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# Evidence of large infrasonic radiation induced by earthquake

## interaction with alluvial sediments.

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#### Abstract

The M<sub>w</sub> 5.9 Ferrara earthquake that struck Northern Italy in May, 20<sup>th</sup>, 2012, was recorded with an infrasound array at a source-to-receiver distance of 300 km. The infrasound record revealed early and late detections characterized by large back-azimuth variations suggesting the existence of an extended area of infrasound radiation. Unlike most of previous studies, the modeled area of maximum infrasound radiation appears to mimic an extended flat area (plain of Po river) with no significant contributions from nearby mountain ranges. The shake map of the earthquake and the map of reported acoustic boom is in good agreement with the modeled area of infrasound radiation suggesting how the transition of seismic waves into acoustic atmospheric waves is efficiently exciting infrasound recorded at far distances from the source. Such a result is in agreement with the significant seismic amplification within the Po plan alluvial sediments.

#### 1. Introduction

Earthquakes are a well-known sources of atmospheric pressure waves [e.g. Mutschlecner and Whitaker, 2005 for a review] that are inferred to be produced by three distinct possible source mechanisms: (*i*) pressure changes due to the vertical displacement of the seismic waves near the infrasound station, (*ii*) the local conversion from seismic waves to the sound pressure near the epicenter area [Olson et al., 2003; Arrowsmith et al., 2009], and (*iii*) pressure waves generated in remote areas by the seismic shaking of the topographic relief [e.g. Le Pichon et al., 2003; Mutschlecner and Withaker, 2005; Le Pichon et al., 2005; Green et al., 2009].

The latter kind of source mechanism, that is usually referred to as *secondary infrasound*, is reported for a wide range of magnitudes, spanning from low magnitude ( $M_L\sim4$ ) shallow events [Green et al., 2009] to mega earthquakes ( $M_L>8$ ), and is usually inferred to be strongly related to steep topographical features [e.g. Le Pichon et al. 2002, 2005]. In particular, Green et al., [2009] modeled the ground-to-air coupling of the 2007 Folkestone earthquake (UK) as being produced by the shaking of the vertical coastal cliffs induced by the nearby (< 5 km epicentral distance) seismic event. Similarly, Arrowsmith et al., [2009] identified from multiple array observations of an earthquake sequence in Nevada a repeating secondary source as being produced by an isolated mountain peak (about 300 m high).

In this study we present infrasound observations of the 2012, ML 5.9, Ferrara earthquake (Italy) as recorded by a small aperture array deployed in the Northwestern Italian Alps at an epicentral distance of  $\sim 300$  km. Infrasound observations are used to infer the location and extent of the

infrasound radiant area, in terms of ground-to-atmosphere coupling of seismic waves, by considering the seismic and infrasonic propagation from the source to the array. Our results suggest the existence of an extended radiant area that appears to match the Po alluvial plan thus providing a new hypothesis on the generation of infrasound from earthquakes.

#### 2. The seismic sequence

On May 20, 2012, at 02:03 UTC, a 5.9  $M_W$  earthquake occurred in the Po plain, Northern Italy, west of the city of Ferrara at a hypocentral depth of 6.3 km (Figure 1: (a) Epicenter of the May, 20th, 2012 Ferrara earthquake (red star) and position of the CHA infrasound array (blue triangle) at a distance of ~300 km from the earthquake epicenter. Raw infrasound data (b) and spectrogram (c) recorded at the CHA array. (d) back-azimuth (ba) and apparent velocity (av) of infrasound detections obtained for the infrasonic signal at the CHA array. Earthquake origin time (red line) and celerity values that would correspond to early and late arrivals (410 and 205 m/s, blue lines) are shown for reference.). The event caused 25 casualties and extensive damage in the area and it was clearly felt in Northern and Central Italy, up to epicentral distances of ~300 km. During the 13 days following the main shock six events with magnitude  $M_W$ >5 occurred, peaking on May 29th, with the 5.8  $M_W$  events ~12 km WSW of the main shock. The seismic sequence of more than 1000 events with a local magnitude MI>3, lasted for ~3 weeks and developed along a south dipping normal fault [Piccinini et al., 2012].

