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AN A POSTERIORI ERROR ESTIMATION FOR THE DISCRETE DUALITY
FINITE VOLUME DISCRETIZATION OF THE STOKES EQUATIONS

Anh Ha LE1, 2 and Pascal OMNES1, 2

Abstract. We derive an a posteriori error estimation for the discrete duality finite volume (DDFV)
discretization of the stationary Stokes equations on very general twodimensional meshes, when a penalty
term is added in the incompressibility equation to stabilize the variational formulation. Two different
estimators are provided: one for the error on the velocity and one for the error on the pressure.
They both include a contribution related to the error due to the stabilization of the scheme, and a
contribution due to the discretization itself. The estimators are globally upper as well as locally lower
bounds for the errors of the DDFV discretization. They are fully computable as soon as a lower bound
for the inf-sup constant is available. Numerical experiments illustrate the theoretical results and we
especially consider the influence of the penalty parameter on the error for a fixed mesh and also of
the mesh size for a fixed value of the penalty parameter. A global error reducing strategy that mixes the
decrease of the penalty parameter and adaptive mesh refinement is described.

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INTRODUCTION

Let Ω be a two-dimensional simply connected polygonal domain with boundary Γ. We consider the Stokes
equations

\[-\Delta \hat{u} + \nabla \hat{p} = f \text{ in } \Omega, \]
\[\nabla \cdot \hat{u} = 0 \text{ in } \Omega, \]
\[\hat{u} = g \text{ on } \Gamma, \]
\[\int_{\Omega} \hat{p}(x) dx = 0, \]

where \(\hat{u}\) is the fluid velocity, \(\hat{p}\) the pressure, \(f\) the body forces per unit mass, and the function \(g\) satisfies
\[\int_{\Gamma} g(\sigma) \cdot \mathbf{n} d\sigma = 0.\] With \(f \in H^{-1}(\Omega)\) and \(g \in H^{1/2}(\Gamma)\), this problem is well-posed (see [18]) due to the so-called

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inf-sup condition: there exists $\beta > 0$ such that:

$$
\beta = \inf_{q \in L^2_0(\Omega)} \sup_{v \in (H^1_0(\Omega))^2} \frac{\int_{\Omega} q \nabla \cdot v(x) \, dx}{\|v\|_{(H^1_0(\Omega))^2} \|q\|_{L^2(\Omega)}},
$$

in which $L^2_0(\Omega)$ is the set of $L^2$ functions over $\Omega$ verifying (4).

Our purpose in this work is to compute an a posteriori error estimation between the exact solution $\hat{u}$, $\hat{p}$ of (1)–(4) and its numerical approximation by the penalized discrete duality finite volume scheme (DDFV) as presented in [23]. Originally developed for linear diffusion equations [17, 20, 21], the DDFV method has been extended to nonlinear diffusion [2, 4, 10], convection-diffusion [11], electro-cardiology [1, 12], drift-diffusion and energy-transport models [6], electro- and magnetostatics [15], electromagnetism [22], and Stokes flows [14, 16, 23]. The originality of the DDFV method is that it is able to treat very general meshes, including very distorted, degenerating, or highly non-conforming meshes (see the numerical tests in [17]). The name of the method comes from the definition of discrete gradient and divergence operators which verify a discrete Green formula.

Like for other equations, the development of a posteriori error estimations for the Stokes problem has followed the a priori investigation of numerical methods. As far as finite elements methods are concerned, R. Verfürth [29] made one the very first contributions by getting two a posteriori error estimations for the mini-element discretization: one is based on a suitable evaluation of the residual, the other is based on the solution of local Stokes problems. Later on, R. Verfürth [30] generalized the first estimator developed in [29] to the non-conforming Crouzeix–Raviart finite element method, neglecting however the consistency error in the estimator. It was shown however in E. Dari et al. [13] that this consistency error may not always be neglected, and, in order to properly take it into account, the authors of [13] use a Helmholtz-Hodge like decomposition (adapted to the Stokes problem) of the velocity error. In the resulting error estimator, this gives rise to terms related to the jumps of the tangential velocity components from one cell to another, in addition to the usual jumps of the normal components of the stress tensor. The case of the non-conforming Fortin–Soulie quadratic elements is also treated in [13].

All the above-cited finite element methods satisfy a uniform discrete inf-sup condition. However, it is often found useful in practice to consider discretizations (especially low-order ones) that do not verify such a uniform discrete inf-sup condition. In this context, C. Bernardi et al. [3] consider the finite element approximation of the Stokes equations when a penalty term is added to stabilize the variational formulation. The a posteriori error estimation they obtain includes two contributions: one related to the discretization on a given mesh, the other related to the penalty term. Based on these two contributions, the mesh refinement and the decrease of the penalty term are linked within an adaptive process.

A very recent contribution by A. Hannukainen et al. [19] sets a general framework for obtaining a posteriori error estimations for the discretization of the Stokes equations. The method is based on the reconstruction of postprocessed $H^1_0$ conforming velocity and $H – \text{div}$ conforming stress tensor fields deduced from the numerical approximation, and it may be applied to various conforming and conforming stabilized finite element methods, the discontinuous Galerkin method, the Crouzeix–Raviart non-conforming finite element method, the mixed finite element method, and a general class of finite volume methods.

However, as far as finite volume methods are concerned, the use of arbitrary meshes in [19] requires first to solve local Stokes problems on a conforming subtriangulation of each control volume, and then to apply the above-cited reconstruction on this subtriangulation. Instead, we would like to obtain error estimates for the solution of the DDFV scheme presented in [23] without having to solve any local problem or to compute any reconstruction. To do this, we shall adapt to the Stokes problem the a posteriori error estimation investigated in [25] for the DDFV discretization of the Laplace equation, using the discrete variational formulation verified by this scheme. The non-conformity of the method is dealt with using the Helmholtz-Hodge like decomposition introduced in [13]. Our estimator also includes a contribution related to the stabilization term in the incompressibility equation, which allows to monitor the amplitude of the penalization coefficient with respect to the mesh refinement process.
Although a DDFV scheme for the Stokes equations was constructed and analyzed in the three dimensional case [24], we limit the discussion in the present article to two space dimensions, because, as emphasized in [24], the 3-D scheme is not a simple extension to three spatial dimensions of the 2-D scheme originally developed in [23], because it is based on a construction for the dual meshes and the diamond mesh that is very specific to 3-D.

This article is organized as follows. Section 1 sets some notations and definitions related to the meshes, to discrete differential operators, and to discrete functions. In section 2, we present the DDFV scheme, and its equivalent variational formula is recalled. In section 3, representations of the errors are elaborated. This is used in section 4 to find a computable upper bound of these errors, provided a lower bound for the inf-sup constant in (5) is known. Such estimations are available in [7–9]. In section 5, the local efficiency of the error estimators is verified. Numerical experiments are presented in section 6.

1. Notations and Definitions

Let Ω be covered by a primal mesh with polygonal cells denoted by $T_i$, $i \in [1, I]$. We associate with each $T_i$ a point $G_i$ located in the interior of $T_i$. With any vertex $S_k$ of the primal mesh, with $k \in [1, K]$, we associate a dual cell $P_k$ by joining points $G_i$ associated with the primal cells surrounding $S_k$ to the midpoints of the edges of which $S_k$ is a node. The notations are summarized in Fig. 1 and 2.

With any primal edge $A_j$ with $j \in [1, J]$, we associate a so-called diamond-cell $D_j$ obtained by joining the vertices $S_{k_1(j)}$ and $S_{k_2(j)}$ of $A_j$ to the points $G_{i_1(j)}$ and $G_{i_2(j)}$ associated with the primal cells that share $A_j$ as a part of their boundaries. When $A_j$ is a boundary edge (there are $J^t$ such edges), the associated diamond-cell is a flat quadrilateral (i.e. a triangle) and we denote by $G_{i_2(j)}$ the midpoint of $A_j$ (thus, there are $J^t$ such additional points $G_i$). The unit normal vector to $A_j$ is $n_j$ and points from $G_{i_1(j)}$ to $G_{i_2(j)}$. We denote by $A_{j_1}'$ (resp. $A_{j_2}'$) the segment joining $G_{i_1(j)}$ (resp. $G_{i_2(j)}$) and the midpoint of $A_j$. Its associated unit normal vector, pointing from $S_{k_1(j)}$ to $S_{k_2(j)}$, is denoted by $n_{j_1}'$ (resp. $n_{j_2}'$). In the case of a boundary diamond-cell, $A_{j_2}'$ reduces to $\{G_{i_2(j)}\}$ and does not play any role. Finally, for any diamond-cell $D_j$, we shall denote by $M_{i_1, k_2 j}$ the midpoints of $[G_{i_1(j)} S_{k_2(j)}]$, with $(\alpha, \beta) \in \{1; 2\}^2$. With $n_j$, $n_{j_1}'$ and $n_{j_2}'$, we associate orthogonal unit vectors $\tau_j$, $\tau_{j_1}'$ and $\tau_{j_2}'$, such that the corresponding orthonormal bases are positively oriented. For any primal cell $T_i$ such that $A_j \subseteq \partial T_i$, we shall define $n_{ji} := n_j$ if $i = i_1(j)$ and $n_{ji} := -n_j$ if $i = i_2(j)$, so that $n_{ji}$ is always exterior to $T_i$. With $n_{ji}$, we associate $\tau_{ji}$ such that $(n_{ji}, \tau_{ji})$ is positively oriented. Similarly, when $A_{j_1}'$ and $A_{j_2}'$ belong to $\partial P_k$, we define $(n'_{jk_1}, \tau'_{jk_1})$ and $(n'_{jk_2}, \tau'_{jk_2})$ so that $n'_{jk_1}$ and $n'_{jk_2}$ are orthogonal to $A'_{j_1}$ and $A'_{j_2}$ and exterior to $P_k$.

**NB:** by a slight abuse of notations, we shall write $i \in \Gamma$ (respectively $j \in \Gamma$ and $k \in \Gamma$) to mean $G_i \in \Gamma$, (resp. $A_j \subseteq \Gamma$ and $S_k \in \Gamma$). The same convention will be used for any set of points other than $\Gamma$; e.g. for $\partial T_i$. We shall write $j \in \partial P_k$ to mean that $A_{j_1}'$ and and $A_{j_2}'$ are subsets of $\partial P_k$.

For $v \in (H^2(\Omega))^2$ with $\mathbf{v} = (v_1, v_2)$, we define

$$\nabla \mathbf{v} = \begin{pmatrix} \frac{\partial v_1}{\partial x} & \frac{\partial v_1}{\partial y} \\ \frac{\partial v_2}{\partial x} & \frac{\partial v_2}{\partial y} \end{pmatrix}, \quad \nabla \times \mathbf{v} = \begin{pmatrix} \frac{\partial v_1}{\partial y} & -\frac{\partial v_1}{\partial x} \\ -\frac{\partial v_2}{\partial y} & \frac{\partial v_2}{\partial x} \end{pmatrix},$$

$$\Delta \mathbf{v} = \begin{pmatrix} \Delta v_1 \\ \Delta v_2 \end{pmatrix}.$$
For future use, we recall Green’s formulae
\[
\int_{\Omega} \Delta v \cdot w \, dx = - \int_{\Omega} \nabla v : \nabla w + \int_{\partial \Omega} (\nabla v \cdot n) \cdot w \, ds, \tag{6}
\]
\[
\int_{\Omega} \nabla v : \nabla \times w \, dx = - \int_{\partial \Omega} (\nabla v \cdot \tau) \cdot w \, ds, \tag{7}
\]
for any \( v \in (H^2(\Omega))^2 \) and \( w \in (H^1(\Omega))^2 \). Here, \( n \) is the outward normal to \( \partial \Omega \) and \( \tau \) is the tangent vector to \( \partial \Omega \) such that \((n, \tau)\) is positively oriented.

**Figure 1.** A non-conforming primal mesh (solid lines) and its associated dual mesh (left, dashed lines) and diamond mesh (right, dotted lines).

**Figure 2.** Notations for an inner diamond-cell (left) and a boundary diamond cell (right).

In the definition of the DDFV scheme, we shall associate the velocity unknowns to the points \( G_i \) and \( S_k \) and the pressure unknowns to the diamond-cells. Moreover the gradient and divergence of the velocity will be defined on the diamond-cells. This leads us to the following definitions.
Definition 1.1. Let $u = (u^T, u^D)$, and $v = (v^T, v^D)$ be in $((\mathbb{R}^2)^I \times (\mathbb{R}^2)^K$. Let $\Phi = (\Phi_j)$ and $\Psi = (\Psi_j)$ be in $((\mathbb{R}^2)^2)^J$. Let $p = (p_j)$ and $q = (q_j)$ be in $\mathbb{R}^j$. We define the following scalar products

\[
(u, v)_{T,P} := \frac{1}{2} \left( \sum_{j \in [I]} |T_j| u^T_i \cdot v^T_i + \sum_{k \in [1,K]} |P_k| u^P_i \cdot v^P_i \right),
\]

\[
(\Phi, \Psi)_D := \sum_{j \in [I]} |D_j| \Phi_j : \Psi_j \text{ and } (p, q)_D = \sum_{j \in [I]} |D_j| p_j q_j.
\]

Definition 1.2. Let $u = (u^T, u^P)$ be in $((\mathbb{R}^2)^I+J^T \times (\mathbb{R}^2)^K$. For any boundary edge $A_j$, with the notations of Fig. 2 (right), we define $\tilde{u}_j$ as the trace of $u$ over $A_j$ by

\[
\tilde{u}_j = \frac{1}{4} \left( u^P_{k_1(j)} + 2u^T_{i_2(j)} + u^P_{k_2(j)} \right).
\]

Let $u = (u^T, u^D)$ be in $((\mathbb{R}^2)^I+J^T \times (\mathbb{R}^2)^K$ and let $w = (w_j)$ be defined on the boundary $\Gamma$. We define the following boundary scalar product

\[
(w, \tilde{u})_{\Gamma_h} := \sum_{j \in \Gamma} |A_j| w_j \cdot \tilde{u}_j.
\]

Definition 1.3. Let $\Phi = (\Phi_j)$ be in $((\mathbb{R}^2)^2)^J$. We define divergences and curls of the tensor field $\Phi$ on the primal and dual cells by

\[
(\nabla^T_h \Phi) \overline{u} := \frac{1}{|T_i|} \sum_{j \in \partial T_i} |A_j| \Phi_j n_{j_i},
\]

\[
(\nabla^P_h \Phi) \overline{u} := \frac{1}{|P_k|} \left( \sum_{j \in \partial P_k} |A_j| \Phi_j n_{j_i} + |A_{j_2}| \Phi_j n_{j_2} + \sum_{j \in \partial P_k \cap \Gamma} |A_j| \Phi_j n_{j} \right),
\]

\[
(\nabla^T_h \Phi) \overline{u} := \frac{1}{|T_i|} \sum_{j \in \partial T_i} |A_j| \Phi_j \tau_{j_i},
\]

\[
(\nabla^P_h \Phi) \overline{u} := \frac{1}{|P_k|} \left( \sum_{j \in \partial P_k} |A_j| \Phi_j \tau_{j_1} + |A_{j_2}| \Phi_j \tau_{j_2} + \sum_{j \in \partial P_k \cap \Gamma} |A_j| \Phi_j \tau_{j} \right).
\]

Definition 1.4. Let $u_1 = ((u_1)^T, (u_1)^P)$ be in $\mathbb{R}^I+J^T \times \mathbb{R}^K$ and $u_2 = ((u_2)^T, (u_2)^P)$ be in $\mathbb{R}^I+J^T \times \mathbb{R}^K$, and $u = (u_1, u_2)$: the discrete gradient $\nabla^D_h u$ and the discrete curl $\nabla^D_h \times u$ are defined by their values in the diamond-cells $D_j$ by

\[
(\nabla^D_h u) \overline{j} = \left( \frac{(\nabla^D_h u_1) \overline{j}}{\nabla^D_h u_2} \right), \quad (\nabla^D_h \times u) \overline{j} = \left( \frac{(\nabla^D_h \times u_1) \overline{j}}{\nabla^D_h \times u_2} \right),
\]

where, for $\phi \in \mathbb{R}^I+J^T \times \mathbb{R}^K$, we define

\[
(\nabla^P_h \phi) \overline{j} := \frac{1}{2|D_j|} \left\{ \phi^P_{k_2(j)} - \phi^P_{k_1(j)} |A_j| n_{j_i} + |A_j| n_{j_2} + |A_{j_2}^T| - \phi^T_{i_2(j)} \right\} |A_j| n_{j_i},
\]

\[
(\nabla^D_h \times \phi) \overline{j} := -\frac{1}{2|D_j|} \left\{ \phi^P_{k_2(j)} - \phi^P_{k_1(j)} |A_j| \tau_{j_1} + |A_j| \tau_{j_2} + |A_{j_2}^T| - \phi^T_{i_2(j)} \right\} |A_j| \tau_{j}.
\]
We also need a discrete divergence of a vector, which is defined using the discrete gradient
\[(\nabla_h^D \cdot \mathbf{u})_j = \text{Trace}((\nabla_h^D \mathbf{u})_j).\]

From basic geometrical arguments, we obtain some properties of the discrete gradient:
\[
\phi^P_{k_2(j)} - \phi^P_{k_1(j)} = (\nabla_h^D \phi)_j \cdot \Delta_{k_2(j)} \Delta_{k_1(j)}, \quad \phi^T_{i_2(j)} - \phi^T_{i_1(j)} = (\nabla_h^D \phi)_j \cdot \Delta_{G_{i_2(j)} G_{i_1(j)}}. \tag{12}
\]

For the penalization of the scheme, we need to define the following (non-consistent) Laplacian-type operator.

**Definition 1.5.** Let \( p = (p_j) \in \mathbb{R}^J \), we define:
\[
(\Delta_h^P p)_j = \frac{1}{|D_j|} \sum_{j' \in \partial D_j} \frac{d_j^2 + d_j'^2}{d_j} (p_{j'} - p_j), \tag{13}
\]
where \( \partial D_j \) is the set of indexes of diamond cells which have a common boundary segment with \( D_j \) and \( d_j = \text{diam}(D_j) \) and \( d_j' = \text{diam}(D_{j'}) \).

**Proposition 1.6.** For \( \Phi \in (\mathbb{R}^{2 \times 2})^J \) and \( \mathbf{u}^T, \mathbf{u}^P \in (\mathbb{R}^2)^{I + J^r} \times (\mathbb{R}^2)^K \) and \( p \in \mathbb{R}^J \), the following discrete Green formula hold:
\[
(\nabla^T_{h \cdot \mathbf{u}} = -(\nabla_h^D \mathbf{u}, \Phi)_D + (\Phi \mathbf{n}, \tilde{\mathbf{u}})_{\Gamma, h}, \tag{14}
(\nabla^T_{h \cdot \mathbf{u}} = \nabla_h^D \mathbf{u}, \Phi)_D + (\Phi \mathbf{\tau}, \tilde{\mathbf{u}})_{\Gamma, h}, \tag{15}
(\nabla^T_{h \cdot \mathbf{u}} = -(\nabla_h^D \mathbf{u}, \Phi)_D + (p \mathbf{n}, \tilde{\mathbf{u}})_{\Gamma, h}, \tag{16}
where \( I_2 \) is the 2 \( \times \) 2 identity matrix and \( \nabla^T_{h \cdot \mathbf{u}} \) and \( \nabla^T_{h \cdot \mathbf{u}} \) stand for \( \nabla_h^T \) and \( \nabla_h^T \) on the primal cells and for \( \nabla_h^T \) and \( \nabla_h^T \) on the dual cells.

Formula (14) is called the discrete Stokes formula and is proved in [23]; its componentwise counterpart can be found in [17]. Formulæ (15) and (16) can be demonstrated in the same way. Proposition 1.7 below may be found componentwise in [15].

**Proposition 1.7.** For all \( \mathbf{u} = (\mathbf{u}_T^I, \mathbf{u}_k^P) \in (\mathbb{R}^2)^{I + J^r} \times (\mathbb{R}^2)^K \), it holds that
\[
(\nabla_h^T \mathbf{u})_i = 0, \quad \forall i \in [1, I], \tag{17}
(\nabla_h^P \mathbf{u})_k = 0, \quad \forall k \in \Gamma. \tag{18}
\]
In addition, for \( k \in \Gamma \), the following equality holds (see Fig. 3 for the notations)

\[
(\nabla_h^P \times (\nabla_h^P u)_k) = \frac{1}{|P_k|} \left[ (u_{i_1(k)}^T - u_{i_2(k)}^T) + \frac{1}{2} (u_{K_1(k)}^P - u_{K_2(k)}^P) \right].
\]

The derivation of the a posteriori error estimates is based on the reformulation of the scheme under a variational form which uses functions associated to the discrete unknowns.

**Definition 1.8.** With any \( u = (u^T_1, u^P_k) \in (\mathbb{R}^2)^{I + J} \times (\mathbb{R}^2)^K \), we associate the function \( u_h \) defined by

\[
(u_h)|_{D_j} \in (P^1(D_j))^2, \quad \forall j \in [1, J],
\]

\[
u_h(M_{\alpha(j)k\beta(j)}) = \frac{1}{2} (u_{\alpha(j)}^T + u_{\beta(j)}^P), \quad \forall j \in [1, J], \quad (\alpha, \beta) \in \{1, 2\}^2.
\]

In addition, for all \( p = (p_j) \in \mathbb{R}^J \), we construct piecewise constant functions corresponding to the approximate pressure, penalty term and to the diamond-cell diameter \( d_j \):

\[
p_h(x) = p_j, \quad \forall x \in D_j, \quad j \in [1, J],
\]

\[
(\nabla_h^P p)_h(x) = (\nabla_h^P p)_j, \quad \forall x \in D_j, \quad j \in [1, J],
\]

\[
d_h(x) = d_j, \quad \forall x \in D_j, \quad j \in [1, J].
\]

The validity of the definition of \( u_h \) by its values in four different points and the proof of the fundamental properties below, which allow to reformulate the scheme under a variational form, may be found in [17].

**Proposition 1.9.** Let \( u = (u^T_1, u^P_k) \in (\mathbb{R}^2)^{I + J} \times (\mathbb{R}^2)^K \) and let \( u_h \) be defined by Definition 1.8. It holds that

\[
(\nabla_h^P u)_j = \nabla u_h|_{D_j}, \quad \forall j \in [1, J],
\]

\[
(\nabla_h^P u)_j = \nabla \cdot u_h|_{D_j}, \quad \forall j \in [1, J].
\]

**Definition 1.10.** Let \( u_h \) be a broken affine function on the diamond mesh: \((u_h)|_{D_j} \in (P^1(D_j))^2, \quad \forall j \in [1, J],\) and \( u_h \) is not necessarily continuous over the interfaces of neighboring diamond-cells. We define its broken gradient and divergence over \( \Omega \) by:

\[
\nabla_h u_h(x) = \nabla u_h|_{D_j}(x) \quad \text{and} \quad \nabla_h \cdot u_h(x) = \nabla \cdot u_h|_{D_j}(x), \quad \forall x \in D_j, \quad j \in [1, J].
\]

2. **The Finite Volume Scheme on General Meshes**

The finite volume scheme used for the numerical approximation of equations (1)–(4) is constructed on the basis of the discrete differential operators defined in section 1. It is very similar to that studied in [23], where proofs of existence, uniqueness, stability with respect to the data and a priori error estimates can be found as soon as \( \varepsilon > 0 \).

\[
(\nabla_h^T \cdot (-\nabla_h^P u + pI_2))_i = f_i^T, \quad \forall i \in [1, I],
\]

\[
(\nabla_h^P \cdot (-\nabla_h^P u + pI_2))_k = f_k^P, \quad \forall k \in [1, K],
\]

\[
(\nabla_h^P \cdot u)_j - \varepsilon d_j^2 (\nabla_h^P p)_j = 0, \quad \forall j \in [1, J],
\]

\[
\frac{u_{k_1(j)}^P + 2u_{i_1(j)}^T + u_{k_2(j)}^P}{4} = g_j, \quad \forall j \in \Gamma,
\]

\[
\sum_{j=1}^{J} |D_j|p_j = 0.
\]
We suppose that $g$ is regular enough, so that we can set $g_j = g(G_{2j})$ in (29), while in (26) and (27), $f_i^T$ and $f_k^P$ are the mean values of $f$ over $T_i$ and $P_k$, respectively:

\[
f_i^T = \frac{1}{|T_i|} \int_{T_i} f(x) dx \quad \text{and} \quad f_k^P = \frac{1}{|P_k|} \int_{P_k} f(x) dx.
\]

(31)

We can prove that the solution of the scheme verifies a discrete variational formulation:

**Proposition 2.1.** Let $u = (u^T_i, u^P_k)$ and $p = (p_j)_{j \in [1,J]}$ be the solution of the scheme (26)-(30). Let $v = (v^T_i, v^P_k)$ such that $\bar{v}_j = 0$ for all $j \in \Gamma$. Let $u_h$ and $v_h$ be the solution associated to $u$ and $v$ by Definition 1.8. Let us set in addition

\[
v^T_h(x) := \frac{1}{2} \left( \sum_{i \in [1, I]} v^T_i \theta^T_i(x) + \sum_{k \in [1, K]} v^P_k \theta^P_k(x) \right),
\]

where $\theta^T_i$ and $\theta^P_k$ are respectively the characteristic function of the cells $T_i$ and $P_k$. Then, it holds that

\[
\sum_j \int_{D_j} \nabla_h u_h : \nabla_h v_h(x)dx - \sum_j \int_{D_j} \nabla_h \cdot v_h p_h(x)dx = \int_\Omega f \cdot v^T_h(x)dx.
\]

(33)

**Proof.** Starting from Eq. (26) and (27), we have

\[-(\nabla^T_h \cdot (\nabla^D_h u_i), v^T_i) + (\nabla^P_k \cdot (pI_2)_k) \cdot v^P_k = f^T_i \cdot v^T_i, \quad \forall i \in [1, I],
\]

(34)

\[-(\nabla^P_k \cdot (\nabla^D_h u))_i \cdot v^P_k + (\nabla^P_k \cdot (pI_2)_k) \cdot v^P_k = f^P_k \cdot v^P_k, \quad \forall k \in [1, K].
\]

(35)

Multiplying (34) by $|T_i|$ and (35) by $|P_k|$ and summing over all $i$ and all $k$, we obtain

\[-(\nabla^T_p \cdot (\nabla^D_h u), v)_{T,P} + (\nabla^T_p \cdot (pI_2), v)_{T,P} = (f, v)_{T,P}.
\]

We can apply (14), (16) and $\bar{v}_j = 0$ for all $j \in \Gamma$ and we obtain

\[(\nabla^D_h u, \nabla^P_h v)_D - (\nabla^D_h \cdot v, p)_D = (f, v)_{T,P}.
\]

Using (23)-(24) and definitions (31)-(32), we obtain (33). □

3. A REPRESENTATION OF THE ERROR

3.1. A representation of the velocity error

The variational formulation of (1) reads:

\[
\int_\Omega \nabla \hat{u} : \nabla vdx - \int_\Omega \hat{p} \nabla \cdot vdx = \int_\Omega f \cdot v(x)dx,
\]

(36)

for all $v \in (H^1_0(\Omega))^2$. We shall estimate the $H^1$ semi norm of the error between the exact solution $\hat{u}$ and the function $u_h$ associated to the solution of the DDFV scheme. For this, we shall denote by $e := \hat{u} - u_h$ and $\epsilon_p := \hat{p} - p_h$ the error in the velocity and pressure, respectively. We have

\[
\| \nabla_h e \|_{L^2(\Omega)} = \left( \sum_j \int_{D_j} |\nabla \hat{u} - \nabla u_h|^2(x)dx \right)^{1/2}.
\]

(37)
Since $\Omega$ is a simply connected domain and since $\nabla_h e = \nabla \hat{u} - \nabla_h u_h$ belongs to $(L^2(\Omega))^2$, we may decompose it in the following way (see Lemma 3.2 in [13]):

$$\nabla_h e = \nabla \hat{\Phi} - qI_2 + \nabla \times \hat{\Psi}, \quad (38)$$

where $q \in L^2_0(\Omega)$, $\hat{\Phi} \in (H^1(\Omega))^2$ with $\nabla \cdot \hat{\Phi} = 0$ and $\hat{\Psi} \in (H^1(\Omega))^2$ with $\int_\Omega \hat{\Psi}(x)dx = 0$, with the following estimations

$$\|\nabla \hat{\Phi}\|_{L^2(\Omega)} \leq \|\nabla_h e\|_{L^2(\Omega)}, \quad (39)$$

$$\|q\|_{L^2(\Omega)} \leq \frac{2 \sqrt{2}}{\beta} \|\nabla_h e\|_{L^2(\Omega)},$$

$$\|\nabla \times \hat{\Psi}\|_{L^2(\Omega)} \leq (1 + \frac{2 \sqrt{2}}{\beta})\|\nabla_h e\|_{L^2(\Omega)},$$

where $\beta$ is defined by (5).

Now, we estimate the velocity error using decomposition (38). First observe that

$$\int_\Omega \nabla_h e : I_2q(x)dx = \int_\Omega \nabla \cdot e q(x)dx = \int_\Omega (\nabla \cdot \hat{u} - \nabla_h \cdot u_h)q(x)dx.$$

From (2) and (28), we have

$$\int_\Omega \nabla_h e : I_2q(x)dx = \varepsilon \int_\Omega d_h^2(x)(\Delta_D p)q(x)dx, \quad (40)$$

where we recall that the function $d_h$ is defined through (22). Multiplying the term $\nabla_h e(x)$ with (38) side by side and integrating over $\Omega$, it holds that

$$\|\nabla_h e\|_{L^2(\Omega)}^2 = \int_\Omega \nabla_h e : (\nabla \hat{\Phi} + \nabla \times \hat{\Psi} - qI_2)dx$$

$$= i_1 + i_2 - \varepsilon \int_\Omega d_h^2(x)(\Delta_D p)q(x)dx, \quad (41)$$

where

$$i_1 = \sum_j \int_{D_j} (\nabla \hat{u} - \nabla u_h) : \nabla \hat{\Phi}(x)dx$$

and

$$i_2 = \sum_j \int_{D_j} (\nabla \hat{u} - \nabla u_h) : \nabla \times \hat{\Psi}(x)dx.$$
but then fixed in what follows, and \[([\nabla_hu_h - p_hI_2)s]_s\] is the jump of the normal component of \(\nabla_hu_h - p_hI_2\) through segment \(s\). Moreover, \(\tau_s\) will be such that \((n_s, \tau_s)\) is a positively oriented orthonormal basis of \(\mathbb{R}^2\) and \([([\nabla_hu_h])_s]_s\) is the jump of the tangential component of \(\nabla_hu_h\) through segment \(s\).

**Proposition 3.2.** Let \(\hat{\Phi}\) be defined in equation (38). Let \(\Phi = (\Phi^T, \Phi^P) \in (\mathbb{R}^2)^{I+J_F} \times (\mathbb{R}^2)^K\) be such that

\[
\hat{\Phi}_j = 0 \text{ for all } j \in \Gamma.
\]  

(42)

The values \((\Phi^T_1, \Phi^P_1)\) are up to now unrelated to \(\hat{\Phi}\), they will be chosen in section 4.2. Then, it holds that

\[
i_1 = \frac{1}{2} \sum_{i \in [1,I]} \int_{T_i} f \cdot (\hat{\Phi} - \Phi^T_i)(x)dx + \frac{1}{2} \sum_{k \in [1,K]} \int_{P_k} f \cdot (\hat{\Phi} - \Phi^P_k)(x)dx
\]

\[
- \frac{1}{2} \sum_{i \in [1,I]} \sum_{s \subset T_i} \int_{s} ((\nabla_hu_h - p_hI_2)s) \cdot (\hat{\Phi} - \Phi^T_i)\sigma d\sigma
\]

\[
- \frac{1}{2} \sum_{k \in [1,K]} \sum_{s \subset P_k} \int_{s} ((\nabla_hu_h - p_hI_2)s) \cdot (\hat{\Phi} - \Phi^P_k)\sigma d\sigma.
\]

(43)

**Proof.** First, since \(\hat{\Phi} \in (H_0^1(\Omega))^2\) and \(\nabla \cdot \hat{\Phi} = 0\), using (36) yields:

\[
i_1 = \sum_j \int_{D_j} \nabla u : \nabla \Phi(x)dx - \sum_j \int_{D_j} \nabla_hu_h : \nabla \hat{\Phi}(x)dx
\]

\[
= \int_\Omega \nabla u : \nabla \Phi(x)dx - \sum_j \int_{D_j} \nabla_hu_h : \nabla \hat{\Phi}(x)dx
\]

\[
= \int_\Omega f \cdot \Phi(x)dx - \sum_j \int_{D_j} \nabla_hu_h : \nabla \Phi(x)dx.
\]

For any \(\Phi = (\Phi^T, \Phi^P)\) satisfying (42), formula (33) leads to

\[
i_1 = \int_\Omega f \cdot (\Phi - \Phi^*_h)(x)dx - \sum_j \int_{D_j} p_h \nabla_h \Phi_h(x)dx
\]

\[
- \sum_j \int_{D_j} \nabla_hu_h : (\nabla \hat{\Phi} - \nabla \Phi_h)(x)dx.
\]

We know that \(p_h \nabla_h \cdot \Phi_h = p_hI_2 : \nabla \Phi_h\) and \(p_hI_2 : \nabla \Phi = 0\) (since \(\nabla \cdot \hat{\Phi} = 0\)), then

\[
i_1 = \int_\Omega f \cdot (\Phi - \Phi^*_h)(x)dx - \sum_j H_1(j),
\]

(44)

where

\[
H_1(j) = \int_{D_j} ((\nabla_hu_h - p_hI_2)s) : (\nabla \hat{\Phi} - \nabla \Phi_h)(x)dx.
\]

(45)
Let us consider a diamond-cell \( D_j \). Since \( \nabla_h u_h - p_h I_2 \) is constant over \( D_j \), we may write, using Green’s formula over \( D_j \),

\[
H_1(j) = \int_{\partial D_j} (\nabla_h u_h - p_h I_2) n_{\partial D_j} \cdot (\hat{\Phi} - \Phi_h)(\sigma) d\sigma,
\]

where \( n_{\partial D_j} \) is the unit normal vector exterior to \( D_j \) on its boundary. Moreover, let \( s \) be any of the four boundary edges of \( D_j \), the function \( \Phi_h \) belongs to \( P^1 \) over \( s \) and the quantity \( \nabla_h u_h - p_h I_2 \) is a constant; the integral of \( (\nabla_h u_h - p_h I_2) n_{\partial D_j} \cdot \Phi_h \) along this edge may thus exactly be computed by the midpoint rule; using the definition of \( \Phi_h \), this function equals \( \frac{1}{2}(\Phi_{T(s)} + \Phi_{k(s)}) \) at the midpoint of \( s \). Thus, it holds that:

\[
H_1(j) = \sum_{s \subset \partial D_j} \int_s (\nabla_h u_h - p_h I_2) n_{s,j} \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) d\sigma,
\]

(46)

where \( n_{s,j} \) is the unit normal vector exterior to \( D_j \) on \( s \).

In (44), in the sum of the \( H_1(j) \) over \( j \in \{1, J\} \), there are two types of edges \( s \) : those in \( S \) and those included in \( \Gamma \). First, each \( s \in S \) is the common edge of two diamond-cells; then, in the sum, there are two corresponding integrals over \( s \), in which we can factorize by \( \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) \). Indeed, the jump of this function through \( s \) vanishes because \( \hat{\Phi} \in (H^1_0(\Omega))^2 \). This implies:

\[
\sum_j \sum_{s \subset \partial D_j} \int_s (\nabla_h u_h - p_h I_2) n_{s,j} \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) d\sigma = \sum_{s \in S} \int_s ((\nabla_h u_h - p_h I_2) n_s) \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) d\sigma.
\]

(47)

Secondly, each diamond-cell \( D_j \) whose boundary intersects \( \Gamma \) has two edges of equal length \( s = [G_{i2}(j)S_{k_\beta(j)}] \) with \( \beta \in \{1, 2\} \) which are included in \( \Gamma \), and their union is exactly \( A_j \). Since \( (\nabla_h u_h - p_h I_2) n_j \) is a constant on \( A_j \), and since \( \sum_{\beta \in \{1, 2\}} \int_{[G_{i2}(j)S_{k_\beta(j)}]} ((\hat{\Phi} - \frac{1}{2} \left( \Phi_{T(i_2(j))} + \Phi_{k_\beta(j)} \right)) (\sigma) d\sigma = \int_{A_j} (\hat{\Phi} - \Phi_h)(\sigma) d\sigma \), we have

\[
\sum_{s \subset \partial D_j \cap \Gamma} \int_s (\nabla_h u_h - p_h I_2) n_{s,j} \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) d\sigma
\]

\[
= \sum_{\beta \in \{1, 2\}} \int_{[G_{i2}(j)S_{k_\beta(j)}]} (\nabla_h u_h - p_h I_2) n_j \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(i_2(j))} + \Phi_{k_\beta(j)} \right) \right) (\sigma) d\sigma
\]

\[
= \int_{A_j} (\nabla_h u_h - p_h I_2) n_j \cdot (\hat{\Phi} - \Phi_h)(\sigma) d\sigma = 0,
\]

(48)

thanks to (42) and to the fact that \( \hat{\Phi} \in (H^1_0(\Omega))^2 \). With (47) and (48), we can write

\[
\sum_{j \in \{1, J\}} H_1(j) = \sum_{s \in S} \int_s ((\nabla_h u_h - p_h I_2) n_s) \cdot \left( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) \right) (\sigma) d\sigma.
\]

(49)

Then, we may write \( \hat{\Phi} - \frac{1}{2} \left( \Phi_{T(s)} + \Phi_{k(s)} \right) = \frac{1}{2} \left( \hat{\Phi} - \Phi_{T(s)} \right) + \left( \hat{\Phi} - \Phi_{k(s)} \right) \). Summing in the right-hand side of (49) the various contributions of \( \Phi_{T(s)} \) for a fixed \( i \) and the various contributions of \( \Phi_{k(s)} \) for a fixed \( k \), we obtain
the following formula

\[
\sum_j H_1(j) = \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset T_i} \int_s \left[ (\nabla_h u_h - p_h I_2) n_s \right] \cdot (\tilde{\Phi} - \Phi_T^T)(\sigma) d\sigma \\
+ \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset P_k} \int_s \left[ (\nabla_h u_h - p_h I_2) n_s \right] \cdot (\tilde{\Phi} - \Phi_k^P)(\sigma) d\sigma.
\]

(50)

Finally, according to (44) and definition (32) of \(\Phi_h^*\), we obtain (43).

Now, we turn to a representation formula for \(i_2\) in (41).

**Proposition 3.3.** Let \(u = (u_i^T, u_k^r)\) be the velocity component of the solution of the scheme (26)–(30) and \(u_h\) the function associated to \(u\) by Def. 1.8. Let \(\hat{\Psi}\) be defined in (38). Let \(\Psi = (\Psi_1^T, \Psi_k^P) \in (\mathbb{R}^2)^{I+J} \times (\mathbb{R}^2)^K\) and \(\Psi_h\) be its associated function. The values \((\Psi_1^T, \Psi_k^P)\) are up to now unrelated to \(\hat{\Psi}\), they will be chosen in section 4.2. Then, the following representation holds

\[
i_2 = \frac{1}{2} \sum_{i \in [1, I]} \sum_{s \subset T_i} \int_s [\nabla_h u_h \tau_s]_s \cdot (\hat{\Psi} - \Psi_T^T)(\sigma) d\sigma \\
+ \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset P_k} \int_s [\nabla_h u_h \tau_s]_s \cdot (\hat{\Psi} - \Psi_k^P)(\sigma) d\sigma \\
- \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} (\nabla g - \nabla_h u_h(\sigma)) \tau_k \cdot (\hat{\Psi}(\sigma) - \Psi_k^P) d\sigma.
\]

(51)

**Proof.** From (41), it holds that

\[
i_2 = \sum_j \int_{D_j} (\nabla \hat{u} - \nabla_h u_h) : \nabla \times \hat{\Psi}(x) dx \\
= \int_{\Omega} \nabla \hat{u} : \nabla \times \hat{\Psi}(x) dx - \sum_j \int_{D_j} \nabla_h u_h : \nabla_h \times \Psi_h(x) dx - \sum_j H_2(j),
\]

(52)

where

\[
H_2(j) = \int_{D_j} \nabla_h u_h : (\nabla \times \hat{\Psi} - \nabla_h \times \Psi_h)(x) dx.
\]

(53)

By application of the continuous Green formula, it holds that

\[
\int_{\Omega} \nabla \hat{u} : \nabla \times \hat{\Psi}(x) dx = - \int_{\partial \Omega} \nabla \hat{u} \cdot \tau \cdot \hat{\Psi}(\sigma) d\sigma = - \int_{\Gamma} \nabla g \cdot \tau \cdot \hat{\Psi}(\sigma) d\sigma.
\]

(54)

We can evaluate the sum of \(H_2(j)\) over \(j\) just like we evaluated the sum of \(H_1(j)\) in Proposition 3.2. There are only two differences. The first is that the gradients of \(\hat{\Phi}\) and \(\Phi_h\) are replaced by the curls of \(\hat{\Psi}\) and \(\Psi_h\), which implies that normal vectors \(n_s\) are replaced by tangent vectors \(-\tau_s\). The second difference is that the boundary integrals do not vanish any more, but can be evaluated like in the discussion that leads to (48). Then, noting that

\[
\sum_{j \in J} \int_{A_j} \nabla_h u_h \tau_j \cdot \hat{\Psi}(\sigma) d\sigma = \sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} \nabla_h u_h \tau_k \cdot \hat{\Psi}(\sigma) d\sigma,
\]
where $\tau_k$ is the tangent vector to $\partial P_k \cap \Gamma$ which is positively oriented with respect to the unit normal vector exterior to $\partial P_k \cap \Gamma$, we obtain the following formula

$$
\sum_j H_2(j) = -\frac{1}{2} \sum_{k \in [1, K]} \sum_{s \in T_k} \int_{s}^\triangledown h u_h \tau_s \cdot (\tilde{\Psi} - \Psi^T_k)(\sigma) d\sigma
$$

$$
-\frac{1}{2} \sum_{k \in [1, K]} \sum_{s \in P_k} \int_{s}^\triangledown h u_h \tau_s \cdot (\tilde{\Psi} - \Psi^P_k)(\sigma) d\sigma
$$

(55)

Moreover, it holds that (see the proof of Lemma 3.5 below):

$$
\text{Lemma 3.4.}
$$

Combining (52), (54), (55), and (56), we obtain (51).

In order to prove (56), we need some technical lemmas related to the $L^2(\Omega)$ scalar product of discrete gradients and curls.

**Lemma 3.4.** Let $\mathbf{u} = (u_k^T, u_k^P)$ be the velocity component of the solution of the scheme (26)–(30) and $\Psi = (\Psi_k^T, \Psi_k^P) \in (\mathbb{R}^2)^{I+J'} \times (\mathbb{R}^2)^K$. It holds that

$$
(\nabla_h^T \mathbf{u}, \Psi)_{T,P} = -\sum_{k \in \Gamma} \int_{\partial P_k \cap \Gamma} (\nabla_h \mathbf{g}(\sigma) - \triangledown h \mathbf{u}_h(\sigma)) \tau_k \cdot \Psi^P_k d\sigma.
$$

(57)

**Proof.** According to Eq. (17) and (18), it holds that

$$
(\nabla_h \times (\nabla_h^P \mathbf{u}))_i = 0, \forall i \in [1, I] \quad \text{and} \quad (\nabla_h^P \times (\nabla_h \mathbf{u}))_k = 0, \forall k \notin \Gamma.
$$

(58)

On the other hand, since the solution of the discrete problem verifies (29), there holds, for $k \in \Gamma$, with the notations of Fig. 3:

$$
\mathbf{u}_k^T(\cdot) = 2\mathbf{g}(G_{I_k}(\cdot)) - \frac{1}{2}(\mathbf{u}_k^P + \mathbf{u}_K^P_{K_1}(\cdot)) \quad \text{and} \quad \mathbf{u}_k^P(\cdot) = 2\mathbf{g}(G_{I_2}(\cdot)) - \frac{1}{2}(\mathbf{u}_k^P + \mathbf{u}_K^P_{K_2}(\cdot)).
$$

(59)

Following (19) and (59), we obtain that

$$
(\nabla_h^P \times (\nabla_h^P \mathbf{u}))_k = \frac{1}{|P_K|} \left[ 2(\mathbf{g}(G_{I_2}(\cdot)) - \mathbf{g}(G_{I_1}(\cdot))) + (\mathbf{u}_K^P_{K_1}(\cdot) - \mathbf{u}_K^P_{K_2}(\cdot)) \right], \forall k \in \Gamma.
$$

(60)

From (58), and using the definition of the scalar product in (8), we obtain

$$
(\nabla_h^T \mathbf{u}, \Psi)_{T,P} = \frac{1}{2} \sum_{k \in \Gamma} |P_K| (\nabla_h^P \times (\nabla_h^P \mathbf{u}))_k \cdot \Psi^P_k.
$$

Using (60) leads to

$$
(\nabla_h^T \mathbf{u}, \Psi)_{T,P} = \sum_{k \in \Gamma} (\mathbf{g}(G_{I_2}(\cdot)) - \mathbf{g}(G_{I_1}(\cdot))) \cdot \Psi^P_k + \frac{1}{2} \sum_{k \in \Gamma} (\mathbf{u}_K^P_{K_1}(\cdot) - \mathbf{u}_K^P_{K_2}(\cdot)) \cdot \Psi^P_k.
$$

(61)
In addition, we have
\[ g(G_{I_2(k)}) - g(G_{I_1(k)}) = (g(G_{I_2(k)}) - g(S_k)) + (g(S_k) - g(G_{I_1(k)})), \]
so that
\[ g(G_{I_2(k)}) - g(G_{I_1(k)}) = -\int_{\partial P_k \cap \Gamma} \nabla g(\sigma) \tau_k d\sigma, \] (62)

In the same way, we have
\[ u^P_{K_1(k)} - u^P_{K_2(k)} = (u^P_{K_1(k)} - u^P_k) + (u^P_k - u^P_{K_2(k)}). \]

Applying property (12) of the discrete gradient, it holds that
\[ u^P_{K_1(k)} - u^P_k = |S_k S_{K_1(k)}| \nabla_h u_h(\sigma) \tau_k = 2|S_k G_{I_1(k)}| \nabla_h u_h(\sigma) \tau_k, \forall \sigma \in [S_k S_{K_1(k)}] \] (63)
and
\[ u^P_k - u^P_{K_2(k)} = |S_{K_2(k)} S_k| \nabla_h u_h(\sigma) \tau_k = 2|S_k G_{I_2(k)}| \nabla_h u_h(\sigma) \tau_k, \forall \sigma \in [S_{K_2(k)} S_k]. \] (64)

As a consequence, it holds that
\[ u^P_{K_1(k)} - u^P_{K_2(k)} = 2\int_{\partial P_k \cap \Gamma} \nabla_h u_h(\sigma) \tau_k d\sigma. \] (65)

Combining (61) with (62) and (65), we obtain (57). \( \square \)

**Lemma 3.5.** Let \( \mathbf{u} = (\mathbf{u}^T, \mathbf{u}^P_k) \) be the velocity component of the solution of the scheme (26)-(30) and \( \Psi = (\Psi^T, \Psi^P_k) \in (\mathbb{R}^2)^{I+J^P} \times (\mathbb{R}^2)^K \). Let \( u_h \) and \( \Psi_h \) be their associated functions through Def. 1.8. Then formula (56) is true.

**Proof.** Applying the discrete Green formula (15), it holds that
\[
\sum_j \int_{\partial D_j} \nabla_h \mathbf{u}_h : \nabla_h \times \Psi_h(x) dx = (\nabla_h^D \mathbf{u}, \nabla_h^D \times \Psi)_D
= (\nabla_h^{T,P} \times (\nabla_h^D \mathbf{u}, \Psi)_{T,P} - (\nabla_h^D \mathbf{u} \cdot \tau, \Psi)_{T,h}. \]

Using the result of Lemma 3.4, we obtain (56). \( \square \)

### 3.2. A representation of the pressure error

We shall estimate the \( L^2 \) norm of the error between the exact solution \( \hat{p} \) and the function associated to the solution of the DDFV scheme by (20). We recall that \( e_p = \hat{p} - p_h \) and \( e = \hat{u} - u_h \).
Proposition 3.6. Let \( \hat{\mathbf{v}} \in (H^1_0(\Omega))^2 \) and \( \mathbf{v} = (\mathbf{v}_i^T, \mathbf{v}_k^P) \in (\mathbb{R}^2)^{l} \times (\mathbb{R}^2)^{K} \) be such that \( \hat{\mathbf{v}}_j = 0 \) for all \( j \in \Gamma \). The values \( (\mathbf{v}_i^T, \mathbf{v}_k^P) \) are up to now unrelated to \( \hat{\mathbf{v}} \), they will be chosen in section 4.3. We have that

\[
\int_{\Omega} e_p \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} = \int_{\Omega} \nabla_h \hat{\mathbf{e}} : \nabla \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \int_{\Omega} \nabla_h \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \frac{1}{2} \sum_{i \in [1, l]} \int_{T_i} f : (\hat{\mathbf{v}} - \mathbf{v}_i^T)(\mathbf{x})d\mathbf{x} - \frac{1}{2} \sum_{k \in [1, K]} \int_{P_k} f : (\hat{\mathbf{v}} - \mathbf{v}_k^P)(\mathbf{x})d\mathbf{x} + \frac{1}{2} \sum_{i \in [1, l]} \sum_{s \subset T_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2)_s] \cdot (\hat{\mathbf{v}} - \mathbf{v}_i^T)(\sigma)d\sigma + \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset P_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2)_s] \cdot (\hat{\mathbf{v}} - \mathbf{v}_k^P)(\sigma)d\sigma.
\]

(66)

Proof. We can use formula (36) to obtain

\[
\int_{\Omega} e_p \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} = \int_{\Omega} \hat{\mathbf{e}} \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \int_{\Omega} p_h \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} = \int_{\Omega} \nabla_h \hat{\mathbf{e}} : \nabla \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \int_{\Omega} \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} + \int_{\Omega} \nabla_h \mathbf{u}_h : \nabla \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \int_{\Omega} p_h \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x}.
\]

(67)

Using Eq. (33), we have

\[
\int_{\Omega} e_p \nabla \cdot \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} = \int_{\Omega} \nabla_h \hat{\mathbf{e}} : \nabla \hat{\mathbf{v}}(\mathbf{x})d\mathbf{x} - \int_{\Omega} \mathbf{f} : (\hat{\mathbf{v}} - \mathbf{v}_h^T)(\mathbf{x})d\mathbf{x} + \int_{\Omega} (\nabla_h \mathbf{u}_h - p_h I_2) : (\nabla \hat{\mathbf{v}} - \nabla_h \mathbf{v}_h)(\mathbf{x})d\mathbf{x}.
\]

(68)

Just like in Proposition 3.2 (see (45) and (50)), we have

\[
\int_{\Omega} (\nabla_h \mathbf{u}_h - p_h I_2) : (\nabla \hat{\mathbf{v}} - \nabla_h \mathbf{v}_h)(\mathbf{x})d\mathbf{x} = \frac{1}{2} \sum_{i \in [1, l]} \sum_{s \subset T_i} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2)_s] \cdot (\hat{\mathbf{v}} - \mathbf{v}_i^T)(\sigma)d\sigma + \frac{1}{2} \sum_{k \in [1, K]} \sum_{s \subset P_k} \int_s [(\nabla_h \mathbf{u}_h - p_h I_2)_s] \cdot (\hat{\mathbf{v}} - \mathbf{v}_k^P)(\sigma)d\sigma.
\]

(69)

From (68) and (69), we obtain (66). \( \square \)

4. A Computable Error Bound

4.1. Preliminaries

In this subsection, we will present some Poincaré-type inequalities which are useful to obtain a computable error bound.

Lemma 4.1. Let \( \omega \) be an open bounded set which is star-shaped with respect to one of its points. Let \( \mathbf{u} \in (H^1(\omega))^2 \) and let \( \mathbf{u}_x \) be the mean-value of \( \mathbf{u} \) over \( \omega \). Then,

\[
\exists C(\omega), \ s.t. \ |\mathbf{u} - \mathbf{u}_x|_{L^2(\omega)} \leq C(\omega) \text{diam}(\omega) \|\nabla \mathbf{u}\|_{L^2(\omega)}.
\]

(70)

Remark 4.2. In what follows, we shall use Lemma 4.1 on primal and dual cells. Since primal cells are (usually) convex, it is useful to note that when \( \omega \) is convex, a universal constant \( C(\omega) \) is given by \( \frac{1}{4} \) (see [26]). On the other hand, a dual cell \( P_k \) may be non-convex, but the way it is constructed implies that it is star-shaped with respect to the vertex \( S_k \) it is associated to, and that it is a union of triangles having \( S_k \) as a vertex. Bounds for \( C(\omega) \) in
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Theorem 4.4. Let $\|u\|_{L^2(E)}^2 \leq \frac{1}{\rho} \left( (2 + \varepsilon^{-2}) \|u\|_{L^2(T)}^2 + \varepsilon^2 \sigma^2 \|
abla u\|_{L^2(T)}^2 \right)$. (71)

4.2. A computable bound for the velocity error

The main result of this Section is the following:

Theorem 4.4. Let $\|\nabla_h e\|_{L^2(\Omega)}$ be defined by (37). Let us define the data oscillations by

\[
\text{osc}(f, T, \Omega) = \left( \sum_{i \in [1,I]} (C(T_i) h_i^T)^2 \|f - f_i^T\|_{L^2(T_i)}^2 \right)^{1/2},
\]

\[
\text{osc}(f, P, \Omega) = \left( \sum_{k \in [1,K]} (C(P_k) h_k^P)^2 \|f - f_k^P\|_{L^2(P_k)}^2 \right)^{1/2}
\]

where $f_i^T$ and $f_k^P$ are defined by (31), $h_i^T := \text{diam}(T_i)$, $h_k^P := \text{diam}(P_k)$ and the constants $C(T_i)$ and $C(P_k)$ are those involved in Lemma 4.1 and Remark 4.2.

We define the local and global error estimators related to the primal mesh:

\[
(\eta_i^T)^2 = \inf_{\mu > 0} \left[ \chi_i(\mu) \sum_{s \in T_i} C_s(\mu) \left( \|\nabla_h u_{h_k} - p h_i I_2 \|_{L^2(s)}^2 \right) \right]
\]

\[
(\eta_i^p)^2 = \sum_{k \in [1,K]} (\eta_k^p)^2
\]

(74)

where the functions $\chi_i$ and $C_s$ are defined by (88) and (87) below.

We define the local and global error estimators related to the dual mesh:

\[
(\eta_k^T)^2 = \inf_{\mu > 0} \left[ \chi_k(\mu) \sum_{s \in P_k} C_s(\mu) \left( \|\nabla_h u_{h_k} - p h_i I_2 \|_{L^2(s)}^2 \right) \right]
\]

\[
(\eta_k^p)^2 = \sum_{k \in [1,K]} (\eta_k^p)^2
\]

(76)
where the function $\chi_k$ is defined by (92) below.

We define the local and global error estimator related to the boundary:

$$
(\zeta_k^P)^2 = \inf_{\mu} \left[ \lambda_k(\mu) \sum_{\alpha=1}^{2} C_{\alpha}(\mu) \|(\nabla g - \nabla u_h)\tau_{J_{\alpha}(k)}\|_{L^2(b_{J_{\alpha}(k)})}^2 \right],
$$

where the functions $C_{\alpha}$ and $\lambda_k$ and the boundary segment $b_{J_{\alpha}(k)}$ are defined in Proposition 4.10 below.

Let us define the indicator related to the penalization:

$$
\zeta_e = \|\nabla_h e\|_{L^2(\Omega)}.
$$

We have

$$
\|\nabla_h e\|_{L^2(\Omega)} \leq \eta = \eta_h + \eta_e.
$$

where

$$
\eta_h = \frac{1}{2} \left( \text{osc}(f,T,\Omega) + \text{osc}(f,P,\Omega) + \eta^T + \eta^P \right) + \left( \frac{1}{2} + \frac{\sqrt{2}}{\beta} \right) \left( \eta^T + \eta^P + 2\zeta^P \right),
$$

$$
\eta_e = \frac{2}{\beta} \zeta_e.
$$

The quantities $\eta$, $\eta_h$, and $\eta_e$ are respectively called the total estimator, the discretization estimator, and the penalization estimator for the velocity.

Proof. The result is obtained using the sum (41), expressions (43) and (51), the bounds (85)–(86), (89)–(90), (93)–(94), (97) and (101) below and the relations (39).

In order to find bounds for the expressions (43) of $i_1$ and (51) of $i_2$, the values of $(\Phi^T_i, \Phi^P_i)$ and $(\Psi^T_i, \Psi^P_i)$ in these expressions have to be specified, since they were up to now arbitrary, except for the boundary midpoint values of $(\Phi^T_i, \Phi^P_i)$ that were chosen so that (42) holds. This is performed in the definition that follows.

Definition 4.5. Since $\hat{\Phi}$, $\hat{\Psi}$ are not necessarily more regular than $(H^1(\Omega))^2$, we choose as an interpolation their $L^2$ projection on the primal and dual cells:

$$
\Phi^T_i = \frac{1}{|T_i|} \int_{T_i} \hat{\Phi}(x)dx, \ \forall i \in [1,I]; \ \Phi^P_i = \frac{1}{|P_k|} \int_{P_k} \hat{\Phi}(x)dx, \ \forall k \in [1,K],
$$

$$
\Psi^T_i = \frac{1}{|T_i|} \int_{T_i} \hat{\Psi}(x)dx, \ \forall i \in [1,I]; \ \Psi^P_i = \frac{1}{|P_k|} \int_{P_k} \hat{\Psi}(x)dx, \ \forall k \in [1,K].
$$

In order to complete the definition of $(\Phi^T_i, \Phi^P_i)$, for any $i \in \Gamma$, the boundary value of $\Phi^T_i$ is given so that (42) holds. Note that it is not necessary to define the value of $\Psi^T_i$ for all $i \in \Gamma$.

Proposition 4.6. Let $h^T_i := \text{diam}(T_i)$, $h^P_k := \text{diam}(P_k)$ and let definitions (72) and (73) hold. Then it holds that

$$
\sum_{i \in [1,I]} \int_{T_i} f \cdot (\hat{\Phi} - \Phi^T_i)(x)dx \leq \text{osc}(f,T,\Omega)\|\nabla \hat{\Phi}\|_{L^2(\Omega)},
$$

$$
\sum_{k \in [1,K]} \int_{P_k} f \cdot (\hat{\Phi} - \Phi^P_k)(x)dx \leq \text{osc}(f,P,\Omega)\|\nabla \hat{\Phi}\|_{L^2(\Omega)}.
$$
Figure 4. For each cell $T_i$ and each vertex $S_k$ of $T_i$, $T_i \cap P_k$ is split in two triangles $t_{ik,1}$ and $t_{ik,2}$.

Proof. Since $\Phi^T_i$ was chosen as the mean-value of $\hat{\Phi}$ over $T_i$ (see (83)), we have
\[
\int_{T_i} f \cdot (\hat{\Phi} - \Phi^T_i)(x) dx = \int_{T_i} (f - f^T_i) \cdot (\hat{\Phi} - \Phi^T_i)(x) dx.
\]
Applying the Cauchy-Schwarz inequality, Lemma 4.1 to $\hat{\Phi}$ over $T_i$ and the discrete Cauchy-Schwarz inequality, we obtain (85). Similarly, we also obtain (86). 

Propositions 4.7 and 4.9 below are proved just like Propositions 5.9 and 5.10 in [25].

Proposition 4.7. For any primal cell $T_i$ and any dual $P_k$ such that $T_i \cap P_k \neq \emptyset$, let $s = [G_iS_k]$ and $t_{ik,1}$ and $t_{ik,2}$ be the triangles defined in Fig. 4 such that $t_{ik,1} \cup t_{ik,2} = T_i \cap P_k$. Let $\rho_{ik,\alpha}$ be the distance from $s$ to the vertex of $t_{ik,\alpha}$ opposite to $s$ and $\sigma_{ik,\alpha}$ be the length of the longest among the two other edges of $t_{ik,\alpha}$. $C(T_i)$ is the constant that appears in (70). For any strictly positive $\theta$, let us define
\[
C_s(\theta) = \left( \frac{1 + \sqrt{1 + \frac{\sigma_{ik,1}^2}{\mu}}}{1 + \sqrt{1 + \frac{\sigma_{ik,2}^2}{\mu}}} \right) \left( \frac{1 + \sqrt{1 + \frac{\sigma_{ik,2}^2}{\mu}}}{1 + \sqrt{1 + \frac{\sigma_{ik,1}^2}{\mu}} \rho_{ik,2} + \left( 1 + \sqrt{1 + \frac{\sigma_{ik,2}^2}{\mu}} \right) \rho_{ik,1} \right),
\]
(87)
\[
\chi_i(\mu) = (C(T_i)h_i^T)^2 + \mu.
\]
(88)

Let $\eta^T$ and $\eta^T_0$ be defined by (74)–(75). With these definitions, it holds that:
\[
\left| \sum_{i \in [1,I]} \sum_{s \in T_i} \int_{s} [(\nabla_h u_h - p_h I_2) n_s \cdot (\hat{\Phi} - \Phi^T_i)(\sigma) \sigma] \right| \leq \eta^T \| \nabla \hat{\Phi} \|_{L^2(\Omega)},
\]
(89)
\[
\left| \sum_{i \in [1,I]} \sum_{s \in T_i} \int_{s} [\nabla_h u_h \tau_s \cdot (\hat{\Psi} - \Psi^T_i)(\sigma) \sigma] \right| \leq \eta^T_0 \| \nabla \hat{\Psi} \|_{L^2(\Omega)}.
\]
(90)
Remark 4.8. The minimization in (74) is performed numerically when we effectively compute the estimators. Although we have no formal proof that the function to be minimized is unimodal, we used in the numerical tests of section 6 the so-called ”golden section search” to perform this minimization (with no claim that this is the most efficient method), starting with the interval $[10^{-2}h_T^2, 10^2h_T^2]$. Moreover, we may already get an idea of the behaviour of $\eta_i^T$ by choosing $\mu = h_T^2$ to evaluate it. By definition of $\sigma_{ik,\alpha}$, this length is lower than the diameter of $T_i$, which implies

$$C_s \left( (h_T^2)^2 \right) \leq \frac{(1 + \sqrt{2})^2}{2(\rho_{ik,1} + \rho_{ik,2})}. \tag{91}$$

Under the hypothesis that the ratios $\rho_{ik,\alpha}$ are all bounded by below by the same constant which does not depend on the mesh, we obtain the following bound

$$\eta_i^T \leq Kh_i^T \sum_{s \in T_i} \|((\nabla_h u_h - p_h I_2) n_s)\|_{L^2(s)},$$

where the constant $K$ does not depend on the mesh. The same remark holds for $\eta_i^T$.

Proposition 4.9. Let us set the same notations as in Prop. 4.7. Let $C_s$ be defined by (87). Let $C(P_k)$ be the constant involved in (70). Let us define for any strictly positive $\mu$,

$$\chi_k(\mu) = (C(P_k)h_k^P)^2 + \mu. \tag{92}$$

Let $\eta^P$ and $\eta^P$ be defined by (76)–(77). With these definitions, it holds that:

$$\left| \sum_{k \in [1, K]} \sum_{s \in P_k} \int_s [((\nabla_h u_h - p_h I_2) n_s) \cdot (\hat{\Phi} - \Phi_k^P)(\sigma) d\sigma \right| \leq \eta^P \|\nabla \hat{\Phi}\|_{L^2(\Omega)}, \tag{93}$$

$$\left| \sum_{k \in [1, K]} \sum_{s \in P_k} \int_s [\nabla_h u_h \tau_s] n_s \cdot (\hat{\Psi} - \Psi_k^P)(\sigma) d\sigma \right| \leq \eta^P \|\nabla \hat{\Psi}\|_{L^2(\Omega)}. \tag{94}$$

The next proposition bounds the boundary terms in the expression (51) of $i_2$.

Proposition 4.10. For any $k \in \Gamma$, let us denote by $D_{j_1(k)}$ and $D_{j_2(k)}$ the two diamond cells whose boundary intersect $\Gamma$ and which have $S_k$ as a vertex. Let $q_{j_1(k)} = P_k \cap D_{j_1(k)}$, $q_{j_2(k)} = P_k \cap D_{j_2(k)}$ and the segment $b_{j_0(k)}$ be the intersection between $\partial q_{j_0(k)}$ and $\Gamma$; see Fig. 5. Let $\rho_{j_0(k)}$ be the distance from $b_{j_0(k)}$ to the vertex of $q_{j_0(k)}$ opposite to $b_{j_0(k)}$ and $\sigma_{j_0(k)}$ be the length of the longest among the two other edges of $q_{j_0(k)}$. Let $C(P_k)$ be the constant that appears in (70). For any strictly positive $\mu$, let us define

$$C_\alpha(\mu) = \frac{2 + \frac{\sigma_{j_0(k)}^2}{\mu + \sqrt{\mu^2 + \mu \sigma_{j_0(k)}^2}}}{\rho_{j_0(k)}},$$

$$\lambda_\alpha(\mu) = (C(P_k)h_k^P)^2 + \mu. \tag{95} \tag{96}$$

Let $\zeta^P$ be defined by (78). With these definitions, it holds that:

$$\sum_{k \in T} \int_{\partial P_k \cap \Gamma} \left| (\nabla g - \nabla_h u_h) \tau_k \cdot (\hat{\Psi} - \Psi_k^P)(\sigma) \right| d\sigma \leq \zeta^P \|\nabla \hat{\Psi}\|_{L^2(\Omega)}. \tag{97}$$
Proof. By application of the Cauchy-Schwarz inequality on each edge $b_{j\alpha}(k)$ and the weighted discrete Cauchy-Schwarz inequality, we obtain for any set of strictly positive discrete real-valued numbers $C^P_{\alpha}$

$$\int_{P_{\alpha}(k)} \left| \langle \nabla g - \nabla h \rangle \cdot (\hat{\psi} - \Psi^P_k) \rangle \right| \, ds = \sum_{\alpha=1}^{2} \int_{b_{j\alpha}(k)} \left| \langle \nabla g - \nabla h \rangle \cdot (\hat{\psi} - \Psi^P_k) \rangle \right| \, ds$$

$$\leq \left[ \sum_{\alpha=1}^{2} C^P_{\alpha} \| \nabla g - \nabla h \| L^2 \right]^{1/2} \left[ \sum_{\alpha=1}^{2} \frac{1}{C^P_{\alpha}} \| \hat{\psi} - \Psi^P_k \| L^2 \right]^{1/2}.$$ \hspace{1cm} (98)

Now, for each segment $b_{j\alpha}(k)$, we can apply the trace inequality (71) on each triangle $q_{j\alpha}(k)$, for all $\alpha \in \{1, 2\}$ and for all strictly positive $\varepsilon_{j\alpha}(k)$. With $C_{1,j\alpha}(k) = \frac{2+\varepsilon_{j\alpha}(k)^2}{P_{j\alpha}(k)}$ and $C_{2,j\alpha}(k) = \frac{\varepsilon_{j\alpha}(k)^2}{P_{j\alpha}(k)}$, we obtain

$$\| \hat{\psi} - \Psi^P_k \| L^2 \leq C_{1,j\alpha}(k)\| \hat{\psi} - \Psi^P_k \| L^2 \leq C_{2,j\alpha}(k)\| \hat{\psi} - \Psi^P_k \| L^2 \leq C_{2,j\alpha}(k)\| \hat{\psi} - \Psi^P_k \| L^2.$$ \hspace{1cm} (99)

Let $\mu > 0$ be arbitrary. For $b_{j\alpha}(k)$ for $\alpha \in \{1, 2\}$, let us choose $\varepsilon_{j\alpha}(k)$ so that

$$\varepsilon_{j\alpha}(k)^2 = \frac{\mu + \sqrt{\mu^2 + \mu^2 \sigma_{j\alpha}(k)^2}}{\sigma_{j\alpha}(k)^2} \iff C_{2,j\alpha}(k) = \mu C_{1,j\alpha}(k)$$

and $C^P_{\alpha} = C_{1,j\alpha}(k)$. It holds that:

$$\sum_{\alpha=1}^{2} \frac{1}{C^P_{\alpha}} \| \hat{\psi} - \Psi^P_k \| L^2 \leq \sum_{\alpha=1}^{2} \left( \| \hat{\psi} - \Psi^P_k \| L^2 + \mu \| \nabla \hat{\psi} \| L^2 \right)$$

$$= \| \hat{\psi} - \Psi^P_k \| L^2 + \mu \| \nabla \hat{\psi} \| L^2 \leq \left[ (C(P_k h_c^2) + \mu \| \nabla \hat{\psi} \| L^2 \right].$$ \hspace{1cm} (100)

Plugging (100) into (98) and applying the discrete Cauchy-Schwarz inequality leads to (97). \hspace{1cm} $\square$
The final term to be evaluated in the right-hand side of (41) is related to the penalization. The proposition that follows is a simple consequence of the Cauchy-Schwarz inequality.

**Proposition 4.11.** Let definition (79) of \( \zeta \) hold. Then, we have

\[
\left| \varepsilon \int_{\Omega} \mathbb{P}_h^{2}(\Delta^P_h p)_{h} q(x) dx \right| \leq \zeta \| q \|_{L^2(\Omega)}.
\] (101)

This ends the various proofs of bounds that were necessary to prove Theorem 4.4.

4.3. A computable bound for the pressure error

**Proposition 4.12.** The following estimate holds:

\[
\| e_p \|_{L^2(\Omega)} \leq \frac{1}{2\beta} \left( 2 \| \nabla_h e \|_{L^2(\Omega)} + \text{osc}(f, T, \Omega) + \text{osc}(f, P, \Omega) + \eta^T + \eta^P \right).
\] (102)

**Proof.** Since \( e_p \in L^2_0(\Omega) \), there exists \( \tilde{v} \in (H^1_0(\Omega))^2 \), such that

\[
\| e_p \|_{L^2(\Omega)} \leq \frac{1}{\beta} \int_{\Omega} e_p \nabla \cdot \tilde{v}(x) dx.
\] (103)

With this given \( \tilde{v} \), we associate \( v = (v^I_t, v^P_t) \in (\mathbb{R}^2)^{I+JT} \times (\mathbb{R}^2)^K \) such that

\[
v^I_t = \frac{1}{|T|} \int_{T} \tilde{v}(x) dx \quad \forall i \in [1, I], \quad v^P_k = \frac{1}{|P_k|} \int_{P_k} \tilde{v}(x) dx \quad \forall k \in [1, K]
\] (104)

and the boundary values of \( v_i, \ i \in \Gamma \) are chosen so that \( \tilde{v}_j = 0 \) for all \( j \in \Gamma \). We may then apply (66) and follow calculations like those involved in propositions 4.6, 4.7 and 4.9. We finally obtain

\[
\left| \int_{\Omega} e_p \nabla \cdot \tilde{v}(x) dx \right| \leq \frac{1}{2} \left( 2 \| \nabla_h e \|_{L^2(\Omega)} + \text{osc}(f, T, \Omega) + \text{osc}(f, P, \Omega) + \eta^T + \eta^P \right) \| \nabla \tilde{v} \|_{L^2(\Omega)}.
\] (105)

Taking (105) into (103) proves (102).

\[ \square \]

5. Efficiency of the estimators

We state the main result of this Section:

**Theorem 5.1.** For any primal cell \( T_i \), let \( h^T_i := \text{diam}(T_i) \) and \( f^T_i \) be the mean-value of \( f \) over \( T_i \). Let \( \eta^T_i \) (resp. \( \eta^P_i \)) be defined in (74) (resp. (75)). For any dual cell \( P_k \), let \( h^P_k := \text{diam}(P_k) \) and \( f^P_k \) be the mean-value of \( f \) over \( P_k \). Let \( \eta^P_k \) (resp. \( \eta^P_k \)) be defined in (76) (resp. (77)). For any boundary dual cell \( P_k \), let \( \xi^P_k \) be defined in (78). Let Hypothesis 5.5 below hold. Then, there exists a constant \( C \) independent of the mesh such that

\[
(\eta^T_i)^2 \leq C \left( \| \nabla_h u_h - \nabla \tilde{u} \|_{L^2(T_i)}^2 + \| p_h - \tilde{p} \|_{L^2(T_i)}^2 \right) + C(h^T_i)^2 \| f - f^T_i \|_{L^2(T_i)}^2,
\] (106)

\[
(\eta^P_i)^2 \leq C \| \nabla_h u_h - \nabla \tilde{u} \|_{L^2(T_i)}^2,
\] (107)

\[
(\eta^P_k)^2 \leq C \left( \| \nabla_h u_h - \nabla \tilde{u} \|_{L^2(P_k)}^2 + \| p_h - \tilde{p} \|_{L^2(P_k)}^2 \right) + C(h^P_k)^2 \| f - f^P_k \|_{L^2(P_k)}^2,
\] (108)

\[
(\eta^P_k)^2 \leq C \| \nabla_h u_h - \nabla \tilde{u} \|_{L^2(P_k)}^2.
\] (109)

Moreover, in the case \( g = 0 \), there exists a constant \( C \) independent of the mesh such that

\[
(\xi^P_k)^2 \leq C \| \nabla_h u_h - \nabla \tilde{u} \|_{L^2(P_k)}^2.
\] (110)
The proof of this theorem is postponed after some technical properties.

In order to prove local efficiency of the estimators, we shall follow the bubble technique as presented in [31]. Since the estimator \( \eta_T \) involves jumps of \( \nabla u_h - p_h I_2 \) through the common edge \( s = [G_1S_k] \) of two neighboring diamond-cells, we shall use functions with a support included in the set \( T_i \cap P_k = \cup_{\alpha = 1,2} t_{ik,\alpha} \), where triangles \( t_{ik,\alpha} \) with \( \alpha = 1 \) or 2 are depicted in Fig. 4. Since we consider a fixed \( s \) in what follows, we simplify the notations to \( s \) and \( T \).

We define the following bubble functions

\[
\begin{align*}
  b_t &= \begin{cases} 
    27\lambda_{t,1}\lambda_{t,2}\lambda_{t,3} \text{ on } t, \\
    0 \text{ elsewhere.}
  \end{cases} \\
  b_s &= \begin{cases} 
    4\lambda_{s,1}\lambda_{s,2} \alpha \in \{1, 2\} \\
    0 \text{ elsewhere.}
  \end{cases}
\end{align*}
\]

(111)

(112)

It holds that \( \omega_t = \text{supp}(b_t) \subset t \) and \( \omega_s := \text{supp}(b_s) = T \cap P_k = t_1 \cup t_2 \). The following propositions are given for example, in [31].

**Proposition 5.3.** It holds that

\[
0 \leq b_t \leq 1, \quad 0 \leq b_s \leq 1, \quad \int_s b_s(\sigma)d\sigma = \frac{2}{3}|s|.
\]

(113)

(114)

**Proposition 5.4.** There exists a constant \( C > 0 \) only depending on the minimal angle in the couple \( (t_1, t_2) \) such that, for \( t = t_1 \) or \( t = t_2 \) and \( h_t = \text{diam}(t) \)

\[
\begin{align*}
  \frac{1}{C} h_t^2 &\leq \int_t b_t(x)dx = \frac{9}{20}|t| \leq Ch_t^2, \\
  \frac{1}{C} |s|^2 &\leq \int_t b_s(x)dx = \frac{1}{3}|t| \leq C|s|^2, \\
  \|\nabla b_t\|_{L^2(t)} &\leq Ch_t^{-1}\|b_t\|_{L^2(t)}, \\
  \|\nabla b_s\|_{L^2(t)} &\leq C|s|^{-1}\|b_s\|_{L^2(t)}.
\end{align*}
\]

(115)

(116)

(117)

(118)

In order to prove the local efficiency of the error estimator we shall make the following hypothesis:

**Hypothesis 5.5.** We suppose that the triangulation of \( \Omega \) composed of all the triangles \( t_{ik,\alpha} \) is regular in the sense that the minimum angles in those triangles are bounded by below independently of the mesh.

From this hypothesis, we derive the following propositions.

**Proposition 5.6.** For any primal cell \( T_i \) and any dual cell \( P_k \) such that \( T_i \cap P_k \neq \emptyset \), let \( s = [G_1S_k] \) and \( t_{ik,1} \) and \( t_{ik,2} \) be the triangles in Fig. 4 such that \( t_{ik,1} \cup t_{ik,2} = T_i \cap P_k \). Let \( h^2_T = \text{diam}(T_i) \), \( h^2_k = \text{diam}(P_k) \) and \( S_{ik} = |T_i \cap P_k| \). Let Hypothesis 5.5 hold. Then, there exists a constant \( C \) independent of the mesh such that

\[
(h^2_T)^2 S_{ik}^{-1} \leq C \text{ and } (h^2_k)^2 S_{ik}^{-1} \leq C.
\]

(119)

Proof. We will only prove the first inequality, since the second one can be treated in the same way.

Let \( \alpha_0 > 0 \) be the lower bound of all the angles of all the triangles \( t_{ik,\alpha} \).
For any $i \in [1,I]$, let $V_i$ be the number of vertices $S_{kl}$, with $\ell \in [1,V_i]$ of the primal cell $T_i$; see Fig. 6 for the notations. First, we note that
\[
V_i \leq V := \frac{2\pi}{2\alpha_0}, \text{ for all } i \in [1,I]. \tag{119}
\]
Let $M_{kl,\ell+1}$ be the midpoint of segment $[S_{kl},S_{kl+1}]$, with the convention that $k_{V_i+1} = k_1$. Then
\[
S_{ikl} = |T_i \cap P_{kl}| = |G_iS_{kl}M_{kl,\ell+1}| + |G_iS_{kl}M_{kl-1,\ell}|. \tag{120}
\]
Now, let us estimate the area of triangle $G_iS_{kl}M_{kl,\ell+1}$. Following Hypothesis 5.5, all the angles of triangle $G_iS_{kl}M_{kl,\ell+1}$ are greater than $\alpha_0$. Let $h_{G_i}$ be the maximum distance from point $G_i$ to the boundary of $T_i$, i.e.,
\[
h_{G_i} := \max\{|G_iS_{kl}|, \ell \in [1,V_i]\}. \tag{121}
\]
We have
\[
|G_iS_{kl}M_{kl,\ell+1}| = \frac{1}{2} \sin(S_{kl}G_iM_{kl,\ell+1})|G_iS_{kl}||G_iM_{kl,\ell+1}|. \tag{122}
\]
By a calculation on triangles $G_iS_{kl}M_{kl,\ell+1}$ and $G_iS_{kl+1}M_{kl,\ell+1}$, it holds that
\[
|G_iM_{kl,\ell+1}| \geq |G_iS_{kl}| \sin \alpha_0 \quad \text{and} \quad |G_iS_{kl+1}| \geq |G_iM_{kl,\ell+1}| \sin \alpha_0. \tag{123}
\]
From (123), we have the recurrence formula
\[
|G_iS_{kl}| \geq (\sin \alpha_0)^2|G_iS_{kl+1}|. \tag{124}
\]
Starting from the vertex $S_k$ which reaches the max in definition (121), the shortest way to go to a given vertex $S_{k_l}$ contains at most $V_i/2$ neighboring vertices for which we may apply (124), and we obtain:

$$|G_i S_k| \geq (\sin \alpha_0)^{V_i} h_{G_i}. \tag{125}$$

Then, from (123), we get that

$$|G_i S_k||G_i M_{k_l,\ell+1}| \geq (\sin \alpha_0)^{2V_i-1} h_{G_i}^2. \tag{126}$$

Combining (119), (122) and (126), and noting that from the definition of $h_{G_i}$, then $h_{G_i} \geq \frac{h_i^T}{2}$, we obtain

$$|G_i S_k M_{k_l,\ell+1}| \geq (\sin \alpha_0)^{2V_i+2} \frac{(h_i^T)^2}{8}. \tag{127}$$

In the same way, we have

$$|G_i S_k M_{k_l-1,\ell}| \geq (\sin \alpha_0)^{2V_i+2} \frac{(h_i^T)^2}{8}. \tag{128}$$

Using (120), (127)–(128), we obtain:

$$S_{ik_l} \geq (\sin \alpha_0)^{2V_i+2} \frac{(h_i^T)^2}{4}. \tag{129}$$

Thus, the inequality is proved with $C = 4 (\sin \alpha_0)^{-\frac{2h_i}{\alpha_0}}^{-2}$. \qed

**Proposition 5.7.** Under Hypothesis 5.5, the positive constants $C(T_i)$ and $C(P_k)$ are bounded independently of the mesh, and the constant $C$ in Prop. 5.4 is bounded by above and by below independently of the mesh.

**Proof.** The constants $C(T_i)$, $C(P_k)$ coming from (70) were bounded explicitly in [32]. From these expressions, it is easily seen that they are bounded if Hyp. 5.5 holds. Moreover, it is proved in [31] that $C$ in Prop. 5.4 depends only on the regularity of the triangles $t_{ik,\alpha}$. \qed

Now, we shall prove the efficiency of the estimators (Proof of Theorem 5.1).

**Proof.** Let us consider an element $T_i$ of the primal mesh and a diamond edge $s$ in $T_i$. Let us recall that by definition, such an edge $s$ does not belong to $\Gamma$. Let us consider the function $w_s = [\nabla \hat{u} - I^2 \hat{p}]b_s$, where $b_s$ is defined by (112). This function belongs to $(H_0^1(\Omega))^2$ and we may thus apply (36), which, taking into account the support of $w_s$, reduces to

$$\int_{\omega_s} (\nabla \hat{u} - I^2 \hat{p}) : \nabla w_s(x)dx = \int_{\omega_s} f \cdot w_s(x)dx. \tag{130}$$

Moreover, $u_h|_{D_j}$ belongs to $(P^1(D_j))^2$, $p_h$ is a constant in each $D_j$ and $w_s$ vanishes on $\Gamma$. Thus, it holds that:

$$\int_{\Omega} (\nabla u_h - I^2 p_h) : \nabla w_s(x)dx = \sum_j \int_{\partial D_j} (\nabla u_h - I^2 p_h) : \nabla w_s(x)dx$$

$$= \sum_j \int_{\partial D_j} ((\nabla u_h - I^2 p_h) n_{\partial D_j}) \cdot w_s(\sigma)d\sigma$$

$$= \sum_{i} \sum_{s' \subset T_i} \int_{s'} ((\nabla u_h - I^2 p_h) n_{s'}) \cdot w_s(\sigma)d\sigma.$$
But since $w_s$ vanishes on all the other edges $s' \neq s$, taking into account the definition of $w_s$ and the property of $b_s$ in (114), it holds that

$$
\int_{\Omega} (\nabla h u_h - I_2 p_h) : \nabla w_s(x) dx = \int_s [ (\nabla h u_h - I_2 p_h) n_s ] : w_s(\sigma) d\sigma \\
= [ (\nabla h u_h - I_2 p_h) n_s ] || w_s(\sigma) d\sigma \\
= 2 \int_s | (\nabla h u_h - I_2 p_h) n_s |^2 \\
= 2 \int_s | (\nabla h u_h - I_2 p_h) n_s |^2. 
$$

(131)

Defining $M := \| (\nabla h u_h - I_2 p_h) n_s \|_{L^2(s)}$, and taking into account (130), we have

$$
M^2 = 2 \int_{\Omega} (\nabla h u_h - I_2 p_h) : \nabla w_s(x) dx = 2 \int_{\omega_s} (\nabla h u_h - I_2 p_h) : \nabla w_s(x) dx \\
= 2 \left[ \int_{\omega_s} (\nabla h u_h - \nabla \hat{u}) : \nabla w_s(x) dx - \int_{\omega_s} (p_h - \hat{p}) \nabla \cdot w_s(x) dx + \int_{\omega_s} f \cdot w_s(x) dx \right]. 
$$

(132)

Using the Cauchy-Schwarz inequality leads to

$$
M^2 \leq 2 \left[ \| \nabla u_h - \nabla \hat{u} \|_{L^2(\omega_s)} + \sqrt{2} \| p_h - \hat{p} \|_{L^2(\omega_s)} \right] \| \nabla w_s \|_{L^2(\omega_s)} + 2 \| f \|_{L^2(\omega_s)} \| w_s \|_{L^2(\omega_s)}. 
$$

(133)

Let us now bound $\| \nabla w_s \|_{L^2(\omega_s)}$ and $\| w_s \|_{L^2(\omega_s)}$. Thanks to (118), it holds that

$$
\| \nabla w_s \|_{L^2(\omega_s)} = \| (\nabla u_s - I_2 p_h) n_s \|_{L^2(\omega_s)} \leq \| (\nabla u_s - I_2 p_h) n_s \| \sqrt{2} \| p_h - \hat{p} \|_{L^2(\omega_s)}, \\
\| w_s \|_{L^2(\omega_s)} = \| (\nabla u_s - I_2 p_h) n_s \|_s \leq \| (\nabla u_s - I_2 p_h) n_s \|_{L^2(\omega_s)}. 
$$

(134) \hspace{1cm} (135)

So there remains to find a bound for $\| b_s \|_{L^2(\omega_s)}$. In order to do this, we first infer from (113) that $b_s^2 \leq b_s$. This implies, using (116)

$$
\| b_s \|_{L^2(\omega_s)} = \left[ \int_{\omega_s} b_s^2(x) dx \right]^{1/2} \leq \int_{\omega_s} b_s^2(x) dx \leq C \| w_s \|_{L^2(\omega_s)}. 
$$

(136)

Taking into account that $\| (\nabla h u_h - I_2 p_h) n_s \| = |s|^{-1/2} M$ and considering (133) to (136), we obtain

$$
M \leq C \left[ |s|^{-1/2} (\| \nabla h u_h - \nabla \hat{u} \|_{L^2(\omega_s)} + \| p_h - \hat{p} \|_{L^2(\omega_s)}) + |s|^{1/2} \| f \|_{L^2(\omega_s)} \right]. 
$$

(137)

One usually expresses $\| f \|_{L^2(\omega_s)}$ as a function of $\| \nabla h u_h - \nabla \hat{u} \|_{L^2(\omega_s)} + \| p_h - \hat{p} \|_{L^2(\omega_s)}$ and of higher order terms. Let $t = t_1$ or $t_2$, and let us denote by $t$ the mean value of $f$ over $t$. Then,

$$
\| f \|_{L^2(t)} \leq \| f - t \|_{L^2(t)} + \| t \|_{L^2(t)}. 
$$

(138)

Then, consider $w_t = f \cdot b_t$, where $b_t$ is defined by (111). The function $w_t$ belongs to $(H^2_0)^2$. Thus, taking into account the support of $b_t$, Eq. (36) reduces to

$$
\int_{\Omega} (\nabla h u_h - I_2 p_h) : \nabla w_t(x) dx = \int_{\Omega} f \cdot w_t(x) dx. 
$$

(139)
Moreover, since $\nabla_h u_h - I_2 p_h$ is a constant over each $t$, and since $w_t$ vanishes on the boundary of $t$, it holds that
\[
\int_t (\nabla_h u_h - I_2 p_h) : \nabla w_t(x) dx = 0.
\] (140)

Since $f_t$ is a constant over $t$, there holds, thanks to (115), (139) and (140),
\[
\|f_t\|_{L^2(t)}^2 = |t| (f_t)^2 = C (f_t)^2 \int_t b_t(x) dx = C \int_t f_t \cdot w_t(x) dx
\]
\[
= C \left[ \int_t f \cdot w_t(x) dx + \int_t (f_t - f) \cdot w_t(x) dx \right]
\]
\[
= C \left[ \int_t (\nabla \hat{u} - \nabla_h u_h) : \nabla w_t(x) dx - \int_t (\hat{p} - p_h) \nabla \cdot w_t(x) dx + \int_t (f_t - f) \cdot w_t(x) dx \right]
\]
\[
\leq C \left( \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(t)} + \sqrt{2} \|\hat{p} - p_h\|_{L^2(t)} \right) \|\nabla w_t\|_{L^2(t)} + C \|f_t - f\|_{L^2(t)} \|w_t\|_{L^2(t)} ,
\] (141)

with $C = 20/9$ in the above expressions. Let us now bound $\|w_t\|_{L^2(t)}$ and $\|\nabla w_t\|_{L^2(t)}$. With (117), it holds that
\[
\|\nabla w_t\|_{L^2(t)} = |f_t| \|\nabla b_t\|_{L^2(t)} \leq |f_t| CH_t^{-1} \|b_t\|_{L^2(t)} ,
\] (142)
\[
\|w_t\|_{L^2(t)} = |f_t| \|b_t\|_{L^2(t)} .
\] (143)

The remaining term that has to be bounded is $\|b_t\|_{L^2(t)}$. For this, we first infer from (113) that $b_t^2(x) \leq b_t(x)$ and then
\[
|f_t| \|b_t\|_{L^2(t)} \leq |f_t| \left( \int_t b_t(x) dx \right)^{1/2} \leq C |f_t| |t|^{1/2} = C \|f_t\|_{L^2(t)} ,
\] (144)
in which $C = \sqrt{9/20}$. Combining (141)–(142)–(143)–(144), we finally get
\[
\|f_t\|_{L^2(t)} \leq C \left( \|f_t - f\|_{L^2(t)} + h_t^{-1} \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(t)} + h_t^{-1} \|\hat{p} - p_h\|_{L^2(t)} \right).
\]

Since $s$ is an edge of $t$, it holds that $|s| \leq h_t$; applying (138), we obtain
\[
\|f\|_{L^2(t)} \leq C \left( \|f_t - f\|_{L^2(t)} + |s|^{-1} \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(t)} + |s|^{-1} \|\hat{p} - p_h\|_{L^2(t)} \right).
\]

Thus, taking into account that $\omega_s = t_1 \cup t_2$, it holds that
\[
\|f\|_{L^2(\omega_s)} \leq \|f\|_{L^2(t_1)} + \|f\|_{L^2(t_2)}
\]
\[
\leq C \left( \|f_{t_1} - f\|_{L^2(t_1)} + \|f_{t_2} - f\|_{L^2(t_2)} \right) + C |s|^{-1} \left( \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(t_1)} \right)
\]
\[
+ \left( \|\hat{p} - p_h\|_{L^2(t_1)} + \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(t_2)} + \|\hat{p} - p_h\|_{L^2(t_2)} \right) \right)
\]
\[
\leq C |s|^{-1} \left( \|\nabla \hat{u} - \nabla_h u_h\|_{L^2(\omega_s)} + \|\hat{p} - p_h\|_{L^2(\omega_s)} \right) + C \|f_{\omega_s} - f\|_{L^2(\omega_s)} .
\] (145)

In this sequence of inequalities, we have used the fact that $f_t$ minimizes $\|f - c\|_{L^2(t)}$ when $c$ runs over $\mathbb{R}^2$; in particular, $\|f_t - f\|_{L^2(t)} \leq \|f_{\omega_s} - f\|_{L^2(t)}$, where $f_{\omega_s}$ is the mean value of $f$ over $\omega_s$. Combining (137) and (145), we obtain
\[
M \leq C |s|^{1/2} \|f - f_{\omega_s}\|_{L^2(\omega_s)} + C |s|^{-1/2} \left( \|\nabla_h u_h - \nabla \hat{u}\|_{L^2(\omega_s)} + \|\hat{p} - p_h\|_{L^2(\omega_s)} \right).
\] (146)
By definition, the local estimator \((\eta^T)^2\) is lower than the value taken by the function in (74) in \(\mu = (h^T)^2\). In (74), we may bound \(C(T)\) by \(1/\pi\) since the primal cells have been supposed to be convex, and with (91) and (146), we obtain

\[
\left(\eta^T\right)^2 \leq C \left(h^T\right)^2 \sum_{s \in T} \frac{|s|^{-1}}{\rho_{ik,1} + \rho_{ik,2}} \left(\|\nabla_h u_h - \nabla\tilde{u}\|_{L^2(\omega_s)}^2 + \|p_h - \tilde{p}\|_{L^2(\omega_s)}^2\right)
+ C \left(h^T\right)^2 \sum_{s \in T} \frac{|s|}{\rho_{ik,1} + \rho_{ik,2}} \|f - f_{\omega_s}\|_{L^2(\omega_s)}^2. \tag{147}
\]

Using Prop. 5.6, and since by definition \(S_{ik} = \frac{1}{2} |s| (\rho_{ik,1} + \rho_{ik,2})\), we have that \(\left(h^T\right)^2 |s|^{-1}(\rho_{ik,1} + \rho_{ik,2})^{-1}\) is bounded by a constant that does not depend on the mesh under Hyp. 5.5. Moreover, under Hyp. 5.5, the ratio \(|s|\rho_{ik,1}^{-1}\) is also bounded by a constant that does not depend on the mesh. So (147) leads to (106).

As far as (107) is concerned, let us consider the function \(v_s = [\nabla_h u_h \tau_s]_{s} b_s\). There obviously holds

\[
\int_{\Omega} \nabla \hat{u} : \nabla \times v_s(x) \, dx = \int_{\omega_s} \nabla \hat{u} : \nabla \times v_s(x) \, dx = 0. \tag{148}
\]

Eq. (148) and the calculations that previously led to (131) may be used to yield

\[
\|\nabla_h u_h \tau_s\|_{L^2(\omega_s)}^2 = \frac{3}{2} \int_{\omega_s} \nabla_h u_h : \nabla \times v_s(x) \, dx
= \frac{3}{2} \int_{\omega_s} (\nabla_h u_h - \nabla\hat{u}) : \nabla \times v_s(x) \, dx
\leq \frac{3}{2} \|\nabla_h u_h - \nabla\hat{u}\|_{L^2(\omega_s)} \|\nabla v_s\|_{L^2(\omega_s)}. \tag{149}
\]

Just like (133) led to (137) and then to (106), inequality (149) leads to (107). The dual inequalities (108), (109) and (110) may be obtained in the same way. The proof of (110) is very similar to that of (107) and (109), but some definitions have to be changed because the segment \(b_{\omega_s(k)}\) in the definition (78) is a boundary segment, and is thus the edge of only one triangle \(t\); the function \(b_s\) is thus defined only in that triangle \(t\).

\section{Numerical results}

First, we study the influence of the parameter \(\varepsilon\) for a fixed mesh and then the influence of the mesh size for a fixed value of the penalty parameter. Secondly, we give an overall process to recursively adapt the value of the penalty parameter and adaptively refine the mesh.

\subsection{Influence of the penalty parameter}

In this subsection, we will work on the domain \(\Omega = [0; 1]^2\). A triangular mesh with rather uniform triangles is used. The exact solution \((\hat{u}, \hat{p})\) is regular with \(\hat{u} = (\partial_y \varphi, -\partial_x \varphi)\) given by

\[
\varphi(x, y) = 100x^2y^2(1-x)^2(1-y)^2 \quad \text{and} \quad \hat{p}(x, y) = 10(x^2 + y^2 - \frac{2}{3}). \tag{150}
\]

Fig. 7 presents the plots of the errors and the estimators when the penalty parameter \(\varepsilon\) goes from \(10^{-5}\) to \(10^2\). They include the actual errors in the \(H^1(\Omega)\) and \(L^2(\Omega)\) norms for the velocity, i.e. the error in the velocity gradient \(\|\nabla\tilde{u} - \nabla_h u_h\|_{L^2(\Omega)}\) and in the velocity \(\|\tilde{u} - u_h\|_{L^2(\Omega)}\), the total estimator, the discretization estimator and the penalization estimator which are given by Theorem 4.4 when we estimate the velocity error. The left
Figure 7. Actual errors in $H^1(\Omega)$ and $L^2(\Omega)$ norms, total estimator, discretization estimator and penalization estimator for the velocity. Left: $h = 10^{-1}$, right: $h = 3.125 \times 10^{-2}$.

(resp. right) figure corresponds to the mesh size $h = 10^{-1}$ (resp. $h = 3.125 \times 10^{-2}$). We see that for a given mesh, the ratio between the penalization estimator and the penalty parameter $\varepsilon$ asymptotically tends to a constant when $\varepsilon$ tends to 0, while the discretization estimator is nearly independent of $\varepsilon$ when $\varepsilon \leq 1$. Moreover, the actual errors decrease with $\varepsilon$ until a certain level. Then, the discretization error is the dominant error and decreasing $\varepsilon$ further does not have any influence on the overall error. It is noticeable that the influence of $\varepsilon$ is more important on the $L^2$ norm than on the $H^1$ semi-norm.

6.2. Influence of the mesh size

On the same square domain $\Omega$ and with the same exact solution as previously, we work with the following values: $\varepsilon = 10^2$, $\varepsilon = 1$ and $\varepsilon = 10^{-2}$. Since the solution is regular, only uniformly refined triangular meshes will be considered. Figure 8 presents the same curves as Fig. 7, but now as a function of $h$, varying from 0.25 to $1.6 \times 10^{-2}$. For the test $\varepsilon = 10^2$, when $h$ is large, $\varepsilon h^2$ is so large that the solution of the penalized model is very different from the solution of the original, non-penalized model. It is then not surprising that the exact error is strongly correlated to the penalization estimator rather than to the total estimator. Moreover, we observe what is called a "numerical locking": for large $h$, the error almost does not decrease when $h$ is decreased; it only starts decreasing when $h$ is small enough. However, the values $\varepsilon = 10^{-2}$ or even $\varepsilon = 1$ are small enough so that no such locking occurs. The penalization estimator is small with respect to the total estimator and we
observe that the actual $H^1$ semi-norm of the error, the total estimator and the discretization estimator decrease roughly like $h$ when $h$ decreases. As above, we notice that the $L^2$ norm of the error is more influenced than the $H^1$ semi-norm when decreasing $\varepsilon$ from 1 to $10^{-2}$.

### 6.3. Adaptive penalty parameter and mesh refinement

We propose the following computational process. We start with a given coarse mesh and an initial value of $\varepsilon$, and we fix some ratio $0 < \gamma \leq 1$.

Then, we compute the numerical solution, and we get $\eta_h$ and $\eta_e$. Then,

- If $\eta_e \geq \gamma \eta_h$, we adapt a new $\varepsilon$ by multiplying the previous value of $\varepsilon$ with the ratio $\frac{3\eta_h}{2\eta_e}$ and keep the same mesh for a new computation. This has the effect of maintaining the error due to the penalization below a certain ratio of the error due to the discretization.
- Otherwise, we adaptively refine the mesh based on the discretization estimator $\eta_h$. For this, on the given mesh, we compute local discretization estimators $\eta_{i,h}$ on each primal cell $T_i$, obtained by adding $\eta_i$, $\eta_i'$ (respectively defined by (74) and (75)) to the contribution of $T_i$ in the oscillation term (72), and by redistributing the sums that appear in the dual estimators (76), (77) and (73) and in the boundary estimators (78) to the cells $T_i$ that have an intersection with $P_k$. More specifically, in (73), the integral over $P_k$ is split into the sum (over all $T_i$ that intersect $P_k$) of the integrals over $T_i \cap P_k$, which are then redistributed to the corresponding $T_i$. Moreover, each local dual estimator (76) and (77) is a sum over diamond edges $s = [G_i S_k]$ of quantities which are thus redistributed to the corresponding $T_i$. Finally, each boundary estimator (78) is a sum of two quantities that can be related each to a given primal cell (see Fig. 5) to which the corresponding contribution in (78) is redistributed. Once all $\eta_{i,h}$ have been computed, we require the refinement of a given primal cell $T_i$ by a factor 4 in terms of area if $\eta_{i,h} \geq (\max_i \eta_{i,h})/2$. Specifically, we use the Triangle mesh generator [27] and its mesh refinement feature which remeshes the domain in a way that user-specified local area constraints are fulfilled. Note that simply dropping the dual and boundary estimators instead of redistributing them onto the primal cells as explained above had almost no influence on the general shape of the convergence curves in Figures 9 and 10 below.

The test we present to illustrate this strategy is also on the domain $\Omega = [0; 1]^2$, The exact solution $(\tilde{u}, \tilde{p})$ is regular with $\tilde{u} = (\varphi_y, -\varphi_x)$, and $\varphi$ is given by

$$\varphi(x, y) = x^2(1-x)^2y^2(1-y)^2 \text{ and } \tilde{p}(x, y) = 5(x^2 + y^2 - \frac{2}{3}).$$

For accuracy reasons, the ratio $\gamma$ may be chosen so that the penalization error is much lower than the discretization error like in the left part of Fig. 9 obtained with $\gamma = 1/500$. We chose this value of $\gamma$ because Fig. 7 showed that the actual $H^1$ error stops decreasing when the penalization estimator is around two orders of magnitude lower than the discretization estimator, while the $L^2$ error stops decreasing when this ratio is around three orders of magnitude. The first four computations on the coarsest mesh are used to tune the value of $\varepsilon$: starting with $\varepsilon = 10^5$, its value is decreased by the adaptive process described above to $2.5 \times 10^{-3}$. We observe (see the overlapping squares on the left part of the $H^1$ semi-norm error curve) that the first two decreasing steps of $\varepsilon$ have the effect of decreasing the actual error. The last decreasing step of $\varepsilon$ from $2.9 \times 10^{-2}$ to $2.5 \times 10^{-3}$ does not affect the total error and we are thus sure that the total error does not suffer from any error introduced by a too high value of $\varepsilon$. Then, the value of $\varepsilon$ does not change any more when we refine the mesh. Along this process, the total estimator is kept almost equal to the discretization estimator.

Moreover, we made a test with $\gamma = 1$ and we present the result on the right part of Fig. 9. This value of $\gamma$ is obviously too high because we observe the numerical locking discussed above.
The next test is inspired from [19]. Our test is in the domain $\Omega = [0, 1]^2$ and the exact couple solution $(\hat{u}, \hat{p})$ is singular with $\hat{u} = (\varphi_y, -\varphi_x)$, and $\varphi$ is

$$\varphi(x, y) = x^7(1 - x)^2y^2(1 - y)^2 \text{ and } \hat{p}(x, y) = \frac{x + y - 1}{10}.$$ 

We observe that the velocity $\hat{u}$ is in $[H^{1/4}(\Omega)]^2$ and there is a boundary singularity on the edge $x = 0$.

We compare the exact $H^1(\Omega)$ semi-norm of the velocity error with the total velocity error estimator for uniform and adaptive refinements. For uniform refinements, only $\varepsilon$ is adapted, while for adaptive refinements, both $\varepsilon$ and the meshes are adapted by the penalty - mesh refinement process described above. In both cases, we choose $\gamma = 1/500$. In Fig. 10, we start with $\varepsilon = 0.1$. The first computation of the adaptive process sets $\varepsilon$ to $1.8 \times 10^{-3}$, and then $\varepsilon$ remains unchanged in the rest of the adaptive process for both adaptive and
uniform refinements. The convergence curve corresponding to the uniform mesh refinement is observed to be asymptotically parallel to the $N^{-1/4}$ curve, while the convergence curve corresponding to the adaptive mesh refinement is parallel to the $N^{-1/3}$ curve. Moreover, the effectivity of the error estimators for both types of refinements is around 25.

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{figure.png}
\caption{Exact errors in the adaptive process using the discretization estimator and exact error.}
\end{figure}

We could have expected the plot of the exact error corresponding to an adaptive mesh refinement to be parallel to the optimal $N^{-1/2}$ straight line, but this is not the case, just like in [19]. We would like to determine whether this problem is caused by our discretization estimator or not. For this, in Fig. 11, we compare the exact errors corresponding to two different adaptive refinement processes: one is driven by our discretization estimator, while the other is driven by the exact error (which we may compute exactly in this test case). Clearly, the exact error is hardly affected by the applied adaptive process.

\begin{thebibliography}{99}
\end{thebibliography}


