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Anh Ha Le, Pascal Omnes

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# AN A POSTERIORI ERROR ESTIMATION FOR THE DISCRETE DUALITY FINITE VOLUME DISCRETIZATION OF THE STOKES EQUATIONS

ANH HA LE<sup>1</sup> AND PASCAL OMNES<sup>1,2</sup>

**Abstract.** We derive an a posteriori error estimation for the discrete duality finite volume (DDFV) discretization of the stationary Stokes equations on very general twodimensional meshes, when a penalty term is added in the incompressibility equation to stabilize the variational formulation. Two different estimators are provided: one for the error on the velocity and one for the error on the pressure. They both include a contribution related to the error due to the stabilization of the scheme, and a contribution due to the discretization itself. The estimators are globally upper as well as locally lower bounds for the errors of the DDFV discretization. They are fully computable as soon as a lower bound for the inf-sup constant is available. Numerical experiments illustrate the theoretical results and we especially consider the influence of the penalty parameter on the error for a fixed mesh and also of the mesh size for a fixed value of the penalty parameter. A global error reducing strategy that mixes the decrease of the penalty parameter and adaptive mesh refinement is described.

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## INTRODUCTION

Let  $\Omega$  be a two-dimensional simply connected polygonal domain with boundary  $\Gamma$ . We consider the Stokes equations

$$-\Delta \hat{\mathbf{u}} + \nabla \hat{p} = \mathbf{f} \text{ in } \Omega, \quad (1)$$

$$\nabla \cdot \hat{\mathbf{u}} = 0 \text{ in } \Omega, \quad (2)$$

$$\hat{\mathbf{u}} = \mathbf{g} \text{ on } \Gamma, \quad (3)$$

$$\int_{\Omega} \hat{p}(x) dx = 0, \quad (4)$$

where  $\hat{\mathbf{u}}$  is the fluid velocity,  $\hat{p}$  the pressure,  $\mathbf{f}$  the body forces per unit mass, and the function  $\mathbf{g}$  satisfies  $\int_{\Gamma} \mathbf{g}(\sigma) \cdot \mathbf{n} d\sigma = 0$ . With  $\mathbf{f} \in H^{-1}(\Omega)$  and  $\mathbf{g} \in H^{1/2}(\Gamma)$ , this problem is well-posed (see [15]) due to the so-called

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<sup>1</sup> CEA, DEN, DM2S-STMF, F-91191 Gif-sur-Yvette Cedex, France. e-mail: [anh-ha.le@cea.fr](mailto:anh-ha.le@cea.fr) & [pascal.omnes@cea.fr](mailto:pascal.omnes@cea.fr)

<sup>2</sup> Université Paris 13, LAGA, CNRS UMR 7539, Institut Galilée, 99, Avenue J.-B. Clément F-93430 Villetaneuse Cedex, France.

inf-sup condition: there exists  $\beta > 0$  such that:

$$\beta = \inf_{q \in L_0^2(\Omega)} \sup_{\mathbf{v} \in (H_0^1(\Omega))^2} \frac{\int_{\Omega} q \nabla \cdot \mathbf{v}(\mathbf{x}) dx}{\|\mathbf{v}\|_{(H_0^1(\Omega))^2} \|q\|_{L^2(\Omega)}}. \quad (5)$$

Our purpose in this work is to compute an a posteriori error estimation between the exact solution  $\hat{\mathbf{u}}, \hat{p}$  of (1)–(4) and its numerical approximation by the penalized discrete duality finite volume scheme (DDFV) as presented in [20]. Originally developed for linear diffusion equations [14, 17, 18], the DDFV method has been extended to nonlinear diffusion [2, 4, 7], convection-diffusion [8], electro-cardiology [1, 9], drift-diffusion and energy-transport models [6], electro- and magnetostatics [12], electromagnetism [19], and Stokes flows [11, 13, 20]. The originality of the DDFV method is that it is able to treat all kind of meshes, including very distorted, degenerating, or highly non-conforming meshes (see the numerical tests in [14]). The name of the method comes from the definition of discrete gradient and divergence operators which verify a discrete Green formula.

Like for other equations, the development of a posteriori error estimations for the Stokes problem has followed the a priori investigation of numerical methods. As far as finite elements methods are concerned, R. Verfürth [23] made one the very first contributions by getting two a posteriori error estimations for the mini-element discretization: one is based on a suitable evaluation of the residual, the other is based on the solution of local Stokes problems. Later on, R. Verfürth [24] generalized the first estimator developed in [23] to the non-conforming Crouzeix–Raviart finite element method, neglecting however the consistency error in the estimator. It was shown however in E. Dari et al. [10] that this consistency error may not always be neglected, and, in order to properly take it into account, the authors of [10] use a Helmholtz-Hodge like decomposition (adapted to the Stokes problem) of the velocity error. In the resulting error estimator, this gives rise to terms related to the jumps of the tangential velocity components from one cell to another, in addition to the usual jumps of the normal components of the stress tensor. The case of the non-conforming Fortin–Soulie quadratic elements is also treated in [10].

All the above-cited finite element methods satisfy a uniform discrete inf-sup condition. However, it is often found useful in practice to consider discretizations (especially low-order ones) that do not verify such a uniform discrete inf-sup condition. In this context, C. Bernardi et al. [3] consider the finite element approximation of the Stokes equations when a penalty term is added to stabilize the variational formulation. The a posteriori error estimation they obtain includes two contributions: one related to the discretization on a given mesh, the other related to the penalty term. Based on these two contributions, the mesh refinement and the decrease of the penalty term are linked within an adaptive process.

A very recent contribution by A. Hannukainen et al. [16] sets a general framework for obtaining a posteriori error estimations for the discretization of the Stokes equations. The method is based on the reconstruction of postprocessed  $H_0^1$  conforming velocity and  $H - div$  conforming stress tensor fields deduced from the numerical approximation, and it may be applied to various conforming and conforming stabilized finite element methods, the discontinuous Galerkin method, the Crouzeix–Raviart non-conforming finite element method, the mixed finite element method, and a general class of finite volume methods.

However, as far as finite volume methods are concerned, the use of arbitrary meshes in [16] requires first to solve local Stokes problems on a conforming subtriangulation of each control volume, and then to apply the above-cited reconstruction on this subtriangulation. Instead, we would like to obtain error estimates for the solution of the DDFV scheme presented in [20] without having to solve any local problem or to compute any reconstruction. To do this, we shall adapt to the Stokes problem the a posteriori error estimation investigated in [21] for the DDFV discretization of the Laplace equation, using the discrete variational formulation verified by this scheme. The non-conformity of the method is dealt with using the Helmholtz-Hodge like decomposition introduced in [10]. Our estimator also includes a contribution related to the stabilization term in the incompressibility equation, which allows to monitor the amplitude of the penalization coefficient with respect to the mesh refinement process.

This article is organized as follows. Section 1 sets some notations and definitions related to the meshes, to discrete differential operators, and to discrete functions. In section 2, we present the DDFV scheme, and its

equivalent variational formula is recalled. In section 3, representations of the errors are elaborated. This is used in section 4 to find a computable upper bound of these errors, provided a lower bound for the inf-sup constant in (5) is known. In section 5, the local efficiency of the error estimators is verified. Numerical experiments are presented in section 6.

## 1. NOTATIONS AND DEFINITIONS

Let  $\Omega$  be covered by a primal mesh with polygonal cells denoted by  $T_i$ ,  $i \in [1, I]$ . We associate with each  $T_i$  a point  $G_i$  located in the interior of  $T_i$ . With any vertex  $S_k$  of the primal mesh, with  $k \in [1, K]$ , we associate a dual cell  $P_k$  by joining points  $G_i$  associated with the primal cells surrounding  $S_k$  to the midpoints of the edges of which  $S_k$  is a node. The notations are summarized in Fig. 1 and 2.

With any primal edge  $A_j$  with  $j \in [1, J]$ , we associate a so-called diamond-cell  $D_j$  obtained by joining the vertices  $S_{k_1(j)}$  and  $S_{k_2(j)}$  of  $A_j$  to the points  $G_{i_1(j)}$  and  $G_{i_2(j)}$  associated with the primal cells that share  $A_j$  as a part of their boundaries. When  $A_j$  is a boundary edge (there are  $J^\Gamma$  such edges), the associated diamond-cell is a flat quadrilateral (i.e. a triangle) and we denote by  $G_{i_2(j)}$  the midpoint of  $A_j$  (thus, there are  $J^\Gamma$  such additional points  $G_i$ ). The unit normal vector to  $A_j$  is  $\mathbf{n}_j$  and points from  $G_{i_1(j)}$  to  $G_{i_2(j)}$ . We denote by  $A'_{j1}$  (resp.  $A'_{j2}$ ) the segment joining  $G_{i_1(j)}$  (resp.  $G_{i_2(j)}$ ) and the midpoint of  $A_j$ . Its associated unit normal vector, pointing from  $S_{k_1(j)}$  to  $S_{k_2(j)}$ , is denoted by  $\mathbf{n}'_{j1}$  (resp.  $\mathbf{n}'_{j2}$ ). In the case of a boundary diamond-cell,  $A'_{j2}$  reduces to  $\{G_{i_2(j)}\}$  and does not play any role. Finally, for any diamond-cell  $D_j$ , we shall denote by  $M_{i_\alpha k_\beta}$  the midpoint of  $[G_{i_\alpha(j)} S_{k_\beta(j)}]$ , with  $(\alpha, \beta) \in \{1; 2\}^2$ . With  $\mathbf{n}_j$ ,  $\mathbf{n}'_{j1}$  and  $\mathbf{n}'_{j2}$ , we associate orthogonal unit vectors  $\boldsymbol{\tau}_j$ ,  $\boldsymbol{\tau}'_{j1}$  and  $\boldsymbol{\tau}'_{j2}$ , such that the corresponding orthonormal bases are positively oriented. For any primal  $T_i$  such that  $A_j \subset \partial T_i$ , we shall define  $\mathbf{n}_{ji} := \mathbf{n}_j$  if  $i = i_1(j)$  and  $\mathbf{n}_{ji} := -\mathbf{n}_j$  if  $i = i_2(j)$ , so that  $\mathbf{n}_{ji}$  is always exterior to  $T_i$ . With  $\mathbf{n}_{ji}$ , we associate  $\boldsymbol{\tau}_{ji}$  such that  $(\mathbf{n}_{ji}, \boldsymbol{\tau}_{ji})$  is positively oriented. Similarly, when  $A'_{j1}$  and  $A'_{j2}$  belong to  $\partial P_k$ , we define  $(\mathbf{n}'_{jk1}, \boldsymbol{\tau}'_{jk1})$  and  $(\mathbf{n}'_{jk2}, \boldsymbol{\tau}'_{jk2})$  so that  $\mathbf{n}'_{jk1}$  and  $\mathbf{n}'_{jk2}$  are orthogonal to  $A'_{j1}$  and  $A'_{j2}$  and exterior to  $P_k$ .

**NB:** by a slight abuse of notations, we shall write  $i \in \Gamma$  (respectively  $j \in \Gamma$  and  $k \in \Gamma$ ) to mean  $G_i \in \Gamma$ , (resp.  $A_j \subset \Gamma$  and  $S_k \in \Gamma$ ). The same convention will be used for any set of points other than  $\Gamma$ ; e.g. for  $\partial T_i$ . We shall write  $j \in \partial P_k$  to mean that  $A'_{j1}$  and  $A'_{j2}$  are subsets of  $\partial P_k$ .

For  $\mathbf{v} \in (H^2(\Omega))^2$  with  $\mathbf{v} = (v_1, v_2)^t$ , we define

$$\nabla \mathbf{v} = \begin{pmatrix} \partial v_1 / \partial x & \partial v_1 / \partial y \\ \partial v_2 / \partial x & \partial v_2 / \partial y \end{pmatrix}, \quad \nabla \times \mathbf{v} = \begin{pmatrix} \partial v_1 / \partial y & -\partial v_1 / \partial x \\ \partial v_2 / \partial y & -\partial v_2 / \partial x \end{pmatrix},$$

$$\Delta \mathbf{v} = \begin{pmatrix} \Delta v_1 \\ \Delta v_2 \end{pmatrix}.$$

If  $A$  and  $B$  are two matrices with dimension  $M$ , we define the inner product

$$A : B = \sum_{i,j=1}^M A_{ij} B_{ij}.$$

For future use, we recall Green's formulae

$$\int_{\Omega} \Delta \mathbf{v} \cdot \mathbf{w} dx = - \int_{\Omega} \nabla \mathbf{v} : \nabla \mathbf{w} + \int_{\partial \Omega} (\nabla \mathbf{v} \mathbf{n}) \cdot \mathbf{w} ds, \quad (6)$$

$$\int_{\Omega} \nabla \mathbf{v} : \nabla \times \mathbf{w} dx = - \int_{\partial \Omega} (\nabla \mathbf{v} \boldsymbol{\tau}) \cdot \mathbf{w} ds, \quad (7)$$

for any  $\mathbf{v} \in (H^2(\Omega))^2$  and  $\mathbf{w} \in (H^1(\Omega))^2$ . Here,  $\mathbf{n}$  is the outward normal to  $\partial\Omega$  and  $\boldsymbol{\tau}$  is the tangent vector to  $\partial\Omega$  such that  $(\mathbf{n}, \boldsymbol{\tau})$  is positively oriented.

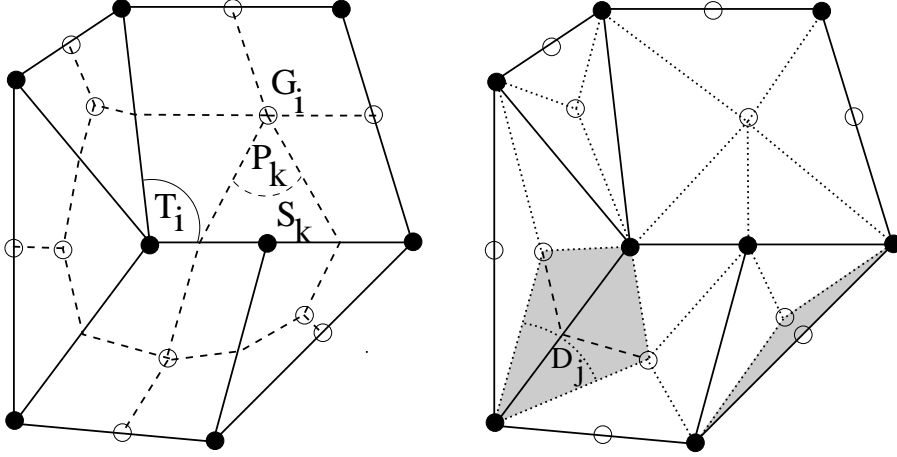


FIGURE 1. A non-conforming primal mesh (solid lines) and its associated dual mesh (left, dashed lines) and diamond mesh (right, dotted lines).

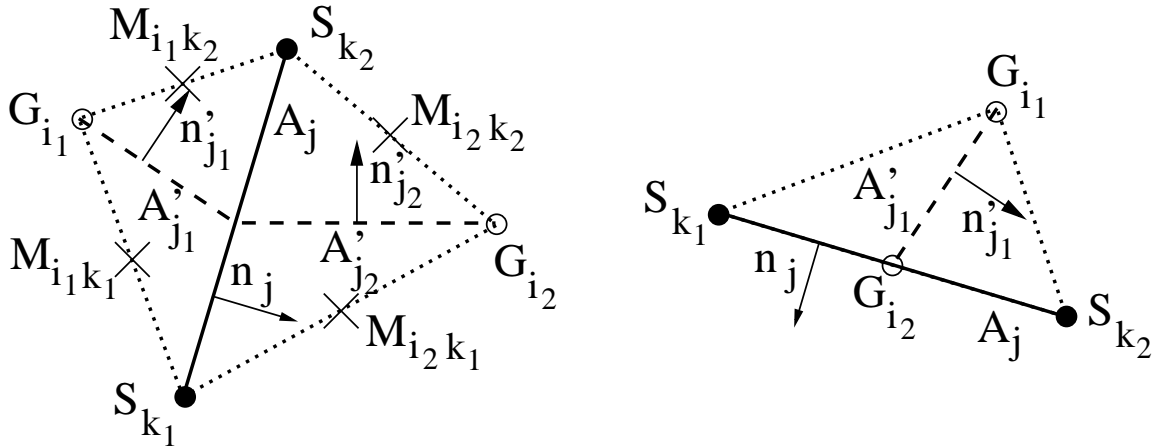


FIGURE 2. Notations for an inner diamond-cell (left) and a boundary diamond cell (right).

