



HAL
open science

Fourier-positivity constraints on QCD dipole models

Bertrand G. Giraud, Robi Peschanski

► **To cite this version:**

Bertrand G. Giraud, Robi Peschanski. Fourier-positivity constraints on QCD dipole models. Physics Letters B, 2016, 760, 26-30, <10.1016/j.physletb.2016.06.033>. <cea-01299882>

HAL Id: cea-01299882

<https://cea.hal.science/cea-01299882v1>

Submitted on 8 Apr 2016

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons CC BY 4.0 - Attribution - International License

Fourier-positivity constraints on QCD dipole models

Bertrand G. Giraud and Robi Peschanski
Institut de Physique Théorique,
CEA, IPhT, F-91191 Gif-sur-Yvette, France
*CNRS, URA 2306 **

Fourier-positivity (\mathcal{F} -positivity), i.e. the mathematical property that a function has a positive Fourier transform, can be used as a constraint on the parametrization of QCD dipole-target cross-sections or Wilson line correlators in transverse position space r . They are Bessel transforms of positive transverse momentum dependent gluon distributions. Using mathematical \mathcal{F} -positivity constraints on the limit $r \rightarrow 0$ behavior of the dipole amplitudes, we identify the common origin of the violation of \mathcal{F} -positivity for various, however phenomenologically convenient, dipole models. It is due to the behavior $r^{2+\epsilon}$, $\epsilon > 0$ softer, even slightly, than color transparency. \mathcal{F} -positivity seems thus to conflict with the present dipole formalism when it includes a QCD running coupling constant $\alpha(r)$.

PACS numbers:

I. ORIGIN OF THE PROBLEM

The introduction of the dipole formalism [1] in QCD predictions for high-energy deep-inelastic scattering (DIS) and other processes at “small x ” has prompted a lot of phenomenological activity with good success during the recent years. Initially, the dipole representation allowed one to express the DIS structure functions off a target in terms of the (imaginary part of) the forward scattering amplitude $\mathcal{N}(Y, r)$ of a colorless quark-antiquark ($q\bar{q}$) pair on a target. Here $q\bar{q}$ is a colorless QCD dipole of transverse size r , elastically scattering on the target at rapidity Y . $\mathcal{N}(Y, r)$, up to a normalization, is the $q\bar{q}$ -target total cross-section, and thus necessarily positive.

The following Fourier-Bessel relations between transverse spatial and momentum coordinates have been written [2],

$$\begin{aligned}\mathcal{N}(Y, r) &= r^2 \int_0^\infty k dk J_0(kr) \tilde{\mathcal{N}}(Y, k) , \\ \tilde{\mathcal{N}}(Y, k) &= \int_0^\infty \frac{dr}{r} J_0(kr) \mathcal{N}(Y, r) ,\end{aligned}\tag{1}$$

where $\tilde{\mathcal{N}}(Y, k)$ is related to the F_2 structure function of DIS off the target. At that time, this connection [3] implied that $\tilde{\mathcal{N}}(Y, k)$ should be positive within a certain approximation. Within this approximation, (1) relate, through a Fourier-Bessel transformation, two physical observables which are both positive. This was our original motivation [4] to discuss Fourier-positivity properties in general, hoping to get constraints on the QCD dipole models.

At present time and in a more general setting, it can be proven that QCD observables at small x like $\mathcal{N}(Y, r)$ and $\tilde{\mathcal{N}}(Y, k)$ are expected to be conjugated *positive* functions through a Bessel transformation. This characteristic feature of positive Fourier partners is indeed more general in the dipole formalism and its various developments than the one that we originally considered [4]. In [5], identical relations (1) are found for $q\bar{q}$ jet correlations in DIS. In that case, $\mathcal{N}(Y, r)$ is obtained from a specific correlator of Wilson lines and $\tilde{\mathcal{N}}(Y, k)$ is the Weizsäcker-Williams transverse momentum dependent (TMD) gluon distribution of the target.

Moreover, in Ref.[6] and in the previous reference [5], a different TMD gluon distribution describes inclusive forward jet production and dijet correlations, respectively, in proton-nucleus (p - A) scattering. In that case, a similar but different Bessel relation can be written between this TMD gluon distribution and the total cross-section of the target with a gluon-gluon (gg) colorless dipole. We will show, later on, that this relation between p - A observables gives rise to the same constraints as (1).