Seismic moment for the main shock was  $7x10^{24}$  dyne-cm and relative source time functions calculated for the event show that a great part of the energy was radiated by a source propagating towards WSW, and this would correspond to an oblique, down-dip rupture propagation [Piccinini et al., 2012]. However, the azimuthal distribution of the relative source time functions (RSTF) amplitude

suggests a more complex pattern of rupture propagation, which was interpreted in terms of a secondary rupture front, which propagated towards the East, roughly parallel to the fault strike [Piccinini et al., 2012].

The observed extensive damage is partly to be explained in terms of seismic site effects in the Po Plain syntectonic alluvial basin. Surrounded by the Alps to the north and by the Appennines mountain ranges to the south, the Po Plain is filled with Plio-Pleistocene sediments, with depths varying from few hundred meters to several kilometers. Amplification of seismic ground-shaking was observed for frequencies between 0.5 and 1.5 Hz, as a consequence of the shallow sediments [Bordoni et al., 2012; Priolo et al., 2012].

#### 3. Infrasound array observation of the M<sub>W</sub> 5.9 earthquake

Infrasound from the main shock was clearly recorded at the CHA infrasound array deployed at a source-to-receiver distance of 294 km from the earthquake epicenter (Figure 1a). The CHA infrasound array is a small aperture ( $\sim$ 140 m), 4-elements array, deployed at an elevation of  $\sim$  2000 m a.s.l. in Champoluc (AO), in the Northwestern Italian Alps (Figure 1a). The array is equipped with 4 Optimic TM 2180 microphones, with a sensitivity of 100 mV/Pa, a low-pass cutoff frequency of 0.5 Hz and instrumental self-noise of -70 dB (relative to 1 Pa²/Hz @ 1 Hz).

On May  $20^{th}$ , 2012, around 02:05 UTC, the four microphones of the CHA array detected the seismic shaking of the ground, that was followed, approximately ten minutes later (02:15 UTC), by a long lasting infrasonic signal (Figure 1b). The seismic wave took approximately 40 seconds to propagate 294 km from the epicenter to the array, with a mean propagation velocity of 7.3 km/s consistent with the Moho discontinuity refracted p wave arrival.

Infrasound produced by the earthquake is recorded between 02:15 and 02:27 UTC as an emergent spindle-shaped signal of long duration, with a peak amplitude of 0.1 Pa in the 1-3 Hz frequency band (Figure 1b,c). The array analysis indicates a continuous arrival of infrasound detected with a stable apparent velocity of  $\sim$  340 m/s and varying back-azimuth (spanning  $\sim$  30°) for the whole emergent long-lasting signal (Figure 1d). The tail of the infrasonic wave-packet is showing an amplitude modulation possibly reflecting multiple sources or arrivals with varying energy from an extended source. This second hypothesis is more consistent with the observed smooth variation of infrasound back-azimuth and stable spectral content.

The timing of the infrasound detections (02:15-02:27 UTC) with respect to the earthquake onset time (02:03 UTC) is not consistent with infrasound being radiated only at the epicenter. At a source-to-receiver distance of 296 km in fact, the corresponding celerity would span between 205 m/s for late detections around 02:27 and 410 m/s for early detections around 02:15 UTC (Figure 1b). Here, only infrasound recorded around 02:20 UTC with back-azimuth of  $\sim$  110°N pointing to the earthquake epicenter appears consistent with a celerity of  $\sim$  300 m/s, thus suggesting stratospheric arrivals of primary infrasound produced at the epicenter. This timing of infrasound detections is rather suggesting the complex generation of infrasound from both primary and secondary sources.

Moreover, while propagation velocity is extremely stable at 340 m/s, back-azimuth is actually showing a large variation up to 30° (Figure 1d) from the real back-azimuth to the earthquake epicenter (110°N). Back-azimuth is observed to increase from the initial value of ~ 90° N (around 02:15 UTC) moving southward up to a maximum value of 110°N (around 02:20 UTC), and to decrease back to ~95°N at the end (around 02:27 UTC) of the emergent phase (Figure 1d). Such a large azimuth variation is the evidence of an extended radiant area for the recorded infrasound.

