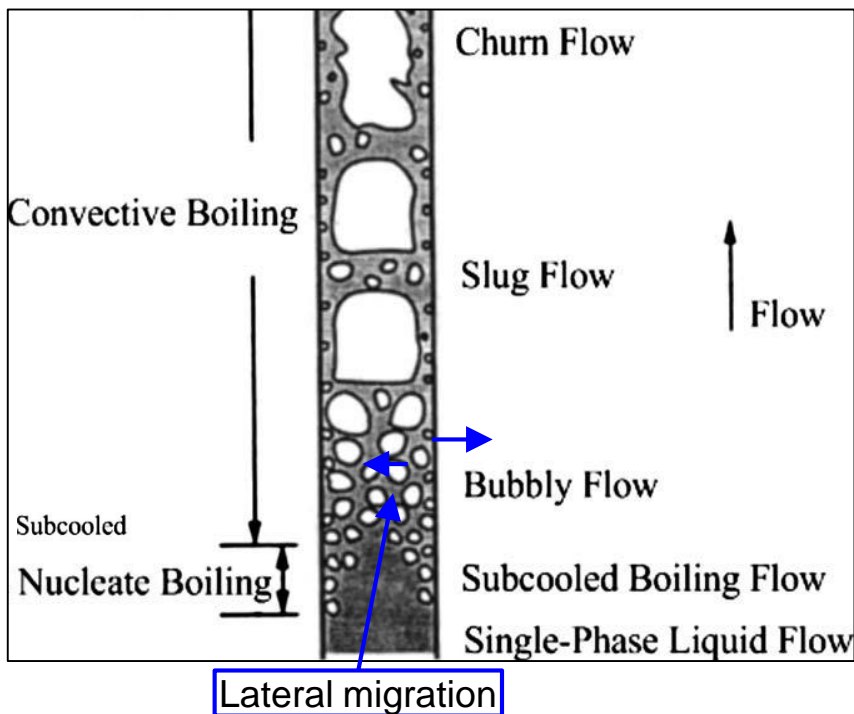
¹ DEN – Service de Thermo-hydraulique et de Mécanique des Fluides
CEA, Université Paris – Saclay, F-91191 Gif-sur-Yvette, France

² PROMES-CNRS UPR 8521
Rambla de la Thermodynamique Tecnosud, 66100 Perpignan, France

International Conference on Multiphase Flow (ICMF)
Rio de Janeiro, Brazil, 19-24 May 2019

guillaume.bois@cea.fr

- Broad application: Critical Heat Flux prediction in PWR conditions (high T&P);
- Focus on the effects of void accumulation at the wall / core due to lateral migration of bubbles;
- Separate effects analysis: adiabatic conditions;
- Up-scaling DNS understanding into CMFD models.



Governing equations (exact, averaged)

Two-fluid model used in CMFD codes

Definition of “two-fluid” variables for CMFD:

- $\overline{\phi_k^k} \triangleq \frac{\overline{\chi_k \phi_k}}{\overline{\chi_k}}$ (weighted **average** by the **phase** indicator)

- Void fraction: $\alpha_v = \overline{\chi_v}$

- Fluctuating velocities: $\mathbf{u}'_k = \mathbf{u}_k - \overline{\mathbf{u}}_k^k$

Phase-averaged momentum equation:

$$\frac{\partial \alpha_k \rho_k \overline{\mathbf{u}}_k^k}{\partial t} + \nabla \cdot [\alpha_k \rho_k (\overline{\mathbf{u}}_k^k \overline{\mathbf{u}}_k^k + \mathbf{R}_{ij}^k)] = -\nabla (\alpha_k \overline{p}_k^k) + \alpha_k \rho_k \mathbf{g} + \nabla \cdot [\alpha_k \overline{\boldsymbol{\tau}}_k^k] - \frac{(\overline{p}_k \mathbf{n}_k - \overline{\boldsymbol{\tau}}_k \cdot \mathbf{n}_k) \cdot \nabla \chi_k}{\overline{\chi_k}}$$

Momentum interfacial transfer defined from local fields:
 $\mathbf{M}_k = \mathbf{M}_k(\chi, u, p)$

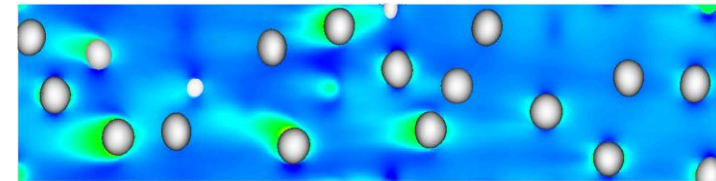
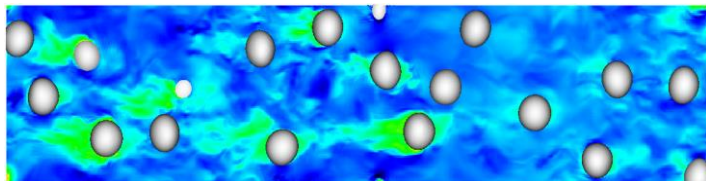
Correlation between fluctuations, Reynolds stresses : $\mathbf{R}_{ij}^k = \overline{u'_i u'_j}^k$

Constitutive equations on \mathbf{R}_{ij}^k and \mathbf{M}_k are required to close the system.

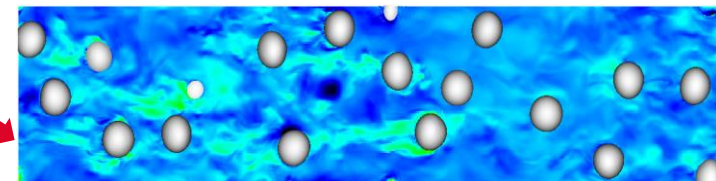
This presentation focuses on lateral motion induced by **migration forces** (\mathbf{M}_k closure)

See abstract 168047. Talk at 11:30am on Thursday for a focus on turbulence modelling.

Modeling of Bubble-Induced Fluctuations in turbulent bubbly up-flow from Direct Numerical Simulations,
A. du Cluzeau, G. Bois, A. Toutant & A. Burlot.