On a mathematical footing, we call “ \mathcal{F} -positivity” that mathematical property of a real positive function whose Fourier transform is itself positive. The aim of our paper is to derive some useful constraints on current QCD dipole

*Electronic address: bertrand.giraud@cea.fr; robi.peschanski@cea.fr

(or related) models coming from \mathcal{F} -positivity. More precisely, let us consider the case of the relations (1). Given a model parameterization for $\mathcal{N}(Y, r) \geq 0$, we want to determine to which conditions it leads to a positive distribution $\tilde{\mathcal{N}}(Y, k) \geq 0$. This problem arising for phenomenological QCD dipole models was already noticed in Refs. [7, 8] and its relation to mathematical properties in [4, 9].

As a mathematical problem, \mathcal{F} -positivity is old and still not completely solved [10]. There exists fundamental theorems [11] ensuring the \mathcal{F} -positivity of real functions, as we shall see later on. However, up to our knowledge, there is no explicit property of the set of all \mathcal{F} -positive functions in 2-(or else) dimensions which would *a priori* suggest the parametrization of $\mathcal{N}(Y, r)$. In the absence of a definite practical scheme, we recently developed [12, 13] a set of practical tools in the form of a hierarchy of necessary conditions. They hold for the general 1-d and 2-d radial cases, which we will now try and apply specifically to the QCD dipole problem.

The plan of our paper is the following. In section II, we describe the mathematical derivation of the specific tools we will apply to the QCD dipole formalism. We insist in particular on the \mathcal{F} -positivity constraints on the small-size limit of the dipole input. In section III, starting with the Fourier partners in Eq.(1), we review examples from the main classes of phenomenologically useful dipole (or related) models and show how and why they often contradict \mathcal{F} -positivity. In section IV, extending our study to the other type of Bessel relations [5, 6], we show that, for both formulations, the QCD dipole or Wilson line correlators models, including a running coupling constant, violate \mathcal{F} -positivity for the same reason, which was yet unknown. The final section V summarizes our results and discusses the origin of the \mathcal{F} -positivity failure, opening possible ways out to reconcile an improved dipole formalism with \mathcal{F} -positivity.

II. MATHEMATICAL ASPECTS AND TOOLS

The fundamental mathematical property characterizing \mathcal{F} -positivity uses the Bochner theorem [11]. Applied for instance to the QCD dipole problem issued from relations (1), it can be expressed as follows. \mathcal{F} -positivity of $\tilde{\mathcal{N}}(Y, k)$ is equivalent to the statement that $\mathcal{N}(Y, r)/r^2$ is not only positive but also *positive-definite*. In the full 2-dimensional transverse coordinate space, positive-definiteness means that for any set of positions in \mathbb{R}^2 , $\{\vec{r}_i, i = 1, \dots, n\}$, for any $n \in \mathbb{N}$ and for any set of numbers $\{u_i, i = 1, \dots, n\}$, the $n \times n$ matrix \mathbb{M} with elements $\mathcal{N}(Y, |\vec{r}_i - \vec{r}_j|)/|\vec{r}_i - \vec{r}_j|^2$ is positive definite. Namely,

$$\sum_{i,j=1}^n u_i \frac{\mathcal{N}(Y, |\vec{r}_i - \vec{r}_j|)}{|\vec{r}_i - \vec{r}_j|^2} u_j > 0, \quad \forall u_i, \forall \vec{r}_i, \forall n. \quad (2)$$

In other terms, the lowest eigenvalue of \mathbb{M} remains positive for all \vec{r}_i, u_i and all values of n .

The Bochner theorem with its applications appears to be still the major tool in the domain. However, to our knowledge, there does not yet exist a mathematical classification of Fourier-positive functions which, for instance, could constrain at once an appropriate parametrization for model building. Testing positive-definiteness (2) cannot be done concretely, due to the infinite number of the constraints. Conversely, a calculation of Fourier transforms for checking Fourier-positivity is obviously possible, but this does not give easy means to select *a priori* appropriate \mathcal{F} -positive sets of functions, or to find the origin of their eventual \mathcal{F} -positivity violation.

Our approach is thus to find new constraints of Fourier-positivity allowing for simple and efficient selection rules on the QCD model parametrizations starting with both *positive* and *positive-definite* input functions whose $\mathcal{N}(Y, r)/r^2$ in Eq.(1) is an example.

In our previous papers [4, 12], we performed a study of large sets of *positive* and *positive-definite* functions in 1 and 2 (radial) dimensions. We used bases of positive polynomials or convexity properties. We found some intricate features of the problem.

In the present work we will use tools based on the Bochner theorem, coming from Ref.[13] where a satisfactory detection of Fourier positivity, thoroughly tested for large sets of functions in one- and radial two dimensions, was obtained. Let us sketch our formulation in the specific case of the (radial) 2-dimensional problem, appropriate for the present study.

