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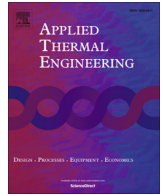
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## In situ investigation of liquid films in pulsating heat pipe



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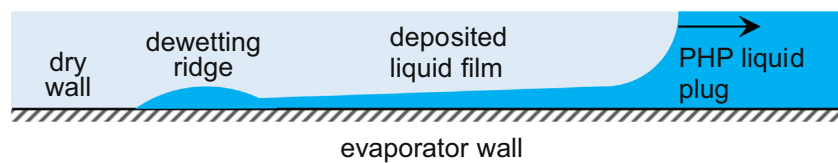
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### HIGHLIGHTS

- 3D shape of the liquid film inside PHP is reconstructed from experimental data.
- Existence of a liquid ridge adjacent to the film contact line is shown.
- Both film shape and thickness are explained theoretically.
- Film profile theory agrees with the experiment.
- Theory for the meniscus oscillation dynamics is proposed and compared to experiment.

### GRAPHICAL ABSTRACT



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### ABSTRACT

To understand functioning of the pulsating (or oscillating) heat pipe (PHP), a liquid film deposited by an oscillating meniscus is studied experimentally inside the simplest, single branch PHP, which is a straight capillary sealed from one end. The PHP capillary is of rectangular section of high aspect ratio. The evaporator is transparent so that the films can be studied by two complementary optical methods: grid deflection method and interferometry. We were able to measure the dynamic film profile during the self induced meniscus oscillations. It has been shown that the PHP films have the same origin as those of the Taylor bubbles; their thickness right after deposition is well described by the classical formulas. The film shape in PHPs differs from the classical wedge-shaped film observed in capillary heat pipes because both of the larger thickness and of the receding triple liquid-vapor-solid contact line. The film slope is very weak, with a growing in time ridge adjacent to the contact line. It is shown that this ridge is the dewetting ridge. Its dynamics is defined mainly by the capillary effects. Such results can be generalized to the conventional multi-branch PHP.

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## 1. Introduction

The pulsating (or oscillating) heat pipe (PHP) is a long capillary tube bent into many branches and partially filled with a two-phase, usually single component, working fluid [1]. During PHP functioning, a moving pattern of multiple vapor bubbles separated

by liquid plugs forms spontaneously inside the tube. Because of their simplicity and high performance, PHPs are often considered as highly promising. Their industrial application is however limited because the functioning of PHPs is not completely understood.

During the last decade, researchers have extensively studied PHPs [2,3]. It has been observed by many researchers that the main flow pattern inside the PHP is the slug flow, i.e. the flow of the “Taylor bubbles” where the gas is surrounded by liquid films. A major part of mass exchange occurs on their interface with the vapor like in the conventional heat pipes. Since the mass exchange provides both a moving force for the oscillations and the heat

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## Nomenclature

$A$	meniscus oscillation amplitude [m]
$b$	singular pressure drop coefficient
$Ca$	dimensionless capillary number
$D$	spacing between portholes [m]
$f$	oscillation frequency [Hz]
$h$	film thickness [m]
$i$	imaginary unit = $\sqrt{-1}$
$J$	evaporation mass flux [ $\text{kg}/(\text{m}^2 \cdot \text{s})$ ]
$k$	liquid heat conductivity [ $\text{W}/(\text{m} \cdot \text{K})$ ]
$L$	latent heat [ $\text{J}/\text{kg}$ ]
$l$	characteristic length scale [m]
$L_r$	added column height [m]
$p$	vapor pressure [Pa]
$T$	temperature [K]
$t$	time [s]
$U$	velocity [m/s]
$w$	dewetting ridge width [m]
$x, y$	coordinates [m]