Space fluctuations of the velocity (WIF)
Time fluctuations of the velocity (WIT)



$$\overline{u'_i u'_i}^l = \overline{u'_i u'_i}^l \Big|_{\text{lam}} + \overline{u'_i u'_i}^l \Big|_{\text{turb}}$$

Instantaneous velocity \rightarrow Reynolds stresses

Momentum interfacial transfers are assessed by DNS:

$$\mathbf{M}_k = \frac{D(\alpha_k \rho_k \bar{\mathbf{u}}_k^k)}{Dt} + \nabla \cdot (\alpha_k \rho_k \mathbf{R}_{ij}^k) + \nabla(\alpha_k \bar{p}_k^k) - \alpha_k \rho_k \mathbf{g} - \nabla \cdot [\alpha_k \bar{\boldsymbol{\tau}}_k^k]$$

They are in disequilibrium because of surface tension:

$$\mathbf{M}_l + \mathbf{M}_v = \overline{\sigma \kappa \nabla \chi_v}$$

Beware, \mathbf{M}_v is not the net force on the gas phase!

Classical interfacial momentum transfer is turned into a trajectory equation for the gas phase:

Net force on the gas phase

$$\alpha_v \rho_v \frac{d\bar{\mathbf{u}}_v^v}{dt} = \mathbf{M}_v^{\text{tot}}$$

$$= \alpha_v (\rho_v - \langle \rho \rangle) \mathbf{g} - \alpha_v \alpha_l (\rho_l - \rho_v) \mathbf{g} - \alpha_v \left[\rho_v \overline{\mathbf{u}'_v \mathbf{u}'_v}^v - \rho_l \overline{\mathbf{u}'_l \mathbf{u}'_l}^l \right] \cdot \nabla \alpha_v + \alpha_v \overline{\sigma \kappa \nabla \chi_v} - \alpha_v \nabla [(\bar{p}_v^v - \bar{p}_l^l) \alpha_v]$$

$$- \alpha_v \nabla \bar{p}_l^l - \alpha_v \left[\alpha_l \nabla \cdot (\rho_l \overline{\mathbf{u}'_l \mathbf{u}'_l}^l) + \alpha_v \nabla \cdot (\rho_v \overline{\mathbf{u}'_v \mathbf{u}'_v}^v) \right]$$

Classically neglected in the standard two-fluid one-pressure model

- Negligible in our conditions (from DNS and for readability): added-mass force, viscous terms

Physical assumptions (classical two-fluid one-pressure modeling):

- Fluctuations are neglected in the vapor phase;
- Differences between pressures are neglected: $\bar{p}_v^v = \bar{p}_v^l = \bar{p}_l^l$
- Lateral forces are made of lift, turbulent dispersion and wall lubrication forces only.

A deeper analysis of this trajectory equation for the gas phase (Newton's second law):

$$\alpha_v \rho_v \frac{d\bar{\mathbf{u}}_v^v}{dt} = \mathbf{M}_v^{\text{tot}}$$

Buoyancy and part of the drag

$$= \alpha_v (\rho_v - \langle \rho \rangle) \mathbf{g} - \alpha_v \alpha_l (\rho_l - \rho_v) \mathbf{g} - \alpha_v \left[\rho_v \overline{\mathbf{u}'_v \mathbf{u}'_v} - \rho_l \overline{\mathbf{u}'_l \mathbf{u}'_l} \right] \cdot \nabla \alpha_v + \alpha_v \overline{\sigma \kappa \nabla \chi_v} - \alpha_v \nabla [(\bar{p}_v^v - \bar{p}_l^l) \alpha_v]$$

Turbulent dispersion

Laminar dispersion

NEW, classically neglected

Two-phases Reynolds stresses:

- Laminar (WIF) $\propto \mathbf{u}_r^2$ NEW
- Turbulent (WIT) $\propto \text{TKE}$

Sub-part of Laminar and turbulent Lift

$$M_{\mathcal{R}}^L|_{\text{lam}} + M_{\mathcal{R}}^L|_{\text{turb}} \propto (C_{LL} + C_{TL}) \rho_l \mathbf{u}_r \nabla \wedge (\mathbf{u}_l)$$

Complex combination of single-phase pressure gradient and lift force due to pressure (induced by surface tension, NEW)

The turbulent part (NEW) can cause core-peaking ($C_{TL} < 0$)
It cannot be assessed on isolated bubbles (like Tomyiama's)
Turbulent lift is still $\propto \rho_l \mathbf{u}_r \nabla \wedge (\mathbf{u}_l)$

