



Fusion and braiding in finite and affine Temperley-Lieb categories

A. M. Gainutdinov, H. Saleur

► To cite this version:

A. M. Gainutdinov, H. Saleur. Fusion and braiding in finite and affine Temperley-Lieb categories. 2016. cea-01468258

HAL Id: cea-01468258

<https://cea.hal.science/cea-01468258>

Preprint submitted on 15 Feb 2017

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Fusion and braiding in finite and affine Temperley–Lieb categories

A.M. Gainutdinov ^{a,b} and H. Saleur ^{c,d}

^a DESY, Theory Group, Notkestrasse 85, Bldg. 2a, 22603 Hamburg, Germany

^b Laboratoire de Mathématiques et Physique Théorique CNRS,
Université de Tours, Parc de Grammont, 37200 Tours, France

^c Institut de Physique Théorique, CEA Saclay, Gif Sur Yvette, 91191, France

^d Department of Physics and Astronomy, University of Southern California,
Los Angeles, CA 90089, USA

ABSTRACT. Finite Temperley–Lieb algebras are “diagram algebra” quotients of (the group algebra of) the famous Artin’s braid group \mathbb{B}_N , while the affine Temperley–Lieb algebras arise as diagram algebras from a generalized version of the braid group. We study asymptotic “ $N \rightarrow \infty$ ” representation theory of these quotients (parametrized by $\mathfrak{q} \in \mathbb{C}^\times$) from a perspective of braided monoidal categories.

Using certain idempotent subalgebras in the finite and affine algebras, we construct infinite “arc” towers of the diagram algebras and the corresponding direct system of representation categories, with terms labeled by $N \in \mathbb{N}$.

The corresponding direct-limit category is our main object of studies.

For the case of the finite Temperley–Lieb algebras, we prove that the direct-limit category is abelian and highest-weight at any $\mathfrak{q} \in \mathbb{C}^\times$ and endowed with braided monoidal structure. The most interesting result is when \mathfrak{q} is a root of unity where the representation theory is non-semisimple. The resulting braided monoidal categories we obtain at different roots of unity are new and interestingly they are not rigid. We observe then a fundamental relation of these categories to a certain representation category of the Virasoro algebra and give a conjecture on the existence of a braided monoidal equivalence between the categories. This should have powerful applications to the study of the “continuum” limit of critical statistical mechanics systems based on the Temperley–Lieb algebra.

We also introduce a novel class of embeddings for the affine Temperley–Lieb algebras (algebras of diagrams drawn without intersections on the surface of a cylinder) and related new concept of fusion or bilinear N -graded tensor product of modules for these algebras. We prove that the fusion rules are stable with the index N of the tower and prove that the corresponding direct-limit category is endowed with an associative tensor product. We also study the braiding properties of this affine Temperley–Lieb fusion. Potential relationship with representations of the product of two Virasoro algebras are left for future work.

1. INTRODUCTION

Let \mathbb{K} be a field and \mathfrak{q} is an invertible element in \mathbb{K} . Graham and Lehrer introduced in [1] a sequence $\mathcal{T}_N^{\mathfrak{q}}$ (with $N = 1, 2, 3, \dots$) of infinite-dimensional algebras over \mathbb{K} as the sets of endomorphisms of the objects N in a certain category of diagrams that depends on \mathfrak{q} . These algebras are called affine Temperley–Lieb and they are extended versions of the Temperley–Lieb quotients of the affine Hecke algebras of type \widehat{A}_{N-1} . They have bases consisting of diagrams drawn without intersections on the surface of a cylinder. A slightly different version of these algebras (the one without the translation generator) was introduced much earlier by Martin and Saleur [2] in the context of 2d statistical mechanics models such as loop models on a cylinder. One can cut the cylinder lengthwise and obtain then a much smaller algebra of diagrams on the strip which is the famous (finite) Temperley–Lieb algebra [3].

The categorical and statistical physics interpretations of the finite and affine Temperley–Lieb (TL) algebras made possible many interesting connections between different parts of mathematics (representation theory, knot theory, low-dimensional topology) and quantum physics. Recent proposals for “fault tolerant” topological quantum computers for instance – which may be realized experimentally using fractional quantum Hall devices – are based on models of non-abelian anyons, whose Hilbert space, interactions and computational properties (e.g. qbit gates) are all expressed in terms of certain representations of the Temperley–Lieb algebra (see e.g. [4, 5], and [6] for a recent review.)

From a different, more physical point of view, it is a well established experimental fact that the properties of two-dimensional statistical mechanics lattice models at their critical point [7] (e.g., those based on representations of the TL algebras with $\mathfrak{q} \in \mathbb{C}^\times$ and when $|\mathfrak{q}| = 1$) can be described using conformal field theories – that is, quantum field theories which are invariant under conformal transformations [8, 9]. This means, for instance, that expectation values of products of local physical observables such as the spin or the energy density can be identified with the Green functions of the conformal field theory – the latter being calculated using abstract algebraic manipulations based on the Virasoro algebra representation theory [9, 10].

It is important to understand that the identification between properties of the lattice model and of the quantum field theory can only hold in the so called *scaling limit* or *the continuum limit*. This requires in particular that all scales (such as the size N of the system) be much larger than the lattice spacing limit. In terms of the finite or affine Temperley–Lieb algebras, this means therefore that one needs to study $N \rightarrow \infty$.

Understanding mathematically the continuum limit of the two-dimensional statistical mechanics lattice models at their critical point remains a challenge. Many works in the mathematical literature have concentrated on the existence of the limit, and investigations of conformal invariance and what it may mean for models defined initially on a discrete lattice. Significant progress was obtained in some cases (such as the Ising model or percolation) spurred in part by developments around the Schramm–Loewner Evolution (see e.g. [11, 12]). For a recent attempt at constructing the conformal stress-energy tensor or the Virasoro algebra on the lattice, see [13, 14].

A different, algebraic route was pioneered by physicists [15], who investigated the potential relations between the Lie brackets of the Virasoro algebra generators and of certain elements in the associative algebras – such as the Temperley–Lieb algebra – underlying the lattice models. This approach saw considerable renewed interest in the last few years as it shows great promise in helping to understand logarithmic conformal field theories.

For recent reviews and references about these theories – from both mathematics and physics point of view see the special volume [16]. For works (in mathematical physics) on relationships between the Temperley–Lieb lattice models and logarithmic conformal field theories in the context of the continuum limit, see [18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28].

From the perspective of the continuum limit, we have thus an interesting problem of understanding the “asymptotic” representation theory of the finite and affine TL algebras when N tends to infinity. There are at least two very distinct parts of the problem: (i) to define appropriately a tower of the algebras and understand its inductive limit as a certain infinite-dimensional operator algebra and (ii) to understand representation-theoretic aspects of the limits, i.e. to define and study towers of the representation categories of the algebras. The first problem (i) seems to be very hard because it requires constructing inductive systems of algebras (and their representations) where the homomorphisms respect also the \mathbb{C} -grading by the Hamiltonian of the lattice model, which is a particular element of the algebra (e.g. TL). Such systems are designed to control the growth of the algebras as the number N of sites tends to infinity and to keep only those states and operators that are eigenstates of the Hamiltonian with *finite* eigenvalue in the limit¹. By the definition, the inductive limit of such an inductive system gives what physicists call *the continuum limit* of lattice models and we use this term as well. Unfortunately, our understanding of the spectrum and eigenstates of the Temperley–Lieb Hamiltonians is very limited, so only rather simple cases within this program were appropriately rigorously treated [28] (see also [25]) and have indeed shown an explicit connection with the conformal field theory on the operator level.

The main goal of our paper is to study the second point (ii) or the “categorical” part of the asymptotic representation theory of the finite and affine TL algebras. Using certain idempotent subalgebras in the finite and affine cases, we construct infinite *arc*-towers of the diagram algebras and the corresponding direct system of representation categories, with terms labeled by $N \in \mathbb{N}$. In more details for the finite TL case, we introduce the “big” representation category $\bigoplus_N \text{TL}_N\text{-mod}$ encapsulating all representations of the TL algebras for all N , and a fixed $\mathfrak{q} \in \mathbb{C}^\times$. We then define naturally the tensor product (bi-functor) in the category using the induction: it is endowed with an *associator* and *braiding* isomorphisms that we prove do satisfy the coherence conditions (the pentagon and hexagon identities).

Inside the \mathbb{N} -graded category of finite TL representations, we construct the “shift” functors $\mathcal{F}_N : \text{TL}_N\text{-mod} \rightarrow \text{TL}_{N+2}\text{-mod}$ associated with the arc-towers of the TL algebras. These functors have a nice property: they map the standard modules to the standard ones of the same weight, projective covers to projective covers, etc. The inductive or direct system of the TL-representation categories thus defined has also certain nice and important properties: the associators are mapped to associators and braiding isomorphism to braiding isomorphism. We then prove that all these structures endow the corresponding direct limit with a braided monoidal category structure, which we denote as \mathcal{C}_∞ . This construction can be considered as a machinery to prove that the TL fusion rules are stable with N for any $\mathfrak{q} \in \mathbb{C}^\times$. When \mathfrak{q} is generic, the category \mathcal{C}_∞ is semi-simple and have left and right duals. Interestingly, the construction is well defined even for \mathfrak{q} a root of unity, which is the most interesting case for the applications,

¹The idea here comes from physics: one should keep at $N \rightarrow \infty$ only relevant “physical” excitations – finite-energy states – above the ground states (that is, the Hamiltonian’s eigenstates of minimum eigenvalue). Note that the Hamiltonian in this discussion has to be properly normalized.



FIGURE 1. The map of the two translation generators $u^{(1)} \in \mathsf{T}_{N_1}^a$ and $u^{(2)} \in \mathsf{T}_{N_2}^a$, with $N_1 = 3$ and $N_2 = 2$, into $\mathsf{T}_{N_1+N_2}^a$ in terms of affine TL diagrams where each crossing/braiding at i -th site has to be replaced by a linear combination of identity and e_i .

in particular to low-dimensional topology and logarithmic conformal field theory. In the root of unity case, we prove that the direct-limit categories we obtain are abelian and highest-weight. They have a *non-simple* tensor unit and in particular not all objects have duals, so the categories are not rigid when \mathfrak{q} is a root of unity.

We also give a rather explicit description of objects in the category C_∞ (subquotient structure of standard and projective objects) and describe the fusion rules. It turns out that this data matches what we know about certain category of the Virasoro algebra representations at critical central charges, so we have in particular an equivalence of abelian categories, described in Prop. 7.1. We formulate (for the first time) an explicit conjecture (in Conj. 7.2) that the two categories are equivalent as braided monoidal categories and give few supporting arguments.

This result provides a mathematical framework to understand the observations in [20, 21] that the representation theory of the Virasoro algebra relevant to the description of the scaling limit of statistical models and the representation theory of the Temperley–Lieb algebra occurring in the lattice formulation of these models “are very similar” (this observation has spurred many developments in the analysis of Virasoro and Temperley–Lieb-like algebras modules – see e.g. [23, 29, 30]).

Note that, since we are dealing with the ordinary Temperley–Lieb algebra, the statistical models here have “open boundary conditions”, which is known to correspond, in the scaling limit, to boundary conformal field theory, involving a single Virasoro algebra [31]. Meanwhile, the physics of critical statistical lattice models away from their boundaries (the so called “bulk” case) is described by two copies of the Virasoro algebra, corresponding to the chiral and anti-chiral dependencies of the correlation functions. The ideas in [19] can then be extended, and involve now, on the lattice side, models with periodic boundary conditions, for which the correct algebra is now the corresponding affine Temperley–Lieb algebra.

Our second result is concentrated on the limit of affine TL representation categories and on a new tensor-product or fusion in the enveloping category (where we call it *the affine TL fusion*) and eventually in the direct-limit category. The affine TL fusion is based on a non-trivial embedding of two affine or periodic TL algebras for N_1 and N_2 into the “big” algebra $\mathsf{T}_{N_1+N_2}^a$ on $N_1 + N_2$ sites. Note that there is an obvious embedding in the finite TL case while in the periodic or affine case it is not a priori clear how to embed the product of two periodic TL algebras into another periodic TL algebra. We introduce a novel diagrammatic way for this in Sec. 3 and describe here shortly only the embedding of the two translation generators $u^{(1)}$ and $u^{(2)}$, as in Fig. 1. (We have thus constructed a novel tower of the affine TL algebras.) The affine TL fusion for the affine TL modules is then defined as the induction from such subalgebra $\mathsf{T}_{N_1}^a \otimes \mathsf{T}_{N_2}^a$ to $\mathsf{T}_{N_1+N_2}^a$.

We then use the machinery elaborated in the finite TL case (in Sec. 5: note it is actually a very general construction that can be applied to a relatively general class of diagram algebras) and prove that the affine TL fusion rules are stable with N and the tensor product is associative.

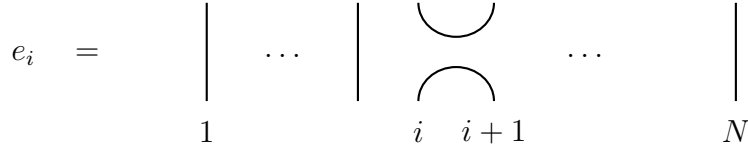
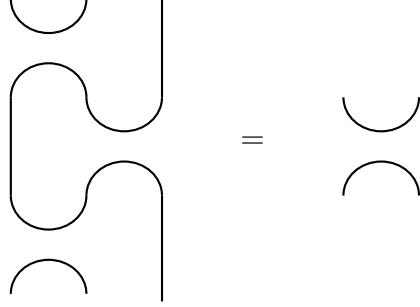
Like in the finite TL case, inside the \mathbb{N} -graded category of affine TL representations, we construct the “shift” functors $\mathcal{F}_N : \mathbb{T}_N^a - \text{mod} \longrightarrow \mathbb{T}_{N+2}^a - \text{mod}$ corresponding to what we call the arc-towers of the affine TL algebras and the corresponding direct limit that is denoted by $\widehat{\mathbb{C}}_\infty$. We have proven that this limit has a tensor product with an associator but there is no braiding (we have computed some of the fusion rules and they are apparently non-commutative). However, we have shown that there is a semi-braiding connecting ‘chiral’ and ‘anti-chiral’ tensor products, the one defined above and the other defined similarly but interchanging over-crossings with under-crossings.

This periodic or “bulk” case is much harder than the “open-boundary” case, and less explored, despite some beautiful results on some simple cases [22, 32, 28, 33, 34, 35]. A key ingredient that was missing so far was how to define for lattice models a fusion that may correspond, in the continuum limit, to the fusion of non chiral fields. The results of the present paper provide a potential way to do this, which will be explored more in our subsequent work [36].

The paper is organized as follows. We first recall in section 2 the definition and well-known facts about the finite Temperley–Lieb (TL) algebra, including the fusion of TL modules, then introduce a tower of finite TL algebras that leads to the concept of an “enveloping” TL representation category endowed with \mathbb{N} -graded bilinear tensor product – an object which, we show later, is crucial for comparison with results from physics. We introduce then a novel tower for the affine TL algebras in section 3 where we also introduce our diagrammatic fusion for affine TL modules. This fusion is further considered in section 4 where a connection to the affine Hecke algebra is discussed. We start in section 5 our discussion of limits of TL representation categories and introduce a braided monoidal structure on the direct limit of the TL categories, with the main result in Thm. 5.8. This section is rather technical and long but provides a general approach for studying the limits associated with towers of algebras. This approach is then used in section 6 in the affine TL case where our main result is formulated in Thm. 6.10. We conclude in section 7 by a discussion of the relation of our work with the Virasoro algebra representation theory, Prop. 7.1, and write down a fundamental conjecture 7.2 about an equivalence of braided tensor categories. In Appendix A, we give several examples of the diagrammatic calculation of affine TL fusion rules.

Conventions: Throughout the paper we fix the field $\mathbb{K} = \mathbb{C}$ for convenience, though most of our results hold for any field \mathbb{K} (we need the characteristic zero only when we discuss the subquotient structure of the TL representations). We also denote by \mathbb{N} the additive semi-group of positive integers $\{1, 2, 3, \dots\}$.

Acknowledgements: We are grateful to Jesper L. Jacobsen for fruitful discussions and general interest in this work. This work was supported in part by the ERC Advanced Grant NuQFT. The work of AMG was also supported by a Humboldt fellowship, DESY and CNRS. AMG thanks IPHT/Saclay for its kind hospitality during last few years where a part of this work was done. AMG also thanks the organizers of the workshop *Conformal Field Theories and Tensor categories*, International Center for Mathematical

FIGURE 2. The diagrammatic representation of e_i .FIGURE 3. The diagrammatic version of the relation $e_i e_{i+1} e_i = e_i$.

Research, Beijing, 2011 where some of our results (those on the direct limit of finite TL categories and the relation to the Virasoro algebra) were presented.

2. FUSION IN TL-MOD CATEGORIES

The (finite) Temperley–Lieb (TL) algebra $\mathrm{TL}_N(m)$ is an associative algebra over \mathbb{C} generated by unit $\mathbf{1}$ and e_j , with $1 \leq j \leq N-1$, satisfying the defining relations

$$\begin{aligned}
 (2.1) \quad e_j^2 &= m e_j, \\
 e_j e_{j \pm 1} e_j &= e_j, \\
 e_j e_k &= e_k e_j \quad (j \neq k, k \pm 1).
 \end{aligned}$$

This algebra has a well-known faithful diagrammatical representation in terms of non-crossing pairings on a rectangle with N points on each of the opposite sides. Multiplication is performed by placing two rectangles on top of each other, and replacing any closed loops by a factor m . While the identity corresponds to the diagram in which each point is directly connected to the point above it, the generator e_i is represented by the diagram, see Fig. 2, where the points i on both sides of the rectangle are connected to the point $i+1$ on the same side, all other points being connected like in the identity diagram. The defining relations are easily checked by using isotopy ambient on the boundary of the rectangle, see Fig. 3.

We will often omit mentioning the parameter m and write simply TL_N as the replacement for $\mathrm{TL}_N(m)$.

2.1. Towers of the TL algebras. The important ingredient of our constructions below are towers of the TL algebras. In terms of the diagrams, we can naturally construct two kinds of towers.

- The first one is standard, it uses *the standard embeddings* of the algebras:

$$(2.2) \quad \mathrm{TL}_N \xrightarrow{\iota} \mathrm{TL}_{N+1}$$

such that

$$(2.3) \quad \iota(e_j) = e_j, \quad 1 \leq j \leq N-1,$$

or in terms of the TL diagrams one adds one vertical string on the right of the diagram considered as an element in \mathbf{TL}_N – this gives an element in \mathbf{TL}_{N+1} .

- The second tower uses what we call the arc-embeddings: a diagram on N sites is enlarged up to the diagram on $N+2$ sites by adding arcs, see precise definition below in Sec. 5.1.

The first type of TL towers is used in the definition of the TL fusion, while the second is used in constructing direct limits of the TL representation categories.

2.2. TL fusion. In this section, we recall *fusion* for modules over the (finite) TL algebras initially introduced in the physics literature [19] and further studied on more formal grounds in [21], which we follow in terms of conventions and notations. The fusion’s definition is based on the standard embeddings (2.3). We then use this fusion construction to define \mathbb{N} -graded tensor-product structure (with an associator) on the direct sum of the categories of TL representations.

Let \mathbf{C}_N denotes the category of finite-dimensional \mathbf{TL}_N -modules (we will usually drop the parameter m for brevity.)

Definition 2.2.1 ([19, 21]). *Let M_1 and M_2 be two modules over \mathbf{TL}_{N_1} and \mathbf{TL}_{N_2} respectively. Then, the tensor product $M_1 \otimes M_2$ is a module over the product $\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2}$ of the two algebras. Using the standard embedding, we consider this product of algebras as a subalgebra in \mathbf{TL}_N , for $N = N_1 + N_2$. The fusion (bi-)functor*

$$(2.4) \quad \times_f : \mathbf{C}_{N_1} \times \mathbf{C}_{N_2} \rightarrow \mathbf{C}_{N_1+N_2}$$

on two modules M_1 and M_2 is then defined as the module induced from this subalgebra, i.e.,

$$(2.5) \quad M_1 \times_f M_2 = \mathbf{TL}_N \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2})} M_1 \otimes M_2 ,$$

where we used the balanced tensor product over $\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2}$.

An explicit calculation of the TL fusion for a large class of indecomposable representations at any non-zero \mathfrak{q} (i.e. including the root of unity cases) is given in [21], see also recent works [29, 27, 30].

We emphasize that the TL fusion is associative, i.e. the functor \times_f in (2.4) is equipped with a family of natural isomorphisms $(M_1 \times_f M_2) \times_f M_3 \cong M_1 \times_f (M_2 \times_f M_3)$, for each triple of \mathbf{TL}_{N_i} modules M_i , $i = 1, 2, 3$, and each triple of the natural numbers (N_1, N_2, N_3) . We leave the proof of the associativity till Sec. 5.4 and give the associator explicitly in Prop. 5.4.1 and prove the pentagon identity in Prop. 5.4.2. On top of it, we also have the braiding, i.e. the tensor product \times_f is commutative, see Sec. 5.5 below.

It is then natural to introduce a large category embracing these structures. We call it the “enveloping” TL representation category (that formally contains TL representations at any N)

$$(2.6) \quad \mathbf{C} = \bigoplus_{N \geq 1} \mathbf{C}_N .$$

The direct sum here means that \mathbf{C} contains \mathbf{C}_N as a full subcategory and there are no morphisms between the full subcategories for different N . The category \mathbf{C} is thus graded

by \mathbb{N} . We will label an object M from \mathbf{C}_N as $M[N]$ to emphasize its grade. We can thus consider \times_f defined in (2.5) as an \mathbb{N} -graded tensor-product functor on \mathbf{C} :

Proposition 2.2.2. *Let \times_f denote the \mathbb{N} -graded bilinear tensor product on \mathbf{C} as defined for each pair $(N_1, N_2) \in \mathbb{N} \times \mathbb{N}$ in (2.4) and (2.5). It is equipped with an associator satisfying the pentagon identity.*

Note that we do not have the tensor unit in \mathbf{C} , as we do not have the zero grade. Let us call such a category (a monoidal category weakened by removing the tensor unit axioms) as a *semi-group category*.

For the later use, we will need the structure of a highest-weight category on \mathbf{C}_N and we recall below few standard facts from the representation theory of $\mathbf{TL}_N(\mathfrak{q} + \mathfrak{q}^{-1})$.

2.3. Standard and projective \mathbf{TL} modules. We also recall [3] *the standard $\mathbf{TL}_N(m)$ modules $\mathcal{W}_j[N]$ of weight $x \leq j \leq N/2$, where $x = \frac{1}{2}(N \bmod 2)$. First, we need to introduce “half-diagrams” (usually called link states) obtained from Temperley-Lieb diagrams (i.e., non-crossing pairings on a rectangle with N points on each of the opposite sides) by cutting these diagrams horizontally in the middle. Each half has N points: some of them are connected by arcs, and some others are not connected to anything. The latter are often called *through-lines* (or defects). The algebra acts in the obvious diagrammatic way by concatenating Temperley-Lieb diagrams with link states, eliminating all loops at the price of multiplying the diagram by m^n , where n is the number of loops, and keeping track of the connectivities using isotopy. It is clear that the number of through-lines cannot increase under the action of the algebra. Standard modules $\mathcal{W}_j[N]$ are obtained by letting the algebra act as usual when the number of through-lines – denoted by $2j$ – is conserved, and setting this action to zero when the number of through-lines decreases. It is well known that these modules are irreducible for \mathfrak{q} generic [3], while their dimension is given by differences of binomial coefficients*

$$(2.7) \quad d_j[N] = \binom{N}{\frac{N}{2} + j} - \binom{N}{\frac{N}{2} + j + 1}.$$

It is well-known also [37] that the category of finite-dimensional \mathbf{TL}_N modules is a highest-weight category, i.e., it has a special class of objects – the standard modules $\mathcal{W}_j[N]$ – with morphisms only in one direction (there is a homomorphism from $\mathcal{W}_k[N]$ to $\mathcal{W}_j[N]$ only if $k \geq j$) and the projective and injective modules have a nice filtration with sections given by the standard and costandard modules, respectively, see more details in [38]. When \mathfrak{q} is generic (not a root of unity) then the category is semi-simple and the structure of a highest-weight category is trivial.

Let $\mathfrak{q} = e^{i\pi/p}$ for integer $p \geq 2$ and set $s \equiv s(j) = (2j + 1) \bmod p$, we then recall the subquotient structure of $\mathcal{W}_j[N]$, which was studied in many works from different perspectives [29, 21, 37, 39] (though we use slightly different conventions here). The standard modules with $s(j) = 0$ are simple. For non-zero $s(j)$ there is a non-trivial homomorphism from $\mathcal{W}_k[N]$ to $\mathcal{W}_j[N]$ only if $k = j$ or $k = j + p - s$, and in the second case:

$$\phi_j : \mathcal{W}_{j+p-s}[N] \rightarrow \mathcal{W}_j[N]$$

with $\ker \phi_j$ given by the socle of $\mathcal{W}_{j+p-s}[N]$ and $\operatorname{im} \phi_j$ is the socle of $\mathcal{W}_j[N]$ and is isomorphic to the head of $\mathcal{W}_{j+p-s}[N]$. Note that ϕ_j exists only if $2(j + p - s) \leq N$. We

thus have the subquotient diagram (for non-zero s)

$$\mathcal{W}_j[N] = \mathcal{X}_j[N] \longrightarrow \mathcal{X}_{j+p-s}[N]$$

where we introduce the notation $\mathcal{X}_j[N]$ for the irreducible quotient of $\mathcal{W}_j[N]$ and set here $\mathcal{X}_k[N] \equiv 0$ if $2k > N$.