In the definition of the DDFV scheme, we shall associate the velocity unknowns to the points  $G_i$  and  $S_k$  and the pressure unknowns to the diamond-cells. Moreover the gradient and divergence of the velocity will be defined on the diamond-cells. This leads us to the following definitions.

**Definition 1.1.** Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$ , and  $\mathbf{v} = (\mathbf{v}_i^T, \mathbf{v}_k^P)$  be in  $(\mathbb{R}^2)^I \times (\mathbb{R}^2)^K$ . Let  $\Phi = (\Phi_j)$  and  $\Psi = (\Psi_j)$  be in  $(\mathbb{R}^{2 \times 2})^J$ . Let  $p = (p_j)$  and  $q = (q_j)$  be in  $\mathbb{R}^J$ . We define the following scalar products

$$(\mathbf{u}, \mathbf{v})_{T,P} := \frac{1}{2} \left( \sum_{i \in [1, I]} |T_i| \mathbf{u}_i^T \cdot \mathbf{v}_i^T + \sum_{k \in [1, K]} |P_k| \mathbf{u}_k^P \cdot \mathbf{v}_k^P \right), \quad (8)$$

$$(\Phi, \Psi)_D := \sum_{j \in [1, J]} |D_j| \Phi_j : \Psi_j \quad \text{and} \quad (p, q)_D = \sum_{j \in [1, J]} |D_j| p_j q_j. \quad (9)$$

**Definition 1.2.** Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  be in  $(\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$ . For any boundary edge  $A_j$ , with the notations of Fig. 2 (right), we define  $\tilde{\mathbf{u}}_j$  as the trace of  $\mathbf{u}$  over  $A_j$  by

$$\tilde{\mathbf{u}}_j = \frac{1}{4} \left( \mathbf{u}_{k_1(j)}^P + 2\mathbf{u}_{i_2(j)}^T + \mathbf{u}_{k_2(j)}^P \right). \quad (10)$$

Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  be in  $(\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$  and let  $\mathbf{w} = (\mathbf{w}_j)$  be defined on the boundary  $\Gamma$ . We define the following boundary scalar product

$$(\mathbf{w}, \tilde{\mathbf{u}})_{\Gamma_h} := \sum_{j \in \Gamma} |A_j| \mathbf{w}_j \cdot \tilde{\mathbf{u}}_j. \quad (11)$$

**Definition 1.3.** Let  $\Phi = (\Phi_j)$  be in  $(\mathbb{R}^{2 \times 2})^J$ . We define divergences and curls of the tensor field  $\Phi$  on the primal and dual cells by

$$\begin{aligned} (\nabla_h^T \cdot \Phi)_i &:= \frac{1}{|T_i|} \sum_{j \in \partial T_i} |A_j| \Phi_j \mathbf{n}_{ji}, \\ (\nabla_h^P \cdot \Phi)_k &:= \frac{1}{|P_k|} \left( \sum_{j \in \partial P_k} (|A'_{j_1}| \Phi_j \mathbf{n}'_{j_1} + |A'_{j_2}| \Phi_j \mathbf{n}'_{j_2}) + \sum_{j \in \partial P_k \cap \Gamma} \frac{|A_j|}{2} \Phi_j \mathbf{n}_j \right), \\ (\nabla_h^T \times \Phi)_i &:= \frac{1}{|T_i|} \sum_{j \in \partial T_i} |A_j| \Phi_j \boldsymbol{\tau}_{ji}, \\ (\nabla_h^P \times \Phi)_k &:= \frac{1}{|P_k|} \left( \sum_{j \in \partial P_k} (|A'_{j_1}| \Phi_j \boldsymbol{\tau}'_{j_1} + |A'_{j_2}| \Phi_j \boldsymbol{\tau}'_{j_2}) + \sum_{j \in \partial P_k \cap \Gamma} \frac{|A_j|}{2} \Phi_j \boldsymbol{\tau}_j \right). \end{aligned}$$

**Definition 1.4.** Let  $u_1 = ((u_1)_i^T, (u_1)_k^P)$  be in  $\mathbb{R}^{I+K+J^\Gamma}$  and  $u_2 = ((u_2)_i^T, (u_2)_k^P)$  be in  $\mathbb{R}^{I+K+J^\Gamma}$ , and  $\mathbf{u} = (u_1, u_2)$ ; the discrete gradient  $\nabla_h^D \mathbf{u}$  and the discrete curl  $\nabla_h^D \times \mathbf{u}$  are defined by their values in the diamond-cells  $D_j$  by

$$(\nabla_h^D \mathbf{u})_j = \begin{pmatrix} (\nabla_h^D u_1)_j^t \\ (\nabla_h^D u_2)_j^t \end{pmatrix}, \quad (\nabla_h^D \times \mathbf{u})_j = \begin{pmatrix} (\nabla_h^D \times u_1)_j^t \\ (\nabla_h^D \times u_2)_j^t \end{pmatrix},$$

where, for  $\phi \in \mathbb{R}^{I+K+J^\Gamma}$ , we define

$$\begin{aligned} (\nabla_h^D \phi)_j &:= \frac{1}{2|D_j|} \left\{ [\phi_{k_2(j)}^P - \phi_{k_1(j)}^P] (|A'_{j_1}| \mathbf{n}'_{j_1} + |A'_{j_2}| \mathbf{n}'_{j_2}) + [\phi_{i_2(j)}^T - \phi_{i_1(j)}^T] |A_j| \mathbf{n}_j \right\}, \\ (\nabla_h^D \times \phi)_j &:= -\frac{1}{2|D_j|} \left\{ [\phi_{k_2(j)}^P - \phi_{k_1(j)}^P] (|A'_{j_1}| \boldsymbol{\tau}'_{j_1} + |A'_{j_2}| \boldsymbol{\tau}'_{j_2}) + [\phi_{i_2(j)}^T - \phi_{i_1(j)}^T] |A_j| \boldsymbol{\tau}_j \right\}. \end{aligned}$$

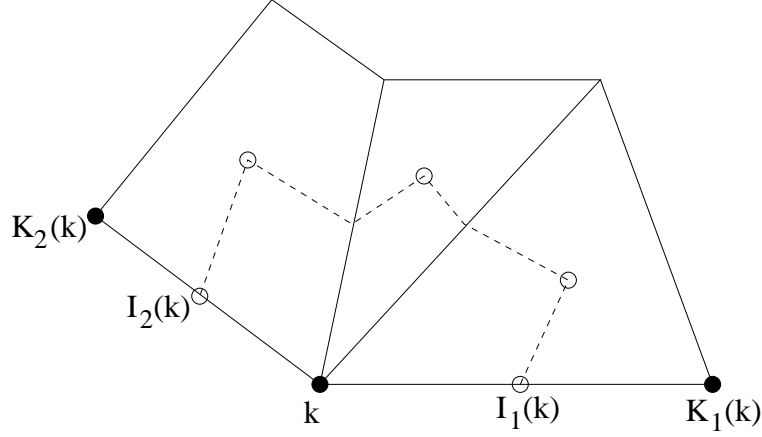


FIGURE 3. Notation for a boundary dual cell in formula (19).

We also need a discrete divergence of a vector, which is defined using the discrete gradient

$$(\nabla_h^D \cdot \mathbf{u})_j = \text{Trace}((\nabla_h^D \mathbf{u})_j).$$

From basic geometrical arguments, we obtain some properties of the discrete gradient:

$$\phi_{k_2(j)}^P - \phi_{k_1(j)}^P = (\nabla_h^D \phi)_j \cdot \overrightarrow{S_{k_1(j)} S_{k_2(j)}} \quad , \quad \phi_{i_2(j)}^T - \phi_{i_1(j)}^T = (\nabla_h^D \phi)_j \cdot \overrightarrow{G_{i_1(j)} G_{i_2(j)}}. \quad (12)$$

For the penalization of the scheme, we need to define the following (non-consistent) Laplacian-type operator.

**Definition 1.5.** Let  $p = (p_j) \in \mathbb{R}^J$ , we define:

$$(\Delta_h^D p)_j = \frac{1}{|D_j|} \sum_{j' \in \partial D_j} (p_{j'} - p_j), \quad (13)$$

where  $\partial D_j$  is the set of indexes of diamond cells which have a common boundary segment with  $D_j$ .

**Proposition 1.6.** For  $\Phi \in (\mathbb{R}^{2 \times 2})^J$  and  $\mathbf{u} = (\mathbf{u}^T, \mathbf{u}^P) \in (\mathbb{R}^2)^{I+J^T} \times (\mathbb{R}^2)^K$  and  $p \in \mathbb{R}^J$ , the following discrete Green formula hold:

$$(\nabla_h^{T,P} \cdot \Phi, \mathbf{u})_{T,P} = -(\nabla_h^D \mathbf{u}, \Phi)_D + (\Phi \mathbf{n}, \tilde{\mathbf{u}})_{\Gamma,h}, \quad (14)$$

$$(\nabla_h^{T,P} \times \Phi, \mathbf{u})_{T,P} = (\nabla_h^D \times \mathbf{u}, \Phi)_D + (\Phi \boldsymbol{\tau}, \tilde{\mathbf{u}})_{\Gamma,h}, \quad (15)$$

$$(\nabla_h^{T,P} \cdot p I_2, \mathbf{u})_{T,P} = -(\nabla_h^D \cdot \mathbf{u}, p)_D + (p \mathbf{n}, \tilde{\mathbf{u}})_{\Gamma,h}, \quad (16)$$

where  $I_2$  is the  $2 \times 2$  identity matrix.

Formula (14) is called the discrete Stokes formula and is proved in [20]; its componentwise counterpart can be found in [14]. Formulae (15) and (16) can be demonstrated in the same way. Proposition 1.7 below may be found componentwise in [12].

**Proposition 1.7.** For all  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P) \in (\mathbb{R}^2)^{I+J^T} \times (\mathbb{R}^2)^K$ , there holds

$$(\nabla_h^T \times (\nabla_h^D \mathbf{u}))_i = 0 \quad , \quad \forall i \in [1, I], \quad (17)$$

$$(\nabla_h^P \times (\nabla_h^D \mathbf{u}))_k = 0 \quad , \quad \forall k \notin \Gamma. \quad (18)$$

In addition, for  $k \in \Gamma$ , the following equality holds (see Fig. 3 for the notations)

$$(\nabla_h^P \times (\nabla_h^D \mathbf{u}))_k = \frac{1}{|P_k|} \left[ (\mathbf{u}_{I_2(k)}^T - \mathbf{u}_{I_1(k)}^T) + \frac{1}{2} (\mathbf{u}_{K_1(k)}^P - \mathbf{u}_{K_2(k)}^P) \right]. \quad (19)$$

The derivation of the a posteriori error estimates is based on the reformulation of the scheme under a variational form which uses functions associated to the discrete unknowns.

**Definition 1.8.** With any  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P) \in (\mathbb{R}^2)^{I+J^T} \times (\mathbb{R}^2)^K$ , we associate the function  $\mathbf{u}_h$  defined by

$$\begin{aligned} (\mathbf{u}_h)|_{D_j} &\in (P^1(D_j))^2, \quad \forall j \in [1, J], \\ \mathbf{u}_h(M_{i_\alpha(j)k_\beta(j)}) &= \frac{1}{2} (\mathbf{u}_{i_\alpha(j)}^T + \mathbf{u}_{k_\beta(j)}^P), \quad \forall j \in [1, J], (\alpha, \beta) \in \{1, 2\}^2. \end{aligned}$$

In addition, for all  $p = (p_j) \in \mathbb{R}^J$ , we construct piecewise constant functions corresponding to the approximate pressure and to the penalty term:

$$p_h(\mathbf{x}) = p_j, \quad \forall \mathbf{x} \in D_j, j \in [1, J], \quad (20)$$

$$(\Delta_h^D p)_h(\mathbf{x}) = (\Delta_h^D p)_j, \quad \forall \mathbf{x} \in D_j, j \in [1, J]. \quad (21)$$

The validity of the definition of  $\mathbf{u}_h$  by its values in four different points and the proof of the fundamental properties below, which allow to reformulate the scheme under a variational form, may be found in [14].

**Proposition 1.9.** Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P) \in (\mathbb{R}^2)^{I+J^T} \times (\mathbb{R}^2)^K$  and let  $\mathbf{u}_h$  be defined by Definition 1.8. There holds

$$(\nabla_h^D \mathbf{u})_j = \nabla \mathbf{u}_h|_{D_j}, \quad \forall j \in [1, J], \quad (22)$$

$$(\nabla_h^D \cdot \mathbf{u})_j = \nabla \cdot \mathbf{u}_h|_{D_j}, \quad \forall j \in [1, J]. \quad (23)$$

**Definition 1.10.** Let  $\mathbf{u}_h$  be in  $(P^1(D_j))^2$ ,  $\forall j \in [1, J]$ , and not necessarily continuous over the interfaces of neighboring diamond-cells. We define its piecewise gradient and divergence over  $\Omega$  by:

$$\nabla_h \mathbf{u}_h(\mathbf{x}) = \nabla \mathbf{u}_h|_{D_j}(\mathbf{x}) \text{ and } \nabla_h \cdot \mathbf{u}_h(\mathbf{x}) = \nabla \cdot \mathbf{u}_h|_{D_j}(\mathbf{x}), \quad \forall \mathbf{x} \in D_j, j \in [1, J]. \quad (24)$$

## 2. THE FINITE VOLUME SCHEME ON GENERAL MESHES

The finite volume scheme used for the numerical approximation of equations (1)–(4) is constructed on the basis of the discrete differential operators defined in section 1. It is very similar to that studied in [20], only the definition of the penalization operator being slightly changed, with no alteration in the existence and uniqueness results that can be proved like in [20], as soon as  $\varepsilon > 0$ .

$$(\nabla_h^T \cdot (-\nabla_h^D \mathbf{u} + pI_2))_i = \mathbf{f}_i^T, \quad \forall i \in [1, I], \quad (25)$$

$$(\nabla_h^P \cdot (-\nabla_h^D \mathbf{u} + pI_2))_k = \mathbf{f}_k^P, \quad \forall k \in [1, K], \quad (26)$$

$$(\nabla_h^D \cdot \mathbf{u})_j - \varepsilon (\Delta_h^D p)_j = 0, \quad \forall j \in [1, J], \quad (27)$$

$$\frac{\mathbf{u}_{k_1(j)}^P + 2\mathbf{u}_{i_2(j)}^T + \mathbf{u}_{k_2(j)}^P}{4} = \mathbf{g}_j, \quad \forall j \in \Gamma, \quad (28)$$

$$\sum_{j=1}^J |D_j| p_j = 0. \quad (29)$$



We suppose that  $\mathbf{g}$  is regular enough, so that we can set  $\mathbf{g}_j = \mathbf{g}(G_{i_2(j)})$  in (28), while in (25) and (26),  $\mathbf{f}_i^T$  and  $\mathbf{f}_k^P$  are the mean values of  $\mathbf{f}$  over  $T_i$  and  $P_k$ , respectively:

$$\mathbf{f}_i^T = \frac{1}{|T_i|} \int_{T_i} \mathbf{f}(x) dx \quad \text{and} \quad \mathbf{f}_k^P = \frac{1}{|P_k|} \int_{P_k} \mathbf{f}(x) dx. \quad (30)$$

We can prove that the solution of the scheme verifies a discrete variational formulation:

**Proposition 2.1.** *Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  and  $p = (p_j)_{j \in [1, J]}$  be the solution of the scheme (25)-(29). Let  $\mathbf{v} = (\mathbf{v}_i^T, \mathbf{v}_k^P)$  such that  $\tilde{\mathbf{v}}_j = 0$  for all  $j \in \Gamma$ . Let  $\mathbf{u}_h$  and  $\mathbf{v}_h$  be the solution associated to  $\mathbf{u}$  and  $\mathbf{v}$  by Definition 1.8. Let us set in addition*

$$\mathbf{v}_h^*(x) := \frac{1}{2} \left( \sum_{i \in [1, I]} \mathbf{v}_i^T \theta_i^T(x) + \sum_{k \in [1, K]} \mathbf{v}_k^P \theta_k^P(x) \right), \quad (31)$$

where  $\theta_i^T$  and  $\theta_k^P$  are respectively the characteristic function of the cells  $T_i$  and  $P_k$ . Then, there holds

$$\sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla_h \mathbf{v}_h(x) dx - \sum_j \int_{D_j} \nabla_h \cdot \mathbf{v}_h p_h(x) dx = \int_{\Omega} \mathbf{f} \cdot \mathbf{v}_h^*(x) dx. \quad (32)$$