Matrices \mathbb{M} from a point lattice. The tool we will use is the set of necessary conditions coming from the positive-definiteness condition (2) for a discrete 2d-lattice of points, $\vec{r}_i = \{n_i^{(1)}r, n_i^{(2)}r\}$, $n_i^{(1),(2)} \in 1, \dots, n$. For example, and further application, let us consider the 3×3 matrix

$$\{\mathbb{M}_{i,j}\}_3 \equiv \left\{ \frac{\mathcal{N}(Y, |\vec{r}_i - \vec{r}_j|)}{|\vec{r}_i - \vec{r}_j|^2} \right\}, \quad \vec{r}_i = \{0, 0\}, \{0, r\}, \{r, 0\}, \quad (3)$$

which, denoting for simplicity $\psi(r) \equiv \mathcal{N}(Y, r)/r^2$, leads to the \mathcal{F} -positivity conditions for the matrix

$$\begin{pmatrix} \psi(0) & \psi(r) & \psi(r\sqrt{2}) \\ \psi(r) & \psi(0) & \psi(r) \\ \psi(r\sqrt{2}) & \psi(r) & \psi(0) \end{pmatrix}. \quad (4)$$

Positive-definiteness implies positivity of the matrix determinant and of its minors along its diagonal, leading, up to a rescaling of r , to the inequalities

$$\psi(0) > \psi(r) > 2 \frac{\psi^2(r/\sqrt{2})}{\psi(0)} - \psi(0). \quad (5)$$

Note that larger lattices including the points of (3) will always lead to condition (5) and will add others [13], forming a hierarchy of necessary conditions for \mathcal{F} -positivity.

An important addendum to the \mathcal{F} -positivity tests for a 2-dimensional radial function $\psi(r)$, as noticed in our Ref.[12], is that they can be extended to the action of the radial Laplacian

$$\Delta_2 [\psi](r) \equiv -\frac{1}{r} \frac{d}{dr} \left(r \frac{d}{dr} \right) [\psi](r) \Rightarrow \Delta_2 [\mathcal{N}/r^2](Y, r) = \int_0^\infty k^3 dk J_0(kr) \tilde{\mathcal{N}}(Y, k). \quad (6)$$

Hence \mathcal{F} -positivity applies also for $\Delta_2 [\psi]$, as well as for its iterations, provided the integral in (6) converges.

III. \mathcal{F} -POSITIVITY TESTS OF THE QCD DIPOLE AMPLITUDE

We shall now apply \mathcal{F} -positivity tests to the amplitude parametrizations of $\mathcal{N}(Y, r)$. For this we review various popular classes of models used in the phenomenology of low x processes.

A. BFKL-type models

A phenomenologically successful model was proposed by [14]. It is based upon an effective BFKL-type amplitude corrected at large transverse size for unitarity. It reads

$$\begin{aligned} \mathcal{N}(Y, r) &= \mathcal{N}_0 \exp \left\{ 2 \left(\gamma_s + \frac{\log(2/rQ)}{\kappa\lambda Y} \right) \log \left(\frac{rQ}{2} \right) \right\} \quad \text{for } rQ \leq 2, \\ \mathcal{N}(Y, r) &= 1 - \exp \{ -a \log^2(brQ) \} \quad \text{for } rQ > 2. \end{aligned} \quad (7)$$

$Q \equiv Q(Y)$ is the rapidity-dependent ‘‘saturation scale’’ and $\gamma_s, \kappa, \lambda$ are characteristic constants of the BFKL amplitude at leading log level [14]. Figure 1 shows the functions \mathcal{N} and $\tilde{\mathcal{N}}$ in the appropriate scaled units, rQ and k/Q respectively, for physical values of $Y = 4, 6, 8$. The left-hand plot displays the dipole amplitude together with its value divided by $(rQ)^2$. The right-hand plot is the gluon amplitude in log-log plot, exhibiting zeroes, showing explicitly the \mathcal{F} -positivity violation of the BFKL-inspired dipole amplitude. A characteristic feature that we retain from the function $\mathcal{N}(Y, r)/(rQ)^2$, which forms the Fourier pair with the distribution $\varphi(Y, k/Q) \equiv \tilde{\mathcal{N}}(Y, k)$ is that it increases from zero at small r . This is a signature of \mathcal{F} -positivity violation since the constraint (5) is not satisfied. Then, already at the first level of our hierarchy of constraints, the amplitude (7) does not obey \mathcal{F} -positivity. This can be traced back to the logarithm $\log(2/rQ)$ in the exponent of the BFKL-inspired expression (7), since it drives the amplitude to zero when $r \rightarrow 0$, for any value of the parameters. In particular, the violation is not related to the presence of non-regular behavior at the transition point $Qr = 2$, as sometimes conjectured.