Let $\mathcal{P}_j[N]$ denotes the projective cover of the simple module $\mathcal{X}_j[N]$. The subquotient structure of the projective covers can be easily deduced due to a reciprocity relation for a highest-weight category. Let $[\mathcal{W}_j : \mathcal{X}_{j'}]$ and $[\mathcal{P}_j : \mathcal{W}_{j'}]$ denote the number of appearances of $\mathcal{X}_{j'}[N]$ in the subquotient diagram for $\mathcal{W}_j[N]$ and the number of appearances of $\mathcal{W}_{j'}[N]$ in the standard filtration for the projective cover $\mathcal{P}_j[N]$, respectively. Then, the reciprocity relation reads

$$(2.8) \quad [\mathcal{P}_j : \mathcal{W}_{j'}] = [\mathcal{W}_{j'} : \mathcal{X}_j],$$

or, in words, the projective modules \mathcal{P}_j are composed of the standard modules that have the irreducible module \mathcal{X}_j as a subquotient. The projective covers $\mathcal{P}_j[N]$ are then simple if $s(j) = 0$, they are equal to $\mathcal{W}_j[N]$ for $0 \leq j \leq \frac{1}{2}(p-2)$ and otherwise have the following structure

$$(2.9) \quad \mathcal{P}_j[N] = \begin{array}{ccc} & \mathcal{X}_j[N] & \\ \swarrow & & \searrow \\ \mathcal{X}_{j-s}[N] & & \mathcal{X}_{j+p-s}[N] \\ \searrow & & \swarrow \\ & \mathcal{X}_j[N] & \end{array}$$

where the nodes are irreducible subquotients and the arrows correspond to \mathbf{TL}_N action, i.e., $\mathcal{X}_j[N]$ in the bottom of the diagram is the irreducible submodule (or socle), while the socle of the quotient $\mathcal{P}_j[N]/\mathcal{X}_j[N]$ is the direct sum $\mathcal{X}_{j-s}[N] \oplus \mathcal{X}_{j+p-s}[N]$ in the middle of the diagram, etc. Also note that the diagram (2.9) has three nodes instead of four if $2(j+p-s) > N$.

3. AFFINE TEMPERLEY–LIEB EMBEDDING

We have seen in the previous section embeddings of the finite Temperley–Lieb algebras $\mathbf{TL}_{N_1} \rightarrow \mathbf{TL}_N$, for $N_1 < N$, that are naturally defined in terms of TL diagrams by adding the vertical strings, or by the use of the standard embeddings $\mathbf{TL}_N \rightarrow \mathbf{TL}_{N+1}$ repeatedly. This standard embedding is the basic step in the definition of the TL fusion. Constructing embeddings of periodic or affine TL algebras is a non-trivial problem that we solve in this section. We first recall the definition of the affine TL algebras (also parametrized by $N \in \mathbb{N}$) and then propose a novel diagrammatical way of defining a tower of these algebras.

3.1. The affine Temperley–Lieb algebras. We recall here two equivalent definitions of the affine Temperley–Lieb algebra - independently introduced and studied in many works [40, 41, 2, 1, 42]. We follow mainly conventions and notations from the work of Graham and Lehrer [1, 43] whenever possible.

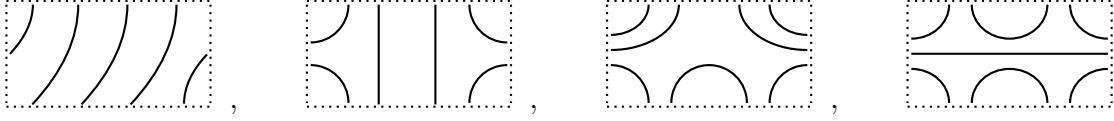


FIGURE 4. Examples of affine diagrams for $N = 4$, with the left and right sides of the framing rectangle identified. The first diagram represents the translation generator u while the second diagram is for the generator $e_4 \in \mathcal{T}_4^a(m)$. The third and fourth ones are examples of $j = 0$ diagrams.

3.1.1. *Definition I: generators and relations.* The affine Temperley–Lieb (aTL) algebra $\mathcal{T}_N^a(m)$ is an associative algebra over \mathbb{C} generated by u , u^{-1} , and e_j , with $j \in \mathbb{Z}/N\mathbb{Z}$, satisfying the defining relations

$$(3.1) \quad \begin{aligned} e_j^2 &= m e_j, \\ e_j e_{j \pm 1} e_j &= e_j, \\ e_j e_k &= e_k e_j \quad (j \neq k, k \pm 1), \end{aligned}$$

which are the standard TL relations but defined now for indices modulo N , and

$$(3.2) \quad \begin{aligned} u e_j u^{-1} &= e_{j+1}, \\ u^2 e_{N-1} &= e_1 \dots e_{N-1}, \end{aligned}$$

where the indices $j = 1, \dots, N$ are again interpreted modulo N .

3.1.2. *Definition II: diagrammatic.* The affine Temperley–Lieb (TL) algebra $\mathcal{T}_N^a(m)$ is an associative algebra over \mathbb{C} spanned by particular diagrams on an annulus with N sites on the inner and N on the outer boundary. The sites are connected in pairs, and only configurations that can be represented using lines inside the annulus that do not cross are allowed. Diagrams related by an isotopy leaving the labeled sites fixed are considered equivalent. We call such (equivalence classes of) diagrams *affine* diagrams. Examples of affine diagrams are shown in Fig. 4 for $N = 4$, where we draw them in a slightly different geometry: we cut the annulus and transform it to a rectangle which we call *framing* so that the sites labeled by ‘1’ are closest to the left and sites labeled by ‘ N ’ are to the right sides of the rectangle. Multiplication $a \cdot b$ of two affine diagrams a and b is defined in a natural way, by joining an inner boundary of a to an outer boundary of the annulus of b , and removing the interior sites. Whenever a closed contractible loop is produced when diagrams are multiplied together, this loop must be replaced by a numerical factor m that we often parametrize by \mathfrak{q} as $m = \mathfrak{q} + \mathfrak{q}^{-1}$.

We also note that the diagrams in this algebra allow winding of through-lines around the annulus any integer number of times, and different windings result in independent algebra elements. Moreover, in the ideal of zero through-lines, any number of non-contractible loops (like in the fourth diagram in Fig. 4) is allowed. The algebra $\mathcal{T}_N^a(m)$ is thus infinite-dimensional. For $N = 1$, it is just the polynomial algebra $\mathbb{C}[u, u^{-1}]$.

3.2. **The affine TL and the affine braid group.** Let \mathbb{CB}_N be the group algebra of the Artin’s braid group. As an associative algebra, it is generated by $g_i^{\pm 1}$, with $1 \leq i \leq N-1$, subject to $g_i g_j = g_j g_i$ for $|i - j| > 1$ and to the standard braid relations:

$$(3.3) \quad g_i g_{i \pm 1} g_i = g_{i \pm 1} g_i g_{i \pm 1}$$

or with the graphical notation

$$g_i = \begin{array}{c} \diagup \quad \diagdown \\ \diagdown \quad \diagup \end{array} \quad g_i^{-1} = \begin{array}{c} \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array}$$

the relations (3.3) can be graphically depicted as

$$(3.4) \quad \begin{array}{c} \diagup \quad \diagdown \\ \diagdown \quad \diagup \\ \diagup \quad \diagdown \\ \diagdown \quad \diagup \end{array} = \begin{array}{c} \diagdown \quad \diagup \\ \diagup \quad \diagdown \\ \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array}$$

It is well-known that the finite TL algebra $\text{TL}_N(\mathbf{q} + \mathbf{q}^{-1})$ is a finite-dimensional quotient of \mathbb{CB}_N where we set

$$(3.5) \quad g_i^{\pm 1} = \pm i(\mathbf{q}^{\pm \frac{1}{2}} \mathbf{1} - \mathbf{q}^{\mp \frac{1}{2}} e_i)$$

and imply the TL relations (2.1).

Let now $\widehat{\mathbb{B}}_N$ be the affine braid group – the group of braids on the surface of a cylinder – it is generated by the translation u (like above) and $g_i^{\pm 1}$, with $0 \leq i \leq N - 1$, subject to $u g_i^{\pm 1} u^{-1} = g_{i+1}^{\pm 1}$ and to the braid relations (3.3) where the index i is now interpreted modulo N . We recall [45] that the affine TL algebra $\mathbb{T}_N^a(\mathbf{q} + \mathbf{q}^{-1})$ can be defined as a quotient of $\widehat{\mathbb{CB}}_N$ where we again set (3.5) and imply the relations (3.1) and (3.2).

Using this connection of \mathbb{T}_N^a with the braid groups, we will sometimes use below the braid generators as a replacement for the TL generators. In the diagrams like (3.4) we emphasize that each under/above crossing of lines should be interpreted as the replacement for the linear combination in (3.5).

3.3. A tower of affine TL algebras. Our approach to the fusion of affine TL modules in the next section relies on the induction functor which associates with any pair of modules over the algebras $\mathbb{T}_{N_1}^a(m)$ and $\mathbb{T}_{N_2}^a(m)$ a module over the bigger algebra $\mathbb{T}_{N_1+N_2}^a(m)$. This functor uses an explicit embedding of the two “small” algebras on N_1 and N_2 sites into the “big” one on $N_1 + N_2$ sites.

We start with defining a “one-step” embedding $\mathbb{T}_N^a(m) \rightarrow \mathbb{T}_{N+1}^a(m)$:

$$(3.6) \quad \begin{aligned} u^{(1)} &\mapsto u g_N^{-1} , \\ e_i^{(1)} &\mapsto e_i, \quad 1 \leq i \leq N - 1 , \\ e_0^{(1)} &\mapsto g_N e_0 g_N^{-1} , \end{aligned}$$

where we label the generators in \mathbb{T}_N^a with the superscript (1) and g_N stands for the combination in (3.5). It is straightforward to check that this map is an algebra map. The kernel of this map is trivial: we have a basis in the image given by placing an extra “vertical” string between the N th and 1st sites of the cylindrical or affine diagram for a basis element in \mathbb{T}_N^a , while each crossing is replaced by the corresponding under-crossing: this gives obviously a bijection between the two bases (explicit diagrams will be given below).

We can use the map (3.6) recursively and define the embedding $\mathbb{T}_N^a(m) \rightarrow \mathbb{T}_{N+k}^a(m)$. Similarly, we can embed the product of two affine TL algebras, $\mathbb{T}_{N_1}^a$ and $\mathbb{T}_{N_2}^a$, into \mathbb{T}_N^a

with $N = N_1 + N_2$. Let us denote the generators in the i th algebra as $u^{(i)}$ and $e_j^{(i)}$, with $i = 1, 2$, and use standard notations for the generators in the “big” algebra \mathbb{T}_N^a . We first define the map on the TL generators $e_j^{(i)}$, where $j \neq 0$, in the standard way

$$(3.7) \quad e_j^{(1)} \mapsto e_j, \quad e_k^{(2)} \mapsto e_{N_1+k}, \quad 1 \leq j \leq N_1 - 1, \quad 1 \leq k \leq N_2 - 1.$$

The translation generators $u^{(1)}$ and $u^{(2)}$ are mapped as (recall, we set $N = N_1 + N_2$)

$$(3.8) \quad u^{(1)} \mapsto u g_{N-1}^{-1} \dots g_{N_1}^{-1}, \quad u^{(2)} \mapsto g_{N_1} \dots g_1 u.$$

In terms of diagrams, these two translation generators are presented as simply as

$$(3.9) \quad u^{(1)} \mapsto \begin{array}{c} \text{Diagram 1: A braid with 5 strands. The top two strands cross twice. The bottom three strands are straight. A dashed box encloses the top two strands and the crossing region.} \end{array} = \begin{array}{c} \text{Diagram 2: A braid with 5 strands. The top two strands are straight. The bottom three strands cross twice. A dashed box encloses the bottom three strands and the crossing region.} \end{array}$$

where we assumed that $N_1 = 3$ and $N_2 = 2$, and for the second translation $u^{(2)}$ we have the diagram

$$(3.10) \quad u^{(2)} \mapsto \begin{array}{c} \text{Diagram 1: A braid with 5 strands. The top three strands are straight. The bottom two strands cross twice. A dashed box encloses the bottom two strands and the crossing region.} \end{array} = \begin{array}{c} \text{Diagram 2: A braid with 5 strands. The top two strands cross twice. The bottom three strands are straight. A dashed box encloses the top two strands and the crossing region.} \end{array}$$

or in words the rightmost string of $u^{(1)}$ (the one that starts at position N_1) passes *above* the N_2 through-lines on the right from it and ends at the position 1, and similarly for $u^{(2)}$ – the leftmost string passes *under* the N_1 through-lines on the left from it. It is then an easy (in terms of diagrams) calculation using the braid relations to check

$$(3.11) \quad u^{(1)} u^{(2)} = u^{(2)} u^{(1)}.$$

Due to the normalization of g_i 's as in (3.5), we have the relations

$$(3.12) \quad g_i g_{i+1} e_i = e_{i+1} e_i, \quad g_i^{-1} g_{i+1}^{-1} e_i = e_{i+1} e_i$$

and many others similar to these. In terms of diagrams, these relations tell us that a TL arc (“half” of the diagram for e_i) can be pulled out under or above any string at the price of the factor 1. We can thus simplify calculations using diagrams with braids and TL arcs only. It is only the twisting that produces a non-trivial factor $i q^{\frac{3}{2}}$:

$$(3.13) \quad e_i g_{i+1} e_i = i q^{\frac{3}{2}} e_i$$

but these relations will not appear in calculations below.

Using the remarks above, we then immediately check the affine TL relations

$$(3.14) \quad (u^{(1)})^2 e_{N_1-1} = e_1 \dots e_{N_1-1}, \quad (u^{(2)})^2 e_{N-1} = e_{N_1+1} \dots e_{N-1}.$$

We then define the map on the periodic TL generators $e_0^{(i)}$ as

$$(3.15) \quad e_0^{(1)} \mapsto g_{N_1} \dots g_{N-1} e_0 g_{N-1}^{-1} \dots g_{N_1}^{-1},$$

$$(3.16) \quad e_0^{(2)} \mapsto g_0^{-1} \dots g_{N_1-1}^{-1} e_{N_1} g_{N_1-1} \dots g_0.$$

Note that in terms of diagrams, this result is natural, as illustrated below (again for $N_1 = 3$ and $N_2 = 2$)

$$(3.17) \quad e_0^{(1)} = \begin{array}{c} \text{Diagram: A vertical line with a cup on the left and a cap on the right, connected to a horizontal line with three vertical lines above it.} \end{array}$$

$$(3.18) \quad e_0^{(2)} = \begin{array}{c} \text{Diagram: A vertical line with a cup on the left and a cap on the right, connected to a horizontal line with three vertical lines below it.} \end{array}$$

where $e_0^{(i)}$ are considered now as the images of the maps (3.15) and (3.16) (we will often use the same notation for the images of elements in the subalgebras.)

Using again the diagrammatical calculation, it is straightforward to check that

$$(3.19) \quad (e_0^{(i)})^2 = (\mathfrak{q} + \mathfrak{q}^{-1}) e_0^{(i)}, \quad i = 1, 2.$$

Further, we also check all the other affine TL relations

$$(3.20) \quad e_0^{(i)} = u^{(i)} e_{N_i-1}^{(i)} (u^{(i)})^{-1} = (u^{(i)})^{-1} e_1^{(i)} u^{(i)}, \quad i = 1, 2,$$

and

$$(3.21) \quad e_0^{(i)} e_1^{(i)} e_0^{(i)} = e_0^{(i)}, \quad e_1^{(i)} e_0^{(i)} e_1^{(i)} = e_1^{(i)},$$

$$(3.22) \quad e_0^{(i)} e_{N_i-1}^{(i)} e_0^{(i)} = e_0^{(i)}, \quad e_{N_i-1}^{(i)} e_0^{(i)} e_{N_i-1}^{(i)} = e_{N_i-1}^{(i)},$$

where we recall that $e_k^{(1)} = e_k$ and $e_k^{(2)} = e_{N_2+k}$. We also see using diagrammatic computation that both the subalgebras $\mathbb{T}_{N_1}^a$ and $\mathbb{T}_{N_2}^a$ indeed commute

$$(3.23) \quad e_0^{(1)} e_0^{(2)} = e_0^{(2)} e_0^{(1)},$$

in addition to (3.11).

So, we have thus constructed a homomorphism of algebras

$$(3.24) \quad \varepsilon_{N_1, N_2} : \quad \mathbb{T}_{N_1}^a \otimes \mathbb{T}_{N_2}^a \longrightarrow \mathbb{T}_N^a,$$

with the image of the generators given in (3.7), (3.8) and (3.15), and (3.16). This homomorphism has trivial kernel (by recursively using the one-step embedding (3.6) that has zero kernel), so we have actually an embedding of algebras.

4. FUSION OF AFFINE TL MODULES



In this section, we introduce *fusion* for modules over the affine TL algebras using the embeddings defined in the previous section. We will use this fusion construction in the next section to define a \mathbb{N} -graded tensor product in the affine TL representation category.

Definition 4.1. *Let M_1 and M_2 be two modules over $\mathsf{T}_{N_1}^a(m)$ and $\mathsf{T}_{N_2}^a(m)$ respectively. Then, the tensor product $M_1 \otimes M_2$ is a module over the product $\mathsf{T}_{N_1}^a(m) \otimes \mathsf{T}_{N_2}^a(m)$ of the two algebras. Using the embedding (3.24), we consider this product of algebras as a subalgebra in $\mathsf{T}_N^a(m)$, for $N = N_1 + N_2$. The (affine) fusion functor $\widehat{\times}_f$ on two modules M_1 and M_2 is then defined as the module induced from this subalgebra, i.e.*

$$(4.1) \quad M_1 \widehat{\times}_f M_2 = \mathsf{T}_N^a \otimes_{(\mathsf{T}_{N_1}^a \otimes \mathsf{T}_{N_2}^a)} M_1 \otimes M_2,$$

where we used the balanced tensor product over $\mathsf{T}_{N_1}^a \otimes \mathsf{T}_{N_2}^a$ and we abuse the notation by writing T_N^a instead of $\mathsf{T}_N^a(m)$.

Below we give explicit examples of the affine TL fusion calculation. Before doing this, let us recall the basic T_N^a -modules called the standard modules.

4.2. Standard T_N^a modules. We introduce here the standard modules $\mathcal{W}_{j,z}[N]$ over $\mathsf{T}_N^a(m)$, which are generically irreducible, and give then several examples of explicit calculations of the fusion. The standard modules are parametrized by pairs (j, z) , with a half-integer j and a non-zero complex number z . In terms of diagrams, the first is the number of through-lines, which we denote by $2j$, $0 \leq j \leq N/2$, connecting the inner boundary of the annulus with $2j$ sites and the outer boundary with N sites. For example, the diagrams  and  correspond to $N = 4$ and $j = 1$, where as usual we identify the left and right sides of the framing rectangles, so the diagrams live on the annulus. We call such diagrams *affine*. The action of an element $a \in \mathsf{T}_N^a(m)$ on $v \in \mathcal{W}_{j,z}$ is then defined by stacking the diagrams: joining the inner boundary of a to the outer boundary of the diagram for v , and removing the interior sites. As usual, a closed contractible loop is replaced by the factor $m = \mathbf{q} + \mathbf{q}^{-1}$ (we will often use this parametrisation by a complex number \mathbf{q}). Whenever the affine diagram thus obtained has a number of through lines less than $2j$, the action is zero. For a given non-zero value of j , it is possible in this action to earn a winding number of the through-lines. In this case, we imply the relation [1]

$$\mu = \mu' \circ u_j^n \equiv z^n \mu',$$

where μ is an affine diagram with $2j$ through lines, μ' is a so-called standard diagram which has no through lines winding the annulus and u_j is the translational operator acting on the $2j$ sites of the inner boundary of μ' . Said differently, whenever $2j$ through-lines wind counterclockwise around the annulus l times, we unwind them at the price of a factor z^{2jl} ; similarly, for clockwise winding, the phase is z^{-2jl} [40, 2]. This is for $j > 0$. If $j = 0$, by the concatenating the diagrams we can produce a non-contractible loop and it has to be replaced by the factor $z + z^{-1}$. Such action gives rise to a generically irreducible $\mathsf{T}_N^a(m)$ module, which we denote by $\mathcal{W}_{j,z}[N]$.

The dimensions of these modules $\mathcal{W}_{j,z}$ are then given by

$$(4.2) \quad \hat{d}_j[N] \equiv \dim \mathcal{W}_{j,z}[N] = \binom{N}{\frac{N}{2} + j}, \quad j \geq 0.$$

Note that these numbers do not depend on z (but modules with different z are not isomorphic).

4.3. Examples of the fusion. We consider here examples of the fusion defined in Def. 4.1 for several pairs of standard modules. We also assume that q is generic, i.e. not a root of unity.

4.3.1. Fusion on $1 + 1$ sites. We begin with a simple example of the fusion of the pair of standard modules $\mathcal{W}_{\frac{1}{2},z}$ on $1 + 1$ sites:

$$(4.3) \quad \mathcal{W}_{\frac{1}{2},z_1}[1] \widehat{\times}_f \mathcal{W}_{\frac{1}{2},z_2}[1] = \text{Ind}_{\mathbb{T}_1^a \otimes \mathbb{T}_1^a}^{\mathbb{T}_2^a} \mathcal{W}_{\frac{1}{2},z_1}[1] \otimes \mathcal{W}_{\frac{1}{2},z_2}[1],$$

where the affine TL on 1 site is the commutative algebra generated by the translation generator, i.e. $\mathbb{T}_1^a = \mathbb{C}[u^{\pm 1}]$. We say that the left \mathbb{T}_1^a (the left component of the tensor product $\mathbb{T}_1^a \otimes \mathbb{T}_1^a$) is generated by $u^{(1)}$ and the right one is generated by $u^{(2)}$. Then, the module $\mathcal{W}_{\frac{1}{2},z_1} \otimes \mathcal{W}_{\frac{1}{2},z_2}$ is one-dimensional and has the basis element v with the action

$$(4.4) \quad u^{(1)}v = z_1v, \quad u^{(2)}v = z_2v.$$

We have to write now all the relations in the module $\mathcal{W}_{\frac{1}{2},z_1} \otimes \mathcal{W}_{\frac{1}{2},z_2}$ using expressions (3.8) for the generators of the “small” algebra $\mathbb{T}_1^a \otimes \mathbb{T}_1^a$ in terms of elements of the “big” algebra \mathbb{T}_2^a . So, using (4.4) and (3.8) we have

$$(4.5) \quad z_1v = ug_1^{-1}v = -i(q^{-\frac{1}{2}}u - q^{\frac{1}{2}}ue_1)v,$$

$$(4.6) \quad \frac{1}{z_1}v = g_1u^{-1}v = i(q^{\frac{1}{2}}u^{-1} - q^{-\frac{1}{2}}e_1u^{-1})v,$$

$$(4.7) \quad z_2v = g_1uv = i(q^{\frac{1}{2}}u - q^{-\frac{1}{2}}e_1u)v,$$

$$(4.8) \quad \frac{1}{z_2}v = u^{-1}g_1^{-1}v = -i(q^{-\frac{1}{2}}u^{-1} - q^{\frac{1}{2}}u^{-1}e_1)v$$

and we note here that $ue_1 = u^{-1}e_1$ and $e_1u^{-1} = e_1u$ because of the relation (3.2) that takes the form $u^2e_1 = e_1$.

Therefore, we have two equations: taking the difference between first and fourth equations we get

$$(4.9) \quad (u - u^{-1})v = iq^{\frac{1}{2}}(z_1 - z_2^{-1})v$$

and second minus third gives

$$(4.10) \quad (u - u^{-1})v = iq^{-\frac{1}{2}}(z_1^{-1} - z_2)v$$

and finally the relation between z_1 and z_2 is

$$(4.11) \quad q(z_1 - z_2^{-1}) = (z_1^{-1} - z_2)$$

that has only two solutions

$$(4.12) \quad z_2 = z_1^{-1} \quad \text{or} \quad z_2 = -qz_1.$$

It tells us that the fusion or the induced module in (4.3) is zero when the condition (4.12) is not satisfied.