*Proof.* Starting from Eq. (25) and (26), we have

$$-(\nabla_h^T \cdot (\nabla_h^D \mathbf{u}))_i \cdot \mathbf{v}_i^T + (\nabla_h^T \cdot (pI_2))_i \cdot \mathbf{v}_i^T = \mathbf{f}_i^T \cdot \mathbf{v}_i^T, \quad \forall i \in [1, I], \quad (33)$$

$$-(\nabla_h^P \cdot (\nabla_h^D \mathbf{u}))_i \cdot \mathbf{v}_k^P + (\nabla_h^P \cdot (pI_2))_k \cdot \mathbf{v}_k^P = \mathbf{f}_k^P \cdot \mathbf{v}_k^P, \quad \forall k \in [1, K]. \quad (34)$$

Multiplying (33) by  $|T_i|$  and (34) by  $|P_k|$  and summing over all  $i$  and all  $k$ , we obtain

$$-(\nabla_h^{T,P} \cdot (\nabla_h^D \mathbf{u}), \mathbf{v})_{T,P} + (\nabla_h^{T,P} \cdot (pI_2), \mathbf{v})_{T,P} = (\mathbf{f}, \mathbf{v})_{T,P}.$$

We can apply (14), (16) and  $\tilde{\mathbf{v}}_j = 0$  for all  $j \in \Gamma$  and we obtain

$$(\nabla_h^D \mathbf{u}, \nabla_h^D \mathbf{v})_D - (\nabla_h^D \cdot \mathbf{v}, p)_D = (\mathbf{f}, \mathbf{v})_{T,P}.$$

Using (22)–(23) and definitions (30)–(31), we obtain (32).  $\square$

### 3. A REPRESENTATION OF THE ERROR

#### 3.1. A representation of the velocity error

The variational formula of (1) reads:

$$\int_{\Omega} \nabla \hat{\mathbf{u}} : \nabla \mathbf{v} dx - \int_{\Omega} \hat{p} \nabla \cdot \mathbf{v} dx = \int_{\Omega} \mathbf{f} \cdot \mathbf{v}(x) dx, \quad (35)$$

for all  $\mathbf{v} \in (H_0^1(\Omega))^2$ . We shall estimate the  $H^1$  semi norm of the error between the exact solution  $\hat{\mathbf{u}}$  and the function  $\mathbf{u}_h$  associated to the solution of the DDFV scheme. For this, we shall denote by  $\mathbf{e} := \hat{\mathbf{u}} - \mathbf{u}_h$  and  $e_p := \hat{p} - p_h$  the error in the velocity and pressure, respectively. We have

$$\|\nabla_h \mathbf{e}\|_{L^2(\Omega)} = \left( \sum_j \int_{D_j} |\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h|^2(x) dx \right)^{1/2}. \quad (36)$$

Since  $\Omega$  is a simply connected domain and since  $\nabla_h \mathbf{e} = \nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h$  belongs to  $(L^2(\Omega))^{2 \times 2}$ , we may decompose it in the following way (see Lemma 3.2 in [10]):

$$\nabla_h \mathbf{e} = \nabla \hat{\Phi} - qI_2 + \nabla \times \hat{\Psi}, \quad (37)$$

where  $q \in L_0^2(\Omega)$ ,  $\hat{\Phi} \in (H_0^1(\Omega))^2$  with  $\nabla \cdot \hat{\Phi} = 0$  and  $\hat{\Psi} \in (H^1(\Omega))^2$  with  $\int_{\Omega} \hat{\Psi}(\mathbf{x}) d\mathbf{x} = 0$ , with the following estimations

$$\begin{aligned} \|\nabla \hat{\Phi}\|_{L^2(\Omega)} &\leq \|\nabla_h e\|_{L^2(\Omega)}, \\ \|q\|_{L^2(\Omega)} &\leq \frac{2}{\beta} \|\nabla_h e\|_{L^2(\Omega)}, \\ \|\nabla \times \hat{\Psi}\|_{L^2(\Omega)} &\leq \left(1 + \frac{2\sqrt{2}}{\beta}\right) \|\nabla_h e\|_{L^2(\Omega)}, \end{aligned} \quad (38)$$

where  $\beta$  is defined by (5).

Now, we estimate the velocity error using decomposition (37). First observe that

$$\int_{\Omega} \nabla_h \mathbf{e} : I_2 q(\mathbf{x}) d\mathbf{x} = \int_{\Omega} \nabla_h \cdot \mathbf{e} q(\mathbf{x}) d\mathbf{x} = \int_{\Omega} (\nabla \cdot \hat{\mathbf{u}} - \nabla_h \cdot \mathbf{u}_h) q(\mathbf{x}) d\mathbf{x}.$$

From (2) and (27), we have

$$\int_{\Omega} \nabla_h \mathbf{e} : I_2 q(\mathbf{x}) d\mathbf{x} = \varepsilon \int_{\Omega} (\Delta_h^D p)_h q(\mathbf{x}) d\mathbf{x}. \quad (39)$$

Multiplying the term  $\nabla_h \mathbf{e}(\mathbf{x})$  with (37) side by side and integrating over  $\Omega$ , there holds

$$\begin{aligned} \|\nabla_h \mathbf{e}\|_{L^2(\Omega)}^2 &= \int_{\Omega} \nabla_h \mathbf{e} : (\nabla \hat{\Phi} + \nabla \times \hat{\Psi} - qI_2) d\mathbf{x} \\ &= i_1 + i_2 - \varepsilon \int_{\Omega} (\Delta_h^D p)_h q(\mathbf{x}) d\mathbf{x}, \end{aligned} \quad (40)$$

where

$$i_1 = \sum_j \int_{D_j} (\nabla \hat{\mathbf{u}} - \nabla \mathbf{u}_h) : \nabla \hat{\Phi}(\mathbf{x}) d\mathbf{x}$$

and

$$i_2 = \sum_j \int_{D_j} (\nabla \hat{\mathbf{u}} - \nabla \mathbf{u}_h) : \nabla \times \hat{\Psi}(\mathbf{x}) d\mathbf{x}.$$

In order to find a suitable representation of  $i_1$  and  $i_2$ , we shall need the following definitions

**Definition 3.1.** The boundary  $\partial D_j$  of any diamond-cell  $D_j$  is composed of the four segments  $[G_{i\alpha(j)} S_{k\beta(j)}]$  with  $(\beta, \alpha) \in \{1, 2\}$  (see Fig. 2). Let us define by  $S$  the set of these edges when  $j$  runs over the whole set of diamond-cells and  $\overset{\circ}{S}$  those edges that do not lie in the boundary  $\Gamma$ . Each  $s \in \overset{\circ}{S}$  is thus a segment that we shall denote by  $[G_{i(s)} S_{k(s)}]$ . We shall also write  $s \in \overset{\circ}{T}_i$  (resp.  $s \in \overset{\circ}{P}_k$ ) if  $s \subset T_i$  (resp.  $s \subset P_k$ ) and  $s \not\subset \Gamma$ . Finally, we shall denote by  $\mathbf{n}_s$  one of the two unit normal vectors to  $s$ , arbitrarily chosen among the two possible choices but then fixed for the sequel, and  $[(\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s$  is the jump of the normal component of  $\nabla \mathbf{u}_h - p_h I_2$  through segment  $s$ . Moreover,  $\boldsymbol{\tau}_s$  will be such that  $(\mathbf{n}_s, \boldsymbol{\tau}_s)$  is a positively oriented orthonormal basis of  $\mathbb{R}^2$  and  $[(\nabla \mathbf{u}_h) \boldsymbol{\tau}_s]_s$  is the jump of the tangential component of  $\nabla \mathbf{u}_h$  through segment  $s$ .

**Proposition 3.2.** Let  $\widehat{\Phi}$  be defined in equation (37). Let  $\Phi = (\Phi_i^T, \Phi_k^P) \in (\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$  be such that

$$\tilde{\Phi}_j = 0 \text{ for all } j \in \Gamma. \quad (41)$$

Then, it holds that

$$\begin{aligned} i_1 &= \frac{1}{2} \sum_{i \in [1, I]} \int_{T_i} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_i^T)(\mathbf{x}) d\mathbf{x} + \frac{1}{2} \sum_{k \in [1, K]} \int_{P_k} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_k^P)(\mathbf{x}) d\mathbf{x} \\ &\quad - \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset \overset{\circ}{T}_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_i^T)(\sigma) d\sigma \\ &\quad - \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset \overset{\circ}{P}_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_k^P)(\sigma) d\sigma. \end{aligned} \quad (42)$$

*Proof.* First, since  $\widehat{\Phi} \in (H_0^1(\Omega))^2$  and  $\nabla \cdot \widehat{\Phi} = 0$ , using (35) yields:

$$\begin{aligned} i_1 &= \sum_j \int_{D_j} \nabla \widehat{\mathbf{u}} : \nabla \widehat{\Phi}(\mathbf{x}) d\mathbf{x} - \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla \widehat{\Phi}(\mathbf{x}) d\mathbf{x} \\ &= \int_{\Omega} \nabla \widehat{\mathbf{u}} : \nabla \widehat{\Phi}(\mathbf{x}) d\mathbf{x} - \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla \widehat{\Phi}(\mathbf{x}) d\mathbf{x} \\ &= \int_{\Omega} \mathbf{f} \cdot \widehat{\Phi}(\mathbf{x}) d\mathbf{x} - \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla \widehat{\Phi}(\mathbf{x}) d\mathbf{x}. \end{aligned}$$

For any  $\Phi = (\Phi_i^T, \Phi_k^P)$  satisfying (41), formula (32) leads to

$$\begin{aligned} i_1 &= \int_{\Omega} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_h^*)(\mathbf{x}) d\mathbf{x} - \sum_j \int_{D_j} p_h \nabla_h \cdot \Phi_h(\mathbf{x}) d\mathbf{x} \\ &\quad - \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : (\nabla \widehat{\Phi} - \nabla_h \Phi_h)(\mathbf{x}) d\mathbf{x}. \end{aligned}$$

We know that  $p_h \nabla_h \cdot \Phi_h = p_h I_2 : \nabla_h \Phi_h$  and  $p_h I_2 : \nabla \widehat{\Phi} = 0$  (since  $\nabla \cdot \widehat{\Phi} = 0$ ), then

$$i_1 = \int_{\Omega} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_h^*)(\mathbf{x}) d\mathbf{x} - \sum_j H_1(j), \quad (43)$$

where

$$H_1(j) = \int_{D_j} (\nabla_h \mathbf{u}_h - p_h I_2) : (\nabla \widehat{\Phi} - \nabla_h \Phi_h)(\mathbf{x}) d\mathbf{x}. \quad (44)$$

Let us consider a diamond-cell  $D_j$ . Since  $\nabla_h \mathbf{u}_h - p_h I_2$  is constant over  $D_j$ , we may write, using Green's formula over  $D_j$ ,

$$H_1(j) = \int_{\partial D_j} (\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_{\partial D_j} \cdot (\widehat{\Phi} - \Phi_h)(\sigma) d\sigma,$$

where  $\mathbf{n}_{\partial D_j}$  is the unit normal vector exterior to  $D_j$  on its boundary. Moreover, let  $s$  be any of the four boundary edges of  $D_j$ , the function  $\Phi_h$  belongs to  $P^1$  over  $s$  and the quantity  $\nabla \mathbf{u}_h - p_h I_2$  is a constant; the integral of  $(\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_{\partial D_j} \cdot \Phi_h$  along this edge may thus exactly be computed by the midpoint rule; using the definition of  $\Phi_h$ , this function equals  $\frac{1}{2}(\Phi_{i(s)}^T + \Phi_{k(s)}^T)$  at the midpoint of  $s$ . Thus, there holds:

$$H_1(j) = \sum_{s \subset \partial D_j} \int_s (\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_{s,j} \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma) d\sigma, \quad (45)$$

where  $\mathbf{n}_{s,j}$  is the unit normal vector exterior to  $D_j$  on  $s$ .

In (43), in the sum of the  $H_1(j)$  over  $j \in [1, J]$ , there are two types of edges  $s$ : those in  $\overset{\circ}{S}$  and those included in  $\Gamma$ . First, each  $s \in \overset{\circ}{S}$  is the common edge of two diamond-cells; then, in the sum, there are two corresponding integrals over  $s$ , in which we can factorize by  $\left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma)$ . Indeed, the jump of this function through  $s$  vanishes because  $\widehat{\Phi} \in (H_0^1(\Omega))^2$ . This implies:

$$\begin{aligned} \sum_j \sum_{\substack{s \subset \partial D_j \\ s \not\subset \Gamma}} \int_s (\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_{s,j} \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma) d\sigma &= \\ \sum_{s \in \overset{\circ}{S}} \int_s [(\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_s] \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma) d\sigma. & \end{aligned} \quad (46)$$

Secondly, each diamond-cell  $D_j$  whose boundary intersects  $\Gamma$  has two edges of equal length  $s = [G_{i_2(j)} S_{k_\beta(j)}]$  with  $\beta \in \{1, 2\}$  which are included in  $\Gamma$ , and their union is exactly  $A_j$ . Since  $(\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_j$  is a constant on  $A_j$ , and since  $\sum_{\beta \in \{1, 2\}} \int_{[G_{i_2(j)} S_{k_\beta(j)}]} (\widehat{\Phi} - \frac{1}{2} (\Phi_{i_2(j)}^T + \Phi_{k_\beta(j)}^P)) (\sigma) d\sigma = \int_{A_j} (\widehat{\Phi} - \tilde{\Phi}_h) (\sigma) d\sigma$ , we have

$$\begin{aligned} & \sum_{s \in \partial D_j \cap \Gamma} \int_s (\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_{s,j} \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma) d\sigma \\ &= \sum_{\beta \in \{1, 2\}} \int_{[G_{i_2(j)} S_{k_\beta(j)}]} (\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_j \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i_2(j)}^T + \Phi_{k_\beta(j)}^P) \right] (\sigma) d\sigma \\ &= \int_{A_j} (\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_j \cdot (\widehat{\Phi} - \tilde{\Phi}_h) (\sigma) d\sigma = 0, \end{aligned} \quad (47)$$

thanks to (41) and to the fact that  $\widehat{\Phi} \in (H_0^1(\Omega))^2$ . With (46) and (47), we can write

$$\sum_{j \in [1, J]} H_1(j) = \sum_{s \in \overset{\circ}{S}} \int_s [(\nabla \mathbf{u}_h - p_h I_2) \mathbf{n}_s] \cdot \left[ \widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) \right] (\sigma) d\sigma. \quad (48)$$

Then, we may write  $\widehat{\Phi} - \frac{1}{2} (\Phi_{i(s)}^T + \Phi_{k(s)}^P) = \frac{1}{2} \left[ (\widehat{\Phi} - \Phi_{i(s)}^T) + (\widehat{\Phi} - \Phi_{k(s)}^P) \right]$ . Summing in the right-hand side of (48) the various contributions of  $\Phi_i^T$  for a fixed  $i$  and the various contributions of  $\Phi_k^P$  for a fixed  $k$ , we obtain

the following formula

$$\begin{aligned} \sum_j H_1(j) &= \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset \overset{\circ}{T}_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_i^T)(\sigma) d\sigma \\ &\quad + \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset \overset{\circ}{P}_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_k^P)(\sigma) d\sigma. \end{aligned} \quad (49)$$

Finally, according to (43) and definition (31) of  $\Phi_h^*$ , we obtain (42).  $\square$

Before we turn to a representation formula for  $i_2$  in (40), we need some technical lemmas related to the  $L^2(\Omega)$  scalar product of discrete gradients and curls.

**Lemma 3.3.** *Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  be the velocity component of the solution of the scheme (25)–(29) and  $\Psi = (\Psi_i^T, \Psi_k^P) \in (\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$ . There holds*

$$(\nabla_h^{T,P} \times (\nabla_h^D \mathbf{u}), \Psi)_{T,P} = - \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} (\nabla \mathbf{g}(\sigma) - \nabla_h \mathbf{u}_h(\sigma)) \boldsymbol{\tau}_k \cdot \Psi_k^P d\sigma, \quad (50)$$

where  $\boldsymbol{\tau}_k$  is the tangent vector to  $\partial P_k \cap \Gamma$  which is positively oriented with respect to the unit normal vector exterior to  $\partial P_k \cap \Gamma$ .