#### 4. Modeling the source area by FDTD analysis

In order to evaluate the area of maximum infrasound radiation, we followed the procedure described by Arrowsmith et al., [2009] and modeled the infrasound source area from the timing and back-azimuth of infrasound detections at the CHA array (Figure 1c), in terms of ground-to-atmosphere coupling of seismic waves. Here, seismic waves are assumed to radiate spherically from the earthquake hypocenter, propagating in the ground at constant velocity and then coupling to the atmosphere to propagate as infrasound. Hypocentral location ( $\lambda_e = 44.896$ ,  $\varphi_e = 11.264$ ,  $h_e = 9.5$  km) and origin time of the event ( $t_e = 02:03:50$  UTC) are given by accurate seismic location by the Italian national seismic network. The infrasound source area is identified with a grid searching procedure, which minimizes the difference between real and modeled back-azimuth and arrival time of infrasound detections.

The searching grid of 71x165 nodes covers Northern Italy, extending 4 degrees in latitude and 8 degrees in longitude and with a grid spacing of 0.05 degrees. We assume each node (i,j) of the grid as being a possible source of infrasound due to ground coupling of the seismic wave within the atmosphere and calculate the expected back-azimuth at the array  $(az_{ij})$  and the seismo-infrasonic travel time  $(t_{ij})$ , with a seismic wave propagating from the hypocenter to the grid node (i,j) at steady velocity and infrasound wave propagating from the grid node to the array along great-distance circles.

The expected back-azimuth of each grid node (i,j) from the array  $(az_{ij})$  is evaluated directly from the great-circle bearing angle for a spherical earth:

$$az_{ij} = \tan^{-1} \left[ \frac{\sin(\varphi_a - \varphi_{ij})\cos(\varphi_a)}{\cos(\lambda_{ij})\sin(\lambda_a) - \sin(\lambda_{ij})\cos(\lambda_a)\cos(\varphi_a - \varphi_{ii})} \right],$$
 (eq.1)

being  $\lambda_{ij}$  and  $\varphi_{ij}$  the latitude and longitude of a given point (i,j) of the searching grid and  $\lambda_a$  and  $\varphi_a$  the coordinates of the central element of the array. Possible ray deflection due to transverse winds is neglected here.

The travel time from the hypocenter to the array  $(t_{ij})$  is calculated as the sum of the seismic travel time from the hypocenter to each node of the searching grid  $(t_s_{ij})$  and the infrasonic travel time from the node to the array  $(t_{ij})$ .

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$$t_{ij} = t_{-}s_{ij} + t_{-}i_{ij}$$
. (eq. 2)

In order to evaluate the seismic travel-time for each node  $(t_s)$  we assumed spherical propagation from a point source located at the earthquake hypocenter  $(\lambda_e, \varphi_e, h_e)$  into a homogeneous half-space with constant velocity of 5.8 km/s. This value is in accordance with first arrivals recorded at seismic stations of the Italian National Seismic Network up to distances of 150 km from the source and represents the mean velocity in the crust [Piccinini et al., 2012]. Here, any directivity of the source or heterogeneity of the medium is neglected. This assumption is however of minor importance, being the seismic propagation velocity one order of magnitude larger than the sound propagation velocity in the atmosphere, thus the uncertainty of the velocity structure in the crust has a limited effect on the modeled timing of the infrasound detection  $(t_{ij})$ .

The infrasonic travel-time from the grid to the CHA array ( $t_{-ij}$ ) is obtained with a 2D FTDT modeling [De-Groot-Hedlin et al., 2011; Lacanna et al., 2014] of the pressure wave propagation in the atmosphere, in order to account for wind effects and atmospheric profile. In particular, the 2D FTDT analysis was applied on 36 profiles centered in the CHA array and spaced by 10 degrees along great circle distances and covering the whole area of the searching grid. Total length of the profiles spans from a minimum of 84 km for the section with back-azimuth of 270 °N to a maximum of 580 km for

section with back-azimuth of 110 °N. Wind, temperature and density for each section from the ground up to an elevation of 70 km are obtained by interpolating ECMWF High-Resolution atmospheric model (HRES) analysis at 91 mean pressure levels up to 0.01 hPa (L91) with a spectral resolution of ~12 km. for the time of occurrence of the event and for the area of analysis.