We then construct a basis for the fusion in the two different cases: (i) $z_2 = -qz_1$ and (ii) $z_2 = z_1^{-1}$. It turns out that in the case (i) the fusion is a one-dimensional \mathbb{T}_2^a -module

while it is two-dimensional in the case (ii). Indeed, assume that $z_2 = -\mathbf{q}z_1$ then we have the relation (using (4.9))

$$(4.13) \quad u^2 \mathbf{v} = i \mathbf{q}^{\frac{1}{2}} (z_1 + \mathbf{q}^{-1} z_1^{-1}) u \mathbf{v} + \mathbf{v}.$$

The relations (4.5)-(4.8) are not the only independent relations in $\mathcal{W}_{\frac{1}{2}, z_1} \otimes \mathcal{W}_{\frac{1}{2}, z_2}$. We have four more

$$(4.14) \quad z_1 z_2 \mathbf{v} = u^{(1)} u^{(2)} \mathbf{v} = u^2 \mathbf{v},$$

$$(4.15) \quad \frac{1}{z_1 z_2} \mathbf{v} = \left(u^{(1)} u^{(2)} \right)^{-1} \mathbf{v} = u^{-2} \mathbf{v},$$

$$(4.16) \quad \frac{z_1}{z_2} \mathbf{v} = u^{(1)} \left(u^{(2)} \right)^{-1} \mathbf{v} = (-\mathbf{q}^{-1} + e_1 + e_0 - \mathbf{q} e_0 e_1) \mathbf{v},$$

$$(4.17) \quad \frac{z_2}{z_1} \mathbf{v} = u^{(2)} \left(u^{(1)} \right)^{-1} \mathbf{v} = (-\mathbf{q} + e_1 + e_0 - \mathbf{q}^{-1} e_1 e_0) \mathbf{v}.$$

Using the first one together with (4.13) we get

$$(4.18) \quad u\mathbf{v} = i\mathbf{q}^{\frac{1}{2}}z_1\mathbf{v}$$

and using then (4.6) we get (recall that $e_0 = ue_1u^{-1}$)

$$(4.19) \quad e_1 \mathbf{v} = 0, \quad e_0 \mathbf{v} = 0.$$

Note that the relations (4.16) and (4.17) are then trivially satisfied. The equation (4.19) actually could be immediately deduced directly from the equation (4.14) that tells that u^2 acts on \mathbf{v} by $-\mathbf{q}z_1^2$ and for generic z_1 it is not possible for the \mathbb{T}_2^a -module $\mathcal{W}_{0,z}$, where u^2 acts as identity. Therefore \mathbf{v} has to belong to $\mathcal{W}_{1,z}$ where e_1 and e_0 act by zero. The value of z here is $i\mathbf{q}^{\frac{1}{2}}z_1$ as follows from (4.18). So, we conclude that because all the generators of \mathbb{T}_2^a but $u^{\pm 1}$ act on \mathbf{v} as zero, the induced module (4.3) is one-dimensional and isomorphic to $\mathcal{W}_{1,i\mathbf{q}^{\frac{1}{2}}z_1}$.

Now we turn to the case (ii) when we fix $z_2 = z_1^{-1}$. In this case, both (4.9) and (4.14) give the same relation

$$(4.20) \quad u^2 \mathbf{v} = \mathbf{v}$$

Using our basic relations (4.5)-(4.6) we obtain the relations

$$(4.21) \quad ue_1\mathbf{v} = \mathfrak{q}^{-1}uv - i\mathfrak{q}^{-\frac{1}{2}}z_1\mathbf{v},$$

$$(4.22) \quad e_1 u v = \mathfrak{q} u v + i \mathfrak{q}^{\frac{1}{2}} z_1^{-1} v$$

and so applying u^{-1} on the both sides of the equations and using (4.20) we get

$$(4.23) \quad e_1 v = \mathfrak{q}^{-1} v - i \mathfrak{q}^{-\frac{1}{2}} z_1 u v,$$

$$(4.24) \quad e_0 \mathbf{v} = \mathbf{q} \mathbf{v} + i \mathbf{q}^{\frac{1}{2}} z_1^{-1} u \mathbf{v}$$

and similar formulas for $ue_0\mathbf{v}$ and $e_0u\mathbf{v}$. Therefore, the action of e_0e_1 , e_1e_0 , etc., on \mathbf{v} is a linear combination of \mathbf{v} and $u\mathbf{v}$. Therefore the induced module (4.3) for $z_2 = z_1^{-1}$ is two-dimensional and irreducible (the irreducibility is easy to check for generic z_1) and thus isomorphic to $\mathcal{W}_{0,z}$. The basis in this module can be chosen as $\{\mathbf{v}_1 = e_1\mathbf{v}, \mathbf{v}_0 =$

$-i\mathbf{q}^{-3/2}z_1e_0\mathbf{v}\}$ which is the standard affine diagrams basis: $\mathbf{v}_1 = \boxed{\text{arc}}^{\text{dashed}}$ and $\mathbf{v}_0 = \boxed{\text{loop}}^{\text{dashed}}$. The weight of the non-contractible loops $z + z^{-1}$ is then computed as

$$(4.25) \quad e_0 \mathbf{v}_1 = (z + z^{-1}) \mathbf{v}_0 \quad \text{and} \quad e_1 \mathbf{v}_0 = (z + z^{-1}) \mathbf{v}_1$$

and a simple calculation gives $z = -i\mathfrak{q}^{-\frac{1}{2}}z_1$. Note that in this case, there is in fact an ambiguity for the sign of z , since the pair (\mathbf{v}_1, z) is defined only up to a sign: in fact, the two modules $\mathcal{W}_{0,\pm z}[2]$ are isomorphic and our choice of the sign in z is just a convention.

Finally, after simple but long calculations we conclude the fusion formula

$$(4.26) \quad \mathcal{W}_{\frac{1}{2},z_1}[1] \hat{\times}_f \mathcal{W}_{\frac{1}{2},z_2}[1] = \begin{cases} \mathcal{W}_{1,i\mathfrak{q}^{\frac{1}{2}}z_1}[2] & \text{when } z_2 = -\mathfrak{q}z_1, \\ \mathcal{W}_{0,-i\mathfrak{q}^{-\frac{1}{2}}z_1}[2] & \text{when } z_2 = z_1^{-1}, \\ 0 & \text{otherwise.} \end{cases}$$

Note that we have the same conditions on non-zero fusion $\mathcal{W}_{\frac{1}{2},z_1}[1] \hat{\times}_f \mathcal{W}_{\frac{1}{2},z_2}[3]$ on $1+3$ sites, as the basic relations have the same form as here. It just takes more calculations to find a basis in the two non-zero cases. It will be proven below (using our construction of direct systems of categories) that the formula (4.26) holds in this case as well.

4.3.2. Fusion on $1+2$ sites. By a direct calculation on $1+2$ sites similar to the previous calculation, we have found the following fusion

$$(4.27) \quad \mathcal{W}_{\frac{1}{2},z_1}[1] \hat{\times}_f \mathcal{W}_{1,z_2}[2] = \begin{cases} \mathcal{W}_{\frac{3}{2},-\mathfrak{q}z_1}[3] & \text{when } z_2 = -i\mathfrak{q}^{3/2}z_1, \\ \mathcal{W}_{\frac{1}{2},-\mathfrak{q}/z_1}[3] & \text{when } z_2 = i\mathfrak{q}^{1/2}z_1^{-1}, \\ 0 & \text{otherwise.} \end{cases}$$

A general formula for the fusion on $N_1 + N_2$ sites for any pair of the standard modules $\mathcal{W}_{j_1,z_1}[N_1]$ and $\mathcal{W}_{j_2,z_2}[N_2]$ is derived in [36], where we also discuss physical implications of the results.

An alternative to the diagrammatical calculation is presented in App. A.3 where we give a few more examples. We use there the relation to the affine Hecke algebra discussed below.

4.4. Remark on the affine Hecke algebra. The affine Temperley Lieb algebra $\mathsf{T}_N^a(m)$ is known to be deeply related to the affine Hecke algebra $\widehat{\mathsf{H}}_N(\mathfrak{q})$ with $m = \mathfrak{q} + \mathfrak{q}^{-1}$. It is useful here to recall some basic facts about the latter [44]. We then show how our definition of the affine TL tower is related to the more standard tower of affine Hecke algebras.

4.4.1. Definition I: The algebra $\widehat{\mathsf{H}}_N(\mathfrak{q})$ is usually defined as follows: it is an associative algebra over \mathbb{C} generated by σ_i , with $1 \leq i \leq N-1$, and $y_j^{\pm 1}$, with $1 \leq j \leq N$, subject to the relations²

$$(4.28) \quad \begin{aligned} (\sigma_i + 1)(\sigma_i - \mathfrak{q}^{-2}) &= 0 \\ [\sigma_i, \sigma_j] &= 0, \quad |i - j| \geq 2 \\ \sigma_i \sigma_{i+1} \sigma_i &= \sigma_{i+1} \sigma_i \sigma_{i+1} \end{aligned}$$

together with

$$(4.29) \quad \begin{aligned} [y_i, y_j] &= 0 \\ [y_j, \sigma_i] &= 0, \quad j \neq i, i+1 \\ \sigma_i y_i \sigma_i &= \mathfrak{q}^{-2} y_{i+1}. \end{aligned}$$

²We use here the standard conventions to facilitate comparison with the literature. A variant - more suitable to the T_N^a quotient is discussed in the appendix.

Note that we have thus the braid generators σ_i (subject to the standard Hecke relations) and the family of commutative generators y_j that have the commutation relations $\sigma_i y_i = \mathbf{q}^{-2} y_{i+1} \sigma_i^{-1} = y_{i+1} \sigma_i + (1 - \mathbf{q}^{-2}) y_{i+1}$. The algebra $\widehat{\mathbf{H}}_N$ can be thus considered as a twisted tensor product $\mathbf{H}_N \otimes \mathbb{C}[y_1^{\pm 1}, \dots, y_N^{\pm 1}]$, where \mathbf{H}_N is the finite Hecke algebra. It is well-known indeed that $\widehat{\mathbf{H}}_N$ is isomorphic as a vector space to the tensor product of the finite Hecke algebra and the algebra of Laurent polynomials in y_j .

4.4.2. *Zelevinsky's tensor product.* The definition of $\widehat{\mathbf{H}}_N$ given above in terms of σ_i and y_j leads us to a (well known) homomorphism of algebras

$$(4.30) \quad \widehat{\mathbf{H}}_{N_1}(\mathbf{q}) \otimes \widehat{\mathbf{H}}_{N_2}(\mathbf{q}) \mapsto \widehat{\mathbf{H}}_{N_1+N_2}(\mathbf{q})$$

where, in the same notations as in (3.7), we have

$$(4.31) \quad \sigma_j^{(1)} \mapsto \sigma_j, \quad \sigma_k^{(2)} \mapsto \sigma_{N_1+k}, \quad 1 \leq j \leq N_1 - 1, \quad 1 \leq k \leq N_2 - 1.$$

and

$$(4.32) \quad y_j^{(1)} \mapsto y_j, \quad y_k^{(2)} \mapsto y_{N_1+k}, \quad 1 \leq j \leq N_1, \quad 1 \leq k \leq N_2.$$

It is thus an embedding of the algebras.

Having the embedding in (4.30), we can now define the affine Hecke fusion $\widehat{\times}_f^{\mathbf{H}}$ as the induced module (see e.g. [44])

$$(4.33) \quad M_1 \widehat{\times}_f^{\mathbf{H}} M_2 = \widehat{\mathbf{H}}_N(\mathbf{q}) \otimes_{(\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2})} M_1 \otimes M_2,$$

where M_1 and M_2 are modules over $\widehat{\mathbf{H}}_{N_1}(\mathbf{q})$ and $\widehat{\mathbf{H}}_{N_2}(\mathbf{q})$, respectively. This fusion was originally introduced by Zelevinsky and since then is usually called *the Zelevinsky's tensor product* of affine Hecke algebra modules.

4.4.3. *Definition II:* Meanwhile, the algebra $\widehat{\mathbf{H}}_N(\mathbf{q})$ admits another definition involving generators σ_i , $i = 1, \dots, N$ (recall i was running only up to $N - 1$ in the previous definition) and a translation generator τ such that

$$(4.34) \quad \begin{aligned} (\sigma_i + 1)(\sigma_i - \mathbf{q}^{-2}) &= 0 \\ [\sigma_i, \sigma_j] &= 0, \quad |i - j| \geq 2 \\ \sigma_i \sigma_{i+1} \sigma_i &= \sigma_{i+1} \sigma_i \sigma_{i+1} \\ \tau \sigma_i \tau^{-1} &= \sigma_{i+1} \end{aligned}$$

and the indices have to be interpreted modulo N . The equivalence of the two definitions follows from the identification

$$(4.35) \quad \tau = y_1 \sigma_1 \dots \sigma_{N-1},$$

see a complete proof in [45]. Note that the relations are invariant under rescaling of τ in the second definition, and y_1 in the first definition.

Let us now see what happens to the homomorphism (4.31)-(4.32) with the second definition of $\widehat{\mathbf{H}}_N$. We have now

$$(4.36) \quad \tau^{(1)} \mapsto y_1 \sigma_1 \dots \sigma_{N_1-1}, \quad \tau^{(2)} \mapsto y_{N_1+1} \sigma_{N_1+1} \dots \sigma_{N_1+N_2-1}$$

Meanwhile, we have also for the algebra on $N_1 + N_2$ sites

$$(4.37) \quad \tau = y_1 \sigma_1 \dots \sigma_{N_1+N_2-1}$$

An easy calculation then leads to

$$(4.38) \quad \tau^{(1)} = \tau \sigma_{N_1+N_2-1}^{-1} \dots \sigma_{N_1}^{-1}, \quad \tau^{(2)} = \mathbf{q}^{2N_1} \sigma_{N_1} \dots \sigma_1 \tau$$

These formulas become identical³ with those we used earlier in (3.8) after taking the proper quotient of $\widehat{\mathbf{H}}_N(\mathbf{q})$ which we describe now.

We now recall that the affine Temperley–Lieb algebra $\mathbf{T}_N^a \equiv \mathbf{T}_N^a(\mathbf{q} + \mathbf{q}^{-1})$ can be obtained from the affine Hecke in two steps [45]: first, we demand the relations

$$(4.39) \quad E_i \equiv 1 + \sigma_i + \sigma_{i+1} + \sigma_i \sigma_{i+1} + \sigma_{i+1} \sigma_i + \sigma_i \sigma_{i+1} \sigma_i = 0, \quad i = 1, \dots, N,$$

or equivalently take a quotient of the affine Hecke algebra $\widehat{\mathbf{H}}_N$ by the two sided ideal generated by E_1 (note that all E_i 's are in the ideal). This quotient – denote it $\widehat{\mathbf{TL}}_N$ – is not in itself the affine Temperley–Lieb algebra \mathbf{T}_N^a because of the second relation in (3.2) which does not follow from (4.39). (This relation follows instead from considering the realization as an algebra of diagrams.) The second step in obtaining \mathbf{T}_N^a is thus, after identifying τ with $(i\mathbf{q}^{-1/2})^{N-1}u$, and σ_i with $i\mathbf{q}^{-1/2}g_i$, to take a quotient by the ideal generated by the element

$$(4.40) \quad \omega \equiv u^2 e_{N-1} - e_1 \dots e_{N-1}.$$

In total, we have

$$(4.41) \quad \mathbf{T}_N^a = \widehat{\mathbf{H}}_N / \mathbf{l}, \quad \text{where } \mathbf{l} = \langle E_1, \omega \rangle$$

where we introduce the two-sided ideal \mathbf{l} generated by E_1 and ω .

4.5. The affine TL fusion from Zelevinsky's tensor product. We will use the following lemma that relates induced modules over an algebra and over its quotient.

Lemma 4.5.1. *Let A be an associative algebra, I a two-sided ideal in A and $B = A/I$ the quotient algebra. Let also C be a subalgebra in A and $I_C = C \cap I$ (I_C is thus an ideal in C). Then, C/I_C is a subalgebra in B and we have an isomorphism of induced modules*

$$(4.42) \quad \text{Ind}_{C/I_C}^B M \cong \text{Ind}_C^A M / (I \cdot \text{Ind}_C^A M),$$

where M is a left C -module with trivial action of I_C , i.e., M is also a module over the quotient algebra C/I_C .

Proof. We first show that $I_C = C \cap I$ is an ideal in C : let $r \in I_C$ then on one side $a \cdot r \cdot b \in C$ for any $a, b \in C$ because $r \in C$ and on the other side $a \cdot r \cdot b$ is also in I because $r \in I$; therefore, $a \cdot r \cdot b$ is in I_C . We consider then the quotient C/I_C which is obviously a subalgebra in $B = A/I$. Therefore, the left hand side of (4.42) is well-defined. We rewrite then the right hand side of (4.42) as

$$A \otimes_C M / (I \cdot A \otimes_C M) \cong (A/I) \otimes_{C/I_C} (M/I_C \cdot M) = B \otimes_{C/I_C} M,$$

where the first isomorphism follows simply from the definition of the balanced tensor product while the second is by definition of B and our assumption on M , which has the trivial I_C action. This finishes the proof of the lemma. \square

³Note that in the affine Hecke algebra, the overall normalization of the y_i generators (in the first definition) or the τ generator (in the second definition) is not fixed, in contrast with the normalization of u in the definition of affine TL (see eq. (3.2)).

In our context, $A = \widehat{\mathbf{H}}_N$ with the ideal $I = \mathbf{l}$ defined in (4.41) and the quotient algebra $B = \mathbf{T}_N^a$, the subalgebra C is the product $\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2}$, with $N_1 + N_2 = N$, and we denote $\mathbf{l}_C = \mathbf{l} \cap C$. Using Lemma 4.5.1, we have for $\mathbf{T}_{N_i}^a$ -modules M_i an isomorphism (reading (4.42) from right to the left)

$$(4.43) \quad M_1 \widehat{\times}_f^{\mathbf{H}} M_2 / (\mathbf{l} \cdot M_1 \widehat{\times}_f^{\mathbf{H}} M_2) \cong \mathbf{T}_N^a \otimes_{\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2} / \mathbf{l}_C} \left(M_1 \otimes M_2 / (\mathbf{l}_C \cdot M_1 \otimes M_2) \right),$$

where $\widehat{\times}_f^{\mathbf{H}}$ is the affine Hecke fusion introduced in (4.33) and the quotient $\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2} / \mathbf{l}_C$ is considered as a subalgebra in \mathbf{T}_N^a . Note that in App. A we actually compute the left-hand side of (4.43) and it agrees with the affine TL fusion computed here and in [36]. Let us formulate the following conjecture:

Conjecture 4.5.2. *We have an isomorphism $\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2} / \mathbf{l}_C \cong \mathbf{T}_{N_1}^a \otimes \mathbf{T}_{N_2}^a$ of algebras, where the ideal \mathbf{l}_C is defined as the intersection of the ideal $\mathbf{l} = \langle E_1, \omega \rangle$ with the subalgebra $\widehat{\mathbf{H}}_{N_1} \otimes \widehat{\mathbf{H}}_{N_2}$ in $\widehat{\mathbf{H}}_{N_1+N_2}$.*

Under Conj. 4.5.2, we obtain the affine TL fusion $M_1 \widehat{\times}_f M_2$ on the right hand side of (4.43). We demonstrate in several examples in App. A that the fusion $\widehat{\times}_f$ for \mathbf{T}_N^a modules indeed can be computed as the quotient $M_1 \widehat{\times}_f^{\mathbf{H}} M_2 / (\mathbf{l} \cdot M_1 \widehat{\times}_f^{\mathbf{H}} M_2)$ of the fusion $\widehat{\times}_f^{\mathbf{H}}$ of the same modules but considered as $\widehat{\mathbf{H}}_N$ -modules, i.e. of the pull-back of the affine TL modules. However, while arbitrary modules of $\widehat{\mathbf{H}}_{N_1}$ and $\widehat{\mathbf{H}}_{N_2}$ can be fused to give a non trivial result, in general, the quotient turns out to be empty, except when some specific “resonance” conditions are satisfied. This is discussed more in the appendix.

5. LIMITS OF TL-MOD CATEGORIES

We now go back to the finite TL representation categories and prepare the machinery that we are going to use in the case (more interesting to us) of affine TL to show that the affine TL fusion does not depends on the choice of the pair (N_1, N_2) and is equipped with an associator. The affine case will be treated in the next section.

5.1. Arc-tower of TL algebras. Recall that in Section 2.1 we introduced a standard tower of TL algebras by using the standard embeddings. These embeddings were used to define the fusion functor for each pair of numbers (N_1, N_2) . Our task here is to connect the \mathbf{TL}_N -mod categories at different N , i.e. to construct an inductive system of \mathbf{TL}_N -mod and to show that the fusion functors induce a monoidal structure in the inductive limit category. For such a connection between different N , we use another tower of the finite Temperley–Lieb algebras given by what we call (*right*) *arc-embeddings*:

$$\mathbf{TL}_N \xrightarrow{\psi} \mathbf{TL}_{N+2}$$

defined in terms of TL diagrams by enlarging the \mathbf{TL}_N diagrams with arcs (instead of the vertical strings) at sites $(N+1, N+2)$ in the top and bottom of the diagram, or in other words

$$\psi(e_j) = \mathbf{e}_{(N)} e_j \mathbf{e}_{(N)},$$

where we introduce the idempotent

$$(5.1) \quad \mathbf{e}_{(N)} = \frac{1}{m} e_{N+1}.$$

It is straightforward to check that ψ defines a homomorphism of algebras with trivial kernel. Such a tower will be called *arc-tower*.

Note also that we have an isomorphism $\mathbf{TL}_N \cong \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$ and can thus consider \mathbf{TL}_N as an idempotent subalgebra in \mathbf{TL}_{N+2} . This allows us to define two functors between the categories of TL modules as follows. Recall that \mathbf{C}_N denotes the category of finite-dimensional \mathbf{TL}_N -modules. We have the localization functor

$$(5.2) \quad \mathcal{L}_N : \mathbf{C}_{N+2} \rightarrow \mathbf{C}_N \quad \text{such that} \quad M \mapsto \mathbf{e}_{(N)} M ,$$

with an obvious map on morphisms, and its right inverse, the so called globalization functor

$$(5.3) \quad \mathcal{G}_N : \mathbf{C}_N \rightarrow \mathbf{C}_{N+2} \quad \text{such that} \quad M \mapsto \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} M ,$$

where $\mathbf{TL}_{N+2} \mathbf{e}_{(N)}$ is considered as a left module over \mathbf{TL}_{N+2} (by the left multiplication) and a right module over the idempotent subalgebra $\mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$ (by the right multiplication), the balanced tensor product is also taken over the idempotent subalgebra $\mathbf{TL}_N = \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$. On morphisms, we have $\mathcal{G}(f) = \text{id} \otimes_{\mathbf{TL}_N} f$. It is a simple exercise to check that the composition $\mathcal{L}_N \circ \mathcal{G}_N$ is naturally isomorphic to the identity functor on \mathbf{C}_N . The reverse composition is not the identity, as the two categories are obviously not equivalent. Instead, we have the following statement.

Proposition 5.2. *The composition $\mathcal{G}_N \circ \mathcal{L}_N$ maps a \mathbf{TL}_{N+2} -module M to $I_e \cdot M$, where I_e is the two-sided ideal generated by $\mathbf{e}_{(N)}$ in \mathbf{TL}_{N+2} .*

Proof. We compute the composition $\mathcal{G}_N \circ \mathcal{L}_N$ as

$$(5.4) \quad \mathcal{G}_N \circ \mathcal{L}_N : \quad M \mapsto \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} (\mathbf{e}_{(N)} \mathbf{TL}_{N+2} \otimes_{\mathbf{TL}_{N+2}} M) ,$$

where we rewrote M as $\mathbf{TL}_{N+2} \otimes_{\mathbf{TL}_{N+2}} M$. Then, we use the associativity of the tensor product over algebras and rewrite the expression in (5.4) as

$$(5.5) \quad \mathcal{G}_N \circ \mathcal{L}_N(M) = (\mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} \mathbf{e}_{(N)} \mathbf{TL}_{N+2}) \otimes_{\mathbf{TL}_{N+2}} M ,$$

where recall that the tensor product is over $\mathbf{TL}_N = \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$. We also note the isomorphism

$$\mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \xrightarrow{\cong} \mathbf{TL}_{N+2} \cdot \mathbf{e}_{(N)} \cdot \mathbf{TL}_{N+2} = I_e$$

of the \mathbf{TL}_{N+2} bimodules (by the left and right multiplication) given by

$$(5.6) \quad a \mathbf{e}_{(N)} \otimes \mathbf{e}_{(N)} b \mapsto a \mathbf{e}_{(N)} b , \quad \text{for } a, b \in \mathbf{TL}_{N+2} .$$

The inverse to this map can be constructed as follows: any element in I_e can be presented as $a \mathbf{e}_{(N)} b$ (though not in a unique way, one can rewrite $a \mathbf{e}_{(N)} b = a' \mathbf{e}_{(N)} d$ with $a' = a \mathbf{e}_{(N)} c$ if $b = c \mathbf{e}_{(N)} d$); take any of these representatives and map $a \mathbf{e}_{(N)} b \mapsto a \mathbf{e}_{(N)} \otimes \mathbf{e}_{(N)} b$. This map is well-defined, i.e., does not depend on the representative because the tensor product is over $\mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$, and obviously inverse to the map (5.6). Finally, using (5.5) together with the isomorphism in (5.6) we obtain an isomorphism of vector spaces

$$\mathcal{G}_N \circ \mathcal{L}_N(M) \xrightarrow{\cong} I_e \otimes_{\mathbf{TL}_{N+2}} M = I_e \cdot M .$$

Note also that the \mathbf{TL}_{N+2} actions are equal on both of sides – they are simply given by the multiplication, so the modules are actually equal and not just isomorphic. \square

Remark 5.2.1. *Note that the ideal I_e generated by $\mathbf{e}_{(N)}$ in \mathbf{TL}_{N+2} is spanned by all TL diagrams except the unit $\mathbf{1}$. This is easy to see in terms of the generators of the subalgebra I_e : all e_i 's are in I_e .*

5.3. The direct limit C_∞ . We use the globalisation functors \mathcal{G}_N , for $N \geq 1$, in constructing certain direct systems of the TL representation categories C_N and eventually their direct limits. By a *direct (or inductive) system* of categories we mean a pair $\{C_i, F_{ij}\}$ of a family of categories C_i indexed by an ordered set I and a family of functors $F_{ij} : C_i \rightarrow C_j$ for all $i \leq j$ satisfying the following properties: (i) F_{ii} is the identity functor on C_i , and (ii) F_{ik} is naturally isomorphic to $F_{jk} \circ F_{ij}$ for all $i \leq j \leq k$. The *direct limit*

$$(5.7) \quad C_\infty \equiv \varinjlim C_i$$

of the direct system $\{C_i, F_{ij}\}$ is defined as the disjoint union $\coprod_i C_i / \sim$ modulo an equivalence relation: two objects $V_i \in C_i$ and $V_j \in C_j$ in the disjoint union are equivalent if and only if there is $k \in I$ such that $F_{ik}(V_i) = F_{jk}(V_j)$; and similarly for morphisms: two morphisms $f_i : V_i \rightarrow W_i$ and $f_j : V_j \rightarrow W_j$ are equivalent if the equality $F_{ik}(f_i) = F_{jk}(f_j)$ holds in $\text{Hom}(F_{ik}(V_i), F_{ik}(W_i))$. We obtain from this definition canonical functors $F_{i\infty} : C_i \rightarrow C_\infty$ mapping each object to its equivalence class. We will use these functors to define additional structures on C_∞ , such as tensor product etc.