- Viscous terms and added mass are neglected

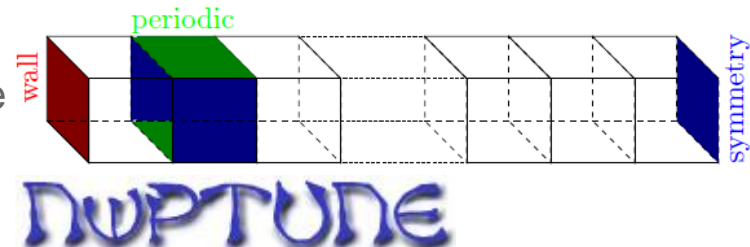
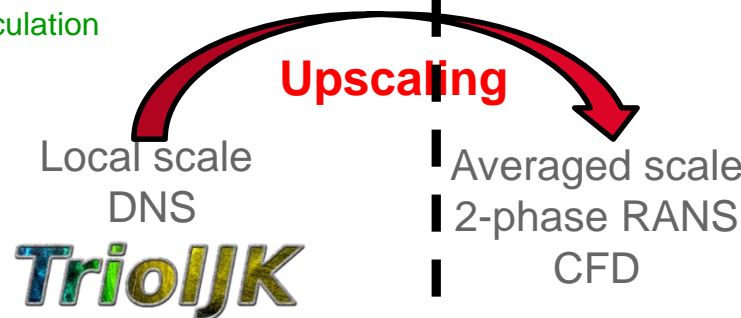
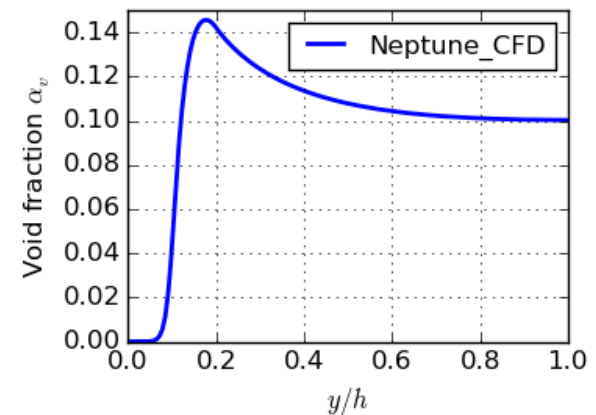
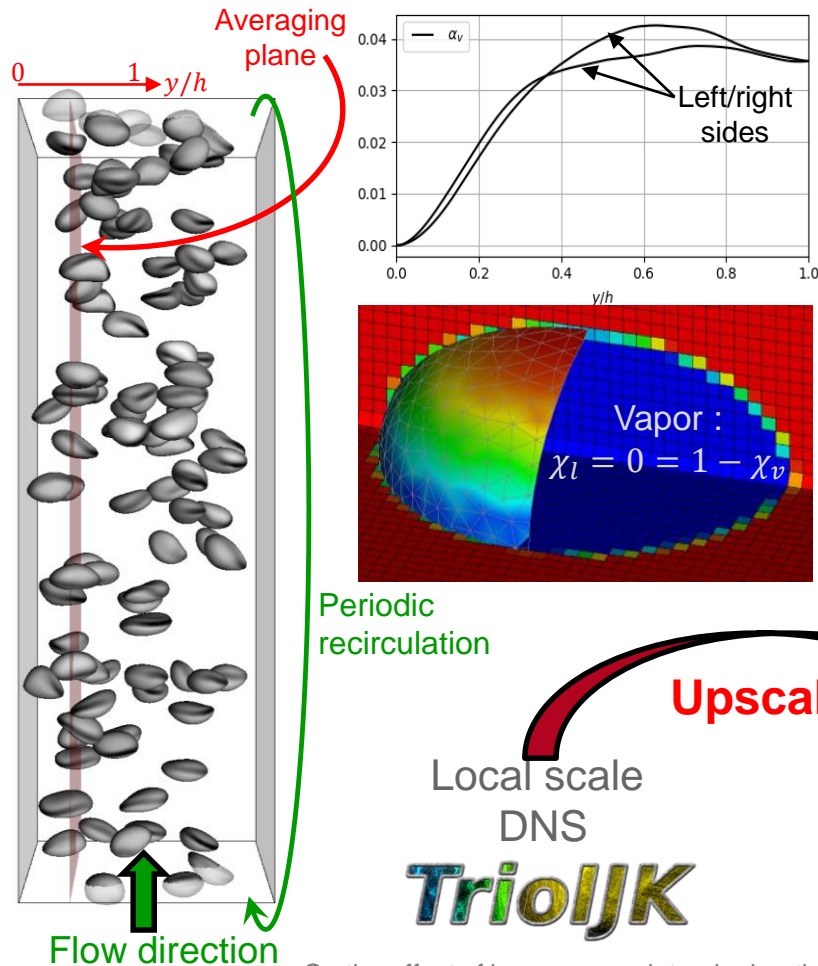
We derived an exact relation with the definition of forces from local variables!

Classical closures pile up assumptions and lack physical understanding of collective effects;

We propose to generate and leverage DNS data to evaluate this equilibrium equation...

■ Interfacial transfers and turbulence in bubbly flows;

- DNS of a bubbly turbulent channel flow with a [Front-Tracking](#) algorithm (TrioCFD/TrioIJK);
- « Numerical experiments »: separated effects unit-tests (adiabatic, no contact line or coalescence);
- Mean profiles and higher-order statistics are computed.



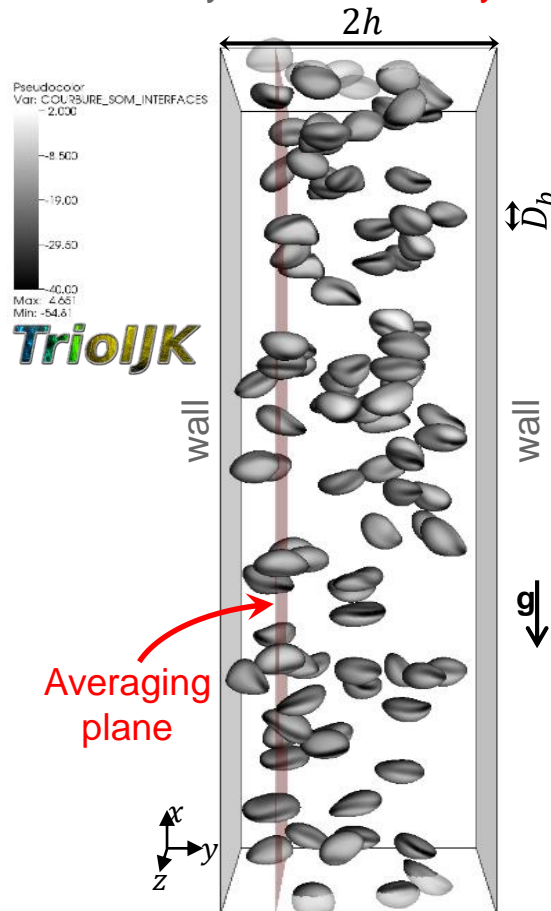
Lu & Tryggvason (Physics of Fluids, 20, 2008).

⇒ Code benchmark for validation;

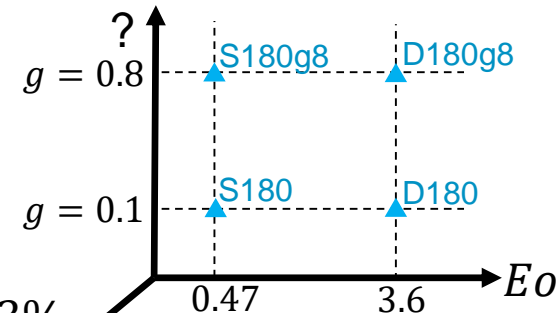
⇒ Additional statistical analysis.

Extended to higher Reynolds numbers;