For each section, a Gaussian-shaped pulse with a frequency of 0.2 Hz was used as the source time function for the 2D FDTD modeling and a lattice grid size of 25 m was applied with a time discretization of 0.0156 seconds, which satisfies the stability conditions of the FDTD. Here, the infrasonic travel-time to the array is calculated every 5 km along each section (Figure 2). Results obtained for all the 36 profiles are then interpolated across the searching grid and the corresponding infrasonic travel-times ( $t_{ij}$ ) are evaluated. The seismo-infrasonic travel time ( $t_{ij}$ ) is eventually calculated over the whole searching grid (2c) by adding the seismic ( $t_{ij}$ ) and infrasonic ( $t_{ij}$ ) counterparts. 2D FDTD modeling predicts direct arrivals for short distances and stratospheric arrivals for longer sections, with maximum seasonal wind blowing to the west at an altitude of about 50 km strongly modifying infrasonic travel times from spherical symmetry along east-west profiles (Figure 2, Figure 3a).

For each infrasonic detection at the CHA array, we measured the back-azimuth  $(az_d)$  and the travel time  $dt_d = t_d - t_e$ , as the difference between the timing of the infrasound detection  $(t_d)$  and the origin time of the earthquake  $(t_e)$  as provided by the independent seismic event location.

In order to evaluate the most probable position of the source of each infrasound detection (d), the normalized difference between the observed and theoretical values of back azimuth ( $dAZ_d(i,j)$ ) and travel time ( $dT_d(i,j)$ ) are evaluated:

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$$\begin{cases} dAZ_{d}(i,j) = \frac{|az_{ij} - az_{d}|}{\max(|az_{ij} - az_{d}|)} \\ dT_{d}(i,j) = \frac{|t_{ij} - dt_{d}|}{\max(|t_{ij} - dt_{d}|)} \end{cases}$$
 (eq. 3)

for each node (i,j) of the grid leading to matrices with values ranging between 0 and 1. When the node (i,j) coincides with the position of the real source the difference between theoretical and measured azimuth  $(az_{ij}-az_d)$  and travel time  $(t_{ij}-dt_d)$  will be zero. Therefore, the node with the lowest values of  $dAZ_d$  and  $dT_d$  will most probably represent the position of the infrasonic source of a given detection (d).

In order to account for both back-azimuth and travel time in the searching procedure, the two matrices are eventually multiplied  $(M_d(i,j) = dAZ_d(i,j) \times dT_d(i,j))$  leading to a matrix with values ranging between 0 and 1, and the possible source position of secondary infrasound is identified in the node of the grid  $(i_{do}, j_{do})$  where the matrice  $M_d(i,j)$  is minimum (Figure 3b).

All the minima evaluated for all the infrasound detections are then combined in a single map, that represents the whole infrasound radiation area (Figure 3b) induced by seismic wave coupling to the atmosphere. The result highlights an extended area of infrasound radiation, that from the epicenter develops for ~100 km to the east along the Po river up to the Po delta, while develops ~100/150 km toward the north-west up to the Garda Lake and the Alpine mountain range (Figure 3b).

#### 5. Modeling the source area by back-projection of infrasound back-azimuth

In addition to the procedure presented above, a location of infrasound detections was performed, considering stratospheric arrivals and transverse wind effects on the propagation path. Detections after 02:25 UT were not considered. Assuming these detections produced from seismic waves originated from the epicenter (not from a spatially extended source), the corresponding locations are found over sea unless considering unrealistic celerity values (<0.2 km/s).

Ray-tracing simulations using the WASP-3D ray theory-based method which account for the the spatiotemporal variations of the horizontal wind terms along the ray paths [Virieux et al., 2004] and the

ECMWF analyses, were performed assuming a source altitude of 1 km, slowness values ranging between 2.6 and 3 s/km. 200 rays were launched with incidence angle equally spaced. Simulations predict a dominant stratospheric waveguide with a refracting height up to 44-54 km altitude, being characterized by a celerity of 0.295 and 0.305 km/s and azimuthal deviation ranging from -4° to 3° depending on the ray trajectory.

The location of the infrasound sources are calculated using an inverse location procedure which combines the seismic source information (epicenter coordinates and origin time), celerity models both infrasound and seismic waves, and the arrival times and wind-corrected azimuth of infrasound waves [e.g. Le Pichon et al., 2002]. The location of each detected infrasonic signal is then back-projected on its back-azimuth with a distance constrained by seismic propagation time from the hypocenter to any source point and travel time though the atmosphere from the source point to CHA. Taking into account the propagation variability due to atmospheric uncertainties and errors in the wave front measurements taking into account the station geometry, a maximum location error of ~ +/-25 km is estimated [Szuberla and Olson, 2004; Le Pichon et al., 2015].