One may object that when this logarithm dominates, the effective BFKL amplitude is no more justified. However, as we now will see, \mathcal{F} -positivity violation due to the small- r behavior of the amplitude is not restricted to BFKL-inspired models.

B. Golec-Biernat-Wusthoff-type models

The Golec-Biernat-Wusthoff model [15] (GBW) is a quite popular model of the dipole amplitude. In a recent paper [16] a generalized expression of this type has been taken for the initial input amplitude for a phenomenological

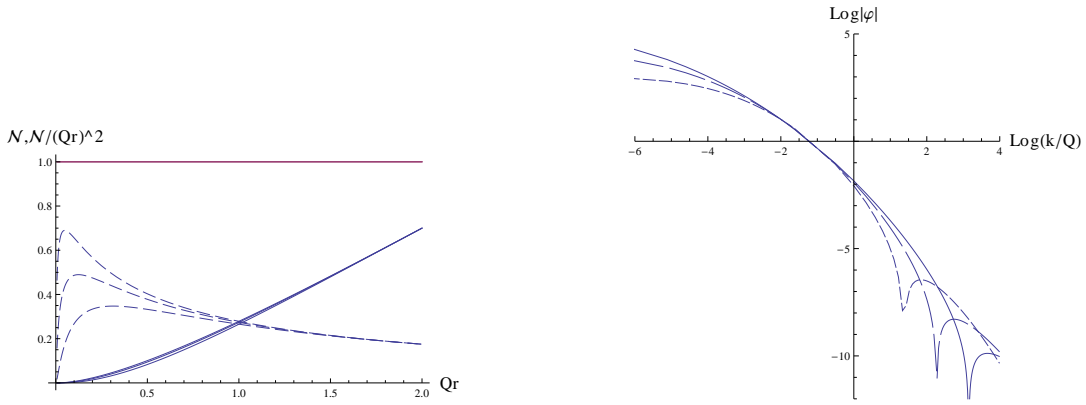


FIG. 1: \mathcal{F} -positivity tests of the BFKL-inspired dipole amplitude. Left, the dipole amplitude: full line, $\mathcal{N}(Y,r)$, dashed line, $\mathcal{N}(Y,r)/(rQ)^2$, as a function of rQ . Right, its Fourier-transform $\varphi(Y,k/Q) \equiv \tilde{\mathcal{N}}(Y,k)$ in a $\log|\varphi|, \log(k/Q)$ plot for $Y = 4, 6, 8$. The dips signal the zeroes of φ after which it is negative. The range of Y runs from bottom to top in the left plot, and from right to left in the right plot.

application of the Balitsky-Kovchegov (BK) evolution equation [17] with rapidity. It reads

$$\mathcal{N}(Y,r) = \left\{ 1 - \exp \left[- \left(\frac{1}{4} (rQ(Y))^2 \right)^p \right] \right\}^{1/p}, \quad (8)$$

where $p = 1$ for the original model [15] while $p > 1$ is clearly preferred in the recent application [16]. One sees from Fig.2 that \mathcal{F} -positivity is still violated even if the inequalities (5) are satisfied. The function $\mathcal{N}/(rQ)^2$, appears to be slightly decreasing with Qr . The situation clarifies if one considers the radial Laplacian operator Δ_2 defined in (6).