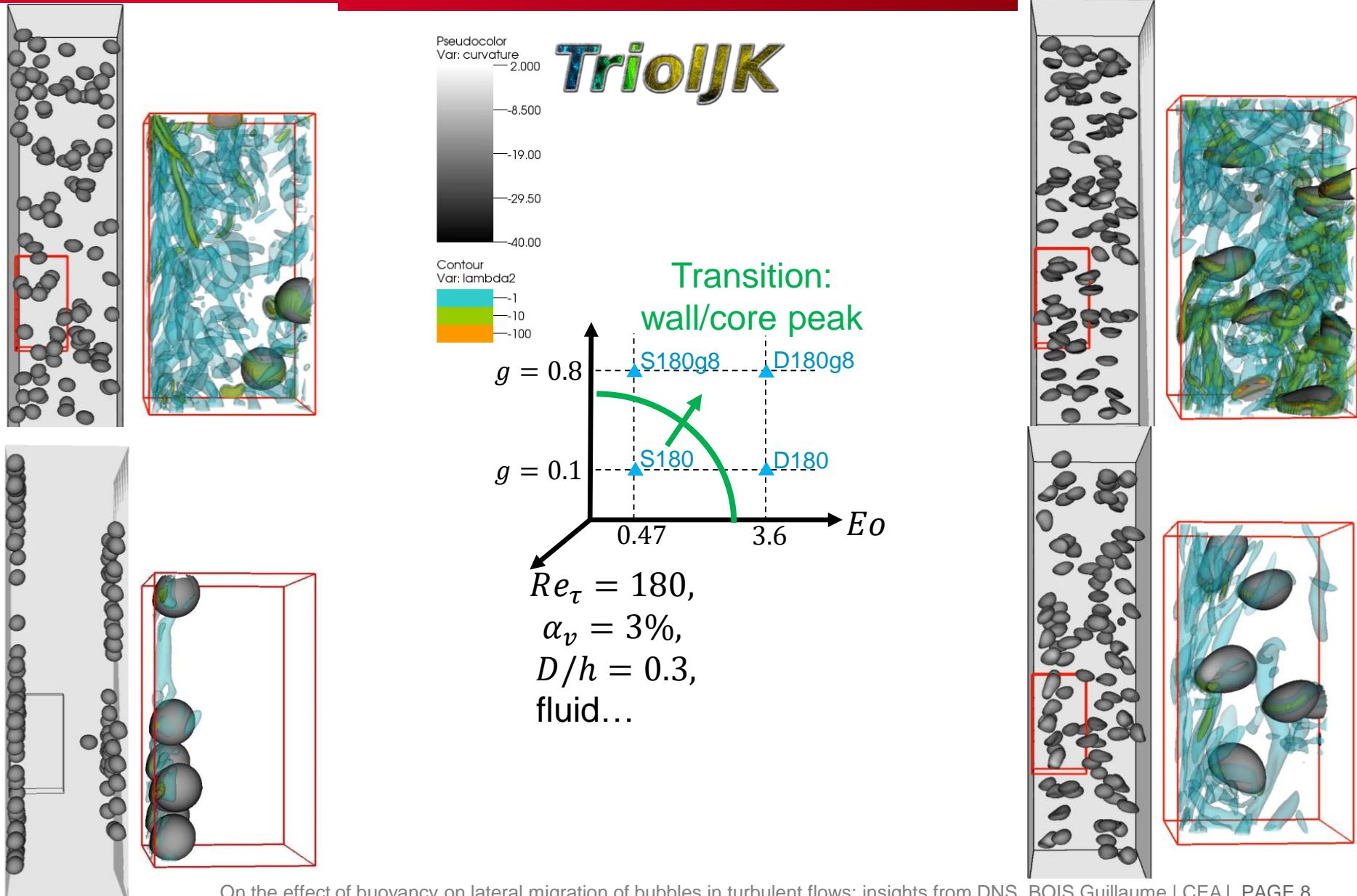
- Study on **deformability** and **gravity** (relative velocity).