Using these simulation results, assuming stratospheric arrivals, a density map of infrasound detections is obtained (Figure 3c). This results into a good agreement with the map obtained with FDTD analysis and identify an extended area of infrasound radiation, which from the earthquake epicenter extends to the east up to the delta of PO river and to the north-west up to the Garda lake and the southern margin of the Alpin chain.

#### 6. Discussion and conclusions

The analysis of infrasound detections at the CHA array, combined with the seismological information available on the location and origin time of the May, 20<sup>th</sup>, 2012 Ferrara earthquake,

allowed to identify a 200 km extended source area of infrasound, spanning from the southern flank of Alpine chain to the north-east ( $\lambda=46^{\circ}$ ,  $\varphi=10^{\circ}$ ), to the delta of Po river ( $\lambda=45^{\circ}$ ,  $\varphi=12.5^{\circ}$ ) to the east (Figure 3) and showing a preferential distribution in the East-West direction from the epicenter to the coastline. Results obtained from FDTD analysis of pressure wave propagation in the atmosphere (Figure 3b) and back-projection of wind-corrected azimuth of infrasound waves (Figure 3c) are highly consistent with each other. Such an extended radiant area is in agreement with the long lasting infrasonic signal and the observed variations of back-azimuth (Figure 1b, d) and might also explain the observed amplitude modulation of the infrasonic wave-packet. Infrasound produced at the epicenter is recorded with a back-azimuth of 110°N around 02:20 UTC, consistent with a celerity of 300 m/s for the epicentral distance of ~300 km. Infrasound radiated from the most western portion (~90°N) of the area is recorded few minutes before the epicentral infrasound (around 02:15 UTC), as a consequence of the shortest distance (~200 km) to the array. In the same way, infrasound radiated east of the epicenter, is recorded with a back-azimuth of ~ 95°N few minutes after the epicentral infrasound (around 02:27 UTC), as a consequence of the longer distance (~400 km) traveled by the infrasonic wave. The radiation area is clearly limited to the coastline to the east, being ground-to-atmosphere coupling of seismic waves unlikely in the sea, and reported up-to-now only for major earthquake [Evers et al., 2014]. This result is a good validation of the proposed modeling. It is worth noting that the area of maximum infrasound radiation modeled from our infrasound

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observation is actually confined within the Po plan, with no significant infrasound radiated by the Alpine and/or Appenines mountain ranges that are actually boarding the Po plain to the north and to the south. We can consider the whole Po plan as almost immediately shacking for the earthquakes, with seismic waves inducing infrasound in the atmosphere. We suggest therefore that the ground shaking of the Po plan is the most likely source mechanism of the secondary infrasound. This hypothesis is corroborated by the spectral content of recorded infrasound (1-2 Hz) peak frequency, that is consistent

with the (0.5-1.5 Hz) peak frequency of the ground-shaking in the PO plan as the effect of the amplification effects due to the soft sediments of the PO valley [Priolo et al., 2012; Bordoni et al., 2012].

This conclusion differs significantly from previous studies [Le Pichon et al., 2003; Arrowsmith et al., 2009; Green et al., 2009], where secondary infrasound was strongly related to topographic relief, while it appears consistent with the work of Walker et al., [2013], where both enhanced topography and flat areas where identified from infrasound observations from the Tohoku earthquake. Anyway, based on our observations we can conclude that the Po plain shaking was definitely the most energetic source of the infrasonic observed detections.

Infrasound detections of the main shock are compared with infrasound detection for the most energetic aftershock (2012/05/29, Ml=5.8) recorded during the 2012 seismic sequence (Figure 4). In both cases infrasound detections show the same pattern of back-azimuth variation with time and delay from the earthquake onsets, confirming the modeled extended source mechanism. Figure 4 suggests also how infrasound detections from source areas far from the earthquake epicenter (i.e. recorded before minute 10 or after min 20 from the earthquake epicenter) might depend on the magnitude and depth of the event.