*Proof.* According to Eq. (17) and (18), there holds

$$(\nabla_h^T \times (\nabla_h^D \mathbf{u}))_i = 0, \quad \forall i \in [1, I] \quad \text{and} \quad (\nabla_h^P \times (\nabla_h^D \mathbf{u}))_k = 0, \quad \forall k \notin \Gamma. \quad (51)$$

On the other hand, since the solution of the discrete problem verifies (28), there holds, for  $k \in \Gamma$ , with the notations of Fig. 3:

$$\mathbf{u}_{I_1(k)}^T = 2\mathbf{g}(G_{I_1(k)}) - \frac{1}{2}(\mathbf{u}_k^P + \mathbf{u}_{K_1(k)}^P) \quad \text{and} \quad \mathbf{u}_{I_2(k)}^T = 2\mathbf{g}(G_{I_2(k)}) - \frac{1}{2}(\mathbf{u}_k^P + \mathbf{u}_{K_2(k)}^P). \quad (52)$$

Following (19) and (52), we obtain that

$$(\nabla_h^P \times (\nabla_h^D \mathbf{u}))_k = \frac{1}{|P_k|} \left[ 2(\mathbf{g}(G_{I_2(k)}) - \mathbf{g}(G_{I_1(k)})) + (\mathbf{u}_{K_1(k)}^P - \mathbf{u}_{K_2(k)}^P) \right], \quad \forall k \in \Gamma. \quad (53)$$

From (51), and using the definition of the scalar product in (8), we obtain

$$(\nabla_h^{T,P} \times (\nabla_h^D \mathbf{u}), \Psi)_{T,P} = \frac{1}{2} \sum_{k \in \Gamma} |P_k| (\nabla_h^P \times (\nabla_h^D \mathbf{u}))_k \cdot \Psi_k^P.$$

Using (53) leads to

$$(\nabla_h^{T,P} \times (\nabla_h^D \mathbf{u}), \Psi)_{T,P} = \sum_{k \in \Gamma} (\mathbf{g}(G_{I_2(k)}) - \mathbf{g}(G_{I_1(k)})) \cdot \Psi_k^P + \frac{1}{2} \sum_{k \in \Gamma} (\mathbf{u}_{K_1(k)}^P - \mathbf{u}_{K_2(k)}^P) \cdot \Psi_k^P. \quad (54)$$

In addition, we have

$$\mathbf{g}(G_{I_2(k)}) - \mathbf{g}(G_{I_1(k)}) = (\mathbf{g}(G_{I_2(k)}) - \mathbf{g}(S_k)) + (\mathbf{g}(S_k) - \mathbf{g}(G_{I_1(k)})),$$

so that

$$\mathbf{g}(G_{I_2(k)}) - \mathbf{g}(G_{I_1(k)}) = - \int_{\partial P_k \cap \Gamma} \nabla \mathbf{g}(\sigma) \boldsymbol{\tau}_k d\sigma, \quad (55)$$

In the same way, we have

$$\mathbf{u}_{K_1(k)}^P - \mathbf{u}_{K_2(k)}^P = (\mathbf{u}_{K_1(k)}^P - \mathbf{u}_k^P) + (\mathbf{u}_k^P - \mathbf{u}_{K_2(k)}^P).$$

Applying property (12) of the discrete gradient, there holds

$$\mathbf{u}_{K_1(k)}^P - \mathbf{u}_k^P = |S_k S_{K_1(k)}| \nabla_h \mathbf{u}_h(\sigma) \boldsymbol{\tau}_k = 2 |S_k G_{I_1(k)}| \nabla_h \mathbf{u}_h(\sigma) \boldsymbol{\tau}_k, \quad \forall \sigma \in [S_k S_{K_1(k)}] \quad (56)$$

and

$$\mathbf{u}_k^P - \mathbf{u}_{K_2(k)}^P = |S_{K_2(k)} S_k| \nabla_h \mathbf{u}_h(\sigma) \boldsymbol{\tau}_k = 2 |S_k G_{I_2(k)}| \nabla_h \mathbf{u}_h(\sigma) \boldsymbol{\tau}_k, \quad \forall \sigma \in [S_{K_2(k)} S_k]. \quad (57)$$

As a consequence, there holds

$$\mathbf{u}_{K_1(k)}^P - \mathbf{u}_{K_2(k)}^P = 2 \int_{\partial P_k \cap \Gamma} \nabla_h \mathbf{u}_h(\sigma) \boldsymbol{\tau}_k d\sigma. \quad (58)$$

Combining (54) with (55) and (58), we obtain (50).  $\square$

**Lemma 3.4.** *Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  be the velocity component of the solution of the scheme (25)–(29) and  $\Psi = (\Psi_i^T, \Psi_k^P) \in (\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$ . Let  $\mathbf{u}_h$  and  $\Psi_h$  be their associated functions through Def. 1.8. There holds*

$$\sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla_h \times \Psi_h(\mathbf{x}) d\mathbf{x} = - \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} (\nabla \mathbf{g}(\sigma) - \nabla_h \mathbf{u}_h(\sigma)) \boldsymbol{\tau}_k \cdot \Psi_k^P d\sigma - (\nabla_h^D \mathbf{u} \boldsymbol{\tau}, \tilde{\Psi})_{\Gamma, h}. \quad (59)$$

*Proof.* Applying the discrete Green formula (15), there holds

$$\begin{aligned} \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla_h \times \Psi_h(\mathbf{x}) d\mathbf{x} &= (\nabla_h^D \mathbf{u}, \nabla_h^D \times \Psi)_D \\ &= (\nabla_h^{T,P} \times (\nabla_h^D \mathbf{u}), \Psi)_{T,P} - (\nabla_h^D \mathbf{u} \boldsymbol{\tau}, \tilde{\Psi})_{\Gamma, h}. \end{aligned}$$

Using the result of lemma 3.3, we obtain (59).  $\square$

With these results, we may now obtain a representation formula for  $i_2$ .

**Proposition 3.5.** *Let  $\mathbf{u} = (\mathbf{u}_i^T, \mathbf{u}_k^P)$  be the velocity component of the solution of the scheme (25)–(29) and  $\mathbf{u}_h$  the function associated to  $\mathbf{u}$  by Def. 1.8. Let  $\hat{\Psi}$  be defined in (37). Let  $\Psi = (\Psi_i^T, \Psi_k^P) \in (\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$  and  $\Psi_h$  be its associated function. Then, the following representation holds*

$$\begin{aligned} i_2 &= \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \in \overset{\circ}{T}_i} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\hat{\Psi} - \Psi_i^T)(\sigma) d\sigma \\ &\quad + \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \in \overset{\circ}{P}_k} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\hat{\Psi} - \Psi_k^P)(\sigma) d\sigma \\ &\quad - \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} (\nabla \mathbf{g}(\sigma) - \nabla_h \mathbf{u}_h(\sigma)) \boldsymbol{\tau}_k \cdot (\hat{\Psi}(\sigma) - \Psi_k^P) d\sigma. \end{aligned} \quad (60)$$

*Proof.* From (40), there holds

$$\begin{aligned} i_2 &= \sum_j \int_{D_j} (\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h) : \nabla \times \hat{\Psi}(\mathbf{x}) d\mathbf{x} \\ &= \int_{\Omega} \nabla \hat{\mathbf{u}} : \nabla \times \hat{\Psi}(\mathbf{x}) d\mathbf{x} - \sum_j \int_{D_j} \nabla_h \mathbf{u}_h : \nabla_h \times \Psi_h(\mathbf{x}) d\mathbf{x} - \sum_j H_2(j), \end{aligned} \quad (61)$$

where

$$H_2(j) = \int_{D_j} \nabla_h \mathbf{u}_h : (\nabla \times \hat{\Psi} - \nabla_h \times \Psi_h)(\mathbf{x}) d\mathbf{x}. \quad (62)$$

By application of the continuous Green formula, there holds

$$\int_{\Omega} \nabla \hat{\mathbf{u}} : \nabla \times \hat{\Psi}(\mathbf{x}) d\mathbf{x} = - \int_{\partial\Omega} \nabla \hat{\mathbf{u}} \boldsymbol{\tau} \cdot \hat{\Psi}(\sigma) d\sigma = - \int_{\Gamma} \nabla \hat{\mathbf{g}} \boldsymbol{\tau} \cdot \hat{\Psi}(\sigma) d\sigma. \quad (63)$$

We can evaluate the sum of  $H_2(j)$  over  $j$  just like we evaluated the sum of  $H_1(j)$  in Proposition 3.2. There are only two differences. The first is that the gradients of  $\hat{\Phi}$  and  $\Phi_h$  are replaced by the curls of  $\hat{\Psi}$  and  $\Psi_h$ , which implies that normal vectors  $\mathbf{n}_s$  are replaced by tangent vectors  $-\boldsymbol{\tau}_s$ . The second difference is that the boundary integrals do not vanish any more, but can be evaluated like in the discussion that leads to (47). Then, noting that

$$\sum_{j \in J^{\Gamma}} \int_{A_j} \nabla_h \mathbf{u}_h \boldsymbol{\tau}_j \cdot \hat{\Psi}(\sigma) d\sigma = \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} \nabla_h \mathbf{u}_h \boldsymbol{\tau}_k \cdot \hat{\Psi}(\sigma) d\sigma,$$

we obtain the following formula

$$\begin{aligned} \sum_j H_2(j) &= -\frac{1}{2} \sum_{i \in [1, I]} \sum_{s \in T_i^{\circ}} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\hat{\Psi} - \Psi_i^T)(\sigma) d\sigma \\ &\quad - \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \in P_k^{\circ}} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\hat{\Psi} - \Psi_k^P)(\sigma) d\sigma \\ &\quad - \sum_{k \in \Gamma} \int_{P_k \cap \Gamma} \nabla_h \mathbf{u}_h \boldsymbol{\tau}_k \cdot \hat{\Psi}(\sigma) d\sigma + (\nabla_h^D \mathbf{u} \boldsymbol{\tau}, \tilde{\Psi})_{\Gamma, h}. \end{aligned} \quad (64)$$

Combining (61), (63), (59) and (64), we obtain (60).  $\square$

### 3.2. A representation of the pressure error

We shall estimate the  $L^2$  norm of the error between the exact solution  $\hat{p}$  and the function associated to the solution of the DDFV scheme by (20). We recall that  $e_p = \hat{p} - p_h$  and  $\mathbf{e} = \hat{\mathbf{u}} - \mathbf{u}_h$ .

**Proposition 3.6.** Let  $\widehat{\mathbf{v}} \in (H_0^1(\Omega))^2$  and  $\mathbf{v} = (\mathbf{v}_i^T, \mathbf{v}_k^P) \in (\mathbb{R}^2)^{I+J^\Gamma} \times (\mathbb{R}^2)^K$  be such that  $\tilde{\mathbf{v}}_j = 0$  for all  $j \in \Gamma$ . We have that

$$\begin{aligned} \int_{\Omega} e_p \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx &= \int_{\Omega} \nabla_h \mathbf{e} : \nabla \widehat{\mathbf{v}}(\mathbf{x}) dx - \frac{1}{2} \sum_{i \in [1, I]} \int_{T_i} \mathbf{f} \cdot (\widehat{\mathbf{v}} - \mathbf{v}_i^T)(\mathbf{x}) dx - \frac{1}{2} \sum_{k \in [1, K]} \int_{P_k} \mathbf{f} \cdot (\widehat{\mathbf{v}} - \mathbf{v}_k^P)(\mathbf{x}) dx \\ &+ \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset \overset{\circ}{T}_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\mathbf{v}} - \mathbf{v}_i^T)(\sigma) d\sigma \\ &+ \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset \overset{\circ}{P}_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\mathbf{v}} - \mathbf{v}_k^P)(\sigma) d\sigma. \end{aligned} \quad (65)$$

*Proof.* We can use formula (35) to obtain

$$\begin{aligned} \int_{\Omega} e_p \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx &= \int_{\Omega} \hat{p} \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx - \int_{\Omega} p_h \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx \\ &= \int_{\Omega} \nabla_h \mathbf{e} : \nabla \widehat{\mathbf{v}}(\mathbf{x}) dx - \int_{\Omega} \mathbf{f} \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx + \int_{\Omega} \nabla_h \mathbf{u}_h : \nabla \widehat{\mathbf{v}}(\mathbf{x}) dx - \int_{\Omega} p_h \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx. \end{aligned} \quad (66)$$

Using Eq. (32), we have

$$\int_{\Omega} e_p \nabla \cdot \widehat{\mathbf{v}}(\mathbf{x}) dx = \int_{\Omega} \nabla_h \mathbf{e} : \nabla \widehat{\mathbf{v}}(\mathbf{x}) dx - \int_{\Omega} \mathbf{f} \cdot (\widehat{\mathbf{v}} - \mathbf{v}_h^*)(\mathbf{x}) dx + \int_{\Omega} (\nabla_h \mathbf{u}_h - p_h I_2) : (\nabla \widehat{\mathbf{v}} - \nabla_h \mathbf{v}_h)(\mathbf{x}) dx. \quad (67)$$

Just like in Proposition 3.2 (see (44) and (49)), we have

$$\begin{aligned} \int_{\Omega} (\nabla_h \mathbf{u}_h - p_h I_2) : (\nabla \widehat{\mathbf{v}} - \nabla_h \mathbf{v}_h)(\mathbf{x}) dx &= \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset \overset{\circ}{T}_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\mathbf{v}} - \mathbf{v}_i^T)(\sigma) d\sigma \\ &+ \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset \overset{\circ}{P}_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\mathbf{v}} - \mathbf{v}_k^P)(\sigma) d\sigma. \end{aligned} \quad (68)$$

From (67) and (68), we obtain (65).  $\square$

## 4. A COMPUTABLE ERROR BOUND

### 4.1. Preliminaries

In this subsection, we will present some Poincaré-type inequalities which are useful to obtain a computable error bound.

**Lemma 4.1.** Let  $\omega$  be an open bounded set which is star-shaped with respect to one of its point. Let  $\mathbf{u} \in (H^1(\omega))^2$  and let  $\bar{\mathbf{u}}_{\omega}$  be the mean-value of  $\mathbf{u}$  over  $\omega$ . Then,

$$\exists C(\omega), \text{ s.t. } \|\mathbf{u} - \bar{\mathbf{u}}_{\omega}\|_{L^2(\omega)} \leq C(\omega) \text{diam}(\omega) \|\nabla \mathbf{u}\|_{L^2(\omega)}. \quad (69)$$

Note that when  $\omega$  is convex, a universal constant  $C(\omega)$  is given by  $\frac{1}{\pi}$  (see [22]). When  $\omega$  is not convex, we may use explicitly computable formulas given, for example, by [5, 26]. The constant  $C(\omega)$  only depends on the shape of  $\omega$ , not on its diameter.



Finally, we will also need a trace inequality (see [21]).

**Lemma 4.2.** *Let  $T$  be a triangle and let  $E$  be one of its edges; let  $\rho$  be the distance from  $E$  to the vertex of  $T$  opposite to  $E$ , and let  $\sigma$  be the longest among the two other sides of  $T$ . Let  $\varepsilon > 0$  be an arbitrary real valued number; then for all  $\mathbf{u} \in (H^1(T))^2$ , there holds*

$$\|\mathbf{u}\|_{L^2(E)}^2 \leq \frac{1}{\rho} \left( (2 + \varepsilon^{-2}) \|\mathbf{u}\|_{L^2(T)}^2 + \varepsilon^2 \sigma^2 \|\nabla \mathbf{u}\|_{L^2(T)}^2 \right). \quad (70)$$

#### 4.2. A computable bound for the velocity error

In the expression (42) of  $i_1$ , the values of  $(\Phi_i^T, \Phi_k^P)$  are arbitrary, except for the boundary midpoint values chosen so that (41) holds. In the expression of  $i_2$  in (60), the values of  $(\Psi_i^T, \Psi_k^P)$  are arbitrary.

**Definition 4.3.** Since  $\widehat{\Phi}$ ,  $\widehat{\Psi}$  are not necessarily more regular than  $(H^1(\Omega))^2$ , we choose as an interpolation their  $L^2$  projection on the primal and dual cells:

$$\Phi_i^T = \frac{1}{|T_i|} \int_{T_i} \widehat{\Phi}(\mathbf{x}) d\mathbf{x}, \quad \forall i \in [1, I]; \quad \Phi_k^P = \frac{1}{|P_k|} \int_{P_k} \widehat{\Phi}(\mathbf{x}) d\mathbf{x}, \quad \forall k \in [1, K], \quad (71)$$

$$\Psi_i^T = \frac{1}{|T_i|} \int_{T_i} \widehat{\Psi}(\mathbf{x}) d\mathbf{x}, \quad \forall i \in [1, I]; \quad \Psi_k^P = \frac{1}{|P_k|} \int_{P_k} \widehat{\Psi}(\mathbf{x}) d\mathbf{x}, \quad \forall k \in [1, K]. \quad (72)$$

In order to complete the definition of  $(\Phi_i^T, \Phi_k^P)$ , for any  $i \in \Gamma$ , the boundary value of  $\Phi_i^T$  is given so that (41) holds. Note that it is not necessary to define the value of  $\Psi_i^T$  for all  $i \in \Gamma$ .