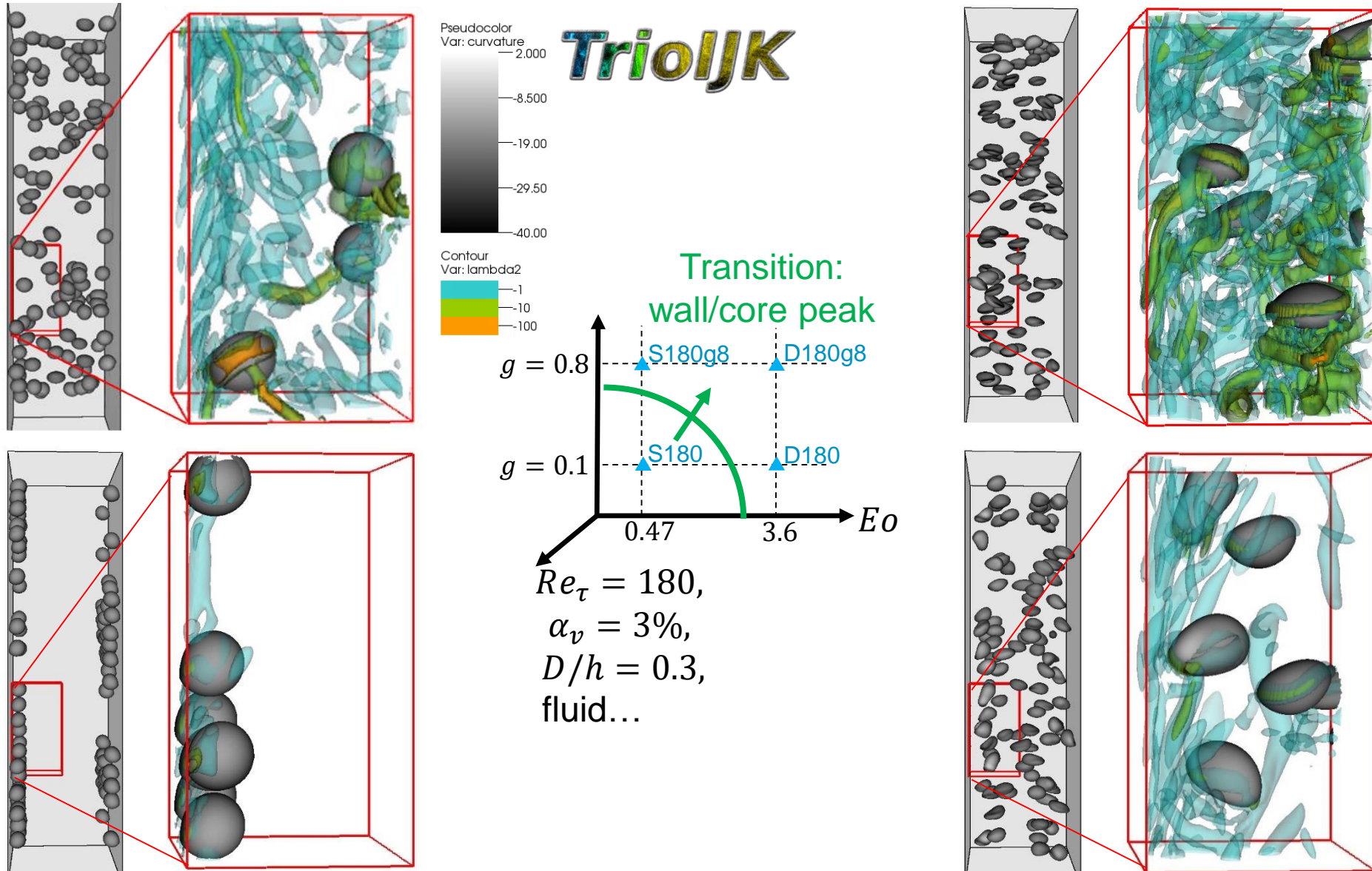


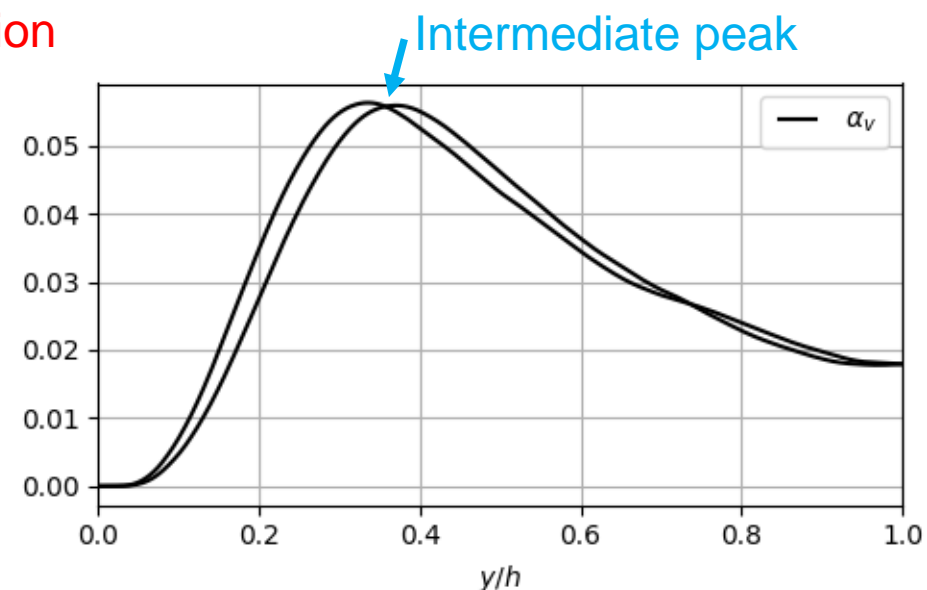
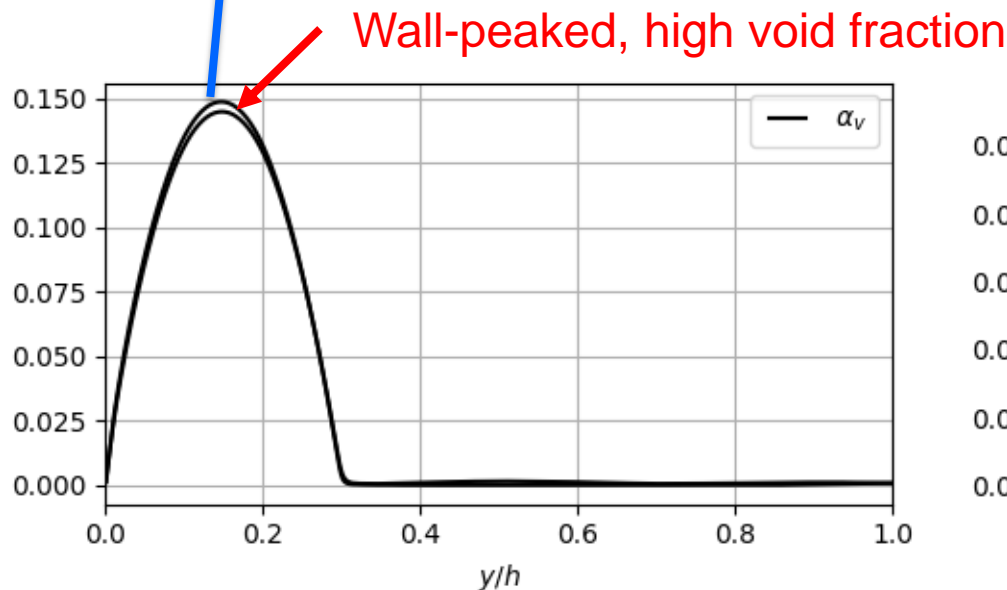
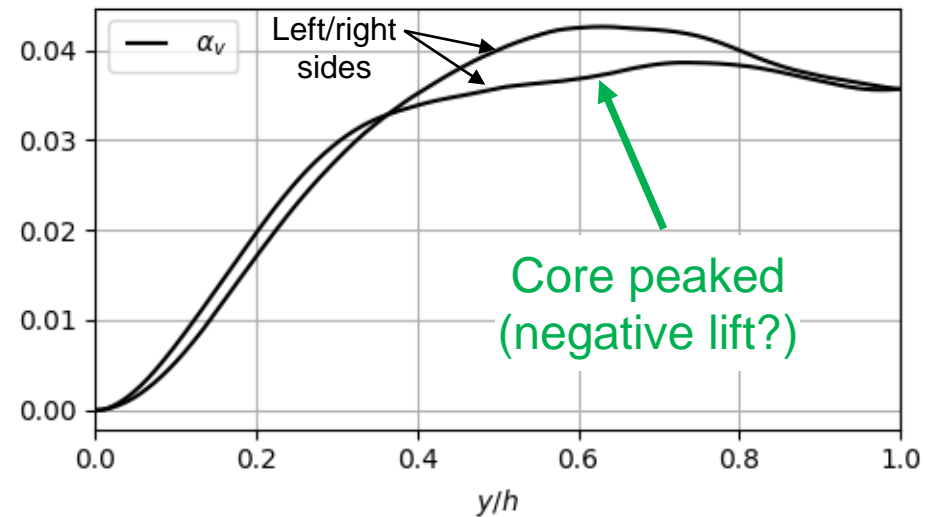
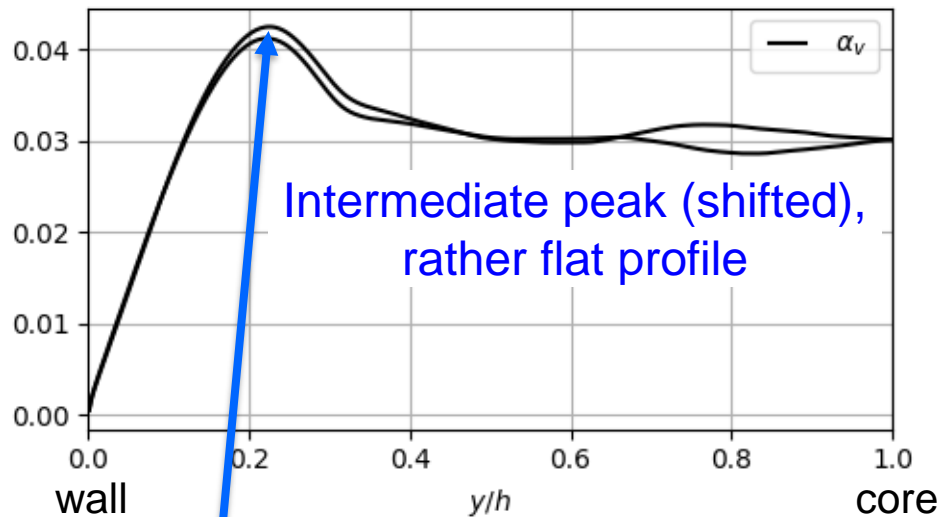
	S	D
$Re_\tau = u_\tau h / \nu_l$	180	
$Eo = \rho g d^2 / \sigma$	0.47	3.6/3.74
α_v (%)	3	
D_b/h	0.3	
N_b	80	
Size	$\pi h \times 2h \times \pi h/2$	
Resolution	$384 \times 1152 \times 192 \approx 85M^{ons}$	
Mesh size	0.3 to $3 h/Re_\tau$	
Fluid	Fictitious	
Gravity	0.1/0.8 (non-dimensional)	