Recall that we introduced the “enveloping” TL category

$$(5.8) \quad \mathbb{C} = \bigoplus_{N \geq 1} C_N$$

which means that \mathbb{C} contains C_N as a full subcategory and there are no morphisms between the full subcategories for different N (it is also known as the disjoint union of C_N ’s). The category \mathbb{C} is thus graded by \mathbb{N} . We will label an object M from C_N as $M[N]$ to emphasize its grade.

Inside \mathbb{C} , we consider two direct systems (recall the definition above, where $C_i = C_i$ and $F_{ij} = \mathcal{G}_{j-2} \circ \dots \circ \mathcal{G}_{i+2} \circ \mathcal{G}_i$, with $i \leq j$)

$$C_1 \xrightarrow{\mathcal{G}_1} C_3 \xrightarrow{\mathcal{G}_3} \dots$$

and

$$C_2 \xrightarrow{\mathcal{G}_2} C_4 \xrightarrow{\mathcal{G}_4} \dots$$

We denote the corresponding direct limits as

$$(5.9) \quad C_\infty^{\text{odd}} = \varinjlim C_{\text{odd}} \quad \text{and} \quad C_\infty^{\text{ev}} = \varinjlim C_{\text{even}} .$$

Then, we define the category

$$(5.10) \quad C_\infty = C_\infty^{\text{ev}} \oplus C_\infty^{\text{odd}} .$$

Note that by the construction the category C_∞ is an abelian \mathbb{C} -linear category for any non-zero value of \mathfrak{q} , i.e., including the roots of unity cases (except $\mathfrak{q} = \pm i$ where our construction is not defined).

Remark 5.3.1. We note that left arc-embeddings can be introduced similarly enlarging diagrams on the left with the pair of the arcs: introduce a new idempotent $\tilde{\mathbf{e}}_{(N)}$ in TL_{N+2} as

$$(5.11) \quad \tilde{\mathbf{e}}_{(N)} = \frac{1}{m} e_1 \in \text{TL}_{N+2}$$

and define left arc-embeddings:

$$\text{TL}_N \xrightarrow{\tilde{\psi}} \text{TL}_{N+2}$$

with

$$\tilde{\psi}(e_j) = \tilde{e}_{(N)} e_{j+2} \tilde{e}_{(N)} .$$

We have also new localization \mathcal{L}_N^l and globalization \mathcal{G}_N^l functors corresponding to the new idempotent subalgebra $\mathbf{TL}_N \cong \tilde{e}_{(N)} \mathbf{TL}_{N+2} \tilde{e}_{(N)}$: they are defined similarly to \mathcal{L}_N and \mathcal{G}_N as in (5.2) and (6.9), respectively. It is clear that \mathcal{L}_N^l and \mathcal{L}_N are naturally isomorphic, as well as their adjoint functors \mathcal{G}_N^l and \mathcal{G}_N . (The isomorphism can be explicitly stated in terms of the TL diagrams.)

The interesting property of the localization functors \mathcal{L}_N is that they map the standard (resp., costandard) module of a weight j on $N+2$ sites to the standard (resp., costandard) module of the same weight j but on N sites. We state this as follows.

Proposition 5.3.2. *For any non-zero \mathbf{q} , the \mathbf{TL}_N -module $\mathbf{e}_{(N)} \mathcal{W}_j[N+2]$ is equal to the standard module $\mathcal{W}_j[N]$ for $j \leq N/2$ and 0 otherwise.*

Proof. We recall that \mathbf{TL}_{N+2} is a quasi-hereditary algebra and the idempotent $\mathbf{e}_{(N)}$ belongs to a heredity chain of \mathbf{TL}_{N+2} , therefore $\mathbf{TL}_N = \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$ is also a quasi-hereditary algebra. The proof of the proposition is then standard (in the theory of quasi-hereditary algebras) and follows for example from [46, Sec. 3.3 and Prop. 3.4].

It is also possible to prove the result by elementary means. First, by adding an arc to the right of any link diagram in $\mathcal{W}_j[N]$, it is clear that we obtain a link diagram in $\mathbf{e}_{(N)} \mathcal{W}_j[N+2]$. Let us call this map by $\rho : \mathcal{W}_j[N] \rightarrow \mathbf{e}_{(N)} \mathcal{W}_j[N+2]$, its kernel is obviously trivial. Conversely, we start from a link diagram in $\mathcal{W}_j[N+2]$. If the last two points are both occupied by through-lines, the action of the idempotent $\mathbf{e}_{(N)}$ is zero. If none of the last two points are occupied by through-lines, the action of $\mathbf{e}_{(N)}$ simply produces, apart from the arc between the $N+1^{th}$ and $N+2^{th}$ points, a new connection between two points among the first N ones, resulting in a link diagram in $\mathcal{W}_j[N]$ after this arc is removed. If only one point is occupied by a through-line, this point is necessarily the $N+2^{th}$ one. The action of $\mathbf{e}_{(N)}$ then moves this through-line to a new position – the point to which the $N+1^{th}$ point was connected by an arc, resulting again, apart from the arc between the $N+1^{th}$ and $N+2^{th}$ points, into a link diagram with the same number of through-lines, that is a link diagram in $\mathcal{W}_j[N]$. We have thus shown by the first map ρ that $\mathcal{W}_j[N] \subset \mathbf{e}_{(N)} \mathcal{W}_j[N+2]$ and the inclusion $\mathbf{e}_{(N)} \mathcal{W}_j[N+2] \subset \mathcal{W}_j[N]$ by the second map. Hence, we have a bijection. Also, both the maps are obviously module maps and we have thus an isomorphism of the \mathbf{TL}_N modules. Moreover, the first map ρ is the identity map – adding the arc on the right is only a convention in terms of the diagrammatical bases in the two spaces – and its inverse is the identity as well. We have thus shown that \mathcal{L}_N sends the standard module $\mathcal{W}_j[N+2]$ to $\mathcal{W}_j[N]$. \square

As a consequence, we obtain the important property of the globalisation functors:

Proposition 5.3.3. *For $x \leq j \leq \frac{N}{2}$ and $x = \frac{1}{2}(N \bmod 2)$, we have*

$$(5.12) \quad \mathcal{G}_N : \quad \mathcal{W}_j[N] \mapsto \mathcal{W}_j[N+2] .$$

Proof. We first compute the composition $\mathcal{G}_N \circ \mathcal{L}_N$ on the module $\mathcal{W}_j[N+2]$ using Prop. 5.2 and obtain

$$(5.13) \quad \mathcal{G}_N \circ \mathcal{L}_N : \quad \mathcal{W}_j[N+2] \mapsto I_e \cdot \mathcal{W}_j[N+2] ,$$

which is $\mathcal{W}_j[N+2]$ for $x \leq j \leq \frac{N}{2}$ and zero otherwise, see Rem. 5.2.1. Then, recall that the composition $\mathcal{L}_N \circ \mathcal{G}_N$ is naturally isomorphic to the identity functor on \mathbf{C}_N . Together with Prop. 5.3.2, we then obtain (5.12). \square

The idea is then to study the direct limit of the standard modules (their subquotient structure) and their fusion rules in the direct limit category. So, our next step is to use the bilinear \mathbb{N} -graded tensor product \times_f on \mathbb{C} , recall Prop. 2.2.2, and to show that it defines a monoidal structure on \mathbb{C}_∞ , with braiding. We start by exploring an associator for \times_f on \mathbb{C} and prove the statement about associativity in Prop. 2.2.2.

5.4. Associativity of TL fusion. The important property of the fusion tensor-product \times_f introduced in (2.5) is the associativity.

Proposition 5.4.1. *Let M_1 , M_2 and M_3 be three modules over $\mathrm{TL}_{N_1}(m)$, $\mathrm{TL}_{N_2}(m)$ and $\mathrm{TL}_{N_3}(m)$, respectively. The tensor product \times_f is equipped with an associator, i.e., we have a family α_{M_1, M_2, M_3} of natural isomorphisms of $\mathrm{TL}_{N_1+N_2+N_3}(m)$ modules*

$$(5.14) \quad \alpha_{M_1, M_2, M_3} : (M_1 \times_f M_2) \times_f M_3 \xrightarrow{\cong} M_1 \times_f (M_2 \times_f M_3)$$

given explicitly as, for any triple of vectors $m_i \in M_i$, with $i = 1, 2, 3$,

$$(5.15) \quad \alpha_{M_1, M_2, M_3} : a \otimes (b \otimes m_1 \otimes m_2) \otimes m_3 \mapsto a \cdot b \otimes (m_1 \otimes (\mathbf{1}_{2+3} \otimes m_2 \otimes m_3)) ,$$

where $a \in \mathrm{TL}_{N_1+N_2+N_3}(m)$, $b \in \mathrm{TL}_{N_1+N_2}(m)$, and $\mathbf{1}_{2+3}$ is the identity in $\mathrm{TL}_{N_2+N_3}$, and $a \cdot b$ stands for the product of the element a with the image of b under the standard embedding of $\mathrm{TL}_{N_1+N_2}$ into $\mathrm{TL}_{N_1+N_2+N_3}$.

Proof. To prove that the map α_{M_1, M_2, M_3} in (5.14) is an isomorphism, we show that it is a composition of two isomorphisms: ℓ and the inverse to r , where ℓ is defined as the composition of the isomorphisms

$$(5.16) \quad \begin{aligned} \ell : (M_1 \times_f M_2) \times_f M_3 &= \mathrm{TL}_{1+2+3} \otimes_{\mathrm{TL}_{1+2} \otimes \mathrm{TL}_3} \left[\left(\mathrm{TL}_{1+2} \otimes_{\mathrm{TL}_1 \otimes \mathrm{TL}_2} M_1 \otimes M_2 \right) \otimes M_3 \right] \\ &\xrightarrow{\cong} \mathrm{TL}_{1+2+3} \otimes_{\mathrm{TL}_{1+2} \otimes \mathrm{TL}_3} \left[\left(\mathrm{TL}_{1+2} \otimes \mathrm{TL}_3 \otimes_{\mathrm{TL}_1 \otimes \mathrm{TL}_2 \otimes \mathrm{TL}_3} (M_1 \otimes M_2 \otimes M_3) \right) \right] \\ &\xrightarrow{\cong} \mathrm{TL}_{1+2+3} \otimes_{\mathrm{TL}_1 \otimes \mathrm{TL}_2 \otimes \mathrm{TL}_3} (M_1 \otimes M_2 \otimes M_3), \end{aligned}$$

where we use the short-hand notations $\mathrm{TL}_i \equiv \mathrm{TL}_{N_i}(m)$, $\mathrm{TL}_{i+j} \equiv \mathrm{TL}_{N_i+N_j}(m)$, etc.; the first line is by definition, in the second line we used that $M_3 \cong \mathrm{TL}_3 \otimes_{\mathrm{TL}_3} M_3$ and an obvious rearrangement of the tensor factors, to establish an isomorphism in the third line we used the associativity of the tensor product over rings — the final result is obvious then. In the third line, we consider the tensor product $M_1 \otimes M_2 \otimes M_3$ as the module over the algebra $\mathrm{TL}_1 \otimes \mathrm{TL}_2 \otimes \mathrm{TL}_3$ which is considered as a subalgebra in TL_{1+2+3} , by the embedding (3.24). Explicitly, the isomorphism ℓ is given by

$$(5.17) \quad \ell : a \otimes (b \otimes m_1 \otimes m_2) \otimes m_3 \mapsto a \cdot b \otimes (m_1 \otimes m_2 \otimes m_3) .$$

Similarly, we introduce the isomorphism r of the right-hand side of (5.14) to the third line in (5.16):

$$(5.18) \quad r : a \otimes (m_1 \otimes (c \otimes m_2 \otimes m_3)) \mapsto a \cdot c \otimes (m_1 \otimes m_2 \otimes m_3) ,$$

where $c \in \mathrm{TL}_{N_2+N_3}(m)$. We have thus the inverse to r , and the map α_{M_1, M_2, M_3} is defined as the composition $r^{-1} \circ \ell$ with the final result in (5.15).

The naturality of α_{M_1, M_2, M_3} is obvious. This finishes our proof of the proposition. \square

Note that \times_f defines the tensor product also for morphisms: for two morphisms in \mathbb{C}_N $f : M \rightarrow M'$ and $g : K \rightarrow K'$, we say $f \times_f g : M \times_f K \rightarrow M' \times_f K'$ for $\mathrm{id}_{\mathrm{TL}_N} \otimes f \otimes g$ where \otimes is in the category of vector spaces.

Proposition 5.4.2. *The family α_{M_1, M_2, M_3} of natural isomorphisms of $\mathbf{TL}_{N_1+N_2+N_3}$ -modules from (5.14) satisfies the pentagon identity*

$$(5.19) \quad \alpha_{M_1, M_2, M_3 \times_f M_4} \circ \alpha_{M_1 \times_f M_2, M_3, M_4} \\ = (\mathrm{id}_{M_1} \times_f \alpha_{M_2, M_3, M_4}) \circ \alpha_{M_1, M_2 \times_f M_3, M_4} \circ (\alpha_{M_1, M_2, M_3} \times_f \mathrm{id}_{M_4})$$

or, equivalently, the “pentagon” diagram

$$\begin{array}{ccc} & (M_1 \times_f M_2) \times_f (M_3 \times_f M_4) & \\ \alpha_{M_1 \times_f M_2, M_3, M_4} \nearrow & & \searrow \alpha_{M_1, M_2, M_3 \times_f M_4} \\ ((M_1 \times_f M_2) \times_f M_3) \times_f M_4 & & M_1 \times_f (M_2 \times_f (M_3 \times_f M_4)) \\ \alpha_{M_1, M_2, M_3} \times_f \mathrm{id}_{M_4} \searrow & & \nearrow \mathrm{id}_{M_1} \times_f \alpha_{M_2, M_3, M_4} \\ (M_1 \times_f (M_2 \times_f M_3)) \times_f M_4 & \xrightarrow{\alpha_{M_1, M_2 \times_f M_3, M_4}} & M_1 \times_f ((M_2 \times_f M_3) \times_f M_4) \end{array}$$

commutes (we use here the abbreviation \times instead of \times_f in the indices of α).

Proof. We denote an element from $((M_1 \times_f M_2) \times_f M_3) \times_f M_4$ as

$$(5.20) \quad a \otimes (b \otimes (c \otimes m_1 \otimes m_2) \otimes m_3) \otimes m_4,$$

for $m_i \in M_i$, $i = 1, 2, 3, 4$, and $a \in \mathbf{TL}_{1+2+3+4}$, $b \in \mathbf{TL}_{1+2+3}$, and $c \in \mathbf{TL}_{1+2}$. (It is clear which tensor products in (5.20) are the balanced tensor products over the \mathbf{TL} subalgebras, so we do not indicate them explicitly.) Using (5.15), we begin with calculating the left-hand side of (5.19) applied to such a vector:

$$(5.21) \quad a \otimes (b \otimes (c \otimes m_1 \otimes m_2) \otimes m_3) \otimes m_4 \\ \mapsto a \cdot b \otimes ((c \otimes m_1 \otimes m_2) \otimes (\mathbf{1}_{3+4} \otimes m_3 \otimes m_4)) \\ \mapsto a \cdot b \cdot c \otimes m_1 \otimes (\mathbf{1}_{2+3+4} \otimes m_2 \otimes (\mathbf{1}_{3+4} \otimes m_3 \otimes m_4)) .$$

On the other hand, the right side of (5.19) gives

$$(5.22) \quad a \otimes (b \otimes (c \otimes m_1 \otimes m_2) \otimes m_3) \otimes m_4 \\ \mapsto a \otimes ((b \cdot c) \otimes m_1 \otimes (\mathbf{1}_{2+3} \otimes m_2 \otimes m_3)) \otimes m_4 \\ \mapsto a \cdot b \cdot c \otimes m_1 \otimes (\mathbf{1}_{2+3+4} \otimes (\mathbf{1}_{2+3} \otimes m_2 \otimes m_3) \otimes m_4) \\ \mapsto a \cdot b \cdot c \otimes m_1 \otimes (\mathbf{1}_{2+3+4} \otimes m_2 \otimes (\mathbf{1}_{3+4} \otimes m_3 \otimes m_4)) ,$$

which equals the third line of (5.21), and so the pentagon diagram indeed commutes. \square

We have thus proven that \mathbf{C} is a semi-group category (note that we have no tensor unit because we do not have the grade zero $N = 0$ subcategory.) Our next step is the introduction of braiding isomorphisms in \mathbf{C} .

5.5. Braiding for TL fusion. Motivated by a construction in [19], we introduce the braiding in \mathbf{C} as follows. Let g_{N_1, N_2} defines the following element in $\mathbf{TL}_{N_1+N_2}$

$$(5.23) \quad g_{N_1, N_2} = (g_{N_2}^{-1} \cdots g_2^{-1} g_1^{-1}) (g_{N_2+1}^{-1} \cdots g_2^{-1}) \cdots (g_{N_2+N_1-1}^{-1} \cdots g_{N_1}^{-1})$$

which passes strings from the left over those from the right (here, $N_1 = 3$, $N_2 = 2$):

$$(5.24) \quad g_{N_1, N_2} \equiv \text{Diagram} = \text{Diagram}$$

and each braid-crossing (or equivalently $g_i^{\pm 1}$) stands for the linear combination (3.5). Note that the element g_{N_1, N_2} defines an automorphism on $\mathbf{TL}_{N_1+N_2}$ by the conjugation $a \mapsto g_{N_1, N_2} \cdot a \cdot g_{N_1, N_2}^{-1}$ which maps the subalgebra $\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2}$ (under the standard embedding) to the subalgebra $\mathbf{TL}_{N_2} \otimes \mathbf{TL}_{N_1}$ as

$$(5.25) \quad a \otimes b \mapsto g_{N_1, N_2} \cdot (a \cdot b) \cdot g_{N_1, N_2}^{-1} = b \otimes a, \quad a \in \mathbf{TL}_{N_1}, \quad b \in \mathbf{TL}_{N_2},$$

where $a \cdot b$ stands for the multiplication of a and b which are considered as the elements in $\mathbf{TL}_{N_1+N_2}$ under the standard embedding, and the last equality is most easily computed in terms of the diagrams. Of course, the conjugation on $\mathbf{TL}_{N_1+N_2}$ is non-trivial, which is easily seen for the generator e_{N_1} .

Because of this flip of the two subalgebras, the action by g_{N_1, N_2} relates the two inductions from modules over these two subalgebras. This allows us to define the family of braiding isomorphisms on the (N_1, N_2) graded components of \mathbf{C} as

$$(5.26) \quad c_{M_1, M_2} : M_1[N_1] \times_f M_2[N_2] \xrightarrow{\cong} M_2[N_2] \times_f M_1[N_1]$$

by the conjugation with g_{N_1, N_2} :

$$(5.27) \quad a \otimes m_1 \otimes m_2 \mapsto g_{N_1, N_2} \cdot a \cdot g_{N_1, N_2}^{-1} \otimes m_2 \otimes m_1,$$

where $a \in \mathbf{TL}_{N_1+N_2}$, and $m_1 \in M_1[N_1]$, $m_2 \in M_2[N_2]$. Here, we write $a \otimes m_1 \otimes m_2$ for a representative in the corresponding class in $M_1[N_1] \times_f M_2[N_2]$. We first check that the map (5.27) is well-defined, i.e., does not depend on a representative in the class. Indeed, assume that $m_1 = b \cdot m'_1$ and $m_2 = c \cdot m'_2$ for some $b \in \mathbf{TL}_{N_1}$, $c \in \mathbf{TL}_{N_2}$ and some $m'_1 \in M_1[N_1]$, $m'_2 \in M_2[N_2]$, and let us compute (5.27) for the other representative (setting here $g \equiv g_{N_1, N_2}$)

$$(5.28) \quad \begin{aligned} a \cdot (b \cdot c) \otimes m'_1 \otimes m'_2 &\mapsto g \cdot abc \cdot g^{-1} \otimes m'_2 \otimes m'_1 = g \cdot a \cdot g^{-1} \cdot (g \cdot bc \cdot g^{-1}) \otimes m'_2 \otimes m'_1 \\ &= g \cdot a \cdot g^{-1} \otimes (c \otimes b) \otimes m'_2 \otimes m'_1 = g \cdot a \cdot g^{-1} \otimes m_2 \otimes m_1, \end{aligned}$$

where we used (5.25) for the second equality, and the first and third \otimes 's in the second line are over the subalgebra $\mathbf{TL}_{N_2} \otimes \mathbf{TL}_{N_1}$. The final result in (5.28) agrees with (5.27).

We emphasize that the $\mathbf{TL}_{N_1+N_2}$ action on the left-hand side of (5.27) is given by the multiplication while an element $a' \in \mathbf{TL}_{N_1+N_2}$ acts by multiplication with $g_{N_1,N_2} \cdot a' \cdot g_{N_1,N_2}^{-1}$ on the right-hand side of (5.27), by the definition of its module structure which is given by applying the automorphism on the algebra. This shows the intertwining property of the map c_{M_1,M_2} . This map is obviously bijective. We have thus proven that c_{M_1,M_2} is an isomorphism in \mathcal{C} .

Note finally that the family c_{M_1,M_2} satisfies the coherence (hexagon) conditions required for the braiding (we use conventions from Kassel's book [47], see the hexagon conditions in eqs. (1.3)-(1.4) in Chapter 13.1.)

Proposition 5.5.1. *The family c_{M_1,M_2} of isomorphisms defined in (5.26)-(5.27) satisfies the hexagon conditions:*

$$(5.29) \quad \alpha_{M_2,M_3,M_1} \circ c_{M_1,M_2 \times_f M_3} \circ \alpha_{M_1,M_2,M_3} = (\mathrm{id}_{M_2} \times_f c_{M_1,M_3}) \circ \alpha_{M_2,M_1,M_3} \circ (c_{M_1,M_2} \times_f \mathrm{id}_{M_3})$$

on the maps from $(M_1 \times_f M_2) \times_f M_3$ to $M_2 \times_f (M_3 \times_f M_1)$ and

$$(5.30) \quad \alpha_{M_3,M_1,M_2}^{-1} \circ c_{M_1 \times_f M_2,M_3} \circ \alpha_{M_1,M_2,M_3}^{-1} = (c_{M_1,M_3} \times_f \mathrm{id}_{M_2}) \circ \alpha_{M_1,M_3,M_2}^{-1} \circ (\mathrm{id}_{M_1} \times_f c_{M_2,M_3})$$

on the maps from $M_1 \times_f (M_2 \times_f M_3)$ to $(M_3 \times_f M_1) \times_f M_2$.

Proof. To show the equalities, we first note that the isomorphism α_{M_1,M_2,M_3} from (5.15) maps a representative to another representative in the same equivalence class corresponding to $a \cdot b \otimes m_1 \otimes m_2 \otimes m_3$. Therefore, the map (5.14) on the set of the equivalence classes, which is the set $(M_1 \times_f M_2) \times_f M_3$, is actually the identity map. The equality (5.29) then follows from the identity $g_{N_1,N_2+N_3} = (\mathbf{1}_{N_2} \otimes g_{N_1,N_3}) \cdot (g_{N_1,N_2} \otimes \mathbf{1}_{N_3})$, where $\mathbf{1}_N$ is the unit on N sites, while the equality (5.30) holds because of the identity $g_{N_1+N_2,N_3} = (g_{N_1,N_3} \otimes \mathbf{1}_{N_2}) \cdot (\mathbf{1}_{N_1} \otimes g_{N_2,N_3})$. \square

We have thus proven the following theorem.

Theorem 5.6. *The category $\mathcal{C} = \oplus_{N>0} \mathcal{C}_N$ is a braided semi-group category: it has the \mathbb{N} -graded tensor product \times_f equipped with the associator $\alpha_{A,B,C}$ defined in (5.14)-(5.15) and satisfying the pentagon condition and equipped with the braiding isomorphisms $c_{A,B}$ defined in (5.26)-(5.27) that satisfy the hexagon conditions.*

5.7. Braided monoidal category \mathcal{C}_∞ . We now recall the category \mathcal{C}_∞ obtained in Sec. 5.1 as the direct limit, see (5.10) with (5.9), of a direct sequence of the categories \mathcal{C}_N inside \mathcal{C} . The direct limit \mathcal{C}_∞ is an abelian \mathbb{C} -linear category by construction (for any \mathfrak{q} which is not 0 or $\pm i$). Our first objective is to study subquotient structure of objects, e.g. the projective objects in \mathcal{C}_∞ . It turns out that there is an interesting correspondence between these projective objects and the so-called staggered representations of the Virasoro algebra (this will be discussed below in Sec. 7). Our second objective is to study different structures on \mathcal{C}_∞ , as tensor product, braiding, dual objects, etc. The idea here is to use the tensor product \times_f in \mathcal{C} , its associator α_{M_1,M_2,M_3} and the braiding c_{M_1,M_2} and to show that these structures can be lifted to the limit category \mathcal{C}_∞ .