**Proposition 4.4.** *Let  $h_i^T := \text{diam}(T_i)$ ,  $h_k^P := \text{diam}(P_k)$ . There exist computable constants  $C(T_i)$  and  $C(P_k)$  such that*

$$\left| \sum_{i \in [1, I]} \int_{T_i} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_i^T)(\mathbf{x}) d\mathbf{x} \right| \leq \text{osc}(\mathbf{f}, T, \Omega) \|\nabla \widehat{\Phi}\|_{L^2(\Omega)}, \quad (73)$$

$$\left| \sum_{k \in [1, K]} \int_{P_k} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_k^P)(\mathbf{x}) d\mathbf{x} \right| \leq \text{osc}(\mathbf{f}, P, \Omega) \|\nabla \widehat{\Phi}\|_{L^2(\Omega)}, \quad (74)$$

where

$$\text{osc}(\mathbf{f}, T, \Omega) = \left( \sum_{i \in [1, I]} (C(T_i) h_i^T)^2 \|\mathbf{f} - \mathbf{f}_i^T\|_{L^2(T_i)} \right)^{1/2}, \quad (75)$$

$$\text{osc}(\mathbf{f}, P, \Omega) = \left( \sum_{k \in [1, K]} (C(P_k) h_k^P)^2 \|\mathbf{f} - \mathbf{f}_k^P\|_{L^2(P_k)} \right)^{1/2} \quad (76)$$

and  $\mathbf{f}_i^T$  and  $\mathbf{f}_k^P$  are defined by (30).

*Proof.* Since  $\Phi_i^T$  was chosen as the mean-value of  $\widehat{\Phi}$  over  $T_i$  (see (71)), we have

$$\int_{T_i} \mathbf{f} \cdot (\widehat{\Phi} - \Phi_i^T)(\mathbf{x}) d\mathbf{x} = \int_{T_i} (\mathbf{f} - \mathbf{f}_i^T) \cdot (\widehat{\Phi} - \Phi_i^T)(\mathbf{x}) d\mathbf{x}.$$

Applying the Cauchy-Schwarz inequality, Lemma 4.1 to  $\widehat{\Phi}$  over  $T_i$  and the discrete Cauchy-Schwarz inequality, we obtain (73). Similarly, we also obtain (74). Note that the primal cells  $T_i$  are usually convex so that  $C(T_i) = \frac{1}{\pi}$  may be used. On the other hand, dual cells  $P_k$  may not always be convex.  $\square$

Propositions 4.5 and 4.7 below are proved just like Propositions 5.9 and 5.10 in [21].

**Proposition 4.5.** *For any primal cell  $T_i$  and any dual  $P_k$  such that  $T_i \cap P_k \neq \emptyset$ , Let  $s = [G_i S_k]$  and  $t_{ik,1}$  and  $t_{ik,2}$  be the triangles defined in Fig. 4 such that  $t_{ik,1} \cup t_{ik,2} = T_i \cap P_k$ . Let  $\rho_{ik,\alpha}$  be the distance from  $s$  to the vertex of  $t_{ik,\alpha}$  opposite to  $s$  and  $\sigma_{ik,\alpha}$  be the length of the longest among the two other edges of  $t_{ik,\alpha}$ .  $C(T_i)$  is the constant that appears in (69). For any strictly positive  $\mu$ , let us define*

$$C_s(\mu) = \frac{\left(1 + \sqrt{1 + \frac{\sigma_{ik,1}^2}{\mu}}\right) \left(1 + \sqrt{1 + \frac{\sigma_{ik,2}^2}{\mu}}\right)}{\left(1 + \sqrt{1 + \frac{\sigma_{ik,1}^2}{\mu}}\right) \rho_{ik,2} + \left(1 + \sqrt{1 + \frac{\sigma_{ik,2}^2}{\mu}}\right) \rho_{ik,1}}, \quad (77)$$

$$\chi_i(\mu) = (C(T_i)h_i^T)^2 + \mu. \quad (78)$$

We define the local and global error estimators related to the primal mesh:

$$(\eta_i^T)^2 = \inf_{\mu > 0} \left[ \chi_i(\mu) \sum_{s \in \overset{\circ}{T}_i} C_s(\mu) \|[(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s\|_{L^2(s)}^2 \right] \quad \text{and} \quad (\eta^T)^2 = \sum_i (\eta_i^T)^2, \quad (79)$$

$$(\eta_i^{\prime T})^2 = \inf_{\mu > 0} \left[ \chi_i(\mu) \sum_{s \in \overset{\circ}{T}_i} C_s(\mu) \|[\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s\|_{L^2(s)}^2 \right] \quad \text{and} \quad (\eta^{\prime T})^2 = \sum_i (\eta_i^{\prime T})^2. \quad (80)$$

With these definitions, there holds:

$$\left| \sum_{i \in [1, I]} \sum_{s \in \overset{\circ}{T}_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_i^T)(\sigma) d\sigma \right| \leq \eta^T \|\nabla \widehat{\Phi}\|_{L^2(\Omega)}, \quad (81)$$

$$\left| \sum_{i \in [1, I]} \sum_{s \in \overset{\circ}{T}_i} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\widehat{\Psi} - \Psi_i^T)(\sigma) d\sigma \right| \leq \eta^{\prime T} \|\nabla \widehat{\Psi}\|_{L^2(\Omega)}. \quad (82)$$

**Remark 4.6.** The minimization in (79) is performed numerically when we effectively compute the estimators. However, we may already get an idea of the behaviour of this quantity by choosing  $\mu = h_{T_i}^2$  to evaluate  $\eta_i^T$ . By definition of  $\sigma_{ik,\alpha}$ , this length is lower than the diameter of  $T_i$ , which implies

$$C_s \left( (h_i^T)^2 \right) \leq \frac{(1 + \sqrt{2})^2}{2(\rho_{ik,1} + \rho_{ik,2})}. \quad (83)$$

Under the hypothesis that the ratios  $\frac{\rho_{ik,\alpha}}{h_i^T}$  are all bounded by below by the same constant which does not depend on the mesh, we obtain the following bound

$$\eta_i^T \leq K h_i^T \sum_{s \in \overset{\circ}{T}_i} \|[(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s\|_{L^2(s)}^2,$$

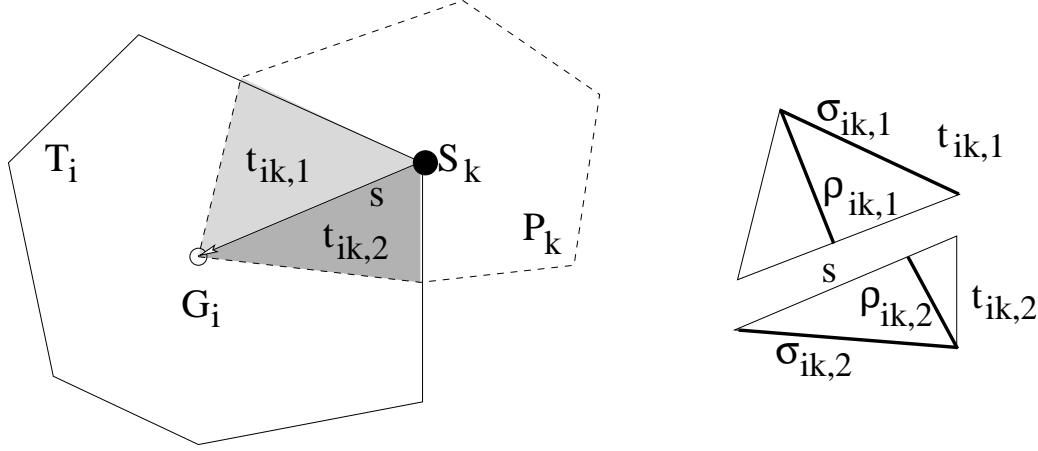


FIGURE 4. For each cell  $T_i$  and each vertex  $S_k$  of  $T_i$ ,  $T_i \cap P_k$  is split in two triangles  $t_{ik,1}$  and  $t_{ik,2}$ .

where the constant  $K$  does not depend on the mesh. The same remark holds for  $\eta'_i{}^T$ .

**Proposition 4.7.** *Let us set the same notations as in Prop. 4.5. Let  $C_s$  be defined by (77). Let  $C(P_k)$  be the constant involved in (69). Let us define for any strictly positive  $\mu$ ,*

$$\chi_k(\mu) = (C(P_k)h_k^P)^2 + \mu. \quad (84)$$

We define the local and global error estimators related to the dual mesh:

$$(\eta_k^P)^2 = \inf_{\mu > 0} \left[ \chi_k(\mu) \sum_{s \in \overset{\circ}{P}_k} C_s(\mu) \|[(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s\|_{L^2(s)}^2 \right] \quad \text{and} \quad (\eta^P)^2 = \sum_k (\eta_k^P)^2, \quad (85)$$

$$(\eta'_k{}^P)^2 = \inf_{\mu > 0} \left[ \chi_k(\mu) \sum_{s \in \overset{\circ}{P}_k} C_s(\mu) \|[\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s\|_{L^2(s)}^2 \right] \quad \text{and} \quad (\eta'^P)^2 = \sum_k (\eta'_k{}^P)^2. \quad (86)$$

With these definitions, there holds:

$$\left| \sum_{k \in [1, K]} \sum_{s \in \overset{\circ}{P}_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2) \mathbf{n}_s]_s \cdot (\widehat{\Phi} - \Phi_k^P)(\sigma) d\sigma \right| \leq \eta^P \|\nabla \widehat{\Phi}\|_{L^2(\Omega)}, \quad (87)$$

$$\left| \sum_{k \in [1, K]} \sum_{s \in \overset{\circ}{P}_k} \int_s [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s \cdot (\widehat{\Psi} - \Psi_k^P)(\sigma) d\sigma \right| \leq \eta'^P \|\nabla \widehat{\Psi}\|_{L^2(\Omega)}. \quad (88)$$

**Proposition 4.8.** *For any  $k \in \Gamma$ , let us denote by  $D_{j_1(k)}$  and  $D_{j_2(k)}$  the two diamond cells whose boundary intersect  $\Gamma$  and which have  $S_k$  as a vertex. Let  $q_{j_1(k)} = P_k \cap D_{j_1(k)}$ ,  $q_{j_2(k)} = P_k \cap D_{j_2(k)}$  and the segment  $b_{j_\alpha(k)}$  be the intersection between  $\partial q_{j_\alpha(k)}$  and  $\Gamma$ ; see Fig. 5. Let  $\rho_{j_\alpha(k)}$  be the distance from  $b_{j_\alpha(k)}$  to the vertex of  $q_{j_\alpha(k)}$  opposite to  $b_{j_\alpha(k)}$  and  $\sigma_{j_\alpha(k)}$  be the length of the longest among the two other edges of  $q_{j_\alpha(k)}$ . Let  $C(P_k)$  be the constant that appears in (69). For any strictly positive  $\mu$ , let us define*

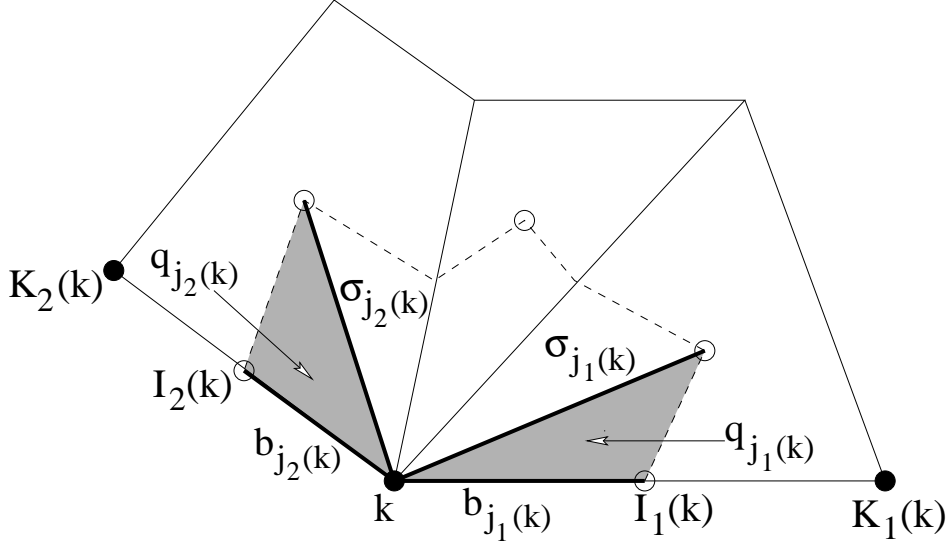


FIGURE 5. For any  $k \in \Gamma$ ,  $S_k$  is the common vertex of  $q_{j_1(k)}$  and  $q_{j_2(k)}$ .

$$C_\alpha(\mu) = \frac{2 + \frac{\sigma_{j_\alpha(k)}^2}{\mu + \sqrt{\mu^2 + \mu\sigma_{j_\alpha(k)}^2}}}{\rho_{j_\alpha(k)}}, \quad (89)$$

$$\lambda_k(\mu) = (C(P_k)h_k^P)^2 + \mu. \quad (90)$$

We define the local and global error estimator related to the boundary:

$$(\zeta_k^P)^2 = \inf_\mu \left[ \lambda_k(\mu) \sum_{\alpha=1}^2 C_\alpha(\mu) \|(\nabla \mathbf{g} - \nabla_h \mathbf{u}_h) \boldsymbol{\tau}_{j_\alpha(k)}\|_{L^2(b_{j_\alpha(k)})}^2 \right] \text{ and } (\zeta^P)^2 = \sum_{k \in \Gamma} (\zeta_k^P)^2. \quad (91)$$

With these definitions, there holds:

$$\sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} |(\nabla \mathbf{g} - \nabla_h \mathbf{u}_h) \boldsymbol{\tau}_k \cdot (\widehat{\Psi} - \Psi_k^P)(\sigma)| d\sigma \leq \zeta^P \|\nabla \widehat{\Psi}\|_{L^2(\Omega)}. \quad (92)$$

*Proof.* By application of the Cauchy-Schwarz inequality on each edge  $b_{j_\alpha(k)}$  and the weighted discrete Cauchy-Schwarz inequality, we obtain for any set of strictly positive real-valued numbers  $C_\alpha^P$

$$\begin{aligned} \int_{\partial P_k \cap \Gamma} |(\nabla \mathbf{g} - \nabla_h \mathbf{u}_h) \boldsymbol{\tau}_k \cdot (\widehat{\Psi} - \Psi_k^P)(\sigma)| d\sigma &= \sum_{\alpha=1}^2 \int_{b_{j_\alpha(k)}} |(\nabla \mathbf{g} - \nabla_h \mathbf{u}_h) \boldsymbol{\tau}_{j_\alpha(k)} \cdot (\widehat{\Psi} - \Psi_k^P)(\sigma)| d\sigma \\ &\leq \left[ \sum_{\alpha=1}^2 C_\alpha^P \|(\nabla \mathbf{g} - \nabla_h \mathbf{u}_h) \boldsymbol{\tau}_{j_\alpha(k)}\|_{L^2(b_{j_\alpha(k)})}^2 \right]^{1/2} \left[ \sum_{\alpha=1}^2 \frac{1}{C_\alpha^P} \|\widehat{\Psi} - \Psi_k^P\|_{L^2(b_{j_\alpha(k)})}^2 \right]^{1/2}. \end{aligned} \quad (93)$$

Now, for each segment  $b_{j_\alpha(k)}$ , we can apply the trace inequality (70) on each triangle  $q_{j_\alpha(k)}$ , for all  $\alpha \in \{1, 2\}$  and for all strictly positive  $\varepsilon_{j_\alpha(k)}$ . With  $C_{1,j_\alpha(k)} = \frac{2+\varepsilon_{j_\alpha(k)}^{-2}}{\rho_{j_\alpha(k)}}$  and  $C_{2,j_\alpha(k)} = \frac{\varepsilon_{j_\alpha(k)}^2 \sigma_{j_\alpha(k)}^2}{\rho_{j_\alpha(k)}}$ , we obtain

$$\|\widehat{\Psi} - \Psi_k^P\|_{L^2(b_{j_\alpha(k)})}^2 \leq C_{1,j_\alpha(k)} \|\widehat{\Psi} - \Psi_k^P\|_{L^2(q_{j_\alpha(k)})}^2 + C_{2,j_\alpha(k)} \|\nabla \widehat{\Psi}\|_{L^2(q_{j_\alpha(k)})}^2.$$