In order to analyze the mechanism of secondary infrasound radiation by the earthquake, we compared the modeled radiant area, obtained both from FDTD analysis (Figure 3b) and back-projection of wind-corrected azimuth of infrasonic waves (Figure 3 c), with shake map and the acoustic effect of the event (Figure 3d). Earthquake booms have been reported for a long time [Michael, 2011 for a review] and are commonly explained as being produced by the refraction of the "p" wave into the atmosphere very close to the listener. The map is obtained interpolating 8766 reports distributed over 1645 municipalities and shows the percentage of earthquake booms felt in a given municipality with respect of the number of reports. The map of earthquake boom shows how the earthquake boom was

clearly felt in a wide area extending from the shoreline to the west, where  $\sim$ 70% of the people reporting the event felt a clear earthquake boom, towards the Alpine mountain chain to the north-east, where  $\sim$ 45% of the people reported the earthquake boom at distances of  $\sim$  200 km from the epicenter. For epicentral distances exceeding 250-300 km this effect is minimal in all directions. Similarly, the shake map of the event, derived from  $\sim$  120 accelerometer record of the Italian National Network (RAN), confirms maximum peak ground acceleration in the epicentral area as well as enhanced ground acceleration North/West and East of the epicenter, in good agreement with the earthquake boom.

We suggest that the good match observed between seismic observation of the event (seismic shake map and map of felt earthquake boom) and the modeled area of maximum infrasound radiation (Figure 3d) is confirming our finding of infrasound being mostly radiated in the alluvial plan. In this context infrasound radiation appears to be controlled by the local shallow geology, with the east-west extension of the infrasonic source area from the epicenter (Figure 3b,c) to the coastline matching the position of the Po and Adige rivers in the alluvial plan, whose recent deposits are sites of maximum seismic amplifications and are characterized by enhanced peak ground acceleration and acoustic effect of the earthquake (Figure 3d). At the same time however, this preferential trend is in agreement with the geometry of the fault plane, thus possibly suggesting a source effect of the geometry and extend of the infrasound radiant area.

#### 7. Data and Resources

Infrasound data of the May, 20<sup>th</sup>, 2012 Ferrara earthquake have been recorded a permanent infrasound array operated in Champoluc (AO, Italy) by the Department of Earth Sciences of the University of Firenze. Raw infrasound data are freely available upon request to the corresponding author. Location and occurrence time of the earthquake is obtained from the Italian seismological instrumental and parametric data-base (ISIDE, http://iside.rm.ingv.it/iside/standard/index.jsp). Data of

earthquake boom are obtained from the database of the Istituto Nazionale di Geofisica e Vulcanologia (<a href="http://www.haisentitoilterremoto.it/">http://www.haisentitoilterremoto.it/</a>). Accelerometer data of the event, used for the shake map, are obtained from the database of the National Accelerometer Network (<a href="http://www.protezionecivile.gov.it/jcms/it/ran.wp">http://www.protezionecivile.gov.it/jcms/it/ran.wp</a>).

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#### 9. References

Arrowsmith, S. J., R. Burlacu, R. Whitaker, and G. Randall (2009). A repeating secondary source of infrasound from the Wells, Nevada, earthquake sequence. Geophys. Res. Lett., 36, L11817. doi:10.1029/2009GL038363.

Bordoni, P., R. M. Azzara, F. Cara, R. Cogliano, G. Cultrera, G. Di Giulio, A. Fodarella, G. Milana, S. Pucillo, G. Riccio, A. Rovelli, P. Augliera, L. Luzi, S. Lovati, M. Massa, F. Facor, R. Puglia, G. Ameri (2012), Preliminary results from EMERISTO, a rapid response network for site-effect studies. *Annals of Geophysics*, *55*(4), 599-607, doi:10.4401/ag-6153.

de Groot-Hedlin, C., M. A. H. Hedlin, and K. Walker (2011), Finite difference of infrasound propagation through a windy, viscous atmosphere: Application to a bolide explosion detected by seismic networks, Geophys. J. Int., 185, 305–320, doi:10.1111/j.1365-246X.2010.04925.x.