5.7.1. Standard and projective objects in \mathcal{C}_∞ . For the first objective, we recall the definition of the standard TL modules given above (2.7) and Prop. 5.3.3 – it suggests the

following numeration of standard objects in the direct limit \mathbf{C}_∞ :

$$(5.31) \quad \mathcal{W}_j \equiv \{ \mathcal{W}_j[2j], \mathcal{W}_j[2j+2], \dots, \mathcal{W}_j[N], \dots \}, \quad j \in \frac{1}{2}\mathbb{N},$$

where the set in the figure brackets is the equivalence class of the objects in \mathbf{C} connected by the globalization functors \mathcal{G}_N for different values of N . In general, for any non-zero \mathbf{TL}_N -module $M[N]$ such that its localisation is zero, we define

$$(5.32) \quad M \equiv \{ M[N], M[N+2], \dots \}, \quad \text{with} \quad M[N+2] \equiv \mathcal{G}_N(M[N]), \quad \text{etc.},$$

as the corresponding equivalence class in the direct limit \mathbf{C}_∞ .

To study the subquotient structure of \mathcal{W}_j and their projective covers, we recall first the definition of the space of morphisms in the direct limit category, see definitions below (5.7): let $M_1, M_2 \in \mathbf{C}_\infty$ then the vector space of morphisms is the set of equivalence classes

$$(5.33) \quad \text{Hom}_{\mathbf{C}_\infty}(M_1, M_2) = \coprod_{N \in \mathbb{N}} \text{Hom}_{\mathbf{C}_N}(M_1[N], M_2[N]) / \sim,$$

where the equivalence relations are defined as $f_1 \sim f_2$ iff f_2 is the image of a sequence of the functors \mathcal{G}_N applied to f_1 . Therefore, to specify a morphism in $\text{Hom}_{\mathbf{C}_\infty}(M_1, M_2)$ it is enough to choose a non-zero representative in $\text{Hom}_{\mathbf{C}_N}(M_1[N], M_2[N])$.

Recall then the description of the abelian categories \mathbf{C}_N in Sec. 2.3. It is clear that when \mathbf{q} is not a root of unity the direct limit \mathbf{C}_∞ is a semi-simple category and isomorphism classes of simple objects are exhausted by \mathcal{W}_j from (5.31).

Let $\mathbf{q} = e^{i\pi/p}$ with integer $p \geq 3$ and recall the notation $s \equiv s(j) = (2j+1) \bmod p$. By the definition of the Hom spaces in \mathbf{C}_∞ we have a morphism from \mathcal{W}_k to \mathcal{W}_j iff $k = j$ or $k = j + p - s$ and the properties of the morphism for $k = j + p - s$ are similar to the finite N case: its kernel is the socle and its image is the socle as well. Then, the subquotient structure of the standard objects \mathcal{W}_j is the following: if $s(j) = 0$ the objects are simple while for non-zero $s(j)$ we have the subquotient structure

$$(5.34) \quad \mathcal{W}_j = \mathcal{X}_j \longrightarrow \mathcal{X}_{j+p-s}$$

where we introduce the notation \mathcal{X}_j for the irreducible quotient of \mathcal{W}_j – it is the equivalence class of the simple \mathbf{TL} modules $\mathcal{X}_j[N]$ in \mathbf{C} (note that now there are no zero conditions on \mathcal{X}_k as we had for finite N cases).

We note further that the projective cover of $\mathcal{X}_j[N+2]$ is $\mathcal{P}_j[N+2] = \mathcal{G}_N(\mathcal{P}_j[N])$. We have thus the projective objects in the direct limit \mathbf{C}_∞ :

$$(5.35) \quad \mathcal{P}_j \equiv \{ \mathcal{P}_j[2j], \mathcal{P}_j[2j+2], \dots, \mathcal{P}_j[N], \dots \}, \quad j \in \frac{1}{2}\mathbb{N},$$

as the equivalence classes of the objects in \mathbf{C} connected by the globalisation functors \mathcal{G}_N for different values of N . Following the description of the projective covers in \mathbf{C}_N around (2.9), the projective objects \mathcal{P}_j are then simple if $s(j) = 0$, they are equal to \mathcal{W}_j

for $0 \leq j \leq \frac{1}{2}(p-2)$ and otherwise have the following structure

$$(5.36) \quad \mathcal{P}_j = \begin{array}{ccc} & \mathcal{X}_j & \\ \swarrow & & \searrow \\ \mathcal{X}_{j-s} & & \mathcal{X}_{j+p-s} \\ \searrow & & \swarrow \\ & \mathcal{X}_j & \end{array}$$

That they are projective covers of the simple objects \mathcal{X}_j is easy to show by the definition of the direct limit (they are projective indecomposable and obviously cover \mathcal{X}_j). Note that now in \mathbf{C}_∞ we do not have the zero conditions as on the nodes for (2.9) and the subquotient structure for \mathcal{P}_j has always four non-zero simple subquotients (if $s(j) \neq 0$ and $2j \geq p$).

We have thus the reciprocity relation

$$(5.37) \quad [\mathcal{P}_j : \mathcal{W}_{j'}] = [\mathcal{W}_{j'} : \mathcal{X}_j]$$

in \mathbf{C}_∞ as well (as in \mathbf{C}_N) and it is a highest-weight category with the standard objects described by (5.31) and (5.34).

We will come back to the \mathcal{W}_j and their projective covers \mathcal{P}_j for \mathfrak{q} a root of unity at the last section 7 where we discuss the connection with the Virasoro algebra representation theory.

5.7.2. Tensor product in \mathbf{C}_∞ . We now turn to our second objective in this section and study the monoidal structure on \mathbf{C}_∞ . Let us define the tensor product $\otimes_{\mathbf{C}_\infty}$ on \mathbf{C}_∞ as

$$(5.38) \quad \otimes_{\mathbf{C}_\infty} : (M_1, M_2) \mapsto M_1 \otimes_{\mathbf{C}_\infty} M_2 = \lim \left[M_1[N_1] \times_f M_2[N_2] \right],$$

where $\lim[\dots]$ stands for taking the direct limit of the object (in this case, from $\mathbf{C}_{N_1+N_2}$) or the corresponding equivalence class (with respect to the functors \mathcal{G}_N). We show now that (5.38) is well-defined and does not depend on representatives in the classes M_1 and M_2 .

We begin by establishing a simple lemma.

Lemma 5.7.3. *Let T be an associative algebra over a field k with unit $\mathbf{1}$, and $I_e = T \cdot e$ is its left ideal generated by an idempotent $e \in T$. We then have*

$$I_e \cdot I_e = I_e.$$

Proof. It is obvious that $I_e \cdot I_e \subset I_e$ because I_e is the left ideal. On the other hand, we also have $I_e \subset I_e \cdot I_e$. Indeed, any element of the form $a \cdot e \in I_e$ can be rewritten in the form $(a \cdot e) \cdot (\mathbf{1} \cdot e) \in I_e \cdot I_e$. This proves the statement in the lemma. \square

And now we use this lemma to prove yet another one.

Lemma 5.7.4. *The TL modules $M_1[N_1] \times_f M_2[N_2]$ and $M_1[N_1] \times_f M_2[N_2 + 2]$ are in the same equivalence class in \mathbf{C}_∞ or, equivalently, we have*

$$(5.39) \quad \mathcal{G}_{N_1+N_2} : M_1[N_1] \times_f M_2[N_2] \mapsto M_1[N_1] \times_f M_2[N_2 + 2].$$

Proof. By definition of \times_f and the functor \mathcal{G}_N , we set $N = N_1 + N_2$ (also $\mathbf{TL}_1 \equiv \mathbf{TL}_{N_1}$, etc.) and we have

$$(5.40) \quad M_1[N_1] \times_f M_2[N_2] = \mathbf{TL}_N \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2] \\ \xrightarrow{\mathcal{G}_N} \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} (\mathbf{TL}_N \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2])$$

and then we establish the following sequence of isomorphisms of the right-hand side of (5.40)

$$(5.41) \quad \text{RHS of (5.40)} \xrightarrow{\text{associativity}} (\mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_N} \mathbf{TL}_N) \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2] \\ \xrightarrow{\mathbf{TL}_N = \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}} \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2] \\ \xrightarrow{\text{Lem. 5.7.3}} \mathbf{TL}_{N+2} \mathbf{e}_{(N)} \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2] ,$$

where the second isomorphism is due to the first balanced tensor product in the first line which is over the subalgebra $\mathbf{TL}_N = \mathbf{e}_{(N)} \mathbf{TL}_{N+2} \mathbf{e}_{(N)}$, while we used Lem. 5.7.3 for the third isomorphism. Note that \mathbf{TL}_2 in the third line stands for $\mathbf{e}_{(N)} \mathbf{TL}_{N_2+2} \mathbf{e}_{(N)}$ and $\mathbf{e}_{(N)}$ is considered as the corresponding idempotent in the subalgebra $\mathbf{TL}_{N_2+2} \subset \mathbf{TL}_{N+2}$, also the module $M_2[N_2]$ has to be considered as the corresponding module over $\mathbf{e}_{(N)} \mathbf{TL}_{N_2+2} \mathbf{e}_{(N)}$. Then, we rewrite \mathbf{TL}_{N+2} as $\mathbf{TL}_{N+2} \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2})} \mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2}$ and establish further isomorphisms

$$(5.42) \quad \text{RHS of (5.41)} \xrightarrow{\cong} \mathbf{TL}_{N+2} \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2})} \left((\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2} \mathbf{e}_{(N)}) \otimes_{(\mathbf{TL}_1 \otimes \mathbf{TL}_2)} M_1[N_1] \otimes M_2[N_2] \right) \\ \xrightarrow{\cong} \mathbf{TL}_{N+2} \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2})} M_1[N_1] \otimes (\mathbf{TL}_{N_2+2} \mathbf{e}_{(N)} \otimes_{\mathbf{TL}_2} M_2[N_2]) ,$$

where we used simple rearrangement of the tensor factors and that $\mathbf{TL}_{N_1} \otimes_{\mathbf{TL}_1} M_1[N_1] = M_1[N_1]$. The right-hand side of (5.42) obviously equals to $M_1[N_1] \times_f (\mathcal{G}_{N_2}(M_2[N_2]))$ which is $M_1[N_1] \times_f M_2[N_2 + 2]$ by the definition of $M_2[N_2 + 2]$. We have thus established an isomorphism between the right-hand sides of (5.40) and (5.42). Finally, note that the \mathbf{TL}_{N+2} actions on these two spaces are equal and not just isomorphic and the composition of the maps is the identity map.

Indeed, the isomorphism from the right-hand side of (5.40) (or $\mathcal{G}_N(M_1[N_1] \times_f M_2[N_2])$) to the right-hand side of (5.42) is given explicitly by the map

$$(5.43) \quad \varkappa : a \mathbf{e}_{(N)} \otimes b \otimes m_1 \otimes m_2 \mapsto a \cdot b \otimes m_1 \otimes (\mathbf{1}_{N_2+2} \mathbf{e}_{(N)} \otimes m_2) ,$$

where $a \in \mathbf{TL}_{N+2}$ and $b \in \mathbf{TL}_N$, and $m_i \in M_i[N_i]$, and we label the generators e_j of \mathbf{TL}_{N_2+2} from $j = N_1 + 1$ to $N + 1$, as usual for the “right” subalgebra, so $\mathbf{e}_{(N)}$ is in the subalgebra \mathbf{TL}_{N_2+2} . Both the elements (in LHS and RHS of (5.43)) are just different representatives of the same equivalence class, which is the element of the balanced tensor product in the third line of (5.41). The map \varkappa is therefore the identity map on the set of the equivalence classes. This finishes our proof. \square

Equivalently, we have the following lemma for the adjoint functors.

Lemma 5.7.5. *Let $N = N_1 + N_2$ and $M_1 \in \mathbf{C}_{N_1}$ and $M_2 \in \mathbf{C}_{N_2+2}$. The localisation functor has the following property:*

$$\mathcal{L}_{N_1+N_2} : M_1[N_1] \times_f M_2[N_2 + 2] \mapsto M_1[N_1] \times_f \mathbf{e}_{(N)} M_2[N_2 + 2] .$$

Proof. We have by definition

$$\mathcal{L}_{N_1+N_2}(M_1[N_1] \times_f M_2[N_2+2]) = \mathbf{e}_{(\mathbf{N})} \mathbf{TL}_{N+2} \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{TL}_{N_2+2})} M_1[N_1] \otimes M_2[N_2+2].$$

Note then that the latter expression is isomorphic to

$$(5.44) \quad \mathbf{e}_{(\mathbf{N})} \mathbf{TL}_{N+2} \mathbf{e}_{(\mathbf{N})} \otimes_{(\mathbf{TL}_{N_1} \otimes \mathbf{e}_{(\mathbf{N})} \mathbf{TL}_{N_2+2} \mathbf{e}_{(\mathbf{N})})} \mathbf{e}_{(\mathbf{N})} (M_1[N_1] \otimes M_2[N_2+2])$$

because we have (by Lem. 5.7.3)

$$\mathbf{e}_{(\mathbf{N})} \mathbf{TL}_N \mathbf{e}_{(\mathbf{N})} \mathbf{TL}_N \cong \mathbf{e}_{(\mathbf{N})} \mathbf{TL}_N.$$

The expression in (5.44) obviously coincides with $M_1[N_1] \times_f \mathbf{e}_{(\mathbf{N})} M_2[N_2+2]$ and thus finishes the proof. \square

Applying repeatedly Lem. 5.7.4, we see that the TL modules $M_1[N_1] \times_f M_2[N_2]$ and $M_1[N_1] \times_f M_2[N_2+2n]$, for integer n , are in the same equivalence class, as an object in \mathbf{C}_∞ . Now, we would like to vary the index N_1 . For this, recall the definition of the left-arc embeddings using the idempotents $\tilde{\mathbf{e}}_{(N)}$ and the corresponding globalisation functors \mathcal{G}_N^l introduced in Rem. 5.3.1. Repeating arguments in the proof of Lem. 5.7.4 we obtain

$$(5.45) \quad \mathcal{G}_{N_1+N_2}^l : M_1[N_1] \times_f M_2[N_2] \mapsto \mathcal{G}_{N_1}^l(M_1[N_1]) \times_f M_2[N_2].$$

Finally, note that the module $\mathcal{G}_{N_1}^l(M_1[N_1])$ is identical to $\mathcal{G}_{N_1}(M_1[N_1])$, as the \mathbf{TL}_{N_1+2} actions are equal (the use of $\tilde{\mathbf{e}}_{(N)}$ instead of $\mathbf{e}_{(\mathbf{N})}$ is a matter of convention). We thus obtain that the TL modules $M_1[N_1] \times_f M_2[N_2]$ and $M_1[N_1+2n] \times_f M_2[N_2]$, for integer n , are in the same equivalence class in \mathbf{C}_∞ . Altogether, we conclude with the following corollary.

Corollary 5.7.6. *The TL modules $M_1[N_1] \times_f M_2[N_2]$ and $M_1[N_1+2n] \times_f M_2[N_2+2m]$, for integer n and m , are in the same equivalence class in the direct limit \mathbf{C}_∞ and thus just different representatives of the same object in \mathbf{C}_∞ . If $m = -n$, the two $\mathbf{TL}_{N_1+N_2}$ modules are even identical in $\mathbf{C}_{N_1+N_2}$. Their equivalence class is by definition the fusion $M_1 \otimes_{\mathbf{C}_\infty} M_2$ of two classes M_1 and M_2 in \mathbf{C}_∞ .*

We next define the associator and the braiding for the tensor product $\otimes_{\mathbf{C}_\infty}$ in \mathbf{C}_∞ . Recall that to specify a morphism in $\text{Hom}_{\mathbf{C}_\infty}(M_1, M_2)$ it is enough to choose a representative in $\text{Hom}_{\mathbf{C}_N}(M_1[N], M_2[N])$.

Definition 5.7.7. *For a triple of objects M_1, M_2, M_3 in \mathbf{C}_∞ we define an isomorphism*

$$(5.46) \quad \alpha_{M_1, M_2, M_3}^{\mathbf{C}_\infty} : (M_1 \otimes_{\mathbf{C}_\infty} M_2) \otimes_{\mathbf{C}_\infty} M_3 \xrightarrow{\cong} M_1 \otimes_{\mathbf{C}_\infty} (M_2 \otimes_{\mathbf{C}_\infty} M_3)$$

as follows: take any triple of positive integers (N_1, N_2, N_3) such that $M_i[N_i]$ are non-zero, for $i = 1, 2, 3$, and take then the corresponding associator $\alpha_{M_1[N_1], M_2[N_2], M_3[N_3]}$ defined in Prop. 5.4.1, then $\alpha_{M_1, M_2, M_3}^{\mathbf{C}_\infty}$ is its equivalence class, i.e. the corresponding element in the quotient (5.33).

Of course, we have to show that our $\alpha^{\mathbf{C}_\infty}$ is well-defined and does not depend on the choice of representatives in the Hom spaces. For this, we first note by Cor. 5.7.6 that $(M_1[N_1] \times_f M_2[N_2]) \times_f M_3[N_3]$ is the same object in \mathbf{C}_N for any choice of (N_1, N_2, N_3) such that $N = N_1 + N_2 + N_3$, and similarly for the other bracketing. The corresponding isomorphisms $\alpha_{M_1[N_1], M_2[N_2], M_3[N_3]}$ for the different choices of N_i are also identical. We then only need to show that the associators for (N_1, N_2, N_3) and (N_1, N_2, N_3+2) are in the same equivalence class, i.e. the second is the image of the first through the functor

\mathcal{G}_N . Recall that $\mathcal{G}_N(f) = \text{id}_{\text{TL}_{N+2}} \times_f f$. To show this, we first calculate the isomorphism $\mathcal{G}_N(\alpha_{M_1[N_1], M_2[N_2], M_3[N_3]})$ on a general element from $\mathcal{G}_N((M_1[N_1] \times_f M_2[N_2]) \times_f M_3[N_3])$ as

$$(5.47) \quad \text{id} \otimes \alpha : \quad c \mathbf{e}_{(N)} \otimes (a \otimes (b \otimes m_1 \otimes m_2) \otimes m_3) \mapsto c \mathbf{e}_{(N)} \otimes a \cdot b \otimes (m_1 \otimes (\mathbf{1}_{N_2+N_3} \otimes m_2 \otimes m_3))$$

and it can be further rewritten as (by applying the identity map $(\text{id} \otimes \varkappa) \circ \varkappa$, see the definition in (5.43))

$$c \cdot a \cdot b \otimes m_1 \otimes (\mathbf{1}_{N_2+N_3+2} \otimes m_2 \otimes (\mathbf{1}_{N_3+2} \mathbf{e}_{(N)} \otimes m_3)) .$$

We then note that the last expression coincides with the image of $\alpha_{M_1[N_1], M_2[N_2], M_3[N_3+2]} \circ \varkappa$ on the left-hand side of (5.47). We have thus shown explicitly (on representatives in the balanced tensor products) the equality

$$\mathcal{G}_N(\alpha_{M_1[N_1], M_2[N_2], M_3[N_3]}) = \alpha_{M_1[N_1], M_2[N_2], M_3[N_3+2]}$$

and this finishes our proof that Def. 5.7.7 is well-defined. It is obvious by the construction that the family of isomorphisms in (5.46) satisfies the pentagon identities because each representative does, recall Prop. 5.4.2.

Definition 5.7.8. In \mathcal{C}_∞ , we similarly define the braiding by the family of isomorphisms

$$(5.48) \quad c_{M_1, M_2}^\infty : \quad M_1 \otimes_{\mathcal{C}_\infty} M_2 \xrightarrow{\cong} M_2 \otimes_{\mathcal{C}_\infty} M_1$$

as the equivalence class corresponding to c_{M_1, M_2} from $\text{Hom}_{\mathcal{C}_N}(M_1[N_1] \times_f M_2[N_2], M_2[N_2] \times_f M_1[N_1])$, which is defined in (5.26), for a choice of N_1 and N_2 such that both $M_1[N_1]$ and $M_2[N_2]$ are non-zero.

Like in the discussion below the definition of the associator in \mathcal{C}_∞ , we show that (5.48) does not depend on the choice of (N_1, N_2) , or Def. 5.7.8 is well-defined. For this, it is enough to show that the two braidings (denote them for brevity as) c_1 and c_2 , one for $M_1[N_1] \times_f M_2[N_2]$ and the second for $M_1[N_1] \times_f M_2[N_2 + 2]$, are in the same equivalence class, or $\mathcal{G}_{N_1+N_2}(c_1) = c_2$. We check the equality by a direct calculation on the representatives in the tensor product $\mathcal{G}_{N_1+N_2}(M_1[N_1] \times_f M_2[N_2])$.

Finally, the braiding isomorphisms (5.48) satisfy the hexagon identities because each representative does, recall Prop. 5.5.1. We have thus proven the following theorem.

Theorem 5.8. The direct-limit category \mathcal{C}_∞ defined in (5.10) with (5.9) is a braided monoidal category with the tensor product $\otimes_{\mathcal{C}_\infty}$ given in (5.38), with the tensor unit \mathcal{W}_0 , with the associator $\alpha_{A, B, C}^\infty$ introduced in Def. 5.7.7 and with the braiding isomorphisms $c_{A, B}^\infty$ defined in Def. 5.7.8.

We only need to comment on the tensor unit \mathcal{W}_0 – it is the equivalence class of the standard TL modules with zero number of through-lines. The tensor unit properties of \mathcal{W}_0 follows from the finite TL fusion (with even N_1 of course)

$$\mathcal{W}_0[N_1] \times_f M[N_2] \cong M[N_2] \times_f \mathcal{W}_0[N_1] \cong M[N_1 + N_2] .$$

It is easy to prove for $N_1 = 2$ that $M[N_2] \times_f \mathcal{W}_0[2]$ actually equals to $\mathcal{G}_{N_2}(M[N_2]) = M[N_2 + 2]$ and then one proceeds by the induction in N_1 using $\mathcal{W}_0[2] \times_f \mathcal{W}_0[2] = \mathcal{W}_0[4]$ and the associativity of the fusion.

5.9. On rigidity and non-rigidity of \mathbf{C}_∞ . We can also introduce (right) duals for the standard objects \mathcal{W}_j in \mathbf{C}_∞ (though not in \mathbf{C} because \mathbf{C} does not have the tensor unit) as follows. We introduce first contragredient objects: for an object $M \in \mathbf{C}_\infty$ let us pick up its representative $M[N]$, we then define the contragredient object M^* as the equivalence class corresponding to the space of linear forms $(M[N])^*$ where the \mathbf{TL}_N action is given with the help of the anti-involution reflecting a TL diagram along the horizontal line in the middle of the diagram. (Note that the definition does not depend on the choice of N .) We then introduce right dual for each standard objects \mathcal{W}_j , with $j \in \frac{1}{2}\mathbb{N}$, as the contragredient object \mathcal{W}_j^* – this duality is equipped with an evaluation map $\text{ev} : \mathcal{W}_j^* \otimes_{\mathbf{C}_\infty} \mathcal{W}_j \rightarrow \mathcal{W}_0$ and a coevaluation map $\text{coev} : \mathcal{W}_0 \rightarrow \mathcal{W}_j \otimes_{\mathbf{C}_\infty} \mathcal{W}_j^*$ (they can be explicitly fixed on representatives at finite N using the diagrammatical formulation of the fusion) satisfying the zig-zag rules. One can similarly introduce right duals for any other objects which are filtered by the standard objects, e.g. for the projective covers. For generic \mathfrak{q} , this gives thus duals for all objects in \mathbf{C}_∞ and the category is actually rigid.

The problem appears at \mathfrak{q} a root of unity: the contragredient objects \mathcal{X}_j^* do not give right duals to \mathcal{X}_j for $0 \leq j \leq \frac{1}{2}(p-2)$. Moreover for these values of j , the simple objects \mathcal{X}_j can not have duals (for our choice of the tensor unit) because the \mathcal{X}_j 's form a tensor ideal that does not contain our tensor unit \mathcal{W}_0 , e.g. for $p = 3$ we compute directly $\mathcal{X}_0^*[N] \times_f \mathcal{X}_0[N] = \mathcal{X}_0[2N]$ which is the simple (one-dimensional) quotient of $\mathcal{W}_0[2N]$. We thus conclude that \mathbf{C}_∞ is not rigid (for $\mathfrak{q} = e^{i\pi/p}$ and $p = 3, 4, 5, \dots$).

We note that the presence of such a tensor ideal that spoils the rigidity is an interesting property which is in common with representation theory of the Virasoro algebra at critical central charges, like $c = 0$, to be discussed in Sec. 7.

We are now going to use the approach elaborated in this section in the case of affine TL algebra representations. Our proofs of statements in this section were designed in such a way that the generalization to the affine case is straightforward.

6. A SEMI-GROUP AFFINE TL CATEGORY

Recall that in the finite TL case we used the idempotent subalgebras to construct direct sequences of TL representation categories and their limit \mathbf{C}_∞ . For the affine TL case, we obtain surprisingly analogous statement by using the same idempotent.