$$Re_\tau = 180, \alpha_v = 3\%, \\ D/h = 0.3, \text{fluid...}$$

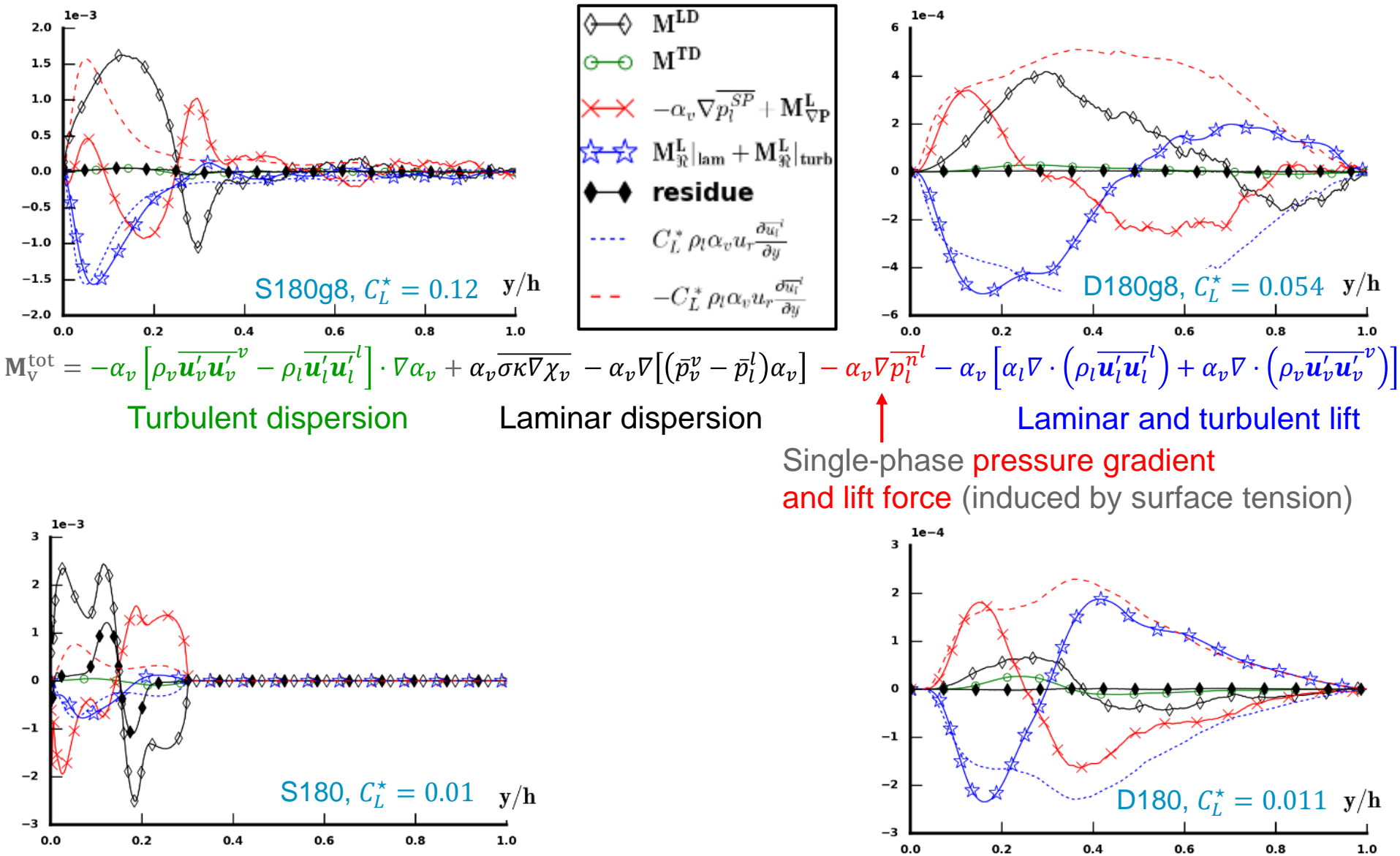






DNS of turbulent bubbly flows – Results (4/4)