329	Evers, L. G., D. Brown, K. D. Heaney, J. D. Assink, P. S. M. Smets, and M. Snellen
330	(2014), Evanescent wave coupling in a geophysical system: Airborne acoustic signals from the
331	Mw 8.1 Macquarie Ridge earthquake, Geophys. Res. Lett., 41, 1644–1650,
332	doi:10.1002/2013GL058801
333	Green, D., J. Guilbert, A. Le Pichon, O. Sebe, and D. Bower (2009), Modelling Ground-
334	to-Air Coupling for the Shallow ML4.2 Folkestone, UK, Earthquake of 2007 April 28. Bull.
335	Seism. Soc. Am.
336	Lacanna, G., M. Ichihara, M. Iwakuni, M. Takeo, M. Iguchi, and M. Ripepe (2014),
337	Influence of atmospheric structure and topography on infrasonic wave propagation, J. Geophys
338	Res. Solid Earth, 119, 2988–3005, doi:10.1002/2013JB010827.
339	Le Pichon, A., J. Guilbert, M. Vallée, J.X. Dessa, and M. Ulziibat (2003), Infrasonic
340	imaging of the Kunlun Mountains during the great 2001 China earthquake. Geophys. Res. Lett.
341	DOI:10.1029/2003GL017581.
342	Le Pichon, A., J. Guilbert, A. Vega, M. Garcés, and N. Brachet (2002), Ground-coupled
343	air waves and diffracted infrasounds from the Arequipa earthquake of June 23, 2001, Geophys.
344	Res. Lett., 1029/2002GL015052.
345	Le Pichon, A., P. Herry, P. Mialle, J. Vergoz, N. Brachet, and M. Garcés (2005),
346	Infrasound associated with 2004-2005 large Sumatra earthquakes and tsunami. Geophys. Res.
347	Lett., 32, L19802, doi:10.1029/2005GL023893.
348	Le Pichon, A., J. D. Assink, P. Heinrich, E. Blanc, A. Charlton-Perez, C. F. Lee, P. Keckhut, A.
349	Hauchecorne, R. Rüfenacht, N. Kämpfer, et al. (2015), Comparison of co-located independent ground

350	based middle atmospheric wind and temperature measurements with numerical weather prediction
351	models, J. Geophys. Res. Atmos., 120, 8318–8331, doi:10.1002/2015JD023273.
352	Michael, A. J. (2011), Earthquake sounds: In Encyclopedia of solid earth geophysics,
353	Springer, pp. 188-192, doi:10.1007/978-90-481-8702-7_201, Dordrecht, The Netherlands.
354	Mutschlecner, J. P., and R. W. Whitaker (2005), Infrasound from earthquakes. J.
355	Geophys. Res., 110, D01108, doi:10.1029/2004JD005067.
356	Olson, J. V., C. R. Wilson, and R. A. Hansen (2003), Infrasound associated with the 2002
357	Denali fault earthquake, Alaska. Geophys. Res. Lett., 30(23) 2195, 1-4,
358	doi:10.1029/2003GL018568.
359	Piccinini, D., N. A. Pino, and G. Saccorotti (2012), Source complexity of the May 20,
360	2012, MW 5.9, Ferrara (Italy) event. <i>Annals of Geophysics</i> , 55(4), doi:10.4401/ag-6111.
361	Priolo, E., M. Romanelli, C. Barnaba, M. Mucciarelli, G. Laurenzano, L. Dall'Olio, N.
362	Abu Zeid, R. Caputo, G. Santarato, L. Vignola, C. Lizza, P. Di Bartolomeo (2012), The Ferrara
363	thrust earthquakes of May-June 2012: preliminary site response analysis at the sites of the OGS
364	temporary network. Annals of Geophysics, 55(4), doi:10.4401/ag-6172.
365	Szuberla, C. A. L., and J. V. Olson, 2004, Uncertainties associated with parameter
366	estimation in atmospheric infrasound arrays, J. Acoust. Soc. Am. 115, 253-258,
367	http://dx.doi.org/10.1121/1.1635407
368	Virieux, J., Garnier, N., Blanc, E. & Dessa, J.X., 2004. Paraxial raytracing for
369	atmospheric wave propagation, Geophys. Res. Lett., 31, doi:10.1029/2004GL020514

### Infrasound by earthquake interaction with alluvial sediment

370	Walker, K. T., A. Le Pichon, T. S. Kim, C. de Groot-Hedlin, IY. Che and M. Garcés
371	(2013), An analysis of ground shaking and transmission loss from infrasound generated by the
372	2011 Tohoku earthquake, J. Geophys. Res. Atmos., 118, 12,831-12,851,
373	doi:10.1002/2013JD020187.
374	