Proposition 6.1. *Let m be non zero. Introduce the idempotent $\mathbf{e}_{(N)} = \frac{1}{m}e_{N+1}$. Then, there is an isomorphism*

$$\psi : \mathbf{T}_N^a \xrightarrow{\cong} \mathbf{e}_{(N)} \mathbf{T}_{N+2}^a \mathbf{e}_{(N)},$$

such that the generators of \mathbf{T}_N^a are mapped as

$$(6.1) \quad u^{\pm 1} \mapsto m \mathbf{e}_{(N)} u^{\pm 1} \mathbf{e}_{(N)},$$

$$(6.2) \quad e_j \mapsto \mathbf{e}_{(N)} e_j, \quad 1 \leq j \leq N-1,$$

$$(6.3) \quad e_N \mapsto m \mathbf{e}_{(N)} e_N e_{N+2} \mathbf{e}_{(N)}.$$

Proof. To prove that ψ is a homomorphism of algebras is a straightforward use of the relations in \mathbb{T}_{N+2}^a . For instance, under the mapping we have

$$\begin{aligned}
 u^2 e_{N-1} &\mapsto m^{-1} e_{N+1} u e_{N+1} u e_{N+1} e_{N-1} \\
 &= m^{-1} e_{N+1} e_{N+2} u^2 e_{N+1} e_{N-1} \\
 &= m^{-1} e_1 e_2 \dots e_{N-1} e_{N+1}
 \end{aligned}
 \tag{6.4}$$

where the first relation in (3.2) (but for \mathbb{T}_{N+2}^a) was used to go from the first to the second line, and similarly the second relation in (3.2) was used to go from the second to the third line. Meanwhile, it is easy to check that

$$e_1 \dots e_{N-1} \mapsto m^{-1} e_1 e_2 \dots e_{N-1} e_{N+1} \tag{6.5}$$

as well, hence checking (3.2) in the image of \mathbb{T}_{N+2}^a .

To prove that the kernel of ψ is zero, we use the graphical representation of the images in (6.1)-(6.3), some of which are shown below for $N = 4$:

$$u = \begin{array}{c} \diagup \quad \diagup \quad \diagup \quad \diagup \end{array} \mapsto \frac{1}{m} \begin{array}{c} \diagup \quad \diagup \quad \diagup \quad \diagup \quad \text{cup} \\ \text{cap} \quad \text{cap} \quad \text{cap} \quad \text{cap} \end{array}
 \tag{6.6}$$

$$e_N = \begin{array}{c} \text{cap} \quad \text{cap} \quad \text{cap} \quad \text{cap} \\ \diagdown \quad \diagdown \quad \diagdown \quad \diagdown \end{array} \mapsto \frac{1}{m} \begin{array}{c} \text{cap} \quad \text{cap} \quad \text{cap} \quad \text{cap} \\ \diagdown \quad \diagdown \quad \diagdown \quad \diagdown \quad \text{cup} \\ \text{cap} \quad \text{cap} \quad \text{cap} \quad \text{cap} \end{array}
 \tag{6.7}$$

□

We can thus consider \mathbb{T}_N^a as the idempotent subalgebra in \mathbb{T}_{N+2}^a . This allows us similarly to Sec. 5.1 to define two functors between the categories of affine TL modules. Let $\widehat{\mathcal{C}}_N$ denotes the category of finite-dimensional \mathbb{T}_N^a -modules. We introduce the localisation functor

$$\widehat{\mathcal{L}}_N : \widehat{\mathcal{C}}_{N+2} \rightarrow \widehat{\mathcal{C}}_N \quad \text{such that} \quad M \mapsto \mathbf{e}_{(\mathbf{N})} M , \tag{6.8}$$

with an obvious map on morphisms, and its right inverse, so called globalisation functor

$$\widehat{\mathcal{G}}_N : \widehat{\mathcal{C}}_N \rightarrow \widehat{\mathcal{C}}_{N+2} \quad \text{such that} \quad M \mapsto \mathbb{T}_{N+2}^a \mathbf{e}_{(\mathbf{N})} \otimes_{\mathbb{T}_N^a} M . \tag{6.9}$$

The composition $\widehat{\mathcal{L}}_N \circ \widehat{\mathcal{G}}_N$ is naturally isomorphic to the identity functor on $\widehat{\mathcal{C}}_N$. Similarly to Prop. 5.2, for the reverse composition we establish the analogous result for \mathbb{T}_N^a (the proof is just a repetition of the proof of Prop. 5.2 replacing TL_N by \mathbb{T}_N^a , etc.)

Proposition 6.2. *The composition $\widehat{\mathcal{G}}_N \circ \widehat{\mathcal{L}}_N$ maps a \mathbb{T}_{N+2}^a -module M to $I_{\mathbf{e}} \cdot M$, where $I_{\mathbf{e}}$ is the two-sided ideal generated by $\mathbf{e}_{(\mathbf{N})}$ in \mathbb{T}_{N+2}^a .*

Remark 6.2.1. *Note that the ideal $I_{\mathbf{e}}$ generated by $\mathbf{e}_{(\mathbf{N})}$ in \mathbb{T}_{N+2}^a is spanned by all affine TL diagrams except the powers of the translation generator u^n , for $n \in \mathbb{Z}$. This is easy to see in terms of the generators of the subalgebra $I_{\mathbf{e}}$: all e_i , with $i \in \mathbb{Z}_{N+2}$, are in $I_{\mathbf{e}}$.*

We study then properties of the two functors $\widehat{\mathcal{L}}_N$ and $\widehat{\mathcal{G}}_N$ with respect to the standard modules introduced in Sec. 4.2. As in the finite TL case, the localisation functors $\widehat{\mathcal{L}}_N$ send the standard (resp., costandard) modules to the standard (resp., costandard) modules of the same weight (j, z) .

Proposition 6.3. *For any non-zero \mathbf{q} , the \mathbb{T}_N^a -module $\mathbf{e}_{(\mathbf{N})}\mathcal{W}_{j,z}[N+2]$ is equal to $\mathcal{W}_{j,z}[N]$ for $j \leq N/2$ and 0 otherwise.*

Proof. We begin with determining $\mathbf{e}_{(\mathbf{N})}\mathcal{W}_{j,z}[N+2]$ as a module over the TL subalgebra $\text{TL}_N \subset \mathbb{T}_N^a$. For this, we use results of [1, Sec. 2.12] and decompose (for a generic value of \mathbf{q})

$$\mathcal{W}_{j,z}[N+2] = \bigoplus_{k=j}^{(N+2)/2} \mathcal{W}_k[N+2]$$

as a module over the TL_{N+2} subalgebra. For \mathbf{q} a root of unity, the direct sum is replaced by the filtration by \mathcal{W}_k modules such that \mathcal{W}_j is a submodule in $\mathcal{W}_{j,z}$ (it is the span of affine TL diagrams of rank-0), \mathcal{W}_{j+1} is a submodule in the quotient $\mathcal{W}_{j,z}/\mathcal{W}_j$ (it is the span of affine TL diagrams of rank-1), etc. Note then that $\mathbf{e}_{(\mathbf{N})}$ is in the subalgebra TL_{N+2} . So, the problem reduces to the finite TL problem: we have to compute $\bigoplus_{k=j}^{(N+2)/2} \mathbf{e}_{(\mathbf{N})}\mathcal{W}_k[N+2]$ (or, equivalently, the action of $\mathbf{e}_{(\mathbf{N})}$ on each section in the filtration by \mathcal{W}_k 's if \mathbf{q} is a root of unity). For this, we use the finite TL result from Prop. 5.3.2: $\mathbf{e}_{(\mathbf{N})}\mathcal{W}_k[N+2] = \mathcal{W}_k[N]$ and obtain

$$\mathbf{e}_{(\mathbf{N})}\mathcal{W}_{j,z}[N+2] = \bigoplus_{k=j}^{N/2} \mathcal{W}_k[N]$$

as the TL_N module (or the corresponding filtration by \mathcal{W}_k 's for \mathbf{q} a root of unity). We have thus $\mathbf{e}_{(\mathbf{N})}\mathcal{W}_{j,z}[N+2] \cong \mathcal{W}_{j,z'}[N]$ for some complex number z' .

The fact that the parameter $z' = z$ can be proven by considering powers of the translation generator. We know by definition that in $\mathcal{W}_{j,z}[N+2]$ there is the relation $u^{N+2} = z^{2j}\mathbf{1}$. Consider now the image of u^N in $\mathcal{W}_{j,z}[N]$. We have

$$\begin{aligned} u^N &\mapsto m^{-1}(e_{N+1}u)^N e_{N+1} \\ &= m^{-1}u e_N (e_{N+1}u)^{N-1} e_{N+1} \\ &= m^{-1}u^2 e_{N-1} e_N (e_{N+1}u)^{N-2} e_{N+1} \\ &= \dots \\ &= m^{-1}u^N e_1 e_2 \dots e_N e_{N+1} \end{aligned} \tag{6.10}$$

where the relation $e_j u = u e_{j-1}$ in \mathbb{T}_{N+2}^a was repeatedly used. We can now replace the product of Temperley–Lieb elements on the right using the second relation in first relation in (3.2) for \mathbb{T}_{N+2}^a , leading to

$$(6.11) \quad u^N \mapsto u^{N+2} \mathbf{e}_{(N)} = z^{2j} \mathbf{e}_{(N)}$$

as required. This finishes our proof of the proposition. \square

Then, we prove the following property of the globalisation functor in the affine case, the proof repeats the one for Prop. 5.3.3 with the use of Prop. 6.3.

Proposition 6.3.1. *For $x \leq j \leq \frac{N}{2}$ and $x = \frac{1}{2}(N \bmod 2)$, we have*

$$(6.12) \quad \widehat{\mathcal{G}}_N : \quad \mathcal{W}_{j,z}[N] \mapsto \mathcal{W}_{j,z}[N+2] .$$

6.4. Associativity of the affine TL fusion. Similarly to the finite TL case, we introduce the “enveloping” affine TL category

$$(6.13) \quad \widehat{\mathcal{C}} = \bigoplus_{N \geq 1} \widehat{\mathcal{C}}_N$$

where $\widehat{\mathcal{C}}_N$ is a full subcategory and there are no morphisms between the full subcategories for different N . The category $\widehat{\mathcal{C}}$ is thus graded by \mathbb{N} . We will label an object M from $\widehat{\mathcal{C}}_N$ as $M[N]$ to emphasize its grade.

As in the previous section, our next step is to introduce a bilinear \mathbb{N} -graded tensor product (bi-functor) on $\widehat{\mathcal{C}}$ and to show that it defines an associative tensor product in the direct limit category $\widehat{\mathcal{C}}_\infty$ introduced below. We will use then Prop. 6.3 for studying the affine TL fusion rules in the direct limit.

Recall that in (4.1) we have defined the affine TL fusion bi-functor

$$(6.14) \quad \widehat{\times}_f : \quad \widehat{\mathcal{C}}_{N_1} \times \widehat{\mathcal{C}}_{N_2} \rightarrow \widehat{\mathcal{C}}_{N_1+N_2}$$

on two modules M_1 and M_2 as the induced module. It obviously respects the \mathbb{N} grading. The important property of the fusion $\widehat{\times}_f$ is the associativity.

Proposition 6.4.1. *Let M_1 , M_2 and M_3 be three modules over $\mathbb{T}_{N_1}^a(m)$, $\mathbb{T}_{N_2}^a(m)$ and $\mathbb{T}_{N_3}^a(m)$, respectively. The tensor product $\widehat{\times}_f$ is equipped with an associator, i.e., we have a family α_{M_1, M_2, M_3} of natural isomorphisms of $\mathbb{T}_{N_1+N_2+N_3}^a(m)$ modules*

$$(6.15) \quad \alpha_{M_1, M_2, M_3} : \quad (M_1 \widehat{\times}_f M_2) \widehat{\times}_f M_3 \xrightarrow{\cong} M_1 \widehat{\times}_f (M_2 \widehat{\times}_f M_3)$$

given explicitly as, for any triple of vectors $m_i \in M_i$, with $i = 1, 2, 3$,

$$(6.16) \quad \alpha_{M_1, M_2, M_3} : \quad a \otimes (b \otimes m_1 \otimes m_2) \otimes m_3 \mapsto a \cdot b \otimes (m_1 \otimes (\mathbf{1}_{2+3} \otimes m_2 \otimes m_3)) ,$$

where $a \in \mathbb{T}_{N_1+N_2+N_3}^a(m)$, $b \in \mathbb{T}_{N_1+N_2}^a(m)$, and $\mathbf{1}_{2+3}$ is the identity in $\mathbb{T}_{N_2+N_3}^a$, and $a \cdot b$ stands for the product of the element a with the image of b under the $\varepsilon_{N_1+N_2, N_3}$ embedding of $\mathbb{T}_{N_1+N_2}^a$ into $\mathbb{T}_{N_1+N_2+N_3}^a$ defined in (3.24).

The isomorphisms α_{M_1, M_2, M_3} satisfy the pentagon identity.

Proof. The proof essentially repeats the poof of Prop. 5.4.1: we replace \mathbb{TL}_N by \mathbb{T}_N^a and \times_f by $\widehat{\times}_f$, and our manipulations in (5.16) with the balanced tensor products are valid for the infinite-dimensional algebras, and define the maps ℓ and r as in (5.17) and (5.18). The map ℓ gives an isomorphism from $(M_1 \widehat{\times}_f M_2) \widehat{\times}_f M_3$ to $\mathbb{T}_{1+2+3}^a \otimes_{\mathbb{T}_1^a \otimes \mathbb{T}_2^a \otimes \mathbb{T}_3^a} (M_1 \otimes M_2 \otimes M_3)$ where $\mathbb{T}_1^a \otimes \mathbb{T}_2^a \otimes \mathbb{T}_3^a \equiv \mathbb{T}_{N_1}^a \otimes \mathbb{T}_{N_2}^a \otimes \mathbb{T}_{N_3}^a$ is considered as the subalgebra in

$\mathbb{T}_{N_1+N_2+N_3}^a$ under the composition $\varepsilon_{N_1+N_2,N_3} \circ (\varepsilon_{N_1,N_2} \otimes \text{id})$, recall the definition (3.24). Similarly the map r gives an isomorphism from the other bracketing $M_1 \hat{\times}_f (M_2 \hat{\times}_f M_3)$ to $\mathbb{T}_{1+2+3}^a \otimes_{\mathbb{T}_1^a \otimes \mathbb{T}_2^a \otimes \mathbb{T}_3^a} (M_1 \otimes M_2 \otimes M_3)$ where now $\mathbb{T}_{N_1}^a \otimes \mathbb{T}_{N_2}^a \otimes \mathbb{T}_{N_3}^a$ is considered as the subalgebra in $\mathbb{T}_{N_1+N_2+N_3}^a$ under a different composition $\varepsilon_{N_1,N_2+N_3} \circ (\text{id} \otimes \varepsilon_{N_2,N_3})$. We finally note that both the compositions have identical images (this is trivial in the finite TL case and a non-trivial but simple check for the affine algebras). Therefore we can define the composition $r^{-1} \circ \ell$ and it gives the associativity isomorphism α_{M_1,M_2,M_3} . The pentagon identity is proven along the same lines as in Prop. 5.4.2. \square

We have thus shown the following.

Proposition 6.4.2. *Let $\hat{\times}_f$ denote the \mathbb{N} -graded bilinear tensor product on $\hat{\mathbb{C}}$ as defined for each pair $(N_1, N_2) \in \mathbb{N} \times \mathbb{N}$ in (6.14) and (4.1). It is equipped with the associator α from Prop. 6.4.1 and it satisfies the pentagon identity. The category $\hat{\mathbb{C}}$ is thus an \mathbb{N} -graded semi-group category.*

6.5. The direct-limit category $\hat{\mathbb{C}}_\infty$. Inside the enveloping category $\hat{\mathbb{C}}$, we consider two direct systems (recall the definition (5.7) with $C_i = \hat{\mathbb{C}}_i$ and $F_{ij} = \hat{\mathcal{G}}_{j-2} \circ \dots \hat{\mathcal{G}}_{i+2} \circ \hat{\mathcal{G}}_i$, $i \leq j$)

$$\hat{\mathbb{C}}_1 \xrightarrow{\hat{\mathcal{G}}_1} \hat{\mathbb{C}}_3 \xrightarrow{\hat{\mathcal{G}}_3} \dots$$

and

$$\hat{\mathbb{C}}_2 \xrightarrow{\hat{\mathcal{G}}_2} \hat{\mathbb{C}}_4 \xrightarrow{\hat{\mathcal{G}}_4} \dots$$

We denote the corresponding direct limits as

$$(6.17) \quad \hat{\mathbb{C}}_\infty^{\text{odd}} = \varinjlim \hat{\mathbb{C}}_{\text{odd}} \quad \text{and} \quad \hat{\mathbb{C}}_\infty^{\text{ev}} = \varinjlim \hat{\mathbb{C}}_{\text{even}} .$$

Then, we define the category

$$(6.18) \quad \hat{\mathbb{C}}_\infty = \hat{\mathbb{C}}_\infty^{\text{ev}} \oplus \hat{\mathbb{C}}_\infty^{\text{odd}} .$$

Note that by the construction the category $\hat{\mathbb{C}}_\infty$ is an abelian \mathbb{C} -linear category for any non-zero value of \mathfrak{q} (except $\mathfrak{q} = \pm i$ where our construction is not defined).

6.6. Objects in $\hat{\mathbb{C}}_\infty$. Prop. 6.3.1 suggests the following numeration of standard objects in the direct limit $\hat{\mathbb{C}}_\infty$:

$$(6.19) \quad \mathcal{W}_{j,z} \equiv \{ \mathcal{W}_{j,z}[2j], \mathcal{W}_{j,z}[2j+2], \dots, \mathcal{W}_{j,z}[N], \dots \} , \quad j \in \frac{1}{2}\mathbb{N}, \quad z \in \mathbb{C}^\times ,$$

where the set in the figure brackets is the equivalence class of the objects in $\hat{\mathbb{C}}$ mapped by the globalisation functors \mathcal{G}_N for different values of N . In general, for any non-zero \mathbb{T}_N^a -module $M[N]$ such that its localisation is zero, we define

$$(6.20) \quad M \equiv \{ M[N], M[N+2], \dots \}, \quad \text{with} \quad M[N+2] \equiv \hat{\mathcal{G}}_N(M[N]), \quad \text{etc.},$$

as the corresponding equivalence class in the direct limit $\hat{\mathbb{C}}_\infty$.

Lemma 6.7. *Let $M_1[N_1]$ and $M_2[N_2]$ be affine TL modules. Then, the affine TL modules $M_1[N_1] \hat{\times}_f M_2[N_2]$ and $M_1[N_1] \hat{\times}_f M_2[N_2+2]$ are in the same equivalence class in $\hat{\mathbb{C}}_\infty$ or, equivalently, we have*

$$(6.21) \quad \hat{\mathcal{G}}_{N_1+N_2} : \quad M_1[N_1] \hat{\times}_f M_2[N_2] \mapsto M_1[N_1] \hat{\times}_f M_2[N_2+2] .$$

The proof of this lemma repeats the proof of the analogous Lem. 5.7.4 where our manipulations with the balanced tensor products are also valid for infinite-dimensional algebras. One has only to replace \mathbf{TL}_N by \mathbf{T}_N^a and the product $a \cdot b$ in the map (5.43) stands now for the product of a and the image of b under our affine TL embedding of \mathbf{T}_N^a into \mathbf{T}_{N+2}^a , recall Sec. 3.3.

Using Lem. 6.7 together with Prop. 6.3.1, we have an immediate application to the calculation of affine TL fusion rules:⁴

$$(6.22) \quad \mathcal{W}_{\frac{1}{2}, z_1}[N_1] \hat{\times}_f \mathcal{W}_{\frac{1}{2}, z_2}[N_2] = \delta_{z_2, -qz_1} \mathcal{W}_{1, iq^{\frac{1}{2}} z_1}[N_1 + N_2] \\ \oplus \delta_{z_2, z_1^{-1}} \mathcal{W}_{0, -iq^{-\frac{1}{2}} z_1}[N_1 + N_2] ,$$

for odd N_1 and N_2 , and

$$(6.23) \quad \mathcal{W}_{\frac{1}{2}, z_1}[N_1] \hat{\times}_f \mathcal{W}_{1, z_2}[N_2] = \delta_{z_2, -iq^{3/2} z_1} \mathcal{W}_{\frac{3}{2}, -qz_1}[N_1 + N_2] \\ \oplus \delta_{z_2, iq^{1/2} z_1^{-1}} \mathcal{W}_{\frac{1}{2}, -q/z_1}[N_1 + N_2] ,$$

for odd N_1 and even N_2 , where we also used the result of the calculation on $1 + 1$ and $1 + 2$ sites in (4.26) and (4.27), respectively. We thus see that the fusion rules are stable with the index N .

We also get the following result, similarly to the finite TL case.

Proposition 6.7.1. *The modules $M_1[N_1] \hat{\times}_f M_2[N_2]$ and $M_1[N_1 + 2n] \hat{\times}_f M_2[N_2 + 2m]$ over the corresponding affine TL algebras, for integer n and m , are in the same equivalence class in the direct limit $\hat{\mathbf{C}}_\infty$. If $m = -n$, the two $\mathbf{T}_{N_1+N_2}^a$ modules are even identical in $\hat{\mathbf{C}}_{N_1+N_2}$. Their equivalence class is by definition the fusion $M_1 \otimes_{\hat{\mathbf{C}}_\infty} M_2$ of two classes M_1 and M_2 in $\hat{\mathbf{C}}_\infty$.*

By this proposition we can introduce the definition of the tensor product in $\hat{\mathbf{C}}_\infty$.

Definition 6.8. *We define the tensor product $\otimes_{\hat{\mathbf{C}}_\infty}$ in $\hat{\mathbf{C}}_\infty$ as*

$$(6.24) \quad \otimes_{\hat{\mathbf{C}}_\infty} : (M_1, M_2) \mapsto M_1 \otimes_{\hat{\mathbf{C}}_\infty} M_2 = \lim \left[M_1[N_1] \hat{\times}_f M_2[N_2] \right] ,$$

where $\lim[\dots]$ stands for taking the direct limit of the object (in this case, from $\hat{\mathbf{C}}_{N_1+N_2}$) or the corresponding equivalence class, for any choice of (N_1, N_2) such that $M_1[N_1]$ and $M_2[N_2]$ are non-zero. (This definition does not depend on such a choice because of Prop. 6.7.1.)

By the definition of $\otimes_{\hat{\mathbf{C}}_\infty}$, note that the fusion obtained in (6.22) and (6.23) allows us to calculate or decompose the tensor products $\mathcal{W}_{\frac{1}{2}, z_1} \otimes_{\hat{\mathbf{C}}_\infty} \mathcal{W}_{\frac{1}{2}, z_2}$ and $\mathcal{W}_{\frac{1}{2}, z_1} \otimes_{\hat{\mathbf{C}}_\infty} \mathcal{W}_{1, z_2}$ in $\hat{\mathbf{C}}_\infty$: we have just to remove the square brackets in the formulas (6.22) and (6.23) replacing $\hat{\times}_f$ by $\otimes_{\hat{\mathbf{C}}_\infty}$.

We next define the associator for the tensor product $\otimes_{\hat{\mathbf{C}}_\infty}$ in $\hat{\mathbf{C}}_\infty$. Recall that to specify a morphism in $\text{Hom}_{\hat{\mathbf{C}}_\infty}(M_1, M_2)$ it is enough to choose a representative in the space $\text{Hom}_{\hat{\mathbf{C}}_N}(M_1[N], M_2[N])$.

⁴To vary the index N_1 , we introduce the globalisation functors $\hat{\mathcal{G}}_N^l$ corresponding to the idempotents $\hat{e}_{(N)}$ as in Rem. 5.3.1 and obtain $\hat{\mathcal{G}}_{N_1+N_2}^l : M_1[N_1] \hat{\times}_f M_2[N_2] \mapsto \hat{\mathcal{G}}_{N_1}^l(M_1[N_1]) \hat{\times}_f M_2[N_2]$. And then note that the $\mathbf{T}_{N_1+2}^a$ module $\hat{\mathcal{G}}_{N_1}^l(M_1[N_1])$ is identical to $\hat{\mathcal{G}}_{N_1}(M_1[N_1])$.

Definition 6.9. For a triple of objects M_1, M_2, M_3 in $\widehat{\mathcal{C}}_\infty$ we define an isomorphism

$$(6.25) \quad \alpha_{M_1, M_2, M_3}^{\widehat{\mathcal{C}}_\infty} : (M_1 \otimes_{\widehat{\mathcal{C}}_\infty} M_2) \otimes_{\widehat{\mathcal{C}}_\infty} M_3 \xrightarrow{\cong} M_1 \otimes_{\widehat{\mathcal{C}}_\infty} (M_2 \otimes_{\widehat{\mathcal{C}}_\infty} M_3)$$

as follows: take any triple of positive integers (N_1, N_2, N_3) such that $M_i[N_i]$ are non-zero, for $i = 1, 2, 3$, and take then the corresponding associator $\alpha_{M_1[N_1], M_2[N_2], M_3[N_3]}$ defined in Prop. 6.4.1, then $\alpha_{M_1, M_2, M_3}^{\widehat{\mathcal{C}}_\infty}$ is its equivalence class, i.e. the corresponding element in the quotient space as in (5.33).