Let  $\mu > 0$  be arbitrary. For  $b_{j_\alpha(k)}$  for  $\alpha \in \{1, 2\}$ , let us choose  $\varepsilon_{j_\alpha(k)}$  so that

$$\varepsilon_{j_\alpha(k)}^2 = \frac{\mu + \sqrt{\mu^2 + \mu \sigma_{j_\alpha(k)}^2}}{\sigma_{j_\alpha(k)}^2} \iff C_{2,j_\alpha(k)} = \mu C_{1,j_\alpha(k)} \quad (94)$$

and  $C_\alpha^P = C_{1,j_\alpha(k)}$ . There holds:

$$\begin{aligned} \sum_{\alpha=1}^2 \frac{1}{C_\alpha^P} \|\widehat{\Psi} - \Psi_k^P\|_{L^2(b_{j_\alpha(k)})}^2 &\leq \sum_{\alpha=1}^2 \left( \|\widehat{\Psi} - \Psi_k^P\|_{L^2(q_{j_\alpha(k)})}^2 + \mu \|\nabla \widehat{\Psi}\|_{L^2(q_{j_\alpha(k)})}^2 \right) \\ &\leq \|\widehat{\Psi} - \Psi_k^P\|_{L^2(P_k)}^2 + \mu \|\nabla \widehat{\Psi}\|_{L^2(P_k)}^2 \\ &\leq [(C(P_k)h_k^P)^2 + \mu] \|\nabla \widehat{\Psi}\|_{L^2(P_k)}^2. \end{aligned} \quad (95)$$

Taking (95) into (93) and applying the discrete Cauchy-Schwarz inequality leads to (92).  $\square$

Before estimating the velocity error, we now define the indicator related to the penalization:

$$\zeta_\varepsilon = \varepsilon \|(\Delta_h^D p)_h\|_{L^2(\Omega)}. \quad (96)$$

In the term in the right-hand side of (40), it is easy to see that

$$\left| \varepsilon \int_{\Omega} (\Delta_h^D p)_h q(x) dx \right| \leq \zeta_\varepsilon \|q\|_{L^2(\Omega)}. \quad (97)$$

**Theorem 4.9.** *Let  $\|\nabla_h \mathbf{e}\|_{L^2(\Omega)}$  be defined by (36), let definitions (75)–(76), (77)–(80), (84)–(86), (89)–(91) and (96) hold. We have*

$$\|\nabla_h \mathbf{e}\|_{L^2(\Omega)} \leq \eta = \eta_h + \eta_\varepsilon. \quad (98)$$

where

$$\eta_h = \frac{1}{2} \left( \text{osc}(f, T, \Omega) + \text{osc}(f, P, \Omega) + \eta^T + \eta^P \right) + \left( \frac{1}{2} + \frac{\sqrt{2}}{\beta} \right) \left( \eta'^T + \eta'^P + 2\zeta^P \right), \quad (99)$$

$$\eta_\varepsilon = \frac{2}{\beta} \zeta_\varepsilon. \quad (100)$$

The quantities  $\eta$ ,  $\eta_h$ , and  $\eta_\varepsilon$  are respectively called the total estimator, the discretization estimator, and the penalization estimator for the velocity.

*Proof.* The result is obtained using (40), (38), (42), (60), (73)–(74), (81)–(82), (87)–(88), (92) and (97).  $\square$

### 4.3. A computable bound for the pressure error

**Proposition 4.10.** *The following estimate holds:*

$$\|e_p\|_{L(\Omega)} \leq \frac{1}{2\beta} (2\|\nabla_h \mathbf{e}\|_{L^2(\Omega)} + \text{osc}(\mathbf{f}, T, \Omega) + \text{osc}(\mathbf{f}, P, \Omega) + \eta^T + \eta^P). \quad (101)$$

*Proof.* Since  $e_p \in L_0^2(\Omega)$ , there exists  $\hat{\mathbf{v}} \in (H_0^1(\Omega))^2$ , such that

$$\|e_p\|_{L^2(\Omega)} \leq \frac{1}{\beta} \frac{\int_{\Omega} e_p \nabla \cdot \hat{\mathbf{v}}(x) dx}{\|\nabla \hat{\mathbf{v}}\|_{L^2(\Omega)}}. \quad (102)$$

With this given  $\hat{\mathbf{v}}$ , we associate  $\mathbf{v} = (\mathbf{v}_i^T, \mathbf{v}_k^P) \in (\mathbb{R}^2)^{I+J^T} \times (\mathbb{R}^2)^K$  such that

$$\mathbf{v}_i^T = \frac{1}{|T_i|} \int_{T_i} \hat{\mathbf{v}}(x) dx \quad \forall i \in [1, I], \quad \mathbf{v}_k^P = \frac{1}{|P_k|} \int_{P_k} \hat{\mathbf{v}}(x) dx \quad \forall k \in [1, K] \quad (103)$$

and the boundary values of  $\mathbf{v}_i$ ,  $i \in \Gamma$  are chosen so that  $\hat{\mathbf{v}}_j = 0$  for all  $j \in \Gamma$ . We may then apply (65) and follow calculations like those involved in propositions 4.4, 4.5 and 4.7. We finally obtain

$$\left| \int_{\Omega} e_p \nabla \cdot \hat{\mathbf{v}}(x) dx \right| \leq \frac{1}{2} (2\|\nabla_h \mathbf{e}\|_{L^2(\Omega)} + \text{osc}(\mathbf{f}, T, \Omega) + \text{osc}(\mathbf{f}, P, \Omega) + \eta^T + \eta^P) \|\nabla \hat{\mathbf{v}}\|_{L^2(\Omega)}. \quad (104)$$

Taking (104) into (102) proves (101).  $\square$

## 5. EFFICIENCY OF THE ESTIMATORS

In order to prove local efficiency of the estimators, we shall follow the bubble function technique as presented in [25]. Since the estimator  $\eta_i^T$  involves jumps of  $\nabla_h \mathbf{u}_h - p_h I_2$  through the common edge  $s = [G_i S_k]$  of two neighboring diamond-cells, we shall use functions with a support included in the set  $T_i \cap P_k = \cup_{\alpha=1,2} t_{ik,\alpha}$ , where triangles  $t_{ik,\alpha}$  with  $\alpha = 1$  or  $2$  are depicted in Fig. 4. Since we consider a fixed  $s$  in what follows, we simplify the notations to  $t_1$  and  $t_2$ . For any triangle  $t$  in  $\{t_1, t_2\}$ , we denote by  $\lambda_{t,\beta}$  the barycentric coordinates associated with the three vertices of  $t$ , with  $\beta \in \{1, 2, 3\}$ . We suppose that the vertices of  $t_1$  and  $t_2$  are locally numbered so that the two nodes of the edge  $s$  are the vertices 1 and 2 of each of the triangles  $t_1$  and  $t_2$ .

**Definition 5.1.** We define the following bubble functions

$$b_t = \begin{cases} 27\lambda_{t,1}\lambda_{t,2}\lambda_{t,3} & \text{on } t, \\ 0 & \text{elsewhere.} \end{cases} \quad (105)$$

$$b_s = \begin{cases} 4\lambda_{t_\alpha,1}\lambda_{t_\alpha,2} & \text{on } t_\alpha, \quad \alpha = \{1, 2\} \\ 0 & \text{elsewhere.} \end{cases} \quad (106)$$

There holds  $\omega_t = \text{supp}(b_t) \subset t$  and  $\omega_s := \text{supp}(b_s) = T_i \cap P_k = t_1 \cup t_2$ . The following propositions are given for example, in [25].

**Proposition 5.2.** *There holds*

$$0 \leq b_t \leq 1, \quad 0 \leq b_s \leq 1, \quad (107)$$

$$\int_s b_s(\sigma) d\sigma = \frac{2}{3} |s|. \quad (108)$$

**Proposition 5.3.** *There exists a constant  $C > 0$  only depending on the minimal angle in the couple  $(t_1, t_2)$  such that, for  $t = t_1$  or  $t = t_2$  and  $h_t = \text{diam}(t)$*

$$\frac{1}{C}h_t^2 \leq \int_t b_t(x)dx = \frac{9}{20}|t| \leq Ch_t^2, \quad (109)$$

$$\frac{1}{C}s^2 \leq \int_t b_s(x)dx = \frac{1}{3}|t| \leq Cs^2, \quad (110)$$

$$\|\nabla b_t\|_{L^2(t)} \leq Ch_t^{-1}\|b_t\|_{L^2(t)}, \quad (111)$$

$$\|\nabla b_s\|_{L^2(t)} \leq Cs^{-1}\|b_s\|_{L^2(t)}. \quad (112)$$

In order to prove the local efficiency of the error estimator we shall make the following hypothesis:

**Hypothesis 5.4.** *We suppose that the triangulation of  $\Omega$  composed of all the triangles  $t_{ik,\alpha}$  is regular in the sense that the minimum angles in those triangles are bounded by below independently of the mesh.*

From this hypothesis, we derive the following propositions.

**Proposition 5.5.** *For any primal cell  $T_i$  and any dual cell  $P_k$  such that  $T_i \cap P_k \neq \emptyset$ , let  $s = [G_i S_k]$  and  $t_{ik,1}$  and  $t_{ik,2}$  be the triangles in Fig. 4 such that  $t_{ik,1} \cup t_{ik,2} = T_i \cap P_k$ . Let  $h_i^T = \text{diam}(T_i)$ ,  $h_k^P = \text{diam}(P_k)$  and  $S_{ik} = |T_i \cap P_k|$ . Let Hypothesis 5.4 hold. Then, there exists a constant  $C$  independent of the mesh such that*

$$(h_i^T)^2 S_{ik}^{-1} \leq C \text{ and } (h_k^P)^2 S_{ik}^{-1} \leq C.$$

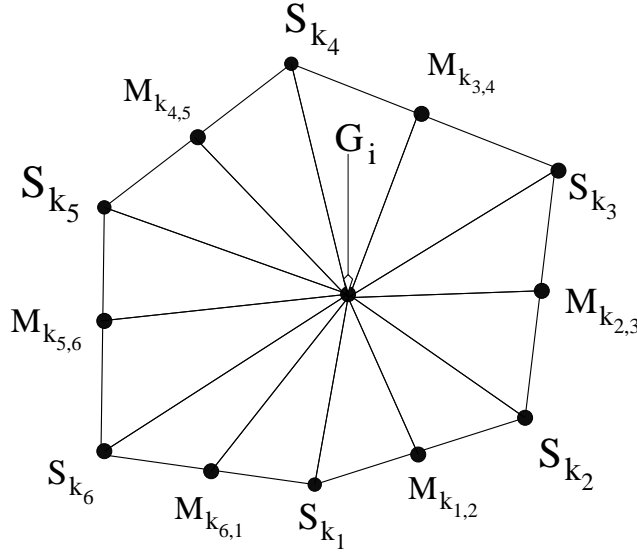


FIGURE 6. Notations of Prop. 5.5.

*Proof.* We will only prove the first inequality, since the second one can be treated in the same way.

Let  $\alpha_0 > 0$  be the lower bound of all the angles of all the triangles  $t_{ik,\alpha}$ .

For any  $i \in [1, I]$ , let  $V_i$  be the number of vertices  $S_{k\ell}$ , with  $\ell \in [1, V_i]$  of the primal cell  $T_i$ ; see Fig. 6 for the notations. First, we note that

$$V_i \leq V := \frac{2\pi}{2\alpha_0}, \text{ for all } i \in [1, I]. \quad (113)$$

Let  $M_{k_\ell, \ell+1}$  be the midpoint of segment  $[S_{k_\ell} S_{k_{\ell+1}}]$ , with the convention that  $k_{V_i+1} = k_1$ . Then

$$S_{ik_\ell} = |T_i \cap P_{k_\ell}| = |G_i S_{k_\ell} M_{k_\ell, \ell+1}| + |G_i S_{k_\ell} M_{k_{\ell-1}, \ell}|. \quad (114)$$

Now, let us estimate the area of triangle  $G_i S_{k_\ell} M_{k_\ell, \ell+1}$ . Following Hypothesis 5.4, all the angles of triangle  $G_i S_{k_\ell} M_{k_\ell, \ell+1}$  are greater than  $\alpha_0$ . Let  $h_{G_i}$  be the the maximum distance from point  $G_i$  to the boundary of  $T_i$ , i.e.,

$$h_{G_i} := \max\{|G_i S_{k_\ell}|, \ell \in [1, V_i]\}. \quad (115)$$

We have

$$|G_i S_{k_\ell} M_{k_\ell, \ell+1}| = \frac{1}{2} \sin(\widehat{S_{k_\ell} G_i M_{k_\ell, \ell+1}}) |G_i S_{k_\ell}| |G_i M_{k_\ell, \ell+1}|. \quad (116)$$

By a calculation on triangles  $G_i S_{k_\ell} M_{k_\ell, \ell+1}$  and  $G_i S_{k_{\ell+1}} M_{k_\ell, \ell+1}$ , there holds

$$|G_i M_{k_\ell, \ell+1}| \geq |G_i S_{k_\ell}| \sin \alpha_0 \quad \text{and} \quad |G_i S_{k_{\ell+1}}| \geq |G_i M_{k_\ell, \ell+1}| \sin \alpha_0. \quad (117)$$

From (117), we have the recurrence formula

$$|G_i S_{k_\ell}| \geq (\sin \alpha_0)^2 |G_i S_{k_{\ell+1}}|. \quad (118)$$

Starting from the vertex  $S_k$  which reaches the max in definition (115), the shortest way to go to a given vertex  $S_{k_\ell}$  contains at most  $V_i/2$  neighboring vertices for which we may apply (118), and we obtain:

$$|G_i S_{k_\ell}| \geq (\sin \alpha_0)^{V_i} h_{G_i}. \quad (119)$$

Then, from (117), we get that

$$|G_i S_{k_\ell}| |G_i M_{k_\ell, \ell+1}| \geq (\sin \alpha_0)^{2V_i+1} h_{G_i}^2. \quad (120)$$

Combining (113), (116) and (120), and noting that from the definition of  $h_{G_i}$ , then  $h_{G_i} \geq \frac{h_i^T}{2}$ , we obtain

$$|G_i S_{k_\ell} M_{k_\ell, \ell+1}| \geq (\sin \alpha_0)^{2V+2} \frac{(h_i^T)^2}{8}. \quad (121)$$

In the same way, we have

$$|G_i S_{k_\ell} M_{k_{\ell-1}, \ell}| \geq (\sin \alpha_0)^{2V+2} \frac{(h_i^T)^2}{8}. \quad (122)$$

Using (114), (121)–(122), we obtain:

$$S_{ik_\ell} \geq (\sin \alpha_0)^{2V+2} \frac{(h_i^T)^2}{4}. \quad (123)$$

Thus, the inequality is proved with  $C = 4 (\sin \alpha_0)^{-\frac{2V}{\alpha_0} - 2}$ .  $\square$

**Proposition 5.6.** *Under Hypothesis 5.4, the positive constants  $C(T_i)$  and  $C(P_k)$  are bounded independently of the mesh, and the constant  $C$  in Prop. 5.3 is bounded by above and by below independently of the mesh.*



*Proof.* The constants  $C(T_i)$ ,  $C(P_k)$  coming from (69) were bounded explicitly in [26]. From these expressions, it is easily seen that they are bounded if Hyp. 5.4 holds. Moreover, it is proved in [25] that  $C$  in Prop. 5.3 depends only on the regularity of the triangles  $t_{ik,\alpha}$ .  $\square$

Now, we will consider the efficiency of the estimators.