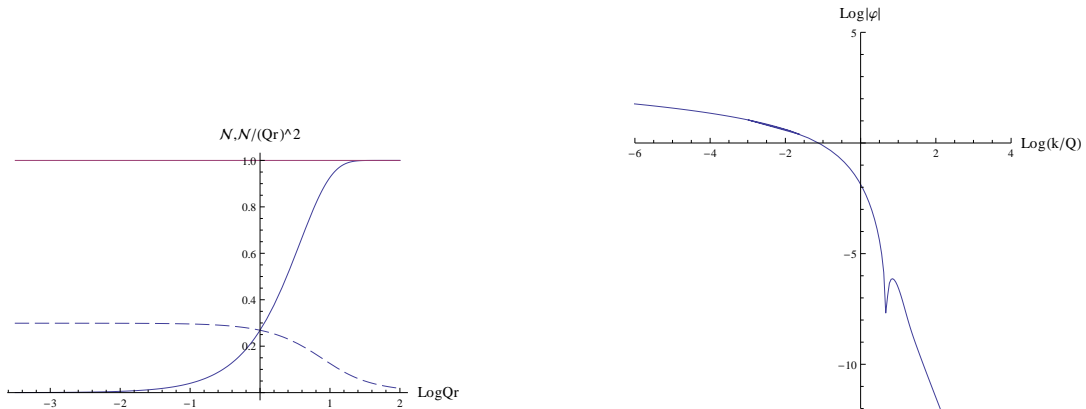


FIG. 2: \mathcal{F} -positivity tests of the GBW-type dipole amplitude. Left, the dipole amplitude: $\mathcal{N}(Y,r)$, full line, $\mathcal{N}(Y,r)/(rQ)^2$, dashed line, in a semi-log plot as a function of $\log(Qr)$. Right, its Fourier-transformed $\varphi(Y,k/Q) \equiv \tilde{\mathcal{N}}(Y,r)$ in a $\log|\varphi|, \log(k/Q)$ plot. The dip corresponds to the first zero of φ . The Fourier pair of functions is displayed for $p=1.148$ in Eq.(8), as found in Ref. [16], table 1.

One finds at small r

$$\Delta_2[\psi](r) = - \left\{ \frac{d^2}{(dr)^2} + \frac{1}{r} \frac{d}{dr} \right\} \left[\frac{\mathcal{N}}{r^2} \right] (Y,r) \approx - \left\{ \frac{d^2}{(dr)^2} + \frac{1}{r} \frac{d}{dr} \right\} \left[\frac{1}{4} - \frac{1}{8p} \left(\frac{r^2}{4} \right)^p \right] = \frac{2p}{4p+1} r^{2(p-1)}. \quad (9)$$

Alongside Eq.(5), \mathcal{F} -positivity requires also $\Delta_2[\psi](0) > \Delta_2[\psi](r)$, which is satisfied only for $p < 1$ in Eq.(9). Hence, any function with $p > 1$ appearing for phenomenological preference in [16] does not satisfy the \mathcal{F} -positivity constraint.

C. Mc Lerran-Venugopalan-type models

The Mc Lerran-Venugopalan model [18] (MV) for the dipole amplitude is well known in low x phenomenology and theory. It gives rise to various recent versions of MV-type, with successful applications to phenomenology. For our present study we will retain first variations of the original model used as a phenomenological input for Y -dependence [19], namely

$$\mathcal{N}(Y, r) = 1 - \exp \left[-\frac{1}{4} (rQ(Y))^{2\gamma} \log \left(e e_c + \frac{1}{r\Lambda} \right) \right], \quad (10)$$

where the constants (γ, e_c) allow a best fit search while they are equal to $(1, 1)$ in the original MV model. Using the condition (5) leads to impose $\gamma \leq 1$. This is satisfied by the MV formulas (10) with $\gamma = 1$ (including the model with $e_c \neq 1$ [19]) but not satisfied for the so-called AAMQS model [20] which has $\gamma = 1.119$.

It is worthwhile to note that a problem of \mathcal{F} -positivity violation related to the MV model has been noticed in Ref.[8] when one uses the original version instead of the approximation by the logarithm in (10). It contains a cut-off $\theta(r\Lambda)$. \mathcal{F} -positivity violation in this case is thus well detected by our test (5) since $0 = \psi(0) < \psi(r)$ for r larger than the cut-off. \mathcal{F} -positivity is restored [8] by replacing the cut-off by a smooth extrapolation. Indeed, the corresponding derivative of $\varphi(r)$ at $r = 0$ is negative, ensuring $\psi(0) > \psi(r)$ as required by \mathcal{F} -positivity. In this case, this criterion is indeed sufficient.

One also considered recently [16] the MV model with running coupling constant, namely [18],

$$\mathcal{N}(Y, r) = \left\{ 1 - \exp \left[- \left(\frac{1}{4} (rQ(Y))^2 \bar{\alpha}(rC) \left[1 + \log \left(\frac{\bar{\alpha}_{sat}}{\bar{\alpha}(rC)} \right) \right] \right)^p \right] \right\}^{1/p}, \quad (11)$$

where $\bar{\alpha}(rC) \propto \log^{-1}(1/(rC))$ is the running QCD coupling in transverse coordinate space and C, α_{sat}, p are phenomenological constants. It is easy to realize that $\mathcal{N}(Y, r)/r^2 \approx \left[\frac{\log \log \frac{1}{r}}{\log \frac{1}{r}} \right] \rightarrow 0$ when $r \rightarrow 0$: \mathcal{F} -positivity is thus violated, in a kind of “weak” logarithmic form, due to the running coupling. We shall see now that this feature is not subsidiary; actually it is related to a deeper theoretical problem of the QCD dipole formalism with a running coupling constant $\alpha(r)$.