### 375 Figure Captions

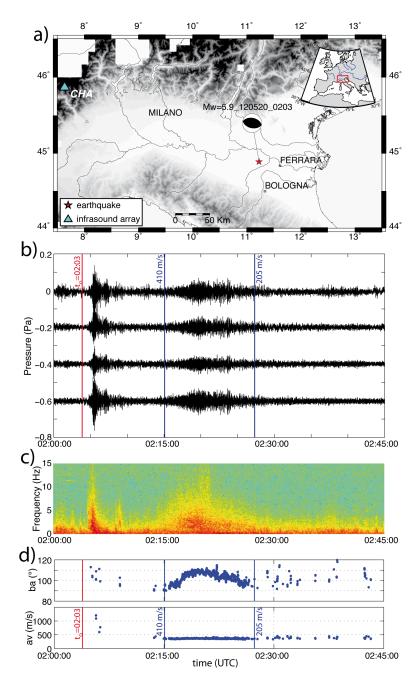
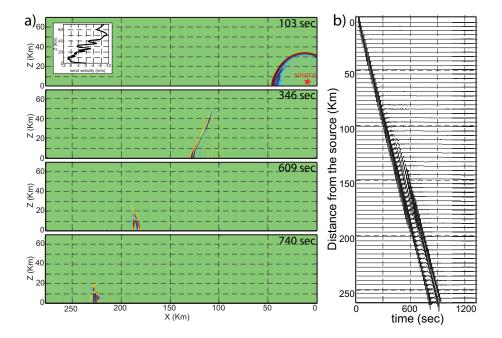


Figure 1: (a) Epicenter of the May, 20<sup>th</sup>, 2012 Ferrara earthquake (red star) and position of the CHA infrasound array (blue triangle) at a distance of ~300 km from the earthquake epicenter. Raw infrasound data (b) and spectrogram (c) recorded at the CHA array. (d) back-azimuth (*ba*) and apparent velocity (*av*) of infrasound detections obtained for the infrasonic signal at the CHA array. Earthquake origin time (red line) and celerity values that would correspond to early and late arrivals (410 and 205 m/s, blue lines) are shown for reference.



source.

Figure 2: FDTD analysis of infrasound propagation along the 180° N atmospheric profile to the array. Snapshots of the infrasound propagation at 103, 346, 609 and 740 seconds respectively (a). Synthetic waveforms evaluated at ground level every 5 km along the profile (b). Waveforms are amplitude normalized in order to enhance the arrival time and show stratospheric arrivals at distances exceeding 130-150 km from the

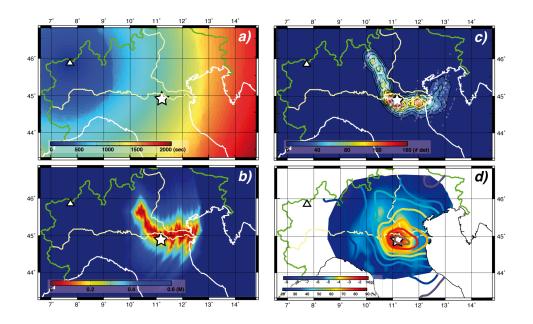


Figure 3: Theoretical seismo-acoustic travel-time (a) of infrasound radiated by secondary sources positioned in the different nodes of the searching grid as it would be recorded by the CHA array. Modeled infrasound radiation area (b) for the main May, 20, 2012, event obtained by combining differences in expected

and observed back-azimuth and travel times (a) for all the infrasound detections of the earthquake. (c) Location (white circles) and density map of infrasonic sources based on ray-tracing. (d) Distribution of earthquake booms felt in northern Italy (contour-lines) and shake map (colored map) for the 2012, May 20<sup>th</sup> event.. Contour-lines represent the percentage of felt boom within the total number of reports. In all subplots position of the earthquake epicenter (white star) and of the CHA array (white triangles) is shown as well as national border (green line) coastline (white line) and the main rivers (yellow lines).

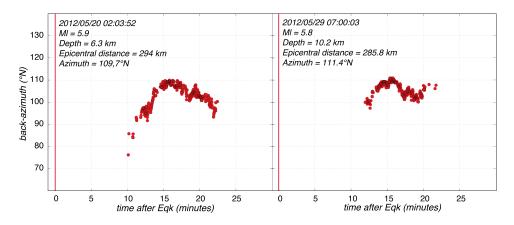


Figure 4: Infrasound detections of the infrasound produced by the May 20<sup>th</sup>, 2012 main shock (MI 5.9) and by the May, 29<sup>th</sup>, 2012 aftershock (MI 5.8). The time of the infrasound detections is expressed in terms of minutes after the earthquake occurrence time.