### Abbreviations

CL	contact line
PHP	pulsating heat pipe
ITO	indium-tin oxide

### Greek symbols

$\Delta T$	internal wall superheating, $T_w - T_{sat}$ [K]
$\mu$	liquid shear viscosity [ $\text{Pa} \cdot \text{s}$ ]
$\nu$	kinematic viscosity [ $\text{m}^2/\text{s}$ ]
$\omega$	angular frequency [ $\text{s}^{-1}$ ]
$\rho$	liquid density [ $\text{kg}/\text{m}^3$ ]
$\sigma$	surface tension [N/m]
$\tau$	shear stress [Pa]
$\theta$	interfacial slope, contact angle

### Subscripts

0	right after film deposition
app	apparent
c	condenser
cl	contact line
i	initial (at $t = 0$ )
L	liquid
m	meniscus
r	reservoir
sat	saturation
V	Voinov
v	visible
w	internal tube wall

exchange, the films are extremely important for the PHP functioning.

In this presentation we concentrate on the film behavior in the evaporator because, for the PHP case, it is more complex than in the condenser where the film thickness grows more or less homogeneously. In the evaporator, the film thinning may lead to its localized drying so that its area changes and the triple liquid-vapor-solid contact line (CL) appears. The most important issues for the PHP theory are the film area and thickness. First we consider the film thickness and its evaporation dynamics and next, the CL effects that turn out to be crucial to understand the film area evolution.

## 2. Experimental

The studies of pure two-phase fluid case in presence of the phase change are much more scarce than in the isothermal case and for the case of different liquid and gas (e.g., water-air), see e.g. [4]. One can mention observations of condensation [5], with application to the LHP condenser. The film evaporation studies are more important for the PHP case where the partial evaporator drying occurs frequently. To our knowledge, the first film observation has been made in the U-turn (two branch) water PHP [6]. Several film observations have been carried out since then [7–9]. The film length has been measured in the single branch PHP [8]. The film profile direct measurements in the capillaries of the circular section are very difficult. For instance, it is impossible to use the confocal microscopy [9] in the cylindrical tubes, although such a method was successful for the capillary HP case. The film profile reconstruction based on the measurements of the spatial temperature distribution with the infra-red camera lacks precision either [10]. It is possible to know only the film length by this method [11].

For the present experiment, we have chosen the two-dimensional capillary, i.e. that of rectangular cross-section with the high aspect ratio,  $2 \times 22 \text{ mm}^2$ . The experiments are carried

out with simplest, single branch PHP. It is a vertical capillary of about 83 cm total length sealed from one (top) end (Fig. 1a). Its bottom end is open to a large closed two-phase reservoir regulated at a constant vapor pressure  $p_r$ . Since the films in different bubbles are independent in the multibranch PHP, such a film study is quite general. The working fluid is the pure ethanol. A special care is taken to extract all the impurities in situ with a special purifying device so that the saturation pressure corresponds to its reference value. The middle part of capillary is the condenser maintained at a constant temperature  $T_c < T_r \equiv T_{sat}(p_r)$  with a thermal bath. The upper part of capillary is the evaporator formed by two transparent sapphire portholes (Fig. 1b) fixed on a copper structure. ITO (indium-tin oxide) layers are deposited at each porthole (outside). They are heated independently with electric current. The evaporator is separated from the condenser by 1 cm-high adiabatic section, made with a thermally insulating material

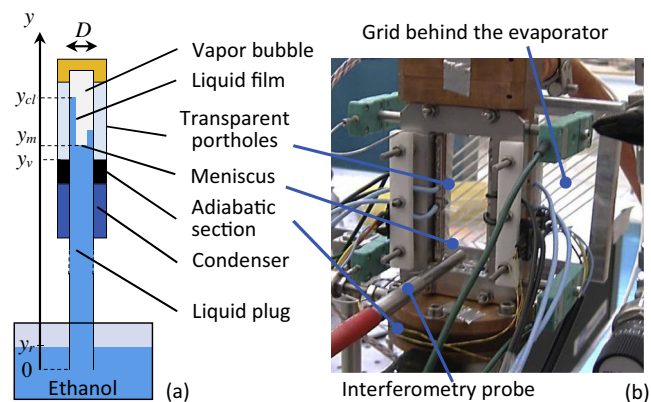


Fig. 1. (a) Sidewise scheme of the whole PHP and (b) a zoom of the PHP upper part (front view of evaporator and adiabatic section). The setup is backlit with the diffuse light and is filmed by a camera from the front.