D. Saturation model with DGLAP evolution

On a theoretical ground, when $r \rightarrow 0$, a correspondence with large k/Q of the observable, the QCD dipole model is faced with the problem of compatibility with the DGLAP evolution equations. Let us examine which constraints this generates on \mathcal{F} -positivity tests. For this sake, and in order to take more general lessons, we shall consider first a well-known model [21] introducing a modification of the saturation model in order to take into account DGLAP evolution. It reads

$$\mathcal{N}(Y, r) = 1 - \exp \left(-\frac{1}{3\sigma_0} \pi^2 r^2 \bar{\alpha}(\mu^2) \times xg(x, \mu^2) \right), \quad x \equiv e^{-Y}. \quad (12)$$

Here $xg(x, \mu^2)$ is the gluon distribution function in the proton considered at momentum fraction x and r -dependent scale $\mu^2 = C/r^2 + \mu_0^2$, with $\alpha(\mu^2) \propto \log \Lambda_{QCD}^2/\mu^2$ and σ_0, C phenomenological constants fitted to the deep-inelastic data.

At small x , the leading behavior of the gluon distribution function can be obtained from the resummation of the double leading logarithms, namely

$$xg(x, \mu^2) \approx \sum_n \frac{\left[\log(1/x) \int_{\mu_0^2}^{\mu^2} \alpha(k^2) dk^2/k^2 \right]^n}{(n!)^2} \sim \exp \sqrt{\log \frac{1}{x} \log \log \frac{\mu^2}{\mu_0^2}}, \quad (13)$$

leading to

$$\mathcal{N}(Y, r) \approx r^2 \frac{1}{\log \frac{1}{r^2}} \exp \sqrt{Y \log \log \frac{1}{r^2}} \rightarrow 0 \quad \text{when } r \rightarrow 0. \quad (14)$$

In this example, we see that the double leading log resummation of the perturbative expansion of the gluon distribution function is unable to fully compensate for the running of the coupling constant in front of (14).

There is a model-independent lesson to be remembered from the example (13) and the investigation of the previous subsection: The QCD asymptotic freedom, appearing in the dipole amplitude in the form of an overall running coupling going logarithmically to zero with the dipole size seems to systematically contradict \mathcal{F} -positivity at large enough transverse momentum. The inverse logarithmic behavior of the running coupling constant is only partially compensated by the resummation at large rapidity Y of the perturbative expansion of the dipole amplitude.

IV. \mathcal{F} -POSITIVITY: A GENERIC PROBLEM.

The various phenomenological models discussed in the previous sections show that there often exists a common problem with \mathcal{F} -positivity at large values of k/Q , i.e. when the gluon momentum k is large w.r.t. the typical (saturation) scale Q of the target. We have seen that the problem is due in all cases to the limit behavior of the amplitude $\mathcal{N}(Y, r)$ when $r \rightarrow 0$. As shown in subsections III C and III D, when the running $\alpha(r)$ of the QCD coupling constant is taken into account, it is responsible for the violation of the \mathcal{F} -positivity test (5). Let us extend now the results obtained with the Fourier partners of (1) to the other type of Bessel relation appearing in Refs [5, 6]. It relates the gg dipole-target forward amplitude $T(Y, r)$ to the TMD gluon distribution $\tilde{\mathcal{T}}(Y, k)$. The Bessel relation in those cases reads

$$\tilde{\mathcal{T}}(Y, k) = \int r dr J_0(kr) \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \mathcal{T}(Y, r) \right) > 0, \quad (15)$$

giving rise again to a \mathcal{F} -positivity constraint.

To analyze this constraint, we introduce a generic parameterization of the gg dipole-target forward when $r \rightarrow 0$, namely,

$$\mathcal{T}(Y, r) \propto r^{2+\epsilon} \quad \text{when } r \rightarrow 0. \quad (16)$$

“Color transparency” corresponds to $\epsilon = 0$. Choosing $0 < \epsilon \ll 1$ allows one to provide an effective description of the effect of the running QCD coupling constant only partially compensated by resummation effects. As we shall now derive, this choice leads to a violation of \mathcal{F} -positivity.

Note that the same behavior (15) applied to $N(Y, r)$ in (1) leads also to the violation of \mathcal{F} -positivity. Indeed one finds $N(Y, r)/r^2 \sim r^\epsilon$, increasing with r and thus violating the constraint (5).

In the integrand of the Bessel transform (15), one recognizes the opposite of the Laplacian operator of (6), namely,

$$\frac{1}{r} \frac{d}{dr} \left(r \frac{d}{dr} \mathcal{T}(Y, r) \right) = -\Delta_2[\mathcal{T}](Y, r) \sim (2 + \epsilon)^2 r^\epsilon \quad \text{when } r \rightarrow 0. \quad (17)$$

Hence, for $\epsilon > 0$, the constraint (5) and thus \mathcal{F} -positivity are violated.