Surprising contributions to the migration forces



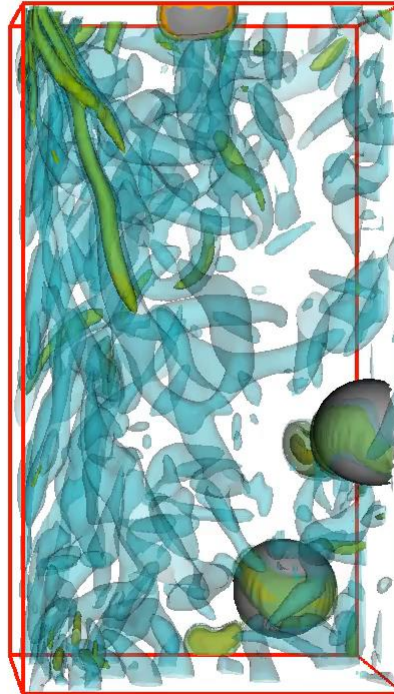
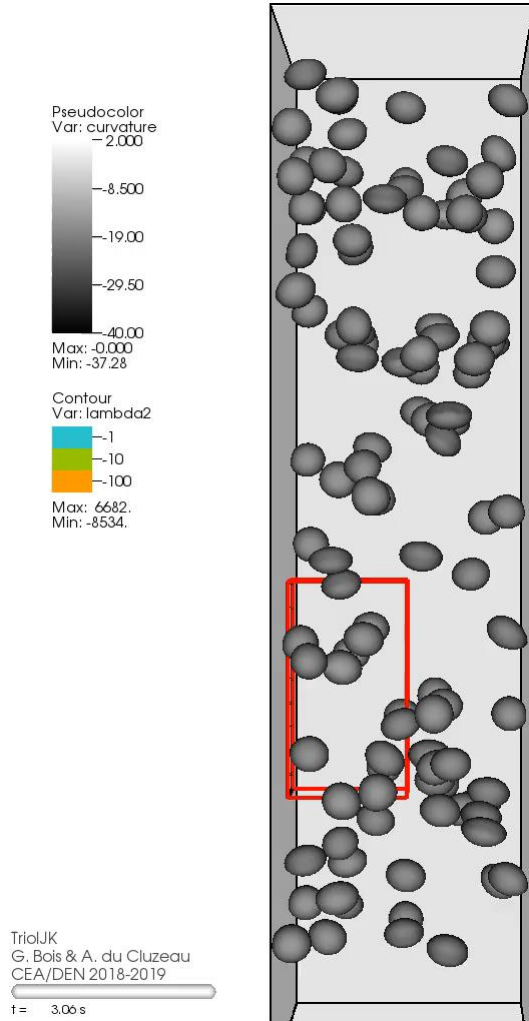
- Core, wall and intermediate peaking are depicted from DNS (through variations of σ and g)
 - Wall-to-core peak transition cannot be reduced to the inversion of the lift coefficient due to the bubble deformability (Eo).
 - Other classical hypotheses are in failure (surface tension effect, laminar Reynolds stresses...)
- Dominant role of the closure relations (interfacial forces and turbulence) on the flow dynamics

⇒ In-depths revision of hypotheses and closures.

 - Bubble-induced **turbulence** level plays a key role in the transition and must be predicted accurately
see *Session XIV – Thursday 11:30am – #168047 – A. Burlot*
Reynolds stresses splitted into classical & bubble-induced turbulence, wake induced fluctuations;
 - **Interfacial momentum transfer:**
 - The new formulation of the trajectory equation is an alternative to the classical particle approach;
 - New forces are essential in the equilibrium (related to surface tension, pressure difference and laminar Reynolds stresses (WIF) effects).
- The database in construction is:
 - Validated against other numerical work (further experimental validation is of interest);
 - Very rich: a lot of statistical information to build all kinds of models;
 - Well-designed for step-by-step analysis: DNS allows One At a Time (OAT) variations.
- Further analysis and model development are under way.
 - To derive and assess models for averaged to 2-phase RANS CFD calculations;
 - Extend the analysis to a broader parametric study and to bubble swarms.

- This work was granted access to the HPC resources of CINES under the allocations x20162b7712, A0022b7712 and A0042b7712 made by GENCI.
- Acknowledgment to CEA/DEN for the HPC ressources allocated on COBALT super-computer at TGCC.





Thank you for
your attention.

Questions?

Dr Guillaume BOIS – Research engineer
DEN/DM2S/STMF/LMSF

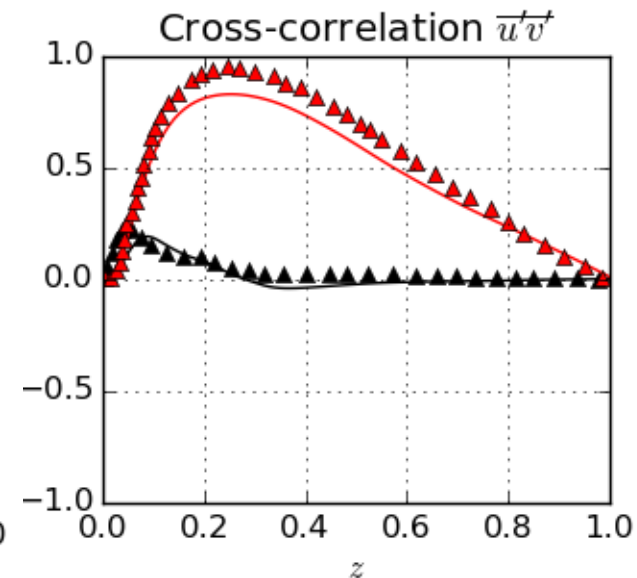
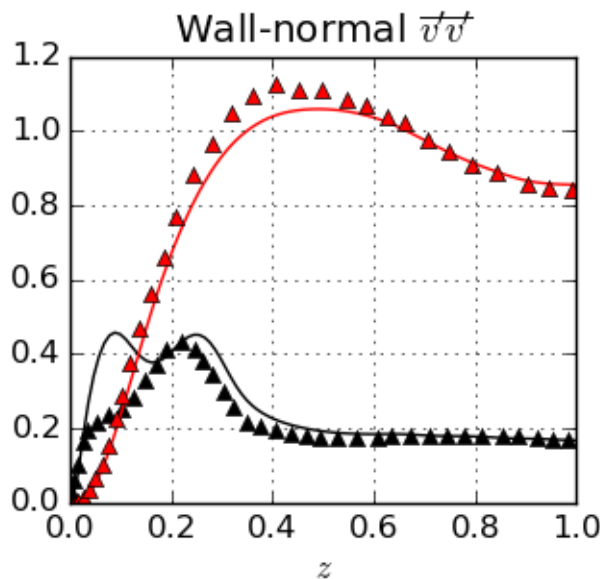
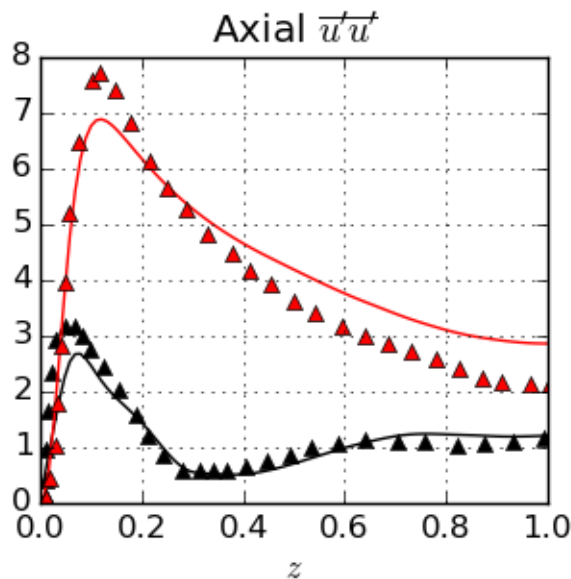
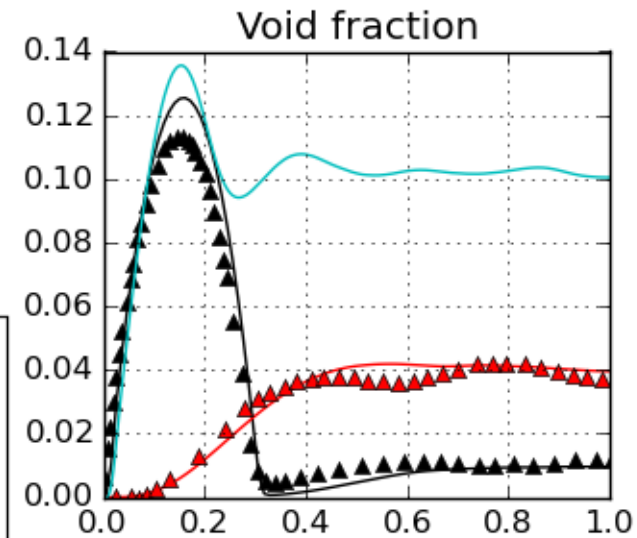
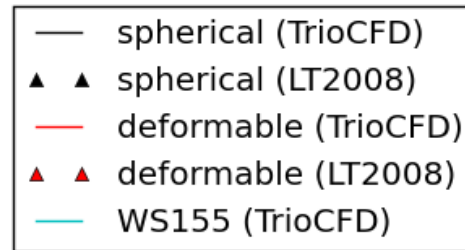
Commissariat à l'énergie atomique et aux énergies alternatives
Centre de Saclay | 91191 Gif-sur-Yvette Cedex
T. +33 (0)1 69 08 69 86 | F. +33 (0)1 69 08 68 86
guillaume.bois@cea.fr