(Vespe<sup>®</sup>). The internal cross-section of different capillary parts is matched with a high precision to avoid as much as possible the CL pinning at their junctions. Two thermocouples are glued externally at the top and at the bottom of each porthole to measure their temperatures. The vapor pressure  $p$  is measured by the fast response pressure transducer located at the top of the evaporator. The evaporator is filmed with a video camera. The portholes are heated with the transparent ITO layers. A much higher power is applied to the rear porthole so that the film on it disappears quickly and the front film can be studied alone.

The meniscus is brought into the evaporator by lowering  $T_c$  (which leads to  $p$  decrease). The liquid meniscus can vertically oscillate around its equilibrium position. These oscillations become self-sustained when an imposed axial thermal gradient causes periodic evaporation and condensation of the fluid [12]. When the meniscus recedes from its topmost position ( $t = 0$  in Fig. 2), the films are deposited at each porthole, their evaporation starts because  $T_w > T_{sat}$  and  $p$  rises. When the meniscus attains a part of the tube where  $T_w < T_{sat}$ , the film condensation occurs at the lower film part [12]. In addition, the vapor expands. Both these effects cause a  $p$  decrease in spite of the continuing evaporation from the upper film part. When the pressure drops low enough, the meniscus begins to rise under the action of the pressure difference  $p_r - p$ . Because of the liquid column inertia, the meniscus rises back into the evaporator and the next oscillation occurs.

### 3. Theoretical reconstruction of the meniscus oscillation

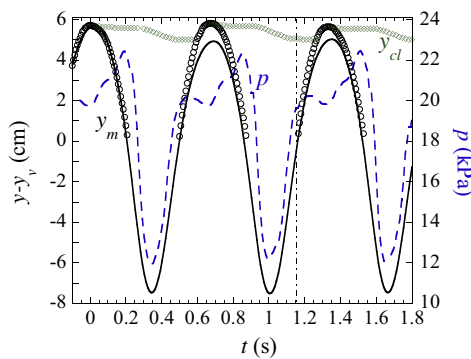
The meniscus equation of motion corresponds to that derived for the meniscus oscillations in the open capillary [13] augmented with the viscous friction term:

$$\rho(y_m + L_r)\dot{U}_m = -\rho g(y_m - y_r) + p_r - p - \frac{1}{2}\rho U_m^2 + \frac{1}{2}b\rho U_m^2 - \frac{2}{D}\tau y_m, \quad (1)$$

where  $L_r$  is the effective length of the liquid in the reservoir that participates in the oscillations, dot means the time derivative,  $U_m = \dot{y}_m$ , and  $y = 0$  corresponds to the bottom of the tube (Fig. 1a). The singular pressure losses coefficient

$$b = \begin{cases} -\frac{1}{2}, & \text{if } U_m > 0, \\ 0, & \text{if } U_m \leq 0 \end{cases}$$

corresponds to a hole in a wall because the tube wall thickness is much larger than  $D = 2$  mm.



**Fig. 2.** Time evolution of the vapor pressure (dashed line), meniscus position (circles) and the topmost CL point in the evaporator measured for  $p_r = 23.7$  kPa,  $T_c = 0$  °C, and heating powers 4 W and 6 W at the front and rear portholes (Fig. 1b), respectively;  $f = 1.5$  Hz. The height  $y_v = 68$  cm corresponds to the bottom of visible fluid (Fig. 1a). The reconstructed theoretically meniscus oscillation is shown with the solid line.