To summarize our discussion, the running $\alpha(r)$ of the QCD coupling constant in front of the dipole amplitude $\mathcal{N}(Y, r)$ or the Wilson line correlator $\mathcal{T}(Y, r)$ lead to a violation of \mathcal{F} -positivity tests in relation with the small- r behavior. It corresponds to a small transverse distance between the Wilson lines or between the two gluons in the gg -dipole, respectively.

V. SUMMARY AND OUTLOOK

We have discussed the constraints on QCD dipole models (and its extensions) arising from \mathcal{F} -positivity. Indeed, the dipole cross-section or the Wilson line correlators in transverse position space are related through Bessel transformations [5] to two different transverse momentum dependent gluon distributions. These can be directly expressed in terms of observables and thus required to be positive. The full mathematical characterization of \mathcal{F} -positive functions is still an open problem. However, in Refs. [4, 12, 13], we have proposed a set of rather simple conditions for \mathcal{F} -positivity. In the present work, we show that they provide quite stringent necessary constraints on QCD dipole models, if not on the dipole formalism itself. Let us quote the main results:

i) Using a set of \mathcal{F} -positivity constraints [12, 13] based on the Bochner theorem [11], we show that dipole models, frequently used in low x physics phenomenology, often do not satisfy \mathcal{F} -positivity. We show that this violation is related to the limit behavior of the model amplitudes when the transverse size $r \rightarrow 0$.

ii) In all cases¹ when \mathcal{F} -positivity is violated, this happens if the behavior of the dipole amplitudes or Wilson line correlators is $\sim r^{2+\epsilon}$, $0 < \epsilon \ll 1$ when the transverse size $r \rightarrow 0$, where “color transparency” corresponds to $\epsilon = 0$.

iii) On a more theoretical level, \mathcal{F} -positivity appears to put a question mark on the QCD dipole (or Wilson correlators) formalism itself. Indeed, when considering a dipole model with a coupling constant running with r , it seems difficult to satisfy the constraint (5). Indeed, asymptotic freedom of $\alpha(r) \sim 1/\log(1/r)$ imposes a decreasing behavior of the amplitude faster than r^2 when $r \rightarrow 0$. It appears only partially compensated by the all-order perturbative resummation of the double logarithms in the low x domain.

To our knowledge, there does not yet exist a satisfactory solution of the \mathcal{F} -positivity constraints on the QCD dipole formalism. One way out of the difficulties would be that the Bessel transform relations could be relaxed at next-leading orders. However then, the question arises how to connect the r -dependent amplitudes with the transverse momentum observables. Another solution [22], would be to modify the Wilson lines formalism starting with the k -dependent running coupling in order to preserve \mathcal{F} -positivity. Indeed, it is conceivable that the inner structure of a QCD dipole, which is after all a composite system of two quarks or two gluons, is resolved at very high transverse momentum of the probe.

Hence the \mathcal{F} -positivity problem of the QCD dipole formalism seems to deserve more study, both from phenomenological and theoretical points of view.

VI. ACKNOWLEDGEMENTS

We want to thank Edmond Iancu, Cyrille Marquet and Gregory Soyez for stimulating discussions.