**Theorem 5.7.** *For any primal cell  $T_i$ , let  $h_i^T := \text{diam}(T_i)$  and  $\mathbf{f}_i^T$  be the mean-value of  $\mathbf{f}$  over  $T_i$ . Let  $\eta_i^T$  (resp.  $\eta_i'^T$ ) be defined in (79) (resp. in (80)). For any dual cell  $P_k$ , let  $h_k^P := \text{diam}(P_k)$  and  $\mathbf{f}_k^P$  be the mean-value of  $\mathbf{f}$  over  $P_k$ . Let  $\eta_k^P$  (resp.  $\eta_k'^P$ ) be defined in (85) (resp. in (86)). And for any boundary dual cell  $P_k$ , let  $\zeta_k^P$  be defined in (91). Let Hypothesis 5.4 hold. Then, there exists a constant  $C$  independent of the mesh such that*

$$(\eta_i^T)^2 \leq C \left( \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(T_i)}^2 + \|p_h - \hat{p}\|_{L^2(T_i)}^2 \right) + C(h_i^T)^2 \|\mathbf{f} - \mathbf{f}_i^T\|_{L^2(T_i)}^2, \quad (124)$$

$$(\eta_i'^T)^2 \leq C \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(T_i)}^2, \quad (125)$$

$$(\eta_k^P)^2 \leq C \left( \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(P_k)}^2 + \|p_h - \hat{p}\|_{L^2(P_k)}^2 \right) + C(h_k^P)^2 \|\mathbf{f} - \mathbf{f}_k^P\|_{L^2(P_k)}^2, \quad (126)$$

$$(\eta_k'^P)^2 \leq C \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(P_k)}^2. \quad (127)$$

Moreover, in the case  $\mathbf{g} = \mathbf{0}$ , there exists a constant  $C$  independent of the mesh such that

$$(\zeta_k^P)^2 \leq C \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(P_k)}^2. \quad (128)$$

*Proof.* Let us consider an element  $T_i$  of the primal mesh and a diamond edge  $s$  in  $\overset{\circ}{T}_i$ . Let us recall that by definition, such an edge  $s$  does not belong to  $\Gamma$ . Let us consider the function  $\mathbf{w}_s = [(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s b_s$ , where  $b_s$  is defined by (106). This function belongs to  $(H_0^1(\Omega))^2$  and we may thus apply (35), which, taking into account the support of  $\mathbf{w}_s$ , reduces to

$$\int_{\omega_s} (\nabla \hat{\mathbf{u}} - I_2 \hat{p}) : \nabla \mathbf{w}_s(x) dx = \int_{\omega_s} \mathbf{f} \cdot \mathbf{w}_s(x) dx. \quad (129)$$

Moreover,  $\mathbf{u}_h$  belongs to  $(P^1(D_j))^2$ ,  $p_h$  is a constant in each  $D_j$  and  $\mathbf{w}_s$  vanishes on  $\Gamma$ . Thus, there holds:

$$\begin{aligned} \int_{\Omega} (\nabla_h \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_s(x) dx &= \sum_j \int_{D_j} (\nabla \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_s(x) dx \\ &= \sum_j \int_{\partial D_j} ((\nabla \mathbf{u}_h - I_2 p_h) \mathbf{n}_{\partial D_j}) \cdot \mathbf{w}_s(\sigma) d\sigma \\ &= \sum_i \sum_{s' \subset \overset{\circ}{T}_i} \int_{s'} [(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_{s'}]_{s'} \cdot \mathbf{w}_s(\sigma) d\sigma. \end{aligned}$$

But since  $\mathbf{w}_s$  vanishes on all the other edges  $s' \neq s$ , taking into account the definition of  $\mathbf{w}_s$  and the property of  $b_s$  in (108), there holds

$$\begin{aligned} \int_{\Omega} (\nabla_h \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_s(x) dx &= \int_s [(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s \cdot \mathbf{w}_s(\sigma) d\sigma \\ &= |[(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s|^2 \int_s b_s(\sigma) \sigma \\ &= \frac{2}{3} |s| |[(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s|^2 \\ &= \frac{2}{3} |[(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s|_{L^2(s)}^2. \end{aligned} \quad (130)$$

Defining  $M := \|[(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s\|_{L^2(\omega_s)}$ , and taking into account (129), we have

$$\begin{aligned} M^2 &= \frac{3}{2} \int_{\Omega} (\nabla_h \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_s(x) dx = \frac{3}{2} \int_{\omega_s} (\nabla_h \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_s(x) dx \\ &= \frac{3}{2} \left[ \int_{\omega_s} (\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}) : \nabla \mathbf{w}_s(x) dx - \int_{\omega_s} (p_h - \hat{p}) \nabla \cdot \mathbf{w}_s(x) dx + \int_{\omega_s} \mathbf{f} \cdot \mathbf{w}_s(x) dx \right]. \end{aligned} \quad (131)$$

Using the Cauchy-Schwarz inequality leads to

$$M^2 \leq \frac{3}{2} \left[ \left( \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)} + \sqrt{2} \|p_h - \hat{p}\|_{L^2(\omega_s)} \right) \|\nabla \mathbf{w}_s\|_{L^2(\omega_s)} \right] + \frac{3}{2} \|\mathbf{f}\|_{L^2(\omega_s)} \|\mathbf{w}_s\|_{L^2(\omega_s)}. \quad (132)$$

Let us now bound  $\|\nabla \mathbf{w}_s\|_{L^2(\omega_s)}$  and  $\|\mathbf{w}_s\|_{L^2(\omega_s)}$ . There holds, thank to (112),

$$\|\nabla \mathbf{w}_s\|_{L^2(\omega_s)} = \|[(\nabla \mathbf{u}_s - I_2 p_h) \mathbf{n}_s]_s\| \|\nabla b_s\|_{L^2(\omega_s)} \leq \|[(\nabla \mathbf{u}_s - I_2 p_h) \mathbf{n}_s]_s\| C |s|^{-1} \|b_s\|_{L^2(\omega_s)}, \quad (133)$$

$$\|\mathbf{w}_s\|_{L^2(\omega_s)} = \|[(\nabla \mathbf{u}_s - I_2 p_h) \mathbf{n}_s]_s\| \|b_s\|_{L^2(\omega_s)}. \quad (134)$$

So there remains to find a bound for  $\|b_s\|_{L^2(\omega_s)}$ . In order to do this, we first infer from (107) that  $b_s^2 \leq b_s$ . This implies, using (110)

$$\|b_s\|_{L^2(\omega_s)} = \left[ \|b_s\|_{L^2(t_1)}^2 + \|b_s\|_{L^2(t_2)}^2 \right]^{1/2} \leq \left[ \int_{t_1 \cup t_2} b_s(x) dx \right]^{1/2} \leq C |s|. \quad (135)$$

Taking into account that  $\|[(\nabla_h \mathbf{u}_h - I_2 p_h) \mathbf{n}_s]_s\| = |s|^{-1/2} M$  and considering (132) to (135), we obtain

$$M \leq C \left[ |s|^{-1/2} (\|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)} + \sqrt{2} \|p_h - \hat{p}\|_{L^2(\omega_s)}) + |s|^{1/2} \|\mathbf{f}\|_{L^2(\omega_s)} \right]. \quad (136)$$

One usually expresses  $\|\mathbf{f}\|_{L^2(\omega_s)}$  as a function of  $\|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)} + \|p_h - \hat{p}\|_{L^2(\omega_s)}$  and of higher order terms. Let  $t = t_1$  or  $t_2$ , and let us denote by  $\mathbf{f}_t$  the mean value of  $\mathbf{f}$  over  $t$ . Then,

$$\|\mathbf{f}\|_{L^2(t)} \leq \|\mathbf{f} - \mathbf{f}_t\|_{L^2(t)} + \|\mathbf{f}_t\|_{L^2(t)}. \quad (137)$$

Then, consider  $\mathbf{w}_t = \mathbf{f}_t b_t$ , where  $b_t$  is defined by (105). The function  $\mathbf{w}_t$  belongs to  $(H_0^1)^2$ . Thus, taking into account the support of  $b_t$ , Eq. (35) reduces to

$$\int_t (\nabla \hat{\mathbf{u}} - I_2 \hat{p}) : \nabla \mathbf{w}_t(x) dx = \int_t \mathbf{f} \cdot \mathbf{w}_t(x) dx. \quad (138)$$

Moreover, since  $\nabla_h \mathbf{u}_h - I_2 p_h$  is a constant over each  $t$ , and since  $\mathbf{w}_t$  vanishes on the boundary of  $t$ , there holds

$$\int_t (\nabla_h \mathbf{u}_h - I_2 p_h) : \nabla \mathbf{w}_t(x) dx = 0. \quad (139)$$

Since  $\mathbf{f}_t$  is a constant over  $t$ , there holds, thanks to (109), (138) and (139),

$$\begin{aligned}
\|\mathbf{f}_t\|_{L^2(t)}^2 &= |t| (\mathbf{f}_t)^2 = C (\mathbf{f}_t)^2 \int_t b_t(\mathbf{x}) d\mathbf{x} = C \int_t \mathbf{f}_t \cdot \mathbf{w}_t(\mathbf{x}) d\mathbf{x} \\
&= C \left[ \int_t \mathbf{f} \cdot \mathbf{w}_t(\mathbf{x}) d\mathbf{x} + \int_t (\mathbf{f}_t - \mathbf{f}) \cdot \mathbf{w}_t(\mathbf{x}) d\mathbf{x} \right] \\
&= C \left[ \int_t (\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h) : \nabla \mathbf{w}_t(\mathbf{x}) d\mathbf{x} - \int_t (\hat{p} - p_h) \nabla \cdot \mathbf{w}_t(\mathbf{x}) d\mathbf{x} + \int_t (\mathbf{f}_t - \mathbf{f}) \cdot \mathbf{w}_t(\mathbf{x}) d\mathbf{x} \right] \\
&\leq C \left( \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(t)} + \sqrt{2} \|\hat{p} - p_h\|_{L^2(t)} \right) \|\nabla \mathbf{w}_t\|_{L^2(t)} + C \|\mathbf{f}_t - \mathbf{f}\|_{L^2(t)} \|\mathbf{w}_t\|_{L^2(t)}, \tag{140}
\end{aligned}$$

with  $C = 20/9$  in the above expressions. Let us now bound  $\|\mathbf{w}_t\|_{L^2(t)}$  and  $\|\nabla \mathbf{w}_t\|_{L^2(t)}$ . With (111), there holds

$$\|\nabla \mathbf{w}_t\|_{L^2(t)} = |\mathbf{f}_t| \|\nabla b_t\|_{L^2(t)} \leq |\mathbf{f}_t| C h_t^{-1} \|b_t\|_{L^2(t)}, \tag{141}$$

$$\|\mathbf{w}_t\|_{L^2(t)} = |\mathbf{f}_t| \|b_t\|_{L^2(t)}. \tag{142}$$

The remaining term that has to be bounded is  $\|b_t\|_{L^2(t)}$ . For this, we first infer from (107) that  $b_t^2(\mathbf{x}) \leq b_t(\mathbf{x})$  and then

$$|\mathbf{f}_t| \|b_t\|_{L^2(t)} \leq |\mathbf{f}_t| \left( \int_t b_t(\mathbf{x}) d\mathbf{x} \right)^{1/2} \leq C |\mathbf{f}_t| |t|^{1/2} = C \|\mathbf{f}_t\|_{L^2(t)}, \tag{143}$$

in which  $C = \sqrt{9/20}$ . Combining (140)–(141)–(142)–(143), we finally get

$$\|\mathbf{f}_t\|_{L^2(t)} \leq C \left( \|\mathbf{f}_t - \mathbf{f}\|_{L^2(t)} + h_t^{-1} \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(t)} + h_t^{-1} \|\hat{p} - p_h\|_{L^2(t)} \right).$$

Since  $s$  is an edge of  $t$ , there holds  $|s| \leq h_t$ ; applying (137), we obtain

$$\|\mathbf{f}\|_{L^2(t)} \leq C \left( \|\mathbf{f}_t - \mathbf{f}\|_{L^2(t)} + |s|^{-1} \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(t)} + |s|^{-1} \|\hat{p} - p_h\|_{L^2(t)} \right).$$

Thus, taking into account that  $\omega_s = t_1 \cup t_2$ , there holds

$$\begin{aligned}
\|\mathbf{f}\|_{L^2(\omega_s)} &\leq \|\mathbf{f}\|_{L^2(t_1)} + \|\mathbf{f}\|_{L^2(t_2)} \\
&\leq C \left( \|\mathbf{f}_{t_1} - \mathbf{f}\|_{L^2(t_1)} + \|\mathbf{f}_{t_2} - \mathbf{f}\|_{L^2(t_2)} \right) + C |s|^{-1} \left[ \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(t_1)} \right. \\
&\quad \left. + \left( \|\hat{p} - p_h\|_{L^2(t_1)} + \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(t_2)} + \|\hat{p} - p_h\|_{L^2(t_2)} \right) \right] \\
&\leq C |s|^{-1} \left( \|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(\omega_s)} + \|\hat{p} - p_h\|_{L^2(\omega_s)} \right) + C \|\mathbf{f}_{\omega_s} - \mathbf{f}\|_{L^2(\omega_s)}. \tag{144}
\end{aligned}$$

In this sequence of inequalities, we have used the fact that  $\mathbf{f}_t$  minimizes  $\|\mathbf{f} - \mathbf{c}\|_{L^2(t)}$  when  $\mathbf{c}$  runs over  $\mathbb{R}^2$ ; in particular,  $\|\mathbf{f}_t - \mathbf{f}\|_{L^2(t)} \leq \|\mathbf{f}_{\omega_s} - \mathbf{f}\|_{L^2(t)}$ , where  $\mathbf{f}_{\omega_s}$  is the mean value of  $\mathbf{f}$  over  $\omega_s$ . Combining (136) and (144), we obtain

$$M \leq C |s|^{1/2} \|\mathbf{f} - \mathbf{f}_{\omega_s}\|_{L^2(\omega_s)} + C |s|^{-1/2} \left( \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)} + \|\hat{p} - p_h\|_{L^2(\omega_s)} \right). \tag{145}$$

By definition, the local estimator  $(\eta_i^T)^2$  is lower than the value taken by the function in (79) in  $\mu = (h_i^T)^2$ . In (79), we may bound  $C(T_i)$  by  $1/\pi$  since the primal cells have been supposed to be convex, and with (83)

and (145), we obtain

$$\begin{aligned} (\eta_i^T)^2 &\leq C (h_i^T)^2 \sum_{s \in \overset{\circ}{T}_i} \frac{|s|^{-1}}{\rho_{ik,1} + \rho_{ik,2}} \left( \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)}^2 + \|p_h - \hat{p}\|_{L^2(\omega_s)}^2 \right) \\ &+ C (h_i^T)^2 \sum_{s \in \overset{\circ}{T}_i} \frac{|s|}{\rho_{ik,1} + \rho_{ik,2}} \|\mathbf{f} - \mathbf{f}_{\omega_s}\|_{L^2(\omega_s)}^2. \end{aligned} \quad (146)$$

Using Prop. 5.5, and since by definition  $S_{ik} = \frac{1}{2}|s|(\rho_{ik,1} + \rho_{ik,2})$ , we have that  $(h_i^T)^2 |s|^{-1}(\rho_{ik,1} + \rho_{ik,2})^{-1}$  is bounded by a constant that does not depend on the mesh under Hyp. 5.4. Moreover, under Hyp. 5.4, the ratio  $|s|\rho_{ik,\alpha}^{-1}$  is also bounded by a constant that does not depend on the mesh. So (146) leads to (124).

As far as (125) is concerned, let us consider the function  $\mathbf{v}_s = [\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s]_s b_s$ . There obviously holds

$$\int_{\Omega} \nabla \hat{\mathbf{u}} : \nabla \times \mathbf{v}_s(\mathbf{x}) \, dx = \int_{\omega_s} \nabla \hat{\mathbf{u}} : \nabla \times \mathbf{v}_s(\mathbf{x}) \, dx = 0. \quad (147)$$

Eq. (147) and the calculations that previously led to (130) may be used to yield

$$\begin{aligned} \|\nabla_h \mathbf{u}_h \boldsymbol{\tau}_s\|_{L^2(s)}^2 &= \frac{3}{2} \int_{\omega_s} \nabla_h \mathbf{u}_h : \nabla \times \mathbf{v}_s(\mathbf{x}) \, dx \\ &= \frac{3}{2} \int_{\omega_s} (\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}) : \nabla \times \mathbf{v}_s(\mathbf{x}) \, dx \\ &\leq \frac{3}{2} \|\nabla_h \mathbf{u}_h - \nabla \hat{\mathbf{u}}\|_{L^2(\omega_s)} \|\nabla \mathbf{v}_s\|_{L^2(\omega_s)}. \end{aligned} \quad (148)$$

Just like (132) led to (136) and then to (124), inequality (148) leads to (125). The dual inequalities (126), (127) and (128) may be obtained in the same way. The proof of (128) is very similar to that of (125) and (127), but some definitions have to be changed because the segment  $b_{j_\alpha(k)}$  in the definition (91) is a boundary segment, and is thus the edge of only one triangle  $t$ ; the function  $b_s$  is thus defined only in that triangle.  $\square$

## 6. NUMERICAL RESULTS

First, we study the influence of the parameter  $\varepsilon$  for a fixed mesh and then the influence of the mesh size for a fixed value of the penalty parameter. Secondly, we give an overall process to recursively adapt the value of the penalty parameter and adaptively refine the mesh.