The arguments showing that $\alpha^{\widehat{\mathcal{C}}_\infty}$ is well-defined and does not depend on the choice of representatives in the Hom spaces are similar to those after Def. 5.7.7. The arguments that the family $\alpha_{M_1, M_2, M_3}^{\widehat{\mathcal{C}}_\infty}$ satisfies the pentagon identity (5.19) are identical to those in the finite TL case.

For the moment, we were not able to deduce a tensor unit in the limit category $\widehat{\mathcal{C}}_\infty$ of affine TL modules, so we have obtained at least a semi-group category. We have thus proven the main theorem of this section.

Theorem 6.10. The category $\widehat{\mathcal{C}}_\infty$ defined in (6.18) with (6.17) is a semi-group category with the tensor product $\otimes_{\widehat{\mathcal{C}}_\infty}$ given in (6.24) with the associator $\alpha_{A, B, C}^{\widehat{\mathcal{C}}_\infty}$ given in (6.25) satisfying the pentagon condition.

The question on existence of the tensor unit in this category will be explored in our next paper [36].

It is now time to explore the braiding properties of the affine TL fusion.

6.11. Non-commutativity of affine TL fusion. In contrast to the finite TL fusion, the affine one $\widehat{\times}_f$ is trivial in most of the cases, recall the result in (4.26) as well as (6.22) and (6.23), except certain resonance conditions on the z -parameters. It makes the fusion $\widehat{\times}_f$ non-commutative. Indeed, we compute the fusion in (4.26) in the two orders:

$$(6.26) \quad \mathcal{W}_{\frac{1}{2}, z}[1] \widehat{\times}_f \mathcal{W}_{\frac{1}{2}, -qz}[1] = \mathcal{W}_{1, iq^{\frac{1}{2}}z}[2] \quad \text{while} \quad \mathcal{W}_{\frac{1}{2}, -qz}[1] \widehat{\times}_f \mathcal{W}_{\frac{1}{2}, z}[1] = 0$$

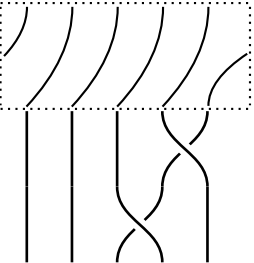
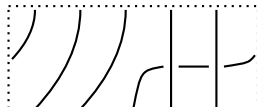
(if q is not ± 1) and

$$\mathcal{W}_{\frac{1}{2}, z}[1] \widehat{\times}_f \mathcal{W}_{\frac{1}{2}, z^{-1}}[1] = \mathcal{W}_{0, -iq^{-\frac{1}{2}}z}[2] \quad \text{while} \quad \mathcal{W}_{\frac{1}{2}, z^{-1}}[1] \widehat{\times}_f \mathcal{W}_{\frac{1}{2}, z}[1] = \mathcal{W}_{0, -iq^{-\frac{1}{2}}z^{-1}}[2] .$$

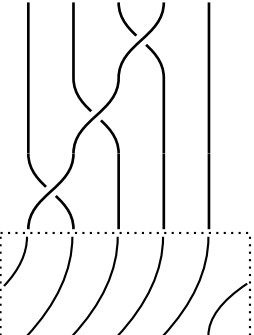
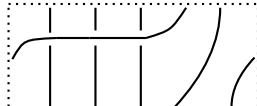
We thus conclude that in contrast to the finite TL fusion \times_f , which has the braiding, our affine TL fusion is non-commutative and there exists no braiding. However, we can introduce another affine TL fusion by replacing g_i by its inverse g_i^{-1} in the definition (3.24) of the embedding $\varepsilon_{N_1, N_2} : \mathbb{T}_{N_1}^a \otimes \mathbb{T}_{N_2}^a \rightarrow \mathbb{T}_{N_1+N_2}^a$ introduced in Sec. 3.3, which diagrammatically corresponds to the interchange between under- and above-crossings. Let us denote such embedding as ε_{N_1, N_2}^- . The definition 4.1 using the embedding ε_{N_1, N_2}^- gives then a different tensor product that we denote as $\widehat{\times}_f^-$. With this new embedding, the two translation generators correspond now, instead of (3.8), to

$$(6.27) \quad \tilde{u}^{(1)} = u g_{N-1} \dots g_{N_1}, \quad \tilde{u}^{(2)} = g_{N_1}^{-1} \dots g_1^{-1} u .$$

In terms of diagrams, we have for instance (with $N_1 = 3$ and $N_2 = 2$), instead of (6.6) and (6.7), the following

(6.28) $\tilde{u}^{(1)} \mapsto$  $=$ 

and

(6.29) $\tilde{u}^{(2)} \mapsto$  $=$ 

It is interesting that there is a braiding-type operation that relates the two affine TL fusions $\widehat{\times}_f$ and $\widehat{\times}_f^-$. Indeed, recall that in Sec. 5.5 we have introduced the braiding c_{M_1, M_2} for the TL fusion given by conjugation (5.26)-(5.27) with the “braid-like” element g_{N_1, N_2} . It is easy to see graphically – or by direct calculation using repeatedly that $ug_i = g_{i+1}u$ – that the following identities hold:

(6.30)
$$\begin{aligned} g_{N_1, N_2} u_{N_1, N_2}^{(1)} &= \tilde{u}_{N_2, N_1}^{(2)} g_{N_1, N_2} , \\ g_{N_1, N_2} u_{N_1, N_2}^{(2)} &= \tilde{u}_{N_2, N_1}^{(1)} g_{N_1, N_2} , \end{aligned}$$

where we temporarily used the notation $u_{N_j, N_k}^{(1,2)}$ and $\tilde{u}_{N_j, N_k}^{(1,2)}$ for images of the translation generators $u^{(1,2)}$ under the homomorphisms ε_{N_j, N_k} and ε_{N_j, N_k}^- , respectively.

Similarly to the finite TL case, we note that the element g_{N_1, N_2} defines an automorphism on $\mathbb{T}_{N_1+N_2}^a$ by the conjugation $a \mapsto g_{N_1, N_2} \cdot a \cdot g_{N_1, N_2}^{-1}$ which maps the subalgebra $\varepsilon_{N_1, N_2}(\mathbb{T}_{N_1}^a \otimes \mathbb{T}_{N_2}^a)$ (i.e., under the first type of the affine TL embedding) to the subalgebra $\varepsilon_{N_2, N_1}^-(\mathbb{T}_{N_2}^a \otimes \mathbb{T}_{N_1}^a)$ (i.e., under the second embedding) as

(6.31)
$$\varepsilon(a \otimes b) \mapsto g_{N_1, N_2} \cdot \varepsilon(a \otimes b) \cdot g_{N_1, N_2}^{-1} = \varepsilon^-(b \otimes a) , \quad a \in \mathbb{T}_{N_1}^a , \quad b \in \mathbb{T}_{N_2}^a .$$

Then, we can introduce a braiding-type relation between $\widehat{\times}_f$ and $\widehat{\times}_f^-$ given by the isomorphism

(6.32)
$$\widehat{c}_{M_1, M_2} : M_1[N_1] \widehat{\times}_f M_2[N_2] \xrightarrow{\cong} M_2[N_2] \widehat{\times}_f^- M_1[N_1]$$

with

(6.33)
$$\widehat{c}_{M_1, M_2} : a \otimes m_1 \otimes m_2 \mapsto g_{N_1, N_2} \cdot a \cdot g_{N_1, N_2}^{-1} \otimes m_2 \otimes m_1 ,$$

where $a \in \mathbb{T}_{N_1+N_2}^a$, and $m_1 \in M_1[N_1]$, $m_2 \in M_2[N_2]$. Recall that we write $a \otimes m_1 \otimes m_2$ here for a representative in the corresponding class in $M_1[N_1] \widehat{\times}_f M_2[N_2]$. The only non-trivial thing to check is that the map (6.33) is well-defined, i.e., does not depend on

a representative in the class. Indeed, assume that $m_1 = b \cdot m'_1$ and $m_2 = c \cdot m'_2$ for some $b \in \mathbb{T}_{N_1}^a$, $c \in \mathbb{T}_{N_2}^a$ and some $m'_i \in M_i[N_i]$, and let us compute (6.33) for the other representative (setting here $g \equiv g_{N_1, N_2}$, $\varepsilon \equiv \varepsilon_{N_1, N_2}$ and $\varepsilon^- \equiv \varepsilon_{N_2, N_1}^-$ for brevity):

$$(6.34) \quad \begin{aligned} a \cdot \varepsilon(b \otimes c) \otimes m'_1 \otimes m'_2 &\mapsto g \cdot a \cdot \varepsilon(b \otimes c) \cdot g^{-1} \otimes m'_2 \otimes m'_1 = g \cdot a \cdot g^{-1} \cdot (g \cdot \varepsilon(b \otimes c) \cdot g^{-1}) \otimes m'_2 \otimes m'_1 \\ &= g \cdot a \cdot g^{-1} \cdot \varepsilon^-(c \otimes b) \otimes m'_2 \otimes m'_1 = g \cdot a \cdot g^{-1} \otimes m_2 \otimes m_1, \end{aligned}$$

where we used (6.31) for the second equality, and note that the \otimes in front of m'_2 on the right-hand side from ‘ \mapsto ’ is over the subalgebra $\varepsilon_{N_2, N_1}^-(\mathbb{T}_{N_2}^a \otimes \mathbb{T}_{N_1}^a)$, as assumed in (6.32). The final result in (6.34) thus agrees with (6.33). Note that if we would use the same tensor product ($\widehat{\times}_f$ or $\widehat{\times}_f^-$) in (6.32)-(6.33), the map would not be well-defined.

The rest of the proof of the isomorphism property repeats the finite TL case discussed in Sec. 5.5: recall that an element $a' \in \mathbb{T}_{N_1+N_2}^a$ acts on the left-hand side of (6.33) by the multiplication with a' while on the right-hand side of (6.33) it acts by the multiplication with $g_{N_1, N_2} \cdot a' \cdot g_{N_1, N_2}^{-1}$. The intertwining property of the map \widehat{c}_{M_1, M_2} in (6.33), i.e., that it commutes with the two $\mathbb{T}_{N_1+N_2}^a$ actions, is then straightforward to check. This map is obviously bijective. We have thus proven that \widehat{c}_{M_1, M_2} is an isomorphism in the category $\widehat{\mathcal{C}}$.

We call the isomorphisms \widehat{c}_{M_1, M_2} as *semi-braiding* associated with the two tensor products $\widehat{\times}_f$ and $\widehat{\times}_f^-$. We finally note that the family of isomorphisms \widehat{c}_{M_1, M_2} defined in (6.32)-(6.33) satisfies an analogue of the coherence (hexagon) conditions (required for the ordinary braiding in a tensor category) but involving the two tensor products $\widehat{\times}_f$ and $\widehat{\times}_f^-$. More properties of the relation between $\widehat{\times}_f$ and $\widehat{\times}_f^-$ will be explored in our forthcoming paper [36].

We finally note that the semi-braiding \widehat{c}_{M_1, M_2} is lifted to the corresponding family of isomorphisms $\widehat{c}_{M_1, M_2}^\infty$ in the direct-limit category $\widehat{\mathcal{C}}_\infty$, similarly to what we have in Def. 5.7.8. This extends Thm. 6.10 by the semi-braiding structure on $\widehat{\mathcal{C}}_\infty$ with respect to the two tensor products – the *chiral* $\otimes_{\widehat{\mathcal{C}}_\infty}$ and the *anti-chiral* $\otimes_{\widehat{\mathcal{C}}_\infty}^-$.

7. OUTLOOK: A RELATION TO VIRASORO ALGEBRA

Let V_p be the Virasoro algebra of central charge

$$(7.1) \quad c(p) = 1 - \frac{6}{p(p-1)}, \quad p \in (1, \infty]$$

i.e., a Lie algebra generated by L_n , with $n \in \mathbb{Z}$, and the central element c with brackets

$$(7.2) \quad [L_n, L_m] = (n-m)L_{n+m} + \frac{c}{12}(n^3 - n)\delta_{n+m, 0}.$$

Verma modules \mathcal{V}_h are highest-weight representations of V_p and completely characterized by the central charge c and the eigenvalue h (so called “conformal weight”) of L_0 on their highest weight vector. These modules \mathcal{V}_h admit singular vectors iff their conformal weights are of the form

$$(7.3) \quad h_{r,s} = \frac{[pr - (p-1)s]^2 - 1}{4p(p-1)},$$

with $r, s \in \mathbb{Z}_+$. The first singular vector appears at level rs , that is, with L_0 eigenvalue given by $h_{r,s} + rs = h_{r,-s}$. For generic central charge c , these modules admit a unique singular vector. The corresponding Kac module is then defined as the quotient of the Verma module with $h_{r,s}$ by its submodule of the weight $h_{r,s} + rs$ (which is also a Verma module): $\mathcal{K}_{r,s} \equiv \mathcal{V}_{h_{r,s}}/\mathcal{V}_{h_{r,-s}}$. This module is irreducible when the central charge is generic. More generally, the Kac module \mathcal{K}_{rs} is defined as the submodule of the Feigin-Fuchs module [48] \mathcal{FF}_{rs} generated by the subsingular vectors of grade strictly less than rs .

We define then a certain abelian \mathbb{C} -linear category of V_p representations which is a highest-weight category with the standard objects given by the Kac modules $\mathcal{K}_{1,n}$, with $n \in \mathbb{N}$, and where indecomposable projective objects admit non-diagonalizable action of L_0 and are filtered by the Kac modules and the length of the filtration is at most two. Here, the projective covers (those which are reducible) are the so-called *staggered* modules [49, 50], which means an extension of two highest-weight modules such that the action of L_0 is non-diagonalizable. In our case, their subquotient structure has a diamond shape, for $j \bmod p \neq \frac{kp-1}{2}$ with $k = 0, 1$,

$$(7.4) \quad \mathcal{P}_{1,2j+1} : \begin{array}{ccc} & h_{1,2j+1} & \\ & \bullet & \\ \swarrow & & \searrow \\ h_{1,1+2(j-s)} & & h_{1,1+2(j+p-s)} \\ \bullet & & \bullet \\ \searrow & & \swarrow \\ & h_{1,2j+1} & \\ & \bullet & \end{array} \quad \text{for } j \geq \frac{p}{2},$$

where we set $s \equiv s(j) = (2j+1) \bmod p$, and the nodes ‘ \bullet ’ together with conformal weights $h_{1,j}$ denote irreducible V_p subquotients. We note that the modules with this subquotient structure and the requirement that the action of L_0 is non-diagonalisable on them are unique up to an isomorphism for the central charges $c(p)$ and integer $p \geq 3$, see [49]. For $j < p/2$, the projective covers $\mathcal{P}_{1,2j+1}$ are $\mathcal{K}_{1,2j+1}$. We denote such an abelian category (generated by the staggered modules $\mathcal{P}_{1,2j+1}$) as \mathbf{Vir}_p .

By a direct comparison of the projective objects in \mathbf{Vir}_p and in \mathbf{C}_∞ for $\mathbf{q} = e^{i\pi/p}$ described in Sec. 5.7.1 we establish an isomorphism between the Hom spaces in both the categories and eventually the following equivalence.

Proposition 7.1. *There exists an equivalence of abelian \mathbb{C} -linear categories \mathbf{C}_∞ for $\mathbf{q} = e^{i\pi/p}$ from Sec. 5.7.1 and \mathbf{Vir}_p such that the simple objects \mathcal{X}_j are identified with irreducible representations of conformal weight $h_{1,2j+1}$, the standard objects \mathcal{W}_j are identified with the Kac modules $\mathcal{K}_{1,1+2j}$, and the projective covers \mathcal{P}_j with the staggered modules $\mathcal{P}_{1,2j+1}$.*

This equivalence has an interesting interpretation from the physics point of view, as already mentioned in the introduction.

When $|m| \leq 2$, statistical mechanics models whose Boltzmann weights are built using representations of the Temperley–Lieb algebra are heuristically known to be critical, and have their continuum limit “described” by conformal field theories, with central charge (7.1) if we parametrize $m = \mathbf{q} + \mathbf{q}^{-1}$ with $\mathbf{q} = e^{i\pi/p}$ and $p \in (1, \infty]$. This statement can be made more precise as follows. The spectrum of eigenvalues of the Hamiltonian $H = -\sum_{i=1}^{N-1} e_i$ (which describes a statistical system with “open boundary conditions”) in modules over the TL algebra has been found to coincide, in the limit $N \rightarrow \infty$ and

after proper rescaling, with the spectrum of the generator $L_0 - \frac{c}{24}$ (note, there is a single Virasoro algebra here because we deal with open boundary conditions, corresponding to boundary conformal field theory) in some corresponding Virasoro modules for the central charge (7.1). When \mathbf{q} is not a root of unity, one can choose without loss of generality to study the TL modules \mathcal{W}_j . The corresponding Virasoro module is then found to be the Kac module with lowest weight given by the conformal weight

$$(7.5) \quad h_{1,1+2j} = \frac{(1 - 2j(p-1))^2 - 1}{4p(p-1)}.$$

Using the definition of the Kac module as the quotient $\mathcal{K}_{1,1+2j} = \mathcal{V}_{h_{1,1+2j}} / \mathcal{V}_{h_{1,-1-2j}}$, the character of this module is

$$(7.6) \quad \text{Tr}_{\mathcal{K}_{1,1+2j}} q^{L_0 - c/24} = q^{-c/24} \frac{q^{h_{1,1+2j}} - q^{h_{1,-1-2j}}}{P(q)}$$

where $P(q) = \prod_{n=1}^{\infty} (1 - q^n)$, and q here is a formal parameter. In the case when \mathbf{q} is a root of unity – that is integer p – the representation theory of both the Temperley-Lieb algebra and the Virasoro algebra become more complicated. The correspondence between these two algebras however continues to hold, after we fix an appropriate category for V_p representations of course. In particular, the projective modules \mathcal{P}_j described in Sec. 5.7.1, see the “diamond shape” diagram in (5.36), can now be put in correspondence with staggered modules $\mathcal{P}_{1,2j+1}$, which are composed of two Kac modules, i.e., have a subquotient structure with a diamond shape again (7.4).

Further, we believe that the category Vir_p has the tensor product structure (or fusion) that we denote by \otimes_{Vir} , and that it is endowed with the tensor unit $\mathcal{K}_{1,1}$ (this part is clear from the Vertex-Operator Algebra realization of V_p), with the associator and the braiding for any p from (7.1): this is based on the rigorous VOA theory of tensor products in [17]. There are then indications from physics (based on computation of the fusion rules) that our result in Prop. 7.1 extends further to the level of monoidal or tensor categories. We make the following conjecture:

Conjecture 7.2. *Let \mathcal{C}_{∞} be the direct limit of TL module categories at any \mathbf{q} such that $\mathbf{q} = e^{i\pi/p}$, which is the braided tensor category from Thm. 5.8, and Vir_p be the braided tensor category of the Virasoro algebra representations at central charge $c(p)$. We have then an equivalence of the braided tensor categories*

$$(7.7) \quad \mathcal{C}_{\infty} \xrightarrow{\sim} \text{Vir}_p$$

such that it reduces to an equivalence from Prop. 7.1 if the categories are considered just as abelian \mathbb{C} -linear categories.

The indications are the following. The category Vir_p or the representation theory of V_p within the category Vir_p is believed to serve as the fundamental description of a class of conformal field theories (CFTs). A fundamental question about CFTs is to determine the operator product expansions (OPEs) of their quantum fields. Thanks to the conformal symmetry, the OPEs are essentially determined by the fusion rules for the corresponding modules over the Virasoro algebra. Going back to our conjecture for generic \mathbf{q} , it is well known indeed that the OPEs of primary fields of conformal field theory associated with

the Kac modules obey the following Virasoro fusion rules [51, 52]

$$(7.8) \quad \mathcal{K}_{1,1+2j_1} \otimes_{\text{Vir}} \mathcal{K}_{1,1+2j_2} = \bigoplus_{j=|j_1-j_2|}^{j_1+j_2} \mathcal{K}_{1,1+2j},$$

which corresponds exactly to the result of fusion in the Temperley-Lieb case, or strictly speaking, to the decomposition for $\otimes_{\mathbb{C}_\infty}$ in our direct-limit category \mathbb{C}_∞ :

$$(7.9) \quad \mathcal{W}_{j_1} \otimes_{\mathbb{C}_\infty} \mathcal{W}_{j_2} = \bigoplus_{|j_1-j_2|}^{j_1+j_2} \mathcal{W}_j.$$

Moreover, in the case \mathfrak{q} a root of unity, fusion of the Temperley-Lieb modules [20, 21, 27] can again be compared with fusion in the corresponding (logarithmic) conformal field theory based on calculations of logarithmic OPEs in [53] and in [21, Sec. 5] (see also more references therein) and in the works [54, 55, 56, 57] that use the so-called Nahm–Gaberdiel–Kausch algorithm [58, 59]. Having the identification from Prop. 7.1 between the modules from both sides, from \mathbb{C}_∞ and Vir_p , we have as well an identification of the fusion rules (or multiplicities of the modules in tensor products of two indecomposables) for all the cases explored so far on the Virasoro side. This agreement motivates Conj. 7.2, to which we hope to get back in subsequent work.

7.3. Non-chiral case and affine TL category $\widehat{\mathbb{C}}_\infty$. The physics of critical statistical lattice models away from their boundaries (the so called “bulk” case) is described by two copies of the Virasoro algebra, corresponding to the chiral and anti-chiral dependencies of the correlation functions. While in the case of rational conformal field theories, most properties in the bulk can be inferred from those near the boundary [60], no such relationship is known to exist in general (see [61] for a discussion) in the case of logarithmic conformal field theories. Meanwhile, lattice models away from their boundaries are obtained by choosing periodic boundary conditions, which corresponds to considering now the affine instead of the finite Temperley-Lieb algebra.

In a series of works on some simple cases, we have begun to explore [22, 32, 28, 35] the relationship between modules of $\mathbb{T}_N^a(m)$ and modules of products of two Virasoro algebras with central charge (7.1). This had led to a deeper understanding of the subquotient structures appearing in logarithmic CFT for values $\mathfrak{q} = i, \mathfrak{q} = e^{i\pi/3}$, (corresponding to central charges $c = -2$ and $c = 0$), and to the introduction of the promising concept of interchiral algebra.

In order to go further in our understanding of LCFTs by using lattice models, it is necessary to understand fusion of non-chiral fields, i.e., the fusion of modules over the product of two Virasoro algebras. Since in the chiral case we observed a full correspondence between fusion in TL and Virasoro fusion rules, it is natural to expect we can learn something about fusion of non-chiral fields by studying fusion of modules of the affine Temperley-Lieb algebra.

It is not so simple to make progress in this direction however. First, we saw that doing a direct calculation of fusion in the periodic case is technically much harder than in the open because of complicated relations between different words in g_i ’s and e_j ’s. We thus need to find another and more constructive way to compute the fusion in periodic systems: this will be studied in our next paper [36]. Second, note that the fusion we have defined in the periodic case does not “reduce” to fusion in the open case once we

restrict to the finite Temperley-Lieb (sub)algebra. Indeed, from the TL_N module

$$(7.10) \quad \mathcal{W}_{j,z}[N] = \bigoplus_{k=j}^{N/2} \mathcal{W}_k[N]$$

we see the TL fusion of the right hand side of this equation using \times_f and thus (7.9) gives a direct sum over many TL modules, while the fusion of the left hand side using $\widehat{\times}_f$ gives in most cases a trivial result. Moreover, it is easy to check that even when this fusion is non trivial, the result does not decompose over TL_N according to the tensor product \times_f . This is a priori different from what one would expect in conformal field theory, where fusion of non-chiral fields is expected to decompose in some simple way in terms of the fusion of the chiral components. It would be also interesting to understand better our semi-braiding in $\widehat{\mathcal{C}}_\infty$ and its relation to non-chiral conformal field theory. We will discuss what happens in our next paper [36].

Meanwhile, we note that it would be interesting to find a duality with a quantum algebra. Recall that in the finite TL case, there is a well known duality with (a finite-dimensional quotient of) the quantum algebra $U_q \mathfrak{sl}(2)$: this duality was actually used in [20] (for $p = 2, 3$ and for projective objects) and then in [21] to compute a large and almost exhaustive list of fusion rules for a pair of indecomposable TL modules at any root of unity. One could then expect that in the periodic case there is a duality but now with (a quotient of) the affine quantum algebra $U_q \widehat{\mathfrak{sl}}(2)$. The idea would be then to use this duality and compute the affine TL fusion using the coproduct in the affine quantum group [44]. We leave this interesting problem for a future work.

APPENDIX A. AFFINE TL FUSION: EXAMPLES

In this section we consider the affine TL fusion from the perspective of the affine Hecke algebra calculations. More precisely, we compute the right-hand side of (4.43) and the results of these calculations agree with the diagrammatical calculation of the fusion in the main text of the paper in Sec. 4.3. This also supports our Conj. 4.5.2.