-
- [1] A. H. Mueller, “Soft gluons in the infinite momentum wave function and the BFKL pomeron,” Nucl. Phys. B **415**, 373 (1994).
 A. H. Mueller and B. Patel, “Single and double BFKL pomeron exchange and a dipole picture of high-energy hard processes,” Nucl. Phys. B **425**, 471 (1994) [hep-ph/9403256].
 A. H. Mueller, “Unitarity and the BFKL pomeron,” Nucl. Phys. B **437**, 107 (1995) [hep-ph/9408245].
- [2] Y. V. Kovchegov, “Unitarization of the BFKL pomeron on a nucleus,” Phys. Rev. D **61**, 074018 (2000) [hep-ph/9905214].
- [3] R. Peschanski, unpublished;
- [4] S. Munier and R. B. Peschanski, “High-energy factorization predictions for the charm structure function F2(c) at HERA,” Nucl. Phys. B **524**, 377 (1998) [hep-ph/9802230].
- [5] B. G. Giraud and R. B. Peschanski, “On positive functions with positive Fourier transforms,” Acta Phys. Polon. B **37**, 331 (2006) [math-ph/0504015].
- [6] F. Dominguez, B. W. Xiao and F. Yuan, “ k_t -factorization for Hard Processes in Nuclei,” Phys. Rev. Lett. **106**, 022301 (2011) [arXiv:1009.2141 [hep-ph]].
 F. Dominguez, C. Marquet, B. W. Xiao and F. Yuan, Phys. Rev. D **83**, 105005 (2011) [arXiv:1101.0715 [hep-ph]].
- [7] C. Marquet, “A QCD dipole formalism for forward-gluon production,” Nucl. Phys. B **705**, 319 (2005) [hep-ph/0409023].
- [8] F. Gélis, unpublished.
- [9] C. Marquet, “Forward inclusive dijet production and azimuthal correlations in pA collisions,” Nucl. Phys. A **796**, 41 (2007) [arXiv:0708.0231 [hep-ph]].
- [10] T. Lappi, “Small x physics and RHIC data,” Int. J. Mod. Phys. E **20**, 1 (2011) [arXiv:1003.1852 [hep-ph]].
- [11] For a first statement on Fourier-positivity:
 P. Lévy: “Fonctions caractéristiques positives” [Positive characteristic functions]. *Comptes Rendus Hebdomadaires des Séances de l’Académie des Sciences, Série A, Sciences Mathématiques* **265** (1967) 249-252. [in French]
 For a comprehensive mathematical presentation,
 I.M. Gel’fand and N.Ya. Vilenkin (1968), *Generalized Functions*, Vol. IV, Academic Press, New-York and London.
 For recent mathematical works:
 P. Jaming, M. Matolcsi and S.G. Révész, “On the extremal rays of the cone of positive, positive definite functions,” J. Fourier Anal. Appl., 1, 561 (2009).
 A. Hinrichs, J. Vybrál, “On positive positive-definite functions and Bochner’s Theorem,” J. Complexity 27, 264, (2011).
- [12] S. Bochner, *Vorlesung über Fouriersche Integrale* (Leipzig Verlag, 1932); English translation: Lectures on Fourier Integrals, *Ann. Math. Stud.*, **42**, (Princeton University Press, 1959).
 I.M. Gel’fand and N.Ya. Vilenkin (1968), *Generalized Functions*, Vol. IV, Academic Press, New-York and London.
- [13] B. G. Giraud and R. Peschanski, “On the positivity of Fourier transforms,” arXiv:1405.3155 [math-ph].

¹ In subsection IIIB a higher rank convexity constraint (9) is violated.

- [13] B. G. Giraud and R. Peschanski, "From "Dirac combs" to Fourier-positivity," arXiv:1509.02373 [math-ph]; to be published in Acta Physica Polonica B, Vol. 47 (2016) No 4.
- [14] E. Iancu, K. Itakura and S. Munier, Phys. Lett. B **590**, 199 (2004) [hep-ph/0310338].
- [15] K. J. Golec-Biernat and M. Wusthoff, "Saturation effects in deep inelastic scattering at low Q^2 and its implications on diffraction," Phys. Rev. D **59**, 014017 (1998) [hep-ph/9807513].
- [16] E. Iancu, J. D. Madrigal, A. H. Mueller, G. Soyez and D. N. Triantafyllopoulos, "Collinearly-improved BK evolution meets the HERA data," Phys. Lett. B **750**, 643 (2015) [arXiv:1507.03651 [hep-ph]].
- [17] I. Balitsky, "Operator expansion for high-energy scattering," Nucl. Phys. B **463**, 99 (1996) [hep-ph/9509348].
Y. V. Kovchegov, "Small x $F(2)$ structure function of a nucleus including multiple pomeron exchanges," Phys. Rev. D **60**, 034008 (1999) [hep-ph/9901281].
- [18] L. D. McLerran and R. Venugopalan, "Computing quark and gluon distribution functions for very large nuclei," Phys. Rev. D **49**, 2233 (1994) [hep-ph/9309289]. *ibidem* Phys. Rev. D **49**, 3352 (1994) [hep-ph/9311205].
- [19] T. Lappi and H. Mntysaari, "Single inclusive particle production at high energy from HERA data to proton-nucleus collisions," Phys. Rev. D **88**, 114020 (2013) [arXiv:1309.6963 [hep-ph]].
- [20] J. L. Albacete, N. Armesto, J. G. Milhano, P. Quiroga-Arias and C. A. Salgado, "AAMQS: A non-linear QCD analysis of new HERA data at small x including heavy quarks," Eur. Phys. J. C **71**, 1705 (2011) [arXiv:1012.4408 [hep-ph]].
- [21] J. Bartels, K. J. Golec-Biernat and H. Kowalski, "A modification of the saturation model: DGLAP evolution," Phys. Rev. D **66**, 014001 (2002) [hep-ph/0203258].
- [22] E. Iancu, private communication.