### 6.1. Influence of the penalty parameter

In this subsection, we will work on the domain  $\Omega = [0; 1]^2$ . A triangular mesh with rather uniform triangles is used. The exact solution  $(\hat{\mathbf{u}}, \hat{p})$  is regular with  $\hat{\mathbf{u}} = (\partial_y \varphi, -\partial_x \varphi)$  given by

$$\varphi(x, y) = 100x^2y^2(1-x)^2(1-y)^2 \text{ and } \hat{p}(x, y) = 10(x^2 + y^2 - \frac{2}{3}). \quad (149)$$

Fig. 7 presents the plots of the errors and the estimators when the penalty parameter  $\varepsilon$  goes from  $10^{-2}$  to  $10^{-8}$ . They include the actual errors in the  $H^1(\Omega)$  and  $L^2(\Omega)$  norms for the velocity, i.e. the error in the velocity gradient  $\|\nabla \hat{\mathbf{u}} - \nabla_h \mathbf{u}_h\|_{L^2(\Omega)}$  and in the velocity  $\|\hat{\mathbf{u}} - \mathbf{u}_h\|_{L^2(\Omega)}$ , the total estimator, the discretization estimator and the penalization estimator which are given by Theorem 4.9 when we estimate the velocity error. The left (resp. right) figure corresponds to the mesh size  $h = 10^{-1}$  (resp.  $h = 3.125 \times 10^{-2}$ ). We see that for a given mesh, the ratio between the penalization estimator and the penalty parameter  $\varepsilon$  asymptotically tends to

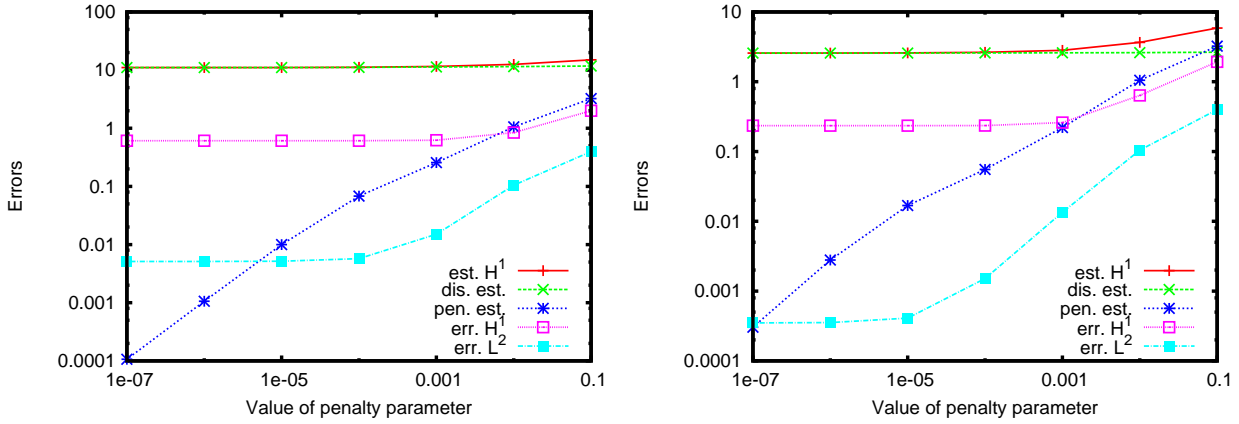


FIGURE 7. Actual errors in  $H^1(\Omega)$  and  $L^2(\Omega)$  norms, total estimator, discretization estimator and penalization estimator for the velocity. Left:  $h = 10^{-1}$ , right:  $h = 3.125 \times 10^{-2}$ .

a constant, while the discretization estimator is nearly independent of  $\varepsilon$ . Moreover, the actual errors decrease with  $\varepsilon$  until a certain level. Then, the discretization error is the dominant error and decreasing  $\varepsilon$  further does not have any influence on the overall error. As expected, when the mesh size is smaller (right part of the Figure), then the value of the penalty parameter for which the errors saturate is also smaller.

### 6.2. Influence of the mesh size

On the same square domain  $\Omega$  and with the same exact solution as previously, we work with a fixed  $\varepsilon = 10^{-2}$  (resp.  $\varepsilon = 10^{-7}$ ) in the left (resp. right) part of Fig. 8. Since the solution is regular, only uniformly refined triangular meshes will be considered. Figure 8 presents the same curves as Fig. 7, but now as a function of  $h$ , varying from 0.25 to  $1.6 \times 10^{-2}$ .

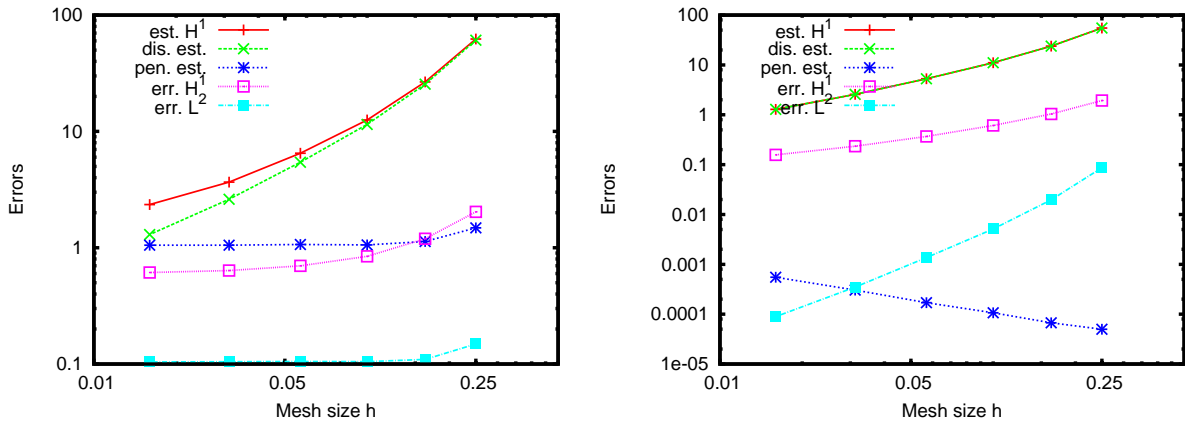


FIGURE 8. Actual errors in  $H^1(\Omega)$  and  $L^2(\Omega)$  norms, total estimator, discretization estimator and penalization estimator for the velocity. Left:  $\varepsilon = 10^{-2}$ , right:  $\varepsilon = 10^{-7}$ .

The actual errors decrease until the mesh size  $h$  is so small that the penalization error will dominate the discretization error, and the total error thus stagnates to a certain level. The penalization estimator is nearly independent of  $h$  in the left figure but behaves roughly like  $h^{-1}$  in the right figure. This behaviour remains

unexplained and further investigations have to be conducted about this. The total estimator and the discretization estimator decrease regularly when  $h$  decreases, roughly like  $h$ , but, then, when  $h$  is small enough the total estimator starts to stagnate because the penalization estimator stops being negligible (better seen in Fig. 9, where  $\varepsilon = 10^{-1}$ ).

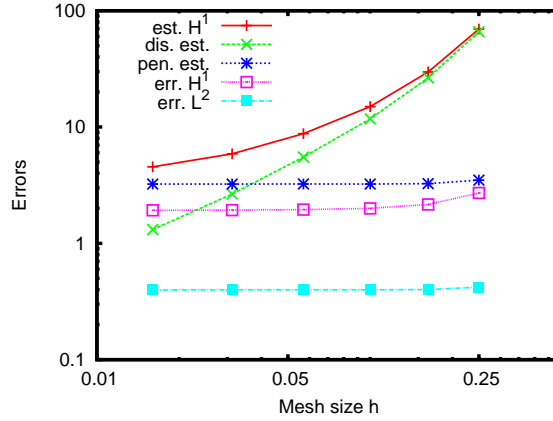


FIGURE 9. Actual errors in  $H^1(\Omega)$  and  $L^2(\Omega)$  norms, total estimator, discretization estimator and penalization estimator for the velocity for  $\varepsilon = 10^{-1}$ .

### 6.3. Adaptive penalty parameter and mesh refinement

We propose the following computational process. We start with a given coarse mesh and an initial value of  $\varepsilon$ , and we fix some ratio  $0 < \gamma \leq 1$ .

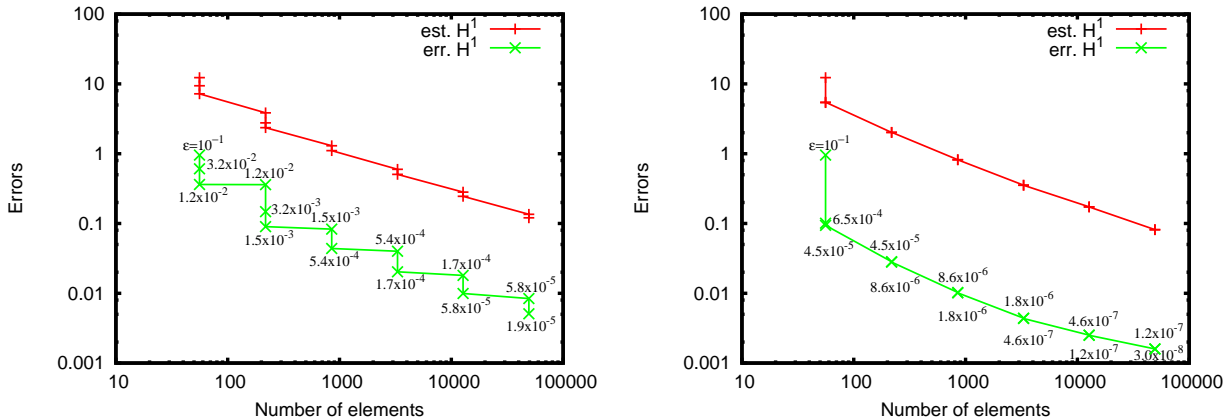


FIGURE 10. Actual errors in  $H^1(\Omega)$  and total estimator for the velocity. Left:  $\gamma = 1/10$ , right:  $\gamma = 1/500$ .

Then, we compute the numerical solution, and we get  $\eta_h$  and  $\eta_\varepsilon$ . Then,

- If  $\eta_\varepsilon \geq \gamma\eta_h$ , we adapt a new  $\varepsilon$  by multiplying the previous value of  $\varepsilon$  with the ratio  $\frac{\gamma\eta_h}{2\eta_\varepsilon}$  and keep the same mesh for a new computation. This has the effect of maintaining the error due to the penalization below a certain ratio of the error due to the discretization.

- Otherwise, we adaptively refine the mesh based on the discretization estimator  $\eta_h$ . For this, on the given mesh, we compute local discretization estimators  $\eta_{i,h}$  on each primal cell  $T_i$ , obtained by adding  $\eta_i, \eta'_i$  (respectively defined by (79) and (80)) to the contribution of  $T_i$  in the oscillation term (75), and by redistributing the sums that appear in the dual estimators (85), (86) and (76) and in the boundary estimators (91) to the cells  $T_i$  that have an intersection with  $P_k$ . Then, we require the refinement of a given primal cell  $T_i$  by a factor 4 in terms of area if  $\eta_{i,h} \geq (\max_i \eta_{i,h})/2$ .

The test we present to illustrate this strategy is also on the domain  $\Omega = [0; 1]^2$ , The exact solution  $(\hat{\mathbf{u}}, \hat{p})$  is regular with  $\hat{\mathbf{u}} = (\varphi_y, -\varphi_x)$ , and  $\varphi$  is given by

$$\varphi(x, y) = x^2(1-x)^2y^2(1-y)^2 \text{ and } \hat{p}(x, y) = 5(x^2 + y^2 - \frac{2}{3}).$$

For accuracy reasons, the ratio  $\gamma$  may be chosen so that the penalization error is much lower than the discretization error like in the right picture of Fig. 10 obtained with  $\gamma = 1/500$ . We observe that the actual error and the total estimator are not affected by the penalty term. Moreover, we made a test with  $\gamma = 1/10$  and we present the result in the left picture of Fig. 10 to show the interplay between the mesh refinement and the decrease of  $\varepsilon$ .

The next test is inspired from [16]. We will compare the exact  $H^1(\Omega)$  error over the velocity, and the total velocity error estimator for uniform and adaptive refinements. We will combine this work with the adaptive penalty - mesh refinement process with  $\gamma = 1/500$ . Our test is in the domain  $\Omega = [0, 1]^2$  and the exact couple solution  $(\hat{\mathbf{u}}, \hat{p})$  is singular with  $\hat{\mathbf{u}} = (\varphi_y, -\varphi_x)$ , and  $\varphi$  is

$$\varphi(x, y) = x^{\frac{7}{4}}(1-x)^2y^2(1-y)^2 \text{ and } \hat{p}(x, y) = \frac{x+y-1}{10}.$$

We observe that the velocity  $\hat{\mathbf{u}}$  is in  $[H^{\frac{5}{4}}(\Omega)]^2$  and there is a boundary singularity on the edge  $x = 0$ .

In Fig. 11, the penalty parameter  $\varepsilon$  decreases from  $10^{-2}$  to  $2 \times 10^{-10}$  for the adaptive refinement (from

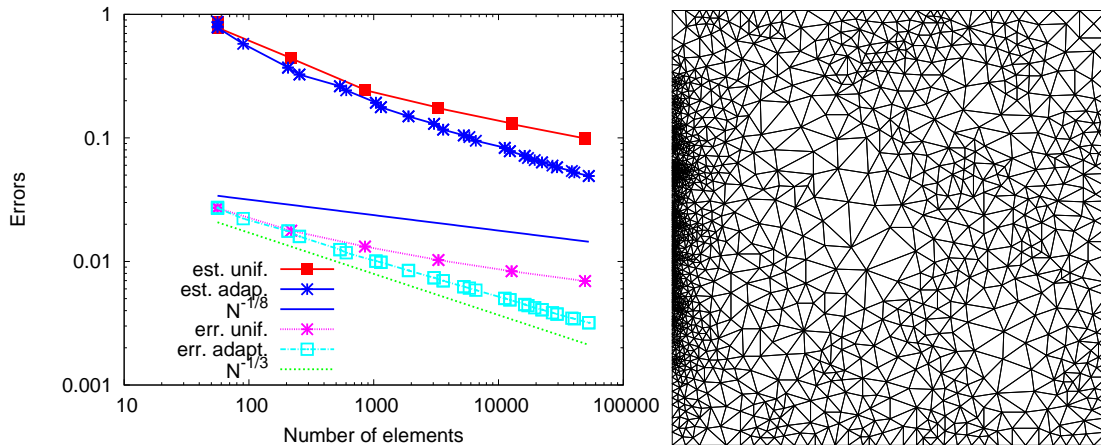


FIGURE 11. Estimated and exact errors in uniform/adaptive refinement (left) and an adapted mesh (right).

$10^{-3}$  to  $4.05 \times 10^{-7}$  for the uniform refinement). The convergence curve corresponding to the uniform mesh refinement is parallel to the  $N^{-1/8}$  straight line, while the convergence curve corresponding to the adaptive

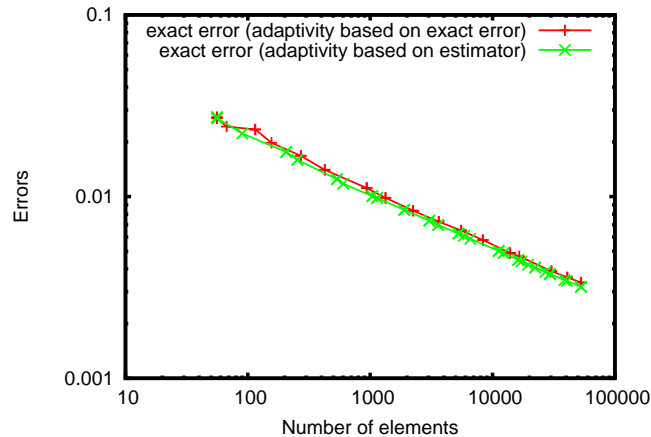


FIGURE 12. Exact errors in the adaptive process using the discretization estimator and exact error.

mesh refinement is parallel to the  $N^{-1/3}$  straight line. Moreover, the effectivity of the error estimators for both types of refinements is around 15.

We could have expected the plot of the exact error corresponding to an adaptive mesh refinement to be parallel to the optimal  $N^{-1/2}$  straight line, but this is not the case, just like in [16]. We would like to determine whether this problem is caused by our discretization estimator or not. For this, in Fig. 12, we compare the exact errors corresponding to two different adaptive refinement process: one is driven by our discretization estimator, while the other is driven by the exact error (which we may compute exactly in this test case). Clearly, the exact error is not mostly affected by the applied adaptive process.

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