Etablissement public à caractère industriel et commercial | R.C.S Paris B 775 685 019

Very good agreements;

Small discrepancies can reasonably be attributed to:

- Higher statistical convergence of our results (longer time);
- Small differences in lagrangian mesh management;
- Discontinuous properties and sharp interfacial force treatment;
- Uniform vs. non-uniform eulerian mesh.



- Local instantaneous description (continuum fluid mechanics)

- Navier-Stokes equations:

$$\frac{\partial \rho_k \mathbf{u}_k}{\partial t} + \nabla \cdot (\rho_k \mathbf{u}_k \mathbf{u}_k) = -\nabla p_k + \rho_k \mathbf{g} + \nabla \cdot \boldsymbol{\tau}_k \quad \text{with} \quad \boldsymbol{\tau}_k \hat{=} \mu_k (\nabla \mathbf{u}_k + \nabla^T \mathbf{u}_k)$$

- Interfacial jump conditions:

- Velocity continuity: $u_1^n = u_2^n$ and $u_1^t = u_2^t$
- Interfacial normal stress balance: $\sum_k (p_k \mathbf{n}_k - \boldsymbol{\tau}_k \cdot \mathbf{n}_k) = -\sigma \kappa \mathbf{n}$

- Extension to full space

- Multiply by phase indicator function χ_k : 1 in phase k, 0 otherwise.

$$\frac{\partial \chi_k \rho_k \mathbf{u}_k}{\partial t} + \nabla \cdot (\chi_k \rho_k \mathbf{u}_k \mathbf{u}_k) = -\nabla (\chi_k p_k) + \chi_k \rho_k \mathbf{g} + \nabla \cdot [\chi_k \mu_k (\nabla \mathbf{u}_k + \nabla^T \mathbf{u}_k)] - (\mathbf{p}_k \mathbf{n}_k - \boldsymbol{\tau}_k \cdot \mathbf{n}_k) \cdot \nabla \chi_k$$

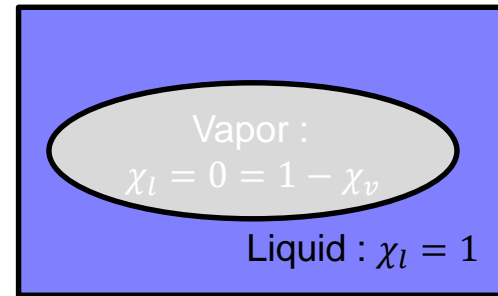
- One-fluid formulation

- Definition of “one-fluid” fields:

$$\phi \hat{=} \sum_k \chi_k \phi_k$$

- Adding up and using jump conditions:

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = -\nabla p + \nabla \cdot [\mu (\nabla \mathbf{u} + \nabla^T \mathbf{u})] + \sigma \kappa \mathbf{n} \delta^i$$



- Combined “one-fluid” formulation valid at any point in the sense of distributions

- Phase-indicator function χ_k is advected by the local velocity field (mixed VOF/FT algorithm)

DNS of turbulent bubbly flows

One-fluid formulation and Front-Tracking

Local instantaneous Navier-Stokes equations...

$$\frac{\partial \rho_k \mathbf{u}_k}{\partial t} + \nabla \cdot (\rho_k \mathbf{u}_k \mathbf{u}_k) = -\nabla p_k + \rho_k \mathbf{g} + \nabla \cdot \boldsymbol{\tau}_k \quad \text{with} \quad \boldsymbol{\tau}_k \hat{=} \mu_k (\nabla \mathbf{u}_k + \nabla^T \mathbf{u}_k)$$

$$\nabla \cdot \mathbf{u}_k = 0$$

... and interfacial jump conditions:

- Velocity continuity: $u_1^n = u_2^n$ and $u_1^t = u_2^t$
- Interfacial normal stress balance: $\sum_k (p_k \mathbf{n}_k - \boldsymbol{\tau}_k \cdot \mathbf{n}_k) = -\sigma \kappa \mathbf{n}$

... are solved by the “one-fluid” formulation valid at any point in the sense of distributions:

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = -\nabla p + \nabla \cdot [\mu (\nabla \mathbf{u} + \nabla^T \mathbf{u})] + \sigma \kappa \mathbf{n} \delta^i$$

$$\nabla \cdot \mathbf{u} = 0$$

“one-fluid” fields: $\phi \hat{=} \sum_k \chi_k \phi_k$ with χ_k the phase indicator function.

The phase-indicator function χ_k is advected by the velocity field (mixed VOF/FT algorithm)

- Interfaces are explicitly tracked;
- All scales are resolved in each phase

Normal, curvature
Phase indicator, velocity, pressure
(MAC arrangement)