The viscous shear stress is defined mainly by the bulk flow, the meniscus effects being negligible for long liquid plugs [14]. One uses usually the Poiseuille expression,  $\tau = 6U_m\mu/D$ . However the oscillatory flow does impact the dissipation. The shear stress is defined in this case by the expression [15]

$$\tilde{\tau} = -\tilde{U}_m\mu\chi \tanh\left(\frac{D\chi}{2}\right) \left[\frac{2 \tanh\left(\frac{D\chi}{2}\right)}{D\chi} - 1\right]^{-1}, \quad (2)$$

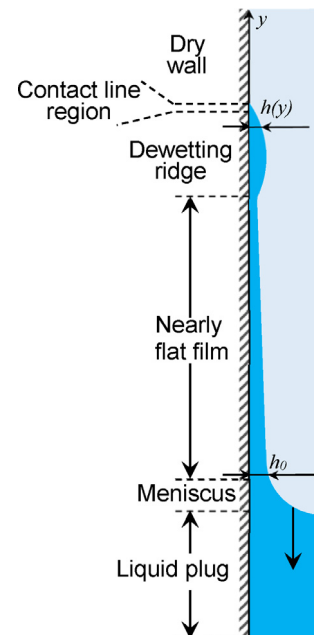
written in the Fourier space, with  $\chi = \sqrt{i\omega/\nu}$ , tilde denoting the Fourier-transformed variables. One mentions that Eq. (2) reduces to the Poiseuille expression in the limit  $\omega t_{visc} \rightarrow 0$ , where  $t_{visc} = D^2/(4\nu)$  is the characteristic viscous relaxation time. In our case, Eq. (2) results in about 1.5 times larger dissipation value than its Poiseuille counterpart.

Since all the parameters in Eq. (1) are known,  $y_m(t)$  evolution can be reconstructed now by using  $p(t)$  measured experimentally (Fig. 2). For the reconstruction,  $L_r = 15$  cm and the ethanol parameters for  $T = 35$  °C are used ( $\rho = 777$  kg/m<sup>3</sup>,  $\mu = 0.9 \cdot 10^{-3}$  Pa·s). By comparing with the part of the curves measured in the transparent evaporator part, one concludes that the reconstruction quality is reasonably good.

### 4. Film thickness

One can distinguish several regions in the evaporator (Fig. 3). First of all, these are liquid plug, dry wall and liquid film regions. The liquid film appears because of its deposition of the receding meniscus (the Landau-Levich phenomenon). The film thickness is controlled by the balance between the viscous friction force that prevents the liquid from flowing along the solid and surface tension in the meniscus region that tends to keep all the fluid inside the plug [16,17]. The thickness of the film deposited by the receding with the velocity  $U_m$  meniscus in a capillary of the diameter  $D$ ,

$$h_0 = 0.67DCa_m^{2/3}, \quad (3)$$



**Fig. 3.** Different regions appearing during the liquid plug receding over a hot wall (in the PHP evaporator).

is defined by the capillary number  $Ca_m = \mu U_m / \sigma$ . The Taylor bubble films have been extensively studied experimentally in isothermal situation. The expression (3) has been corrected to account for many secondary effects like liquid inertia but still remains applicable. In our experiment, the film thickness has been measured by interferometry at a fixed height similarly to [18], by using the Ocean Optics HR2000+ spectrometer. The white light is sent perpendicularly to the film surface by the apparatus probe and the reflected light is recovered by the same probe (Fig. 1). The rays reflected from both film interfaces (with the solid and with the vapor) interfere, which creates multiple maxima in the light spectrum measured by apparatus. The local film thickness is directly deduced from the wavelengths corresponding to the maxima. The thickness measured during the same run at different probe heights (that correspond to different  $U_m$ ) is presented as a function of  $Ca_m$  in Fig. 4. One can see that the measurements are coherent with Eq. (3).

### 5. Slope of the drying film

Under isothermal situation, the deposited film remains flat. The reason for this is that the film is immobile in the wall reference while the liquid flow occurs only at two film ends: in the region where the film is connected to the meniscus and near the dried out solid wall portion (Fig. 3). The film profile is not perturbed by the gravity either: the film thickness is much smaller than the capillary length, which is usually millimetric. However the liquid film thickness may vary because of evaporation [10]. To account for this effect, one can neglect the flow along the film so that the film thickness  $h(y, t)$  changes because of evaporation only:

$$\dot{h} = -\frac{J}{\rho}. \quad (4)$$

The film that has stayed longer after its deposition by the receding meniscus becomes thinner so at a given time moment the film should be thicker near the meniscus and thinner far from it. To obtain the film shape, consider the energy balance at the film surface. It is always at  $T_{sat}$  corresponding to the vapor pressure (perhaps, except at the areas where the evaporation is so strong that molecular kinetic effects become important) because the fluid is pure. By neglecting the heat flux leaving the surface for the vapor with respect to the latent heat, one can write the energy balance as [19,20]

$$J = \frac{k\Delta T}{hL} \quad (5)$$

and use it in Eq. (4). By assuming invariable  $\Delta T$  in the evaporator section (due to its high heat conductivity), one gets the solution of Eq. (4),

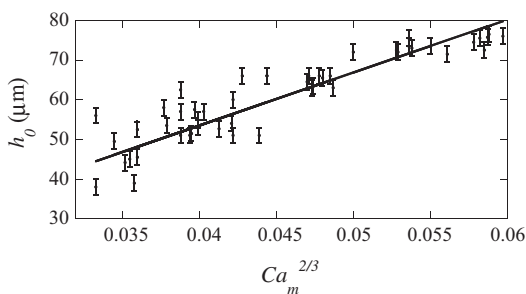


Fig. 4. Film thickness measured by interferometry (characters) just after the film deposition at different heights to obtain different meniscus velocities. It is compared to Eq. (3) (line).

$$h(y, t) = \sqrt{h_0^2 - 2 \frac{k\Delta T}{L\rho} (t - t_0)}, \quad (6)$$

where  $t_0 < t$  is the time moment of the meniscus passage through the point  $y$ ;  $h_0 = h(t = t_0)$  is the thickness given by Eq. (3), which depends on  $U_m = U_m(t_0)$ . The variables  $y$  and  $t_0$  are related through  $y = y_m(t_0)$  where  $y_m(t)$  is the meniscus motion law. For the slope calculation, one can approximate the actual meniscus trajectory with the harmonic oscillation

$$y_m = y_i + A[\cos(\omega t) - 1], \quad (7)$$

where  $y_i = y_m(t = 0)$ . One obtains

$$t_0 = \frac{1}{\omega} \arccos\left(\frac{y - y_i}{A} + 1\right), \quad (8)$$

to be used in Eq. (6).

The film thickness evolution calculated with Eq. (6) is compared to that measured at several fixed heights after the film deposition in Fig. 5. This figure describes the nearly flat film portion schematized in Fig. 3. It is indeed nearly flat: its experimental slope of  $0.6^\circ$  (Fig. 5b) is well reproduced by the theoretical curve (Fig. 5a).

Both theory and experiment presented in this section have their limitations: they are not capable to resolve the dewetting ridge schematized in Fig. 3. The interferometry measurements are limited because of large interface slopes. To overcome this limitation, the grid deflection technique will be applied below. To describe the dewetting ridge theoretically, the CL effects need to be considered.

### 6. Contact line effects

It is known since long ago [21] that evaporation causes the apparent contact angle increase. In the liquid wedge geometry characteristic for the CL region (Fig. 3), such an effect is explained by the strong evaporation (cf. Eq. (5)) in the vicinity of the CL where  $h(y) \rightarrow 0$  when  $y \rightarrow y_{cl}$ . The theory of this phenomenon has been developed by a number of researchers, see e.g. [22–26], and is summarized here very briefly. A liquid mass loss localized in the microscopic vicinity of the CL requires a liquid flow towards it, which creates the viscous pressure drop along  $y$ , i.e. the pressure  $p_L$  in the liquid decreases with  $y$ . Since the vapor pressure is homogeneous, both the interfacial pressure jump  $p - p_L(y)$  and the local interface curvature defined through the Laplace equation increase with  $y$ , which leads to the interface slope decrease with  $y$ . As the

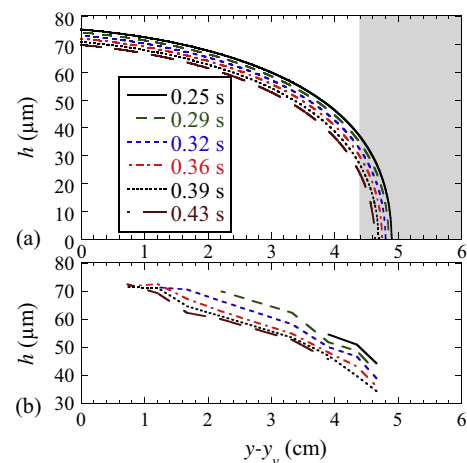


Fig. 5. Film thinning caused by evaporation: (a) calculated with  $A = 6$  cm,  $\Delta T = 10^\circ$  and the meniscus curvature radius of  $0.8$  mm to conform to  $h_0$  value obtained with the interferometry measurements (b). The region influenced by the CL effects is shadowed in (a).



interface slope at the CL  $\theta$  is fixed by the intermolecular forces, the slope  $\theta_V$  at some microscopic distance from the CL depends on evaporation rate and may attain a large value  $\theta_V \gg \theta$ . The  $\theta_V$  value increases with  $\Delta T$  and depends weakly on a microscopic length scale  $l_V \simeq 100$  nm defined in practice by the hydrodynamic slip length divided by the interface slope. The vapor recoil effect [27] also contributes to the  $\theta_V$  increase. Note that the Voinov angle can be seen as an effective contact angle caused by the substrate superheating.

If the CL was immobile, this angle could be measured in a hypothetical setup where the interface shape is not imposed by the macroscopic constraints. In reality, the CL moves and one needs to account for the related flow and slope change. It turns out that in the presence of the CL motion the apparent contact angle (i.e. that obtained by extrapolation of the macroscopic shape toward the CL) is expressed by the Cox-Voinov-type expression [26]

$$\theta_{app}^3 = \theta_V^3 + 9Ca_{cl} \ln \frac{w}{el_V}, \quad (9)$$

where  $e \simeq 2.71$ ,  $Ca_{cl} = \mu U_{cl} / \sigma$  and  $w$  is a relevant macroscopic scale (dewetting ridge width in the present case [28]).

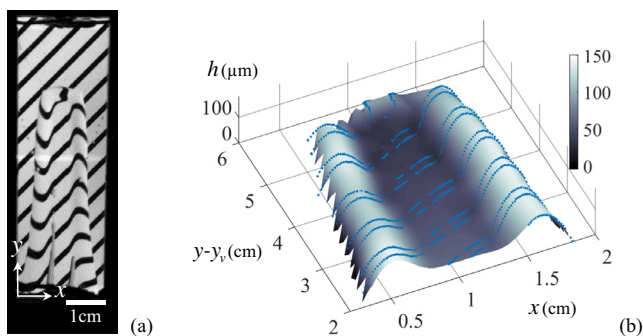
One knows that the liquid film is unstable when the solid surface is not wetted. Once a CL is formed, it recedes spontaneously. The liquid gathered by receding cannot flow inside the thin film because of the strong shear. For this reason it is accumulated in a “dewetting” ridge [29] formed along the film edge. When the ridge is small, its shape is controlled by the surface tension; its cross-section should be circular.

It is quite important to note that since the ridge shape is circular,  $\theta_{app}$  and  $w$  are related to each other once the ridge volume is given. As the ridge accumulates the fluid, its volume is defined by  $U_{cl}$ . One obtains finally that  $U_{cl}$  is defined only by  $\theta_V$  [29,30]. Within a logarithmic prefactor calculated in [28],  $Ca_{cl} \sim \theta_V^3 / 9$ .

One can now analyze the experimental data to find the described above phenomena experimentally.

## 7. Film shape reconstruction

A non-intrusive, optical grid deflection technique [30–32] has been used to capture the film shape in situ, during the PHP functioning. A wire grid (formed by parallel wires) is placed between the diffuse light source and the cell, cf. Fig. 1b. The wire direction is inclined to  $45^\circ$  to avoid it to be parallel to the CL. The incident light rays are refracted at the inclined film interface. The local film slope can be deduced from distortion of the grid image (Fig. 6a). A local  $h$  measurement is needed to reconstruct its spatial variation



**Fig. 6.** Raw image (a) and the corresponding 3D reconstruction of the film profile (b). The surface is obtained by interpolation between reconstructed points (characters) along the images of grid edges. Note that the lateral ridge parts appear because of the rectangular tube and heating mode; they are absent in the circular tube PHP.

along the single wire image; the discussed above interferometry data have been used for this purpose.

An obtained image and corresponding to it reconstruction of the film shape are shown in (Figs. 6). The time moment corresponding to the image is shown by the vertical dash-dotted line in Fig. 2. The central film portion is obtained from the interferometry measurements, while the dewetting ridge shape is calculated using the grid deflection analysis.

To obtain the apparent contact angle one needs to cut the 3D fit of the film surface by  $y = \text{const}$  plane. The fits show that the dewetting ridges are indeed circular, in agreement with the above theory. Since the experimental data are impossible to obtain very close to the contact line, one needs to use the fits to determine  $\theta_{app}$ .

From  $\theta_{app}$  and  $U_{cl}$ , the contact angle can be calculated with Eq. (9). Note that  $U_{cl} < 0$  and  $\theta_V > \theta_{app}$ . In the present experiment,  $\theta_V$  can reach  $30^\circ$  [33].

## 8. Discussion and conclusions

One can understand now the CL dynamics (Fig. 2). Since the thermal inertia of the sapphire portholes is high, one can reasonably assume that the temperature variation of their internal surfaces is small so that the  $\Delta T$  variation is inverse of that of  $p$ . Let us remind that  $\theta_V$  is the increasing function of  $\Delta T$ . At  $t = 0$ ,  $p$  is high ( $\theta_V$  is low) so that CL is pinned on surface defects. The  $p$  drop leads to a sharp  $\theta_V$  increase so that the dewetting starts: CL recedes. Once  $p$  rises again,  $\theta_V$  drops and becomes pinned again. The meniscus washes up the film, brings CL to the upper position and the oscillation starts again.

One can see that because of the contact angle increase appearing due to intensive evaporation in a microscopically small region, a thicker ridge forms at the film edge. This locally decreases the film evaporation so that the film edge receding dynamics is defined rather by dewetting effect than by the overall film evaporation as assumed in the state-of-the-art PHP simulations [34]. The physics of the CL receding is understood. The CL velocity is defined mainly by the so called Voinov contact angle that is defined mainly by the superheating of internal PHP walls.

The discovered film shape contradicts previous ideas [8,10] about the wedge-like films in the PHP. In reality, the film is nearly flat (which conforms to the hypothesis made in 1D PHP simulations [34]) with a ridge at its end.

Because the films inside the bubbles of conventional, multi-branch PHP evolve independently of each other, the findings obtained here for the single branch PHP apply for the multi-branch case. The particularity of the flat geometry is the trapezoidal film shape, which is absent in the circular cross-section tubes. However the capillary ridge and, more generally, film dynamics observed here holds in the general case. Since the film area is one of the key factors influencing the PHP dynamics, these findings are very important to be implemented in the future both single and multi-branch PHP simulations.

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