A.1. $j_1 = j_2 = 1/2$. We consider again the fusion of $\mathcal{W}_{\frac{1}{2},z_1}[1]$ and $\mathcal{W}_{\frac{1}{2},z_2}[1]$ discussed in Sec. 4.3.1 but now from the point of view of the affine Hecke algebra, i.e. we are going to analyze $\mathcal{W}_{\frac{1}{2},z_1}[1] \widehat{\times}_f^H \mathcal{W}_{\frac{1}{2},z_2}[1]$ where $\widehat{\times}_f^H$ is introduced in (4.33). With the normalizations adopted in this paper, we use g_i, x_i instead of σ_i, y_i and modify the relations (4.29) into

$$(A.1) \quad g_i x_i g_i = x_{i+1}$$

together with

$$(A.2) \quad u = x_1 g_1 \dots g_{N-1}$$

instead of (4.35). Recall then (4.4):

$$u^{(1)} \mathbf{v} = z_1 \mathbf{v}, \quad u^{(2)} \mathbf{v} = z_2 \mathbf{v},$$

we therefore have

$$(A.3) \quad x_1 \mathbf{v} = z_1 \mathbf{v}, \quad x_2 \mathbf{v} = z_2 \mathbf{v}.$$

We introduce $\mathbf{w} = g_1 \mathbf{v}$ and note that the two vectors \mathbf{v} and \mathbf{w} form a basis in the induced module $\mathcal{W}_{\frac{1}{2},z_1}[1] \widehat{\times}_f^H \mathcal{W}_{\frac{1}{2},z_2}[1]$. This is easy to see using the relations $x_i g_i = x_{i+1} g_i^{-1}$,

recall (A.1). We use then these affine Hecke relations to obtain the matrix representation of the generators in the space $\mathbb{C}\mathbf{v} \oplus \mathbb{C}\mathbf{w}$:

$$(A.4) \quad x_1 = \begin{pmatrix} z_1 & -i\mathbf{q}^{-1/2}z_2(1 - \mathbf{q}^2) \\ 0 & -\mathbf{q}z_2 \end{pmatrix}, \quad x_2 = \begin{pmatrix} z_2 & i\mathbf{q}^{-3/2}(\mathbf{q}^2 - 1)z_2 \\ 0 & -\mathbf{q}^{-1}z_1 \end{pmatrix}.$$

We have meanwhile

$$(A.5) \quad g_1 = \begin{pmatrix} 0 & -\mathbf{q}^{-1} \\ 1 & -i\mathbf{q}^{1/2}(\mathbf{q}^{-2} - 1) \end{pmatrix}$$

and similarly for g_2 with a bit more complicated matrix. We check then that e_1 and e_2 defined via $e_i = \mathbf{q} + i\mathbf{q}^{1/2}g_i$ do satisfy the Temperley–Lieb relations. Moreover, $u = x_1g_1$ obeys $u^2 = z_1z_2\mathbf{1}$ (recall that u^N is central). Finally, we find

$$(A.6) \quad \begin{aligned} e_1e_2e_1 &= (z + z^{-1})^2e_1 \\ e_2e_1e_2 &= (z + z^{-1})^2e_2 \end{aligned}$$

with $z = \pm i\mathbf{q}^{-1/2}\sqrt{z_1z_2^{-1}}$. For $N = 2$, the second relation in (3.2) becomes simply $u^2e_1 = e_1$ and combining with $u^2 = z_1z_2\mathbf{1}$ this gives the condition $z_1z_2 = 1$. So, the result of the fusion is therefore zero unless $z_2 = z_1^{-1}$, and thus $z = -i\mathbf{q}^{-1/2}z_1$ (with the same sign convention as discussed in the main text in Sec. 4.3.1). Hence, we have found

$$(A.7) \quad \mathcal{W}_{\frac{1}{2}, z_1}[1] \hat{\times}_f \mathcal{W}_{\frac{1}{2}, z_2}[1] = \mathcal{W}_{0, -i\mathbf{q}^{-\frac{1}{2}}z_1}[2] \quad \text{when } z_2 = z_1^{-1}.$$

The result (A.7) is obtained assuming \mathbf{w} is linearly independent of \mathbf{v} . Otherwise, one gets $g_1\mathbf{v} = \lambda\mathbf{v}$ where λ is a constant easily determined – using the relation between g_i^2 and g_i – to be $\lambda = i\mathbf{q}^{1/2}$. It follows that

$$(A.8) \quad \mathcal{W}_{\frac{1}{2}, z_1}[1] \hat{\times}_f \mathcal{W}_{\frac{1}{2}, z_2}[1] = \mathcal{W}_{1, i\mathbf{q}^{\frac{1}{2}}z_1}[2] \quad \text{when } z_2 = -\mathbf{q}z_1.$$

In other words, the induced affine Hecke module admits an invariant subspace at $z_2 = -\mathbf{q}z_1$ and the quotient by this submodule allows the action of $\mathbb{T}_2^a(\mathbf{q} + \mathbf{q}^{-1})$. This result can be interpreted as follows: for generic values of z_1 and z_2 the induced module $\mathcal{W}_{\frac{1}{2}, z_1}[1] \hat{\times}_f^H \mathcal{W}_{\frac{1}{2}, z_2}[1]$ does not admit the action of \mathbb{T}_2^a because the ideal \mathbf{l} from (4.41) generates the whole module, so the right hand side of (4.43) is zero, while for $z_2 = z_1^{-1}$ the ideal \mathbf{l} acts as zero and (4.43) gives (A.7), and for $z_2 = -\mathbf{q}z_1$ it generates a one-dimensional invariant subspace and (4.43) gives (A.8), and these are all possible cases.

Combining (A.7) with (A.8), we see that our affine Hecke calculation of the affine TL fusion (under the result in (4.43) and Conj. 4.5.2) is in agreement with the previous diagrammatical calculation in the main text resulted in (4.26).

A.2. $j_1 = 0$ and $j_2 = 1/2$. The definition of fusion $\hat{\times}_f$ holds for all modules of course, not just the standard ones. As an example, we consider here the case of $\mathcal{W}_{0, \mathbf{q}}[N]$, which is well known [1] to be reducible with a submodule isomorphic to $\mathcal{W}_{1, 1}[N]$, and admits a simple (for generic \mathbf{q}) quotient $\mathcal{W}_{0, \mathbf{q}}/\mathcal{W}_{1, 1} \equiv \overline{\mathcal{W}}_{0, \mathbf{q}}[N]$ of dimension $\hat{d}_0[N] - \hat{d}_1[N]$, see (4.2).

Restricting to the simplest case $N = 2$, the module $\overline{\mathcal{W}}_{0, \mathbf{q}}[2]$ is one-dimensional and has the basis vector \mathbf{u} with the action of the affine Hecke generators: $g_1\mathbf{u} = -i\mathbf{q}^{-3/2}\mathbf{u}$, and $u^{(1)}\mathbf{u} = \mathbf{u}$, and $x_1\mathbf{u} = -i\mathbf{q}^{-3/2}\mathbf{u}$. We consider then the fusion $\hat{\times}_f^H$ with a one dimensional module $\mathcal{W}_{\frac{1}{2}, z}[1]$. Within the induced module, the vectors $g_2\mathbf{u} \equiv \mathbf{v}$ and $g_1\mathbf{v} \equiv \mathbf{w}$ are

linearly independent and form a basis with \mathbf{u} . We then find in the $\mathbf{u}, \mathbf{v}, \mathbf{w}$ basis the matrix representation:

$$(A.9) \quad g_1 = \begin{pmatrix} i\mathbf{q}^{-3/2} & 0 & 0 \\ 0 & 0 & -\mathbf{q}^{-1} \\ 0 & 1 & i\mathbf{q}^{1/2}(1 - \mathbf{q}^{-2}) \end{pmatrix}, \quad g_2 = \begin{pmatrix} 0 & -\mathbf{q}^{-1} & 0 \\ 1 & i\mathbf{q}^{1/2}(1 - \mathbf{q}^{-2}) & 0 \\ 0 & 0 & -i\mathbf{q}^{-3/2} \end{pmatrix}.$$

It is easy to check however that the generators $e_i \equiv \mathbf{q} + i\mathbf{q}^{1/2}g_i$ now do not satisfy the required relation $e_i e_{i+1} e_i = e_i$ for $i = 1, 2$. It is necessary to take a quotient implying the linear relation $\mathbf{u} = -i\mathbf{q}^{-1/2}\mathbf{v} + \mathbf{q}^{-1}\mathbf{w}$. In the \mathbf{v}, \mathbf{w} basis for instance, one finds now

$$(A.10) \quad g_1 = \begin{pmatrix} 0 & -\mathbf{q}^{-1} \\ 1 & i\mathbf{q}^{1/2}(1 - \mathbf{q}^{-2}) \end{pmatrix}$$

and

$$(A.11) \quad g_2 = \begin{pmatrix} i\mathbf{q}^{1/2} & 0 \\ -\mathbf{q}^{-2} & -i\mathbf{q}^{-3/2} \end{pmatrix}$$

together with

$$(A.12) \quad x_1 = \begin{pmatrix} i\mathbf{q}^{3/2} & 0 \\ 0 & z \end{pmatrix}, \quad x_2 = \begin{pmatrix} -z\mathbf{q}^{-1} & iz\mathbf{q}^{-5/2}(\mathbf{q}^2 - 1) \\ iz\mathbf{q}^{-3/2}(\mathbf{q}^2 - 1) & -i\mathbf{q}^{1/2} - z\mathbf{q}^{-3}(\mathbf{q}^2 - 1)^2 \end{pmatrix}$$

and a slightly more complicated expression for x_3 . The point however is that, in this restricted space, $[x_i, x_j] = 0$ if and only if $z = i\mathbf{q}^{3/2}$. We find then

$$(A.13) \quad u = \begin{pmatrix} i\mathbf{q}^{-3/2} & -\mathbf{q}^{-1} \\ 1 - \mathbf{q}^2 - \mathbf{q}^{-2} & i(\mathbf{q}^{1/2} - \mathbf{q}^{-3/2}) \end{pmatrix}$$

and check that, in the quotient, we have

$$(A.14) \quad u^2 e_2 = e_1 e_2$$

while $u^3 = i\mathbf{q}^{3/2}\mathbf{1}$. In other words, the ideal \mathbf{l} from (4.41) does not generate the whole module but only a proper invariant subspace only at $z = i\mathbf{q}^{3/2}$. The resulting quotient-module being two dimensional is a quotient of the standard (three dimensional) \mathbf{T}_3^a -module $\mathcal{W}_{\frac{1}{2}, z}[3]$ when $z = i\mathbf{q}^{3/2}$: it is known that $\mathcal{W}_{\frac{1}{2}, i\mathbf{q}^{3/2}}$ is reducible and admits a two dimensional irreducible quotient-module $\overline{\mathcal{W}}_{1/2, i\mathbf{q}^{3/2}}[3]$, and it is the only two-dimensional irreducible module for \mathbf{T}_3^a at generic \mathbf{q} . Therefore, we obtain the affine TL fusion

$$(A.15) \quad \overline{\mathcal{W}}_{0, \mathbf{q}}[2] \widehat{\times}_f \mathcal{W}_{\frac{1}{2}, z}[1] = \begin{cases} \overline{\mathcal{W}}_{1/2, z}[3] & z = i\mathbf{q}^{3/2}, \\ 0 & \text{otherwise} \end{cases}.$$

again using the result in (4.43) and Conj. 4.5.2. We finally note that the result (A.15) will be also confirmed in our next paper [36].

A.3. $j_1 = 1/2$ and $j_2 = 1$. Finally we consider the fusion $\mathcal{W}_{\frac{1}{2}, z_1}[1] \widehat{\times}_f \mathcal{W}_{1, z_2}[2]$ with the first module on one site the second on two sites. Starting with the only basis element \mathbf{u} in the ordinary tensor product $\mathcal{W}_{\frac{1}{2}, z_1}[1] \otimes \mathcal{W}_{1, z_2}[2]$ we again generate the two more basis elements $g_1\mathbf{u} = \mathbf{v}$ and $g_2\mathbf{v} = \mathbf{w}$ in the induced module. Using relations in the module \mathcal{W}_{1, z_2} , we have $u^{(2)}\mathbf{u} = z_2\mathbf{u}$ and on the other hand $u^{(2)}\mathbf{u} = i\mathbf{q}^{1/2}x_2\mathbf{u}$ (because $u^{(2)} = x_2g_2$)

and therefore $x_2 \mathbf{u} = -i\mathbf{q}^{-1/2} z_2 \mathbf{u}$, while $x_1 \mathbf{u} = z_1 \mathbf{u}$. The defining relations lead, in the $\mathbf{u}, \mathbf{v}, \mathbf{w}$ basis, to

$$(A.16) \quad x_1 = \begin{pmatrix} z_1 & z_2(\mathbf{q} - \mathbf{q}^{-1}) & i\mathbf{q}^{1/2} z_2(\mathbf{q} - \mathbf{q}^{-1}) \\ 0 & i\mathbf{q}^{1/2} z_2 & 0 \\ 0 & 0 & i\mathbf{q}^{1/2} z_2 \end{pmatrix}$$

and similar matrix expressions for the other generators. The generators g_1 and g_2 give rise to generators e_1, e_2 that satisfy the Temperley–Lieb relations:

$$(A.17) \quad e_1 = \begin{pmatrix} \mathbf{q} & -i\mathbf{q}^{-1/2} & 0 \\ i\mathbf{q}^{1/2} & \mathbf{q}^{-1} & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad e_2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \mathbf{q} & -i\mathbf{q}^{-1/2} \\ 0 & i\mathbf{q}^{1/2} & \mathbf{q}^{-1} \end{pmatrix}$$

One also finds $u^3 = z_1 z_2^2 \mathbf{1}$. However, the T_3^a relation $u^2 e_2 = e_1 e_2$ implies the constraint $z_1 z_2 = i\mathbf{q}^{1/2}$, and so $u^3 = i\mathbf{q}^{1/2} z_2 \mathbf{1}$. This leads to the three dimensional module $\mathcal{W}_{\frac{1}{2}, i\mathbf{q}^{1/2} z_2}[3] = \mathcal{W}_{\frac{1}{2}, -\mathbf{q} z_1^{-1}}[3]$, or in other words, we have shown that the ideal \mathbf{l} from (4.41) acts by zero only at $z_1 = i\mathbf{q}^{1/2} z_2^{-1}$.

We can also consider the case where the ideal \mathbf{l} generates an invariant subspace. It happens indeed when \mathbf{v} is proportional to \mathbf{u} , which implies that \mathbf{w} is proportional to \mathbf{u} as well. After taking the corresponding quotient, one finds then $e_1 = e_2 = 0$ and $u = -\mathbf{q} z_1 \mathbf{1}$, while $z_2 = -i\mathbf{q}^{3/2} z_1$. It follows that the quotient is the (one-dimensional) standard module $\mathcal{W}_{3/2, -\mathbf{q} z_1} = \mathcal{W}_{3/2, -i\mathbf{q}^{-1/2} z_2}$.

We thus conclude that the affine TL fusion is zero at all values of z_1 and z_2 except the following cases:

$$(A.18) \quad \mathcal{W}_{\frac{1}{2}, z_1}[1] \hat{\times}_f \mathcal{W}_{1, z_2}[2] = \mathcal{W}_{\frac{1}{2}, i\mathbf{q}^{1/2} z_2}[3], \quad z_1 = i\mathbf{q}^{1/2} z_2^{-1},$$

$$(A.19) \quad \mathcal{W}_{\frac{1}{2}, z_1}[1] \hat{\times}_f \mathcal{W}_{1, z_2}[2] = \mathcal{W}_{3/2, -i\mathbf{q}^{-1/2} z_2}[3], \quad z_1 = i\mathbf{q}^{-3/2} z_2,$$

which is in agreement with (4.27).

REFERENCES

- [1] J.J. Graham and G.I. Lehrer, *The representation theory of affine Temperley-Lieb algebras*, L'Ens. Math. 44, 173 (1998).
- [2] P.P. Martin and H. Saleur, *On an algebraic approach to non planar statistical mechanics*, Comm. Math. Phys. 158, 155 (1993).
- [3] P. P. Martin, *Potts models and related problems in statistical mechanics*, World Scientific, Singapore (1991)
- [4] C. Nayak, S. Simon, A. Stern, M. Freedman and S. Das Sarma, *Non abelian anyons and topological quantum computation*, Rev. Mod. Phys. 80, 1083 (2008).
- [5] M.H. Freedman, M.J. Larsen and Z. Wang, *A modular functor which is universal for quantum computation*, Comm. Math. Phys. 228 (2002) 177.
- [6] Z. Wang, *Quantum computation*, available on <http://www.math.ucsb.edu/~zhenghwa/>
- [7] R. J. Baxter, *Exactly solved models in statistical mechanics*, Academic Press (1982).
- [8] M. Schottenloher, *A mathematical introduction to conformal field theory*, Springer (1997)
- [9] P. di Francesco, P. Mathieu and D. Senechal, *Conformal Field Theory*, Springer.
- [10] K. Gawedzky, *Conformal field theory*, Seminaire N. Bourbaki 704 (1988) 95.
- [11] G. F. Lawler, O. Schramm, and W. Werner, *On the scaling limit of planar self-avoiding walk*. In Fractal geometry and applications: a jubilee of Benoit Mandelbrot, Part 2, volume 72 of Proc. Sympos. Pure Math., pages 339-364. Amer. Math. Soc., Providence, RI, 2004.
- [12] Stanislav Smirnov, *Towards conformal invariance of 2D lattice models*. In International Congress of Mathematicians. Vol. II, pages 1421-1451. Eur. Math. Soc., Zurich, 2006.

- [13] D. Chelkak, A. Glazman and S. Smirnov, *Discrete stress-energy tensor in the loop $O(n)$ model*, arXiv:1604.06339
- [14] C. Hongler, F. Johansson Viklund, K. Kytola, *Lattice representations of the Virasoro algebra I: discrete Gaussian free field*, arXiv:1307.4104
- [15] W.M. Koo and H. Saleur, *Representations of the Virasoro algebra from lattice models*, Nucl. Phys. B426, 459 (1994).
- [16] A. M. Gainutdinov, D. Ridout and I. Runkel Eds, *Logarithmic conformal field theory*, J. Phys. A46 (2013).
- [17] Y.-Z. Huang, J. Lepowsky, and L. Zhang, *Logarithmic tensor category theory for generalized modules for a conformal vertex algebra, I-VIII*, arXiv:1012.4193; arXiv:1012.4196; arXiv:1012.4197; arXiv:1012.4198; arXiv:1012.4199; arXiv:1012.4202, arXiv:1110.1929, arXiv:1110.1931.
- [18] P. A. Pearce, J. Rasmussen and J. B. Zuber, *Logarithmic minimal models*, J. Stat. Mech.0611:P11017 (2006).
- [19] N. Read and H. Saleur, *Enlarged symmetry algebras of spin chains, loop models, and S-matrices*, Nucl. Phys. B 777, 263 (2007).
- [20] N. Read and H. Saleur, *Associative-algebraic approach to logarithmic conformal field theories*, Nucl. Phys. B 777, 316 (2007).
- [21] A.M. Gainutdinov and R. Vasseur, *Lattice fusion rules and logarithmic operator product expansions*, Nucl. Phys. B 868, 223-270 (2013).
- [22] A.M. Gainutdinov, N. Read and H. Saleur, *Continuum limit and symmetries of the periodic $gl(1|1)$ spin chain*, Nucl. Phys. B871 (2013) 245.
- [23] A. M. Gainutdinov, J. L. Jacobsen, H. Saleur and R. Vasseur, *A physical approach to the classification of indecomposable Virasoro representations from the blob algebra*, Nucl. Phys. B873, 614 (2013).
- [24] P.A. Pearce, J. Rasmussen, S. Villani, *Infinitely extended Kac table of solvable critical dense polymers*, J. Phys. A: Math. Theor. 46 (2013) 175202.
- [25] A. M. Gainutdinov, H. Saleur and I. Yu. Tipunin, *Lattice W -algebras and logarithmic CFTs*, J. Phys. A: Math. Theor. 47 (2014) 495401.
- [26] A. Morin-Duchesne, J. Rasmussen, D. Ridout, *Boundary algebras and Kac modules for logarithmic minimal models*, arXiv:1503.07584.
- [27] J. Belletete, *Fusion rules for the Temperley-Lieb algebra and its dilute generalizations*, arXiv:1505.02112.
- [28] A. M. Gainutdinov, N. Read and H. Saleur, *Associative algebraic approach to logarithmic CFT in the bulk: the continuum limit of the $gl(1|1)$ periodic spin chain, Howe duality and the interchiral algebra*, Comm. Math. Phys., 341, 35 (2016).
- [29] D. Ridout and Y. St Aubin, *Standard modules, induction and the Temperley-Lieb algebra*, arXiv:1204.4505
- [30] J. Belletete, D. Ridout and Y. St Aubin, *Restriction and induction of indecomposable modules over the Temperley-Lieb algebra*, arXiv:1605.05159
- [31] J. Cardy, *Conformal invariance and surface critical behavior*, Nucl. Phys. B240 (1984) 514.
- [32] A. M. Gainutdinov, N. Read and H. Saleur, *Bimodule structure in the periodic $gl(1|1)$ spin chain*, Nucl. Phys. B 871 (2013) 289-329
- [33] A. Morin-Duchesne, P. Pearce and J. Rasmussen, *Modular invariant partition function of critical dense polymers*, arXiv:1303.4895
- [34] A. Morin-Duchesne and Y. St Aubin, *Jordan cells of periodic loop models*, arXiv:1302.5483
- [35] A. M. Gainutdinov, N. Read, H. Saleur and R. Vasseur, *The periodic $sl(2|1)$ alternating spin chain and its continuum limit as a bulk LCFT at $c = 0$* , JHEP (2015) 114.
- [36] A.M. Gainutdinov, J.L. Jacobsen, H. Saleur, in preparation.
- [37] B.W. Westbury, *The representation theory of the Temperley-Lieb algebras*, Math. Z. 219, 539-565 (1995).
- [38] V. Dlab and C.M. Ringel, *A construction of quasi-hereditary algebras*, Compos. Math. 70, 155-175 (1989).
- [39] P.P. Martin and D. McAnally, *On Commutants, Dual Pairs and Non-Semisimple Algebras from Statistical Mechanics*, Int. J. Mod. Phys. A 7, Supp. 1B, 675 (1992).

- [40] P.P. Martin and H. Saleur, *The blob algebra and the periodic Temperley-Lieb algebra*, Lett. Math. Phys. 30, 189 (1994).
- [41] V. Jones, *A quotient of the affine Hecke algebra in the Brauer algebra*, L'enseignement Mathématique 40, 313 (1994).
- [42] R.M. Green, *On representations of affine Temperley-Lieb algebras*, Algebras and Modules II, CMS Conference Proceedings, vol. 24, Amer. Math. Soc., Providence, RI, 1998, pp. 245–261.
- [43] J.J. Graham and G. I. Lehrer, *The two-step nilpotent representations of the extended Affine Hecke algebra of type A*, Compositio Mathematica 133 (2002) 173.
- [44] V. Chari and A. Pressley, *Quantum affine algebras and affine Hecke algebras*, Pac. Jour. Math. 174, 295 (1996).
- [45] J.J. Graham and G.I. Lehrer, *Diagram algebras, Hecke algebras, and decomposition numbers at roots of unity*, Ann. Sci. Ecole Norm. Sup. 36, 479–524 (2003).
- [46] P.P. Martin and D. Woodcock, *On the structure of the blob algebra*, J. Algebra 225, 957–988 (2000).
- [47] C. Kassel, *Quantum Groups*, Springer (1995)
- [48] B Feigin and D Fuchs, *Representations of the Virasoro algebra*, Adv. Stud. Contemp. Math., 7:465554, 1990.
- [49] K. Kytölä and D. Ridout, *On Staggered Indecomposable Virasoro Modules*, J. Math. Phys. 50, 123503 (2009).
- [50] F. Rohsiepe, *On reducible but indecomposable representations of the Virasoro algebra*, hep-th/9611160.
- [51] A.A. Belavin, A.M. Polyakov and A.B. Zamolodchikov, *Infinite conformal symmetry in two-dimensional quantum field theory*, Nucl. Phys. B 241, 333 (1984).
- [52] I. Frenkel, M. Zhu, *Vertex operator algebras associated to modified regular representations of the Virasoro algebra*, arXiv:1012.5443.
- [53] V. Gurarie and A.W.W. Ludwig, *Conformal Field Theory at central charge $c=0$ and Two-Dimensional Critical Systems with Quenched Disorder*, hep-th/0409105.
- [54] H. Eberle and M. Flohr, *Virasoro representations and fusion for general augmented minimal models*, J. Phys. A 39, 15245–15286 (2006).
- [55] P. Mathieu and D. Ridout, *From Percolation to Logarithmic Conformal Field Theory*, Phys. Lett. B 657, 120 (2007).
- [56] P. Mathieu and D. Ridout, *Logarithmic $M(2,p)$ Minimal Models, their Logarithmic Couplings, and Duality*, Nucl. Phys. B 801, 268 (2008).
- [57] M. R. Gaberdiel, I. Runkel, S. Wood, *Fusion rules and boundary conditions in the $c = 0$ triplet model*, J. of Phys. A: Mathematical and Theoretical, V 42, N 32 (2009).
- [58] W. Nahm, *Quasirational fusion products*, Int. J. Mod. Phys. B8, 3693 (1994).
- [59] M. Gaberdiel and H. Kausch, *Indecomposable Fusion Products*, Nucl. Phys. B 477, 293 (1996).
- [60] R. Behrend, P. Pearce, V. Petkova and J.B. Zuber, *On the Classification of Bulk and Boundary Conformal Field Theories*, Phys.Lett. B444 (1998) 163.
- [61] I. Runkel, M. Gaberdiel and S. Wood, *Logarithmic bulk and boundary conformal field theory and the full centre construction*, contribution to the proceedings of 'Conformal field theories and tensor categories' (Beijing, June 2011), arXiv:1201.6273