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Penning and associative ionisation of highly excited rubidium atoms

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Abstract. Highly excited and resonant states are created in a rubidium vapour in order to study the collisional ionisation with the ground state and between these different excited species. A mass analysis allows us to separate the atomic and molecular channels. Penning ionisation with a resonant state is shown to have a great efficiency (rate coefficient greater than $10^{-8} \text{ cm}^3 \text{ s}^{-1}$) for all the highly excited levels under study (6d, 8s, 7d). The associative ionisation rate coefficient between these excited atoms and a ground-state atom lies in the expected range, 10^{-10} to $5 \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ and seems to be independent of l and j values.

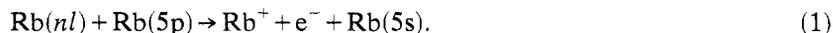
1. Introduction

Mass analysis of ions produced in a Rb vapour where atomic excited states are created by a two-step photoexcitation technique allows a convenient means for the identification of the main reaction channels producing these ions. In a previous paper (Chéret *et al* 1981) we have presented a measurement of the rate coefficient for associative ionisation when Rb(9s) excited atoms collide with ground-state atoms. In a second paper (Chéret *et al* 1982) it has been shown that atomic ions are produced with a large rate coefficient ($\sim 10^{-8} \text{ cm}^3 \text{ s}^{-1}$) by collision between two short-lifetime excited atoms, one in the 8s state the other in a 5p state.

The results are interesting for ionised gases or high-density excited vapour studies (Roussel *et al* 1980, 1981, Lucatorto and McIlrath 1976, Allegrini *et al* 1981, Le Gouët *et al* 1982).

To complement this work a systematic study involving other excited states has been performed. Two important collisional processes have been pointed out.

Penning ionisation (rate coefficient $k_1(nl)$)



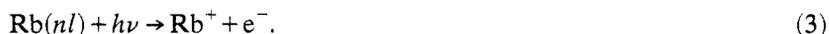
Associative ionisation (rate coefficient $k_2(nl)$)



The results obtained for the rate coefficients for these two mechanisms are presented for the 6d, 8s and 7d levels.

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In our experimental conditions it is shown that photoionisation processes (cross section $\sigma(nl, \lambda_i)$) play an important role



The experimental technique is presented and the results are discussed. It is shown that other mechanisms producing atomic or molecular ions are negligible under our experimental conditions.

2. Experimental set-up

The reaction chamber is a flat stainless steel cylinder 25 cm in diameter and 8 cm high (figure 1). Four diameters are used for photoionisation and optical measurements. A mass spectrometer set on the vertical axis of the cell in a separated differentially pumped chamber performs ion analysis. The ions are extracted from the reaction chamber by weak electric fields (a few V cm^{-1}) through a small thin hole (with diameter 1 mm). The temperature of the whole apparatus is 450 K and a heated reservoir filled with rubidium metal is used to maintain a constant pressure of alkali vapour in the range from 10^{-5} to 10^{-3} Torr. Under these conditions the Rb concentration lies in the range 2×10^{11} to $2 \times 10^{13} \text{ cm}^{-3}$.

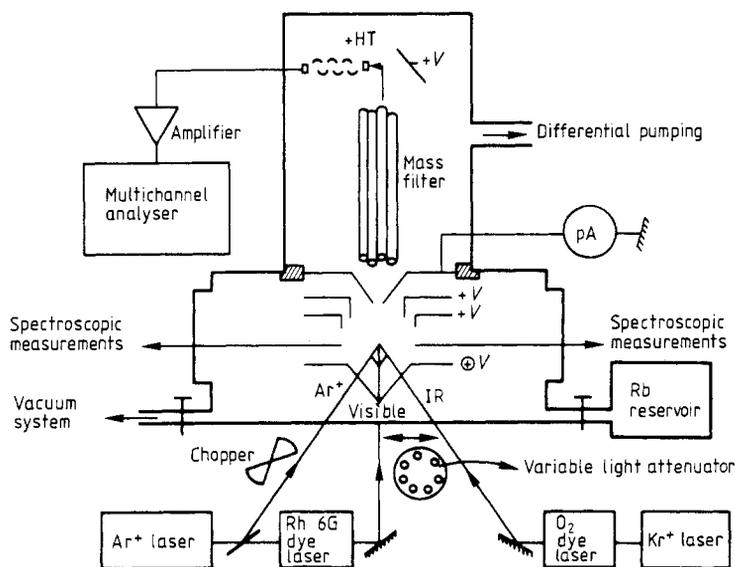


Figure 1. Scheme of the experimental device.

The first cw infrared dye laser beam (25 GHz width, laser power in the centre of the cell P_1 , beam section S_1 , wavelength λ_1) excites Rb atoms to one of the first resonant doublets of Rb ($5p^2P_{3/2}$) at $\lambda_1 = 780.2 \text{ nm}$ or Rb ($5p^2P_{1/2}$) at $\lambda_1 = 795.0 \text{ nm}$. A second multimode cw dye laser beam (40 GHz, P_2 , S_2 , λ_2) creates s or d atomic excited states in the centre of the cell. The overlapping volume of the two laser beams is estimated to be $V = 1.0 \times 10^{-2} \text{ cm}^{-3}$.

According to reactions (1) and (2) these excited states produce atomic and molecular ions. According to reaction (3) photoionisation by the two dye laser beams produces atomic ions. The Rb^+ and Rb_2^+ ion currents are determined, relative to one another, by the mass spectrometer. The total ion current, which is the sum of these two partial currents, is measured in absolute value on a polarised plate inside the cell. The transmission through our system is assumed to be equal for the two ion species. The electron multiplier is working in the analogue mode.

3. Equations governing the ion currents

In spite of a selective photoexcitation, cascading radiative de-excitation of the directly pumped level (nl) gives an appreciable population of the lower states ($n'l'$). Their densities are correlated to the nl concentration by the following linear system using the Einstein coefficients $A_{i \rightarrow j}$:

$$[n'l'] = \left(\sum_p A_{p \rightarrow n'l'} [p] \right) \left(\sum_q A_{n'l' \rightarrow q} \right)^{-1}. \quad (4)$$

The square brackets indicate a density and p (or q) the levels above (or below) the ($n'l'$) state considered. The Einstein coefficients have been calculated by Gounand (1979) using the Bates–Damgaard theoretical results.

The calculated relative densities obtained for each (nl) level studied are summarised in table 1. They have been checked experimentally by fluorescence measurements showing that collisional excitation transfer does not play an important role in the pressure range studied. This result is not surprising as the pressure is low and the energy gaps between the different excited states are large.

Table 1. Population densities due to radiative de-excitation.

Level	7d	8p	5f	8s	6d	7p	4f	7s	5d	6p	6s	4d
$\frac{n'l'}{nl} (\times 10^2)$					100	8.8	1.6	0.8	3.2	2.9	1.0	3.3
				100	0	34.2	0	3.3	11.8	19.9	1.1	5.3
	100	5.4	1.9	0.3	1.0	1.0	0.7	0.2	2.0	3.6	0.8	3.0

It is obvious that the atomic and molecular ions can be produced by the pumped level and by the radiatively populated states. For a given level ($n'l'$) the corresponding currents are:

$$I_1'(n'l') = eVk_1(n'l')[5p][n'l'] \quad \text{for Penning ionisation} \quad (5)$$

$$I_2'(n'l') = eVk_2(n'l')[5s][n'l'] \quad \text{for associative ionisation} \quad (6)$$

$$I_\nu'(n'l', \lambda_i) = eV \frac{P_i \lambda_i}{S_i hc} \sigma(n'l', \lambda_i)[n'l'] \quad \text{for photoionisation.} \quad (7)$$

In these expressions, e , c , h are the usual notations for the elementary charge, the speed of light and the Planck constant. V is the interaction volume which is equal to the overlapping volume of the laser beams. $k_1(n'l')$, $k_2(n'l')$ and $\sigma(n'l', \lambda_i)$ are respectively the rate coefficients and the photoionisation cross section at λ_i for the level ($n'l'$).

The measured atomic and molecular ion currents $I_1^+(nl)$ and $I_2^+(nl)$, corresponding to the laser excitation of the (nl) level, are the sum of the partial currents $I_1^+(n'l')$, $I_2^+(n'l')$ and $I_v^+(n'l', \lambda_i)$ for all the levels and wavelengths. They can be written as

$$I_1^+(nl) = eV \left([5p] \sum_{n'l'} k_1(n'l') \frac{[n'l']}{[nl]} + \sum_{n'l', \lambda_i} \sigma(n'l', \lambda_i) \frac{P_i \lambda_i [n'l']}{S_i hc [nl]} \right) [nl] \quad (8)$$

$$I_2^+(nl) = eV [5s] \sum_{n'l'} k_2(n'l') \frac{[n'l']}{[nl]} [nl]. \quad (9)$$

Introducing effective rate coefficients $K_i(nl)$ and effective photoionisation cross sections $Q(nl, \lambda_i)$:

$$K_1(nl) = \sum_{n'l'} k_1(n'l') \frac{[n'l']}{[nl]} \quad (10)$$

$$Q(nl, \lambda_i) = \sum_{n'l'} \sigma(n'l', \lambda_i) \frac{[n'l']}{[nl]} \quad (11)$$

correlated to the radiative equilibrium obtained in the gas, equations (8) and (9) reduce, respectively, to:

$$I_1^+(nl) = eV \left(K_1(nl) [5p] + \sum_{\lambda_i} \frac{P_i \lambda_i}{S_i hc} Q(nl, \lambda_i) \right) [nl] \quad (12)$$

$$I_2^+(nl) = eV K_2(nl) [5s] [nl]. \quad (13)$$

4. Experimental procedure

$Q(nl, \lambda_i)$ are determined from density ratios presented in table 1 and from calculations of photoionisation cross sections by Aymar *et al* (1976). These results are summarised in table 2. The determination of $K_1(nl)$ and $K_2(nl)$ from (12) and (13) only requires density measurements of the 5p and nl concentrations.

These quantities are deduced from the ion current resulting from the photoionisation induced by the Ar ion laser. For this purpose a part of its beam is sent in the interaction volume on a third diameter of the cell (see figure 1).

In a first step the IR dye laser working alone produces the 5p state and the Ar^+ is tuned at 472.7 nm in the case of the $5p^2P_{1/2}$ or at 476.5 nm for the $5p^2P_{3/2}$ (this laser beam is characterised by P_3 , S_3 , λ_3 and the overlapping volume is $V_3 = 1.2 \times 10^{-2} \text{ cm}^3$). The resulting photoionisation current is

$$I^+ = eV_3 \frac{P_3 \lambda_3}{S_3 hc} \sigma(5p, \lambda_3) [5p] \quad (14)$$

where $\sigma(5p, \lambda_3)$ is known from the experimental results of Moskvin (1963) in good agreement with the theoretical calculations of Aymar *et al* (1976), $[5p]$ is deduced from I^+ with a good accuracy.

In a second step the Ar^+ laser (P_4 , S_4 , λ_4 , overlapping volume with the two dye laser beams, $V_4 = 7 \times 10^{-3} \text{ cm}^3$) tuned at $\lambda_4 = 514.5 \text{ nm}$, pumps the visible dye and photoionises all the excited levels except the 5p state. The atomic ion current I_1^+ due to Penning ionisation and to photoionisation by the two dye laser beams is increased by ΔI^+ . From the effective photoionisation cross section $Q(nl, \lambda_4)$ associated to the

Table 2. (a) Individual photoionisation cross sections, and, (b) effective photoionisation cross sections, both in units of 10^{-18} cm^2 .

(a)					
Levels	σ (7802 Å)	σ (630 nm)	σ (616 nm)	σ (572 nm)	σ (517 nm)
7d	5.7			2.3	2.05
8p	2.1			2.0	1.7
5f	1.0			0.32	0.25
8s	0.015		$<10^{-3}$	$<10^{-3}$	$<10^{-3}$
6d	8.0	6.4	6.0	5.0	3.3
7p	7.0	5.5	5.0	4.0	3.15
4f	2.0	0.7	0.55	0.45	0.35
7s	0.1	0.03	0.015	$<10^{-2}$	$<10^{-3}$
5d	16.0	12.0	10.1	9.8	8.0
6p	15.0	13.0	10.8	9.5	7.5
6s	0.0	0.25	0.20	0.18	0.15
4d	0.0	27.0	25.0	18.0	15.0

(b)			
Levels	$Q(nl, 780)$	$Q(nl, \lambda_2)$	$Q(nl, 515)$
6d	9.5	7.5 ($\lambda_2 = 630 \text{ nm}$)	4.8
8s	6.5	6.4 ($\lambda_2 = 616 \text{ nm}$)	4.7
7d	6.7	3.8 ($\lambda_2 = 572 \text{ nm}$)	3.3

nl state ΔI^+ can be written:

$$\Delta I^+ = eV_4 \frac{P_4 \lambda_4}{S_4 hc} Q(nl, \lambda_4)[nl]. \quad (15)$$

ΔI^+ is typically a tenth of I_1^+ , both are measured simultaneously by chopping (not synchronously) the IR and the Ar^+ laser beams, the ion current being recorded on a multichannel analyser.

Under these conditions, all the parameters required to determine $K_1(nl)$ and $K_2(nl)$ from (12) and (13) are known. Nevertheless, these basic formulae assume that the three mechanisms considered are dominant.

The variations of I_1^+ and I_2 when plotted against the laser powers P_1 and P_2 and against the pressure check this hypothesis. It is shown experimentally that $[5p]$ varies as P_1 and $[nl]$ varies as $P_1 P_2$ for laser powers up to 320 mW. Consequently, the molecular ion current and the various partial atomic ion currents, i.e. Penning collisional current I_1 , IR photoionisation current $I_\nu(nl, \lambda_1)$ and visible photoionisation current $I_\nu(nl, \lambda_2)$, vary with P_1 and P_2 in the following manner:

$$\begin{aligned} I_2(nl) &= eVK_2(nl)[5s][nl] \sim P_1 P_2 \\ I_1(nl) &= eVK_1(nl)[5p][nl] \sim P_1 \times P_1 P_2 \\ I_\nu(nl, \lambda_1) &= eV \frac{P_1 \lambda_1}{S_1 hc} Q(nl, \lambda_1)[nl] \sim P_1 \times P_1 P_2 \\ I_\nu(nl, \lambda_2) &= eV \frac{P_2 \lambda_2}{S_2 hc} Q(nl, \lambda_2)[nl] \sim P_2 \times P_1 P_2. \end{aligned} \quad (16)$$

Therefore, if one of the powers is kept constant, for instance P_2 , the variation of

$$I_1^i(nl) = I_1(nl) + I_v(nl, \lambda_1) + I_v(nl, \lambda_2)$$

against P_1 should be the sum of a linear and a quadratic term. The linear part L_1 comes from the photoionisation by the visible radiation, the quadratic term Q_1 comes from the Penning ionisation and from the photoionisation induced by the IR beam. At constant P_1 , the Penning ionisation and the IR photoionisation provide a linear term L_2 , the visible photoionisation a quadratic term Q_2 : when the constant power is changed the linear and quadratic terms exchange their mechanisms. For a given set of powers P_1 and P_2 $I_1^i(nl)$ is

$$I_1^i(nl) = P_1 + Q_1 = l_1(P_2)P_1 + q_1(P_2)P_1^2 \quad (17)$$

or

$$I_1^i(nl) = P_2 + Q_2 = l_2(P_1)P_2 + q_2(P_1)P_2^2$$

where l_i and q_i are parameters to be determined.

Experimental investigation of these dependences confirms the validity of the mechanisms considered and rules out other reactions which have not the proper variation with the two laser powers. These possible processes are discussed further.

Experimentally, at different pressures, the laser powers are varied using density filters fixed on a rotating wheel and the ion currents are recorded synchronously. This technique allows us to determine the l_i and q_i values from a least-squares fit method with a good correlation factor. The ratios between the linear and the quadratic parts of the atomic ion current for given laser powers can be written

$$\frac{Q_1}{L_1} = \frac{L_2}{Q_2} = \frac{P_1}{P_2} \frac{S_2}{S_1} \frac{\lambda_1}{\lambda_2} \frac{Q(nl, \lambda_1)}{Q(nl, \lambda_2)} + K_1(nl)[5p] \left(\frac{P_2}{S_2} \frac{Q(nl, \lambda_2)\lambda_2}{hc} \right)^{-1} \quad (18)$$

At a given pressure, the shape of the atomic ion current curve and the linear and quadratic term exchange show that the only possible mechanisms are (1), (2) and (3). The experimental results give on one side the photoionisation current induced by the visible laser and on the other side the IR photoionisation current plus the Penning ionisation current. The relative importance of these two last currents is determined from the theoretical photoionisation cross sections. The study of the pressure dependence confirms the validity of the cross sections selected: at a given laser power $I_v(nl, \lambda_1)$ is proportional to nl , and $I_1(nl)$ to $[nl][5p]$, so when the pressure increases and consequently $[5p]$, $I_1(nl)$ increases as the square of $I(nl, \lambda_1)$. This last study demonstrates the validity of the hypothesis.

5. Determination of the rate coefficients

Measurements of the atomic and molecular ion currents as a function of the laser powers and of the pressure allow us to determine $K_1(nl)$ and $K_2(nl)$ unambiguously from (12) and (13). The experimental procedure gives the rate coefficients by different methods.

(a) From (12) and (15) the following form of $K_1(nl)$ can be deduced

$$K_1(nl) = \left(\frac{I_1^i}{\Delta I^+} - \alpha_1 \frac{P_1}{P_4} - \alpha_2 \frac{P_2}{P_4} \right) \frac{\beta P_4}{[5p]} \quad (19)$$

with

$$\alpha_i = \frac{V}{V_4} \frac{\lambda_i}{\lambda_4} \frac{Q(nl, \lambda_i)}{Q(nl, \lambda_4)} \frac{S_4}{S_i}$$

and

$$\beta = \frac{V_4}{V} \frac{\lambda_4}{hcS_4} Q(nl, \lambda_4).$$

Here, the main experimental quantity is $I_1/\Delta I^+$. It is obtained by a chopping technique and a simultaneous measurement of the currents. This experimental procedure correlates the photoionisation currents $I_\nu(nl, \lambda_i)$ induced by the infrared and the visible laser beams to the photoionisation current ΔI^+ given by the argon ion laser tuned at 514.5 nm:

$$I_\nu(nl, \lambda_i) = \alpha_i \Delta I^+ P_i / P_4.$$

It is an experimental check of the importance of the photoionisation mechanisms.

(b) $K_1(nl)$ can be also determined from (18):

$$K_1(nl) = \left(\frac{Q_1}{L_1} - \gamma \frac{P_1}{P_2} \right) \frac{\delta P_2}{[5p]} \quad (20)$$

with

$$\gamma = \frac{\lambda_1}{\lambda_2} \frac{Q(nl, \lambda_1)}{Q(nl, \lambda_2)} \frac{S_2}{S_1} \quad \text{and} \quad \delta = \frac{\lambda_2}{hcS_2} Q(nl, \lambda_2).$$

Here, the main experimental quantity is Q_1/L_1 (or L_2/Q_2). This is determined by a modulation technique and a least squares fit. The elementary processes and the effective rate coefficient are obtained simultaneously. It should be noticed that the interaction volume between the IR and the visible laser beams does not appear in formula (20).

The accuracy of the two methods is comparable. Their results are in good agreement and complementary. The experimental results presented in this paper for $K_1(nl)$ are the average values obtained by these two approaches.

(c) The associative rate coefficient $K_2(nl)$ can be written from (13) and (15) as

$$K_2(nl) = \frac{I_2}{\Delta I^+ P_4} \frac{\beta}{[5p]} \quad (21)$$

The linear dependence of Rb_2^+ on the laser powers is checked by the modulation technique.

6. Power measurements

The determination of $K_1(nl)$ and $K_2(nl)$ needs the knowledge of the laser powers P_i at different wavelengths λ_i in the interaction zone. λ_2 , λ_3 and λ_4 do not correspond to transitions to the ground state and consequently are slightly absorbed in the gas. The main absorption is on the entrance and the exit window of the cell: the measurement of the input E_i and output O_i powers allows the calculation of the laser power P_i :

$$P_i = (1 - A_i)E_i = (E_i O_i)^{1/2}$$

assuming the same absorption A_i for each window.

For the resonance line λ_1 the measurement is more difficult. A strong absorption occurs at λ_1 in the first few millimeters into the cell (the interaction area is located 12.5 cm away from the entrance window). However, the photoionisation is produced by the whole laser line which is wider (25 GHz) than the absorption line (Doppler width 1 GHz). Consequently, the influence of the resonant absorption is not very important. The laser power, in the interaction zone, is determined by taking into account the absorption due to the window and correcting the entrance power from the resonance absorption A_R . For this purpose, entrance and exit powers are measured at the resonance (E_1, O_1) and out of resonance (E', O'). E_1, O_1, E' and O' are related by:

$$O_1 = (1 - A_R)(1 - A_1)^2 E_1 \quad O' = (1 - A_1)^2 E'$$

and give

$$P_1 = (1 - A_R)(1 - A_1)E_1 = O_1(E'/O')^{1/2}.$$

7. Experimental results

The $6d \ ^2D_{3/2}$, $6d \ ^2D_{5/2}$, $8s \ ^2S_{1/2}$ and $7d \ ^2D_{5/2}$ levels have been excited from the $5p \ ^2P_{3/2}$ level, and $6d \ ^2D_{3/2}$ from $5p \ ^2P_{1/2}$. The $7d \ ^2D_{3/2}$ has an absorption coefficient which is too low and is energetically too close to the $7d \ ^2D_{5/2}$ level to be excited alone. The d and s levels are pumped using Rh 6G, the 5p using oxazine 2.

Density measurements and rate coefficient determinations have been performed in the pressure range 10^{-5} to 10^{-3} Torr. Figure 2 represents 5p and 8s concentrations plotted against pressure at constant laser powers. It should be noticed that [8s] varies as [5p] and that the shape of the [5p] curve differs from a linear dependence. This

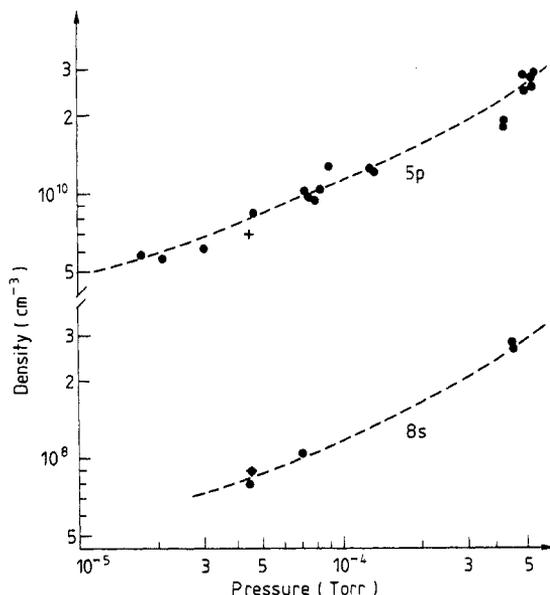


Figure 2. $5p \ ^2P_{3/2}$ and $8s \ ^2S_{1/2}$ concentrations plotted against pressure at $P_1 = 164$ mW and $P_2 = 216$ mW.

can be explained by the strong radiation trapping of the resonance line. Density measurements avoid the difficult theoretical treatment of excitation by a resonance line.

Density measurements have been performed at different powers and different pressures to check the dependences upon P_1 and P_2 . As the density ratios $[5p]/[5s]$ (calculated from figure 2) and $[nl]/[5p]$ are low, i.e. far from saturation values, it is not surprising that $[5p]$ varies as P_1 and $[nl]$ as P_1P_2 . The following values have been obtained for the $[nl]/[5p]$ ratios: 1:390 for $6d^2D_{3/2}$ at $P_2 = 120$ mW, 1:54 for $6d^2D_{5/2}$ (120 mW), 1:78 for $8s^2S_{1/2}$ (216 mW), 1:53 for $7d^2D_{5/2}$ (108 mW) pumped from the $5p^2P_{3/2}$ level, and 1:48 for the $6d^2D_{3/2}$ pumped from the $5p^2P_{1/2}$ ($P_2 = 152$ mW).

The experimental variations of I_1 and I_2 plotted against P_1 and P_2 are consistent with the prediction of formulae (16). Figure 3 represents a typical variation obtained for the $7d^2D_{5/2}$ level at intermediate pressure, $p = 4 \times 10^{-4}$ Torr. The behaviour of the Rb^+ curves corresponds to the sum of a linear and a quadratic term which exchange when the constant power is changed. The whole procedure can be applied to determine $K_1(nl)$ and $K_2(nl)$.

The effective rate coefficients as calculated from (19) or (20) and from (21) are summarised in table 3. The differences between $K_1(nl)$ calculated from (19) or (20)

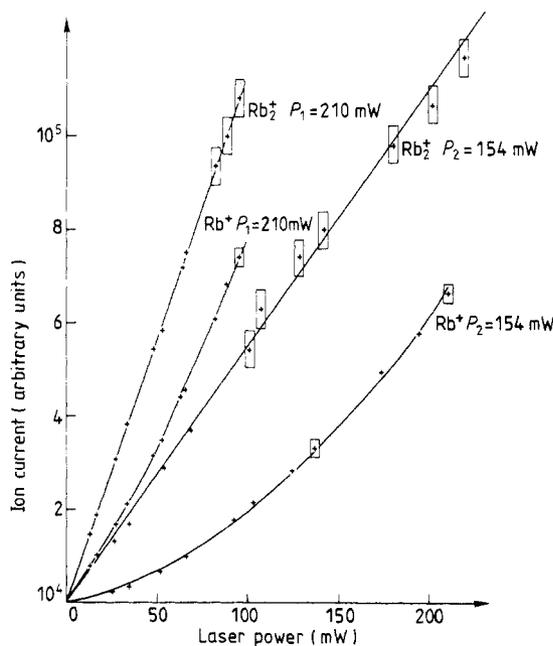


Figure 3. Ion currents plotted against P_1 and P_2 for the $7d^2D_{5/2}$ at 4.7×10^{-5} Torr.

Table 3. Effective collisional rate coefficients $K_i(nl)$ and individual rate coefficients $k_i(nl)$.

Level	$K_1(nl)$ ($10^{-8} \text{ cm}^3 \text{ s}^{-1}$)	$K_2(nl)$ ($10^{-10} \text{ cm}^3 \text{ s}^{-1}$)	$k_1(nl)$ ($10^{-8} \text{ cm}^3 \text{ s}^{-1}$)	$k_2(nl)$ ($10^{-10} \text{ cm}^3 \text{ s}^{-1}$)
$6d^2D_{3/2}$	4.9 ± 1.0	2.0 ± 0.5	4.2 ± 0.9	2.0 ± 0.5
$6d^2D_{5/2}$	4.8 ± 1.0	1.7 ± 0.4	4.1 ± 0.9	1.7 ± 0.4
$8s^2S_{1/2}$	3.7 ± 1.2	2.5 ± 0.6	2.2 ± 0.7	2.5 ± 0.6
$7d^2D_{5/2}$	2.6 ± 1.0	4.4 ± 0.9	2.2 ± 0.9	4.2 ± 0.9

are not significant. This strongly supports the calculated values of the photoionisation cross sections. No variation of the rate coefficients with the pressure has been observed. Figure 4 gives for the 8s level the collisional atomic current $I_1(8s)$, the molecular current $I_2(8s)$ and the total collisional current $I_1(8s) + I_2(8s)$ plotted against pressure. Full points and crosses are deduced from measurements, the full and broken curves are calculated with the same values of $K_1(nl)$ and $K_2(nl)$ at different pressures with the densities deduced from figure 2. Changing $K_i(nl)$ induces a parallel shift of the curves. The same rate coefficients explain all the measurements.

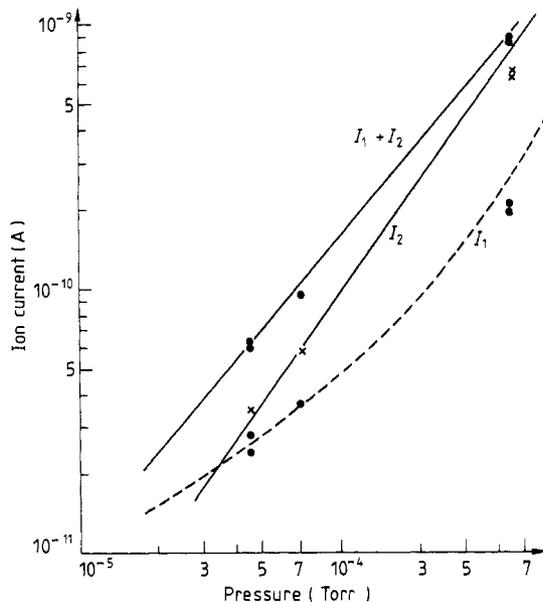


Figure 4. Ion collisional currents plotted against pressure for the $8s\ ^2S_{1/2}$ level at $P_1 = 154\text{ mW}$ and $P_2 = 210\text{ mW}$.

8. Individual rate coefficients

It is easy to deduce the individual associative rate coefficients $k_2(nl)$ for the pumped level from the effective coefficients $K_2(nl)$.

In figure 5 the Rb_2^+ limit corresponding to a potential well of 5300 cm^{-1} for this ion is reported, as calculated by Valance (1978). The 6d levels are the lowest for which associative ionisation is energetically possible: consequently $K_2(6d)$ is equal to $k_2(6d)$. As the $8s \rightarrow 6d$ transition is optically forbidden and as the collisional transfers are negligible there is no level below 8s which contributes to associative ionisation: $K_2(8s)$ is also equal to $k_2(8s)$.

When the 7d level is excited the 8p, 5f, 8s and 6d levels also contribute to the molecular ion current. $k_2(6d)$ and $k_2(8s)$ have been determined. Klucharev *et al* (1980) gives an experiential rate coefficient for the 8p level close to the value presented here for the 6d and 8s levels. These different measurements indicate that there is no effect of the orbital quantum number for highly excited atoms. This last conclusion has been extended to the 5f level. An interpolation against the effective quantum

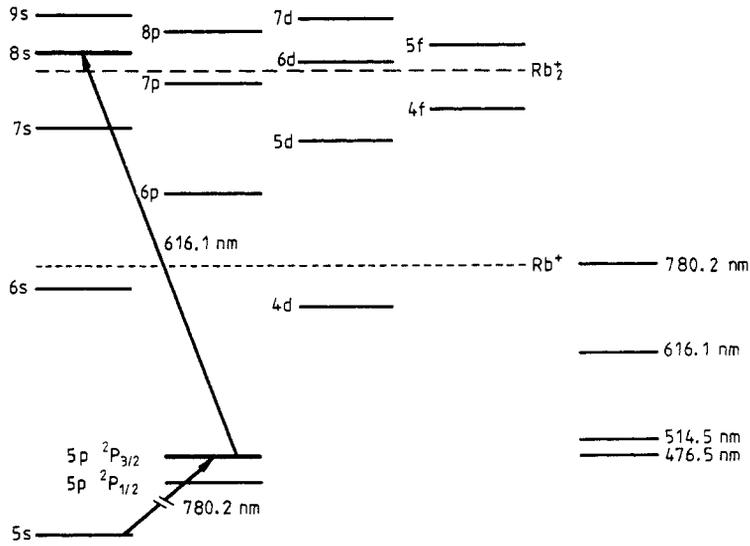


Figure 5. Energy diagram of the Rb levels showing a two-step excitation of the 8s level. The upper broken line represents the lowest energy limit for associative ionisation and the lower broken line the limit for Penning ionisation with a 5p ²P_{3/2} level. The full lines on the right-hand side represent the lowest photoionisation energy limit at the corresponding wavelengths.

number n^* between $k_2(8s)$ and $k_2(8p)$ has been done to determine $k_2(5f)$ (figure 6). In fact the 5f density is quite low and the correction is weak. The $k_2(nl)$ values are summarised in table 3.

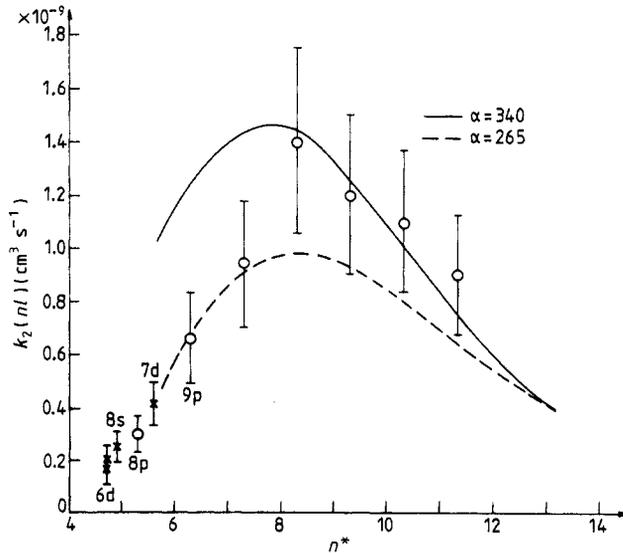


Figure 6. The $k_2(nl)$ variation plotted against n^* . The full and broken curves are the Mihajlov and Janev (1981) theoretical curves for two values of the Rb polarisability α : —, $\alpha = 340$; ---, $\alpha = 265$. The experimental data are; ○, Klucharev *et al* (1980); ×, this work.

All the levels above the 6s level can contribute to the Penning ion current and no rate coefficient is available except the present $K_1(nl)$ values. Consequently, $k_1(nl)$ values are more difficult to obtain. Below the directly pumped level the most populated are the p states (table 1). They contribute strongly to the population equilibrium when the 8s level is pumped and more slightly in the case of a d level. As there is no significant difference between $K_1(8s)$, $K_1(6d)$ or $K_1(7d)$, $k_1(np)$ is certainly not much greater than $k_1(nd)$. The same remark can be done for the 4f level taking into account $K_1(6d)$, $K_1(7d)$ and the relative concentrations. As no level seems to play a dominant role $K_1(6d)$, $K_1(8s)$ and $K_1(7d)$ are corrected from the population distribution (table 1) to determine $k_1(6d)$, $k_1(8s)$ and $k_1(7d)$. These values are reported in table 3.

The neighbouring levels of the 6d or 7d states have low relative concentrations and as a matter of fact $k_1(nd)$ differs slightly from $K_1(nd)$. On the other hand, the neighbouring levels are very important for the 8s level. If $k_1(8s)$ and $k_1(7s)$ are assumed negligible a common value of $5.6 \times 10^{-8} \text{ cm}^3 \text{ s}^{-1}$ is obtained for $k_1(7p)$, $k_1(5d)$ and $k_1(6p)$, whereas $k_1(6d)$ and $k_1(7d)$ are weakly modified. As this large value is not inconsistent with the 6d and 7d measurements, an unambiguous determination of $k_1(ns)$ needs complementary measurements.

9. Accuracy of the method

Error bars corresponding to measurement dispersion are indicated in table 3. Different parameters are used to determine K_1 and K_2 . Some of these, the photoionisation cross sections and the Einstein coefficient are calculated from theoretical works. Only the ratios of theoretical data are used except for the 5p photoionisation cross section which is also known from measurements. The fluorescence intensity measurements and the photoionisation currents confirm the Einstein coefficient ratios and the variation of the photoionisation cross section with wavelength. The rate coefficients can be considered as reliable from that point of view.

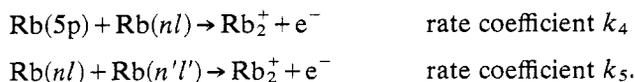
Other parameters such as interaction volumes and laser powers are correlated to experimental diagnostics. The determination of [5p] requires knowledge of the interaction volume between the IR and the Ar^+ laser beams. This volume has been checked by moving the Ar^+ laser beam transversally through the IR beam. The 5p distribution has been found to coincide with the IR beam diameter. Only interaction volume ratios are used in formulae (19) and (21). Volumes do not appear in formula (20). Under these conditions the apparatus can be considered as giving absolute values of the rate coefficients.

10. Neglected mechanisms

The experimental variations of I_1 and I_2 with P_1 and P_2 show that various mechanisms of ion production are negligible under our experimental conditions.

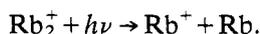
(A) The linear shape of Rb_2^+ rules out the following (and similar) mechanisms with a quadratic behaviour with the powers.

(a) associative ionisation between excited states:



An upper limit of the reaction rates can be estimated by assuming the linearity to be better than 10%. k_4 should be less than $2 \times 10^{-9} \text{ cm}^3 \text{ s}^{-1}$ and k_5 than $2 \times 10^{-7} \text{ cm}^3 \text{ s}^{-1}$. This last value is too large to have a physical meaning but the upper limit obtained for k_4 shows that associative ionisation between two excited states is less than a tenth of the Penning ionisation. Such a result has often been obtained in Penning ionisation studies (Manus *et al* 1977, Pesnelle *et al* 1981).

(b) Photodissociation of Rb_2^+ :

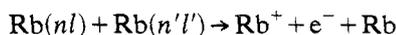


An upper limit of $5 \times 10^{-17} \text{ cm}^2$ can be determined for the cross section of this mechanism. This is not at variance with the experimental results of Bearman and Leventhal (1978) who estimate the photodissociation cross section of Na_2^+ to be 10^{-17} cm^2 .

The lack of photodissociation has also been checked with the Ar^+ laser beam.

(B) The variation of the atomic ion current with the visible and the IR powers allows us to discard the following (and similar) mechanisms.

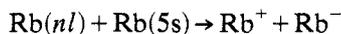
(a) Penning ionisation between two highly excited states:



which varies quadratically with the two laser powers.

The upper limit of the rate coefficient is estimated to be $10^{-7} \text{ cm}^3 \text{ s}^{-1}$.

(b) Ion pair formation:

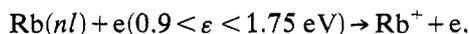


which varies linearly with the two laser powers.

From the electron affinity of the negative rubidium ion (Frey *et al* 1978) this reaction is not energetically possible for the 6d and 8s levels. The rate coefficient for the 7d level is estimated to be less than $8 \times 10^{-12} \text{ cm}^3 \text{ s}^{-1}$.

Other mechanisms can be ruled out by calculations or experimental observations.

(C) Ionisation by hot electrons:



With the laser powers used, these electrons are mainly formed by photoionisation. A calculation, for the 8s level, using the Percival theory (1966) gives an electron impact ionisation cross section of $3 \times 10^{-14} \text{ cm}^2$ at an electron energy of 1.5 eV. Such a cross section provides ion currents three orders of magnitude less than the total measured atomic current.

(D) The ion currents collected with only the IR laser beam tuned at the resonance are also at least three orders of magnitude less than the measured currents. No current is detected when this laser is detuned.

11. Discussion and conclusion

The experimental associative rate coefficients obtained for s or d levels are of the same order of magnitude as the coefficients measured by Klucharev (1980) for p levels. This seems to indicate a negligible influence of the orbital quantum number on the cross section value. Figure 6 gives the set of experimental results together

with the theoretical curves obtained by Mihajlov and Janev (1981) for two values of the Rb polarisability.

However, it should be noticed that the calculated curves are only valid for p and d levels. The Mihajlov model gives results depending on the photoionisation cross section near the threshold. The Kramer approximation used by Mihajlov is reasonable for the calculation of the p or d photoionisation cross sections near the threshold but gives results about ten times too high for the s levels. Actually, the theoretical value of $k_2(ns)$ should rather be ten times lower than $k_2(np)$ or $k_2(nd)$, in disagreement with the experimental observation.

Lower values of associative ionisation cross sections for s states have been obtained theoretically and experimentally in helium (Cohen 1976): $k(3s)/k(3p) = 1:30$ and $k(3s)/k(3d) = 1:50$ but also $k(4s)/k(4d) = 1:3$ and $k(4s)/k(4d) = 1:7$. These data might indicate that when n increases associative ionisation cross sections for an s state and a p or d state come close together.

Recently Kushawaha and Leventhal (1982) determined rate coefficients of 7.8×10^{-11} and $3.3 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1}$ for the associative ionisation of the 4d and 5s sodium atoms colliding with a ground-state Na atom (3s). The first coefficient lies in the range of the present measurements and the second seems to agree with the preceding theoretical remarks. The 5s and 4d levels are the lowest states of Na for which the associative ionisation with a ground-state atom is energetically possible which prevents any correction for cascading de-excitation. However, this mechanism gives a concentration for the 4p level which is half that of the 5s state. Under these conditions it is highly surprising that the authors do not see any Na^+ ions. The photoionisation of the 4p level should give an atomic ion current of the same order of magnitude as the measured molecular ion current according to the published 5s density and laser powers. They also considered the Penning ionisation coefficient to be very low: this is in disagreement with the results of Le Gouët *et al* (1982). A Penning rate coefficient of the order of $10^{-8} \text{ cm}^3 \text{ s}^{-1}$ is used by these authors to account for electron energy distribution in excited Na vapour.

The present experimental work indicates clearly the importance of Penning ionisation when two excited short-lifetime Rb atoms collide. It would be of a great interest to have a specific theoretical interpretation for (i) Penning ionisation including two excited states, and (ii) associative ionisation especially involving ns